

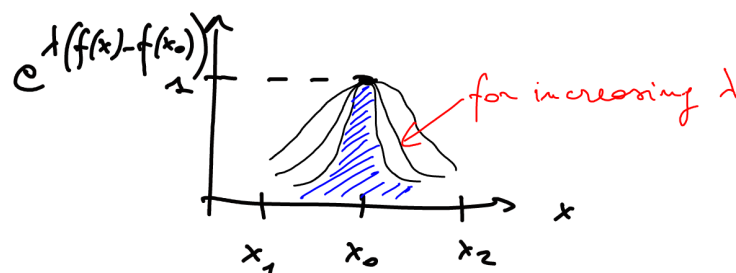
The Laplace method

In several situations one wish to evaluate complicated integrals which have a form

$$(21) \quad I(\lambda) = \int_{x_1}^{x_2} dx \, g(x) e^{\lambda f(x)} \quad \lambda \in \mathbb{R}$$

where g and f are continuous and differentiable functions. Although $I(\lambda)$ cannot be calculated for any arbitrary λ , it happens that it can be well approximated (under appropriate conditions) as $\lambda \rightarrow \infty$.

The core idea of Laplace's method is that the major contribution to the integral in (21) as $\lambda \rightarrow \infty$ comes from the neighborhood of the point in $[x_1, x_2]$ where $f(x)$ gets its maximum value, which we call x_0 .



There are essentially three cases:

- 1) x_0 is an interior maximum, $x_1 < x_0 < x_2$ and $f'(x_0) = 0$. We assume that $f''(x_0) < 0$ (actually if $f'(x_0) = f''(x_0) = f'''(x_0) = 0$ but $f^{(iv)}(x_0) < 0$ we can apply very similar ideas and the calculations are not more difficult). Also $g(x_0) \neq 0$.

and it is finite.

As we expect that the dominant contributions comes from the neighborhood of x_0 and $f(x) = f(x_0) + \frac{(x-x_0)^2}{2} f''(x_0) + \Theta(|x-x_0|^3)$ as $x \rightarrow x_0$, we obtain

$$I(\lambda) \approx \int_{x_1}^{x_2} g(x) e^{\lambda \left(f(x_0) + \frac{(x-x_0)^2}{2} f''(x_0) \right)} dx \underset{\text{at lead. order}}{\approx} g(x_0) e^{\lambda f(x_0)} \underbrace{\int_{x_1}^{x_2} e^{-\lambda \frac{|f''(x_0)|}{2} (x-x_0)^2} dx}_{\text{leading order}}$$

We change var. $s = (x-x_0) \sqrt{\frac{|f''(x_0)|}{2} \lambda}$ so

$$I(\lambda) = g(x_0) e^{\lambda f(x_0)} \underbrace{\int_{(x_1-x_0) \sqrt{\frac{|f''(x_0)|}{2} \lambda}}^{(x_2-x_0) \sqrt{\frac{|f''(x_0)|}{2} \lambda}} \frac{2}{\lambda |f''(x_0)|} e^{-s^2} ds}_{\text{as } \lambda \rightarrow \infty \rightarrow \int_{-\infty}^{+\infty} e^{-s^2} ds = \sqrt{\pi}}$$

(22)

$$I(\lambda) = g(x_0) e^{\lambda f(x_0)} \sqrt{\frac{2\pi}{\lambda |f''(x_0)|}}$$

as $\lambda \rightarrow +\infty$
(leading order)

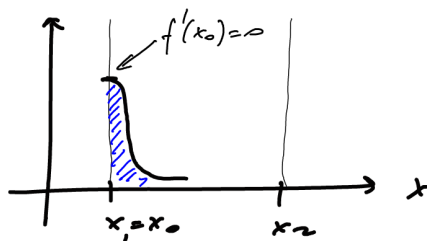
one should prove that this is the leading order (we did not).

Exercise: show that the next to leading order of $I(\lambda)$ in eq. (22) is given by

$$I(\lambda) = e^{\lambda f(x_0)} \sqrt{\frac{2\pi}{\lambda |f''(x_0)|}} \left(g(x_0) + \frac{c}{\lambda} \right) \quad \text{as } \lambda \rightarrow \infty$$

where c is a constant that depends on the derivatives of f up to 4th order (at $x=x_0$) and on $g(x_0)$ and $g'(x_0)$.

2) $x_0 = x_1$ or $x_0 = x_2$ (x_0 is a flat endpoint) and $f'(x_0) = 0$.



It is easy to show that the leading order formula is

$$(23) \quad I(\lambda) = g(x_0) e^{\lambda f(x_0)} \sqrt{\frac{\pi}{2\lambda |f''(x_0)|}}$$

Example

The modified Bessel function of second kind is a special function that occurs in many applications. It is given by

$$K_\nu(x) = \int_0^\infty e^{-x \cosh t} \cosh(\nu t) dt, \quad x > 0$$

We wish to estimate $K_\nu(x)$ as $x \rightarrow +\infty$ for fixed ν .

Since $\cosh'(t) = \sinh(t)$, $t > 0$, the max of $e^{-x \cosh t}$ as a function of t ($x > 0$) occurs at $t=0$, with zero derivative.

So we can apply eq. (23) ($\cosh t \approx 1 + \frac{t^2}{2} + \dots$)
 $f(t) = \cosh t$

$$K_\nu(x) = e^{-x} \sqrt{\frac{\pi}{2x}} \quad \text{as } x \rightarrow \infty$$

at leading order

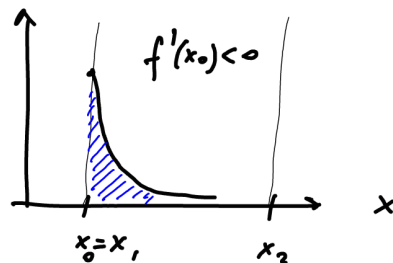
notice that it does not depend on ν !

Exercise: show that

$$K_\nu(x) = \sqrt{\frac{\pi}{2x}} e^{-x} \left(1 + \frac{c(\nu)}{x} + \dots \right)$$

and $c(\nu) = \frac{4\nu^2 - 1}{8}$.

3) $x_0 = x_1$ or $x_0 = x_2$ (x_0 is an endpoint) but $f'(x_0) \neq 0$
 Without loss of generality we take $x_0 = x_1$ and $f'(x_0) < 0$



In this case we get $f(x) = f(x_0) + (x - x_0)f'(x_0) + \dots$
 hence

$$\begin{aligned} I(\lambda) &\approx \int_{x_1=x_0}^{x_2} g(x) e^{\lambda [f(x_0) + (x-x_0)f'(x_0)]} dx = \\ &= g(x_0) e^{\lambda f(x_0)} \int_{x_0}^{x_2} e^{\lambda (x-x_0)f'(x_0)} dx = \end{aligned}$$

$$= g(x_0) e^{\lambda f(x_0)} \int_0^{x_2 - x_0} e^{\lambda f'(x_0) s} ds = \frac{g(x_0) e^{\lambda f(x_0)}}{\lambda f'(x_0)} \left[e^{\lambda s f'(x_0)} \right]_0^{x_2 - x_0}$$

\uparrow
 $s = x - x_0$

$$= \frac{g(x_0) e^{\lambda f(x_0)}}{\lambda f'(x_0)} \left(e^{\lambda (x_2 - x_0) f'(x_0)} - 1 \right)$$

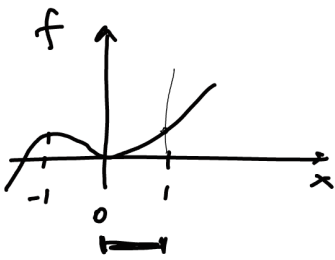
as $\lambda \rightarrow \infty$ and $f'(x_0) < 0$ we get

$$(24) \quad I(\lambda) \cong g(x_1) \frac{e^{\lambda f(x_1)}}{\lambda |f'(x_1)|} \quad \text{as } \lambda \rightarrow \infty$$

Example:

Obtain the leading order approximation of

$$I(\lambda) = \int_0^1 x^m e^{\lambda(3x^2 + 2x^3)} dx \quad \text{as } \lambda \rightarrow +\infty$$



The max is at $x=1$, $f(1)=5$, $f'(1)=12$
 $g(1)=1$. We can apply eq. (24)

$$I(\lambda) \simeq \frac{e^{5\lambda}}{12\lambda} \quad \text{at leading order}$$

(no dependence on m)

Exercise:

Calculate the leading order approx. when the integral is

$$\int_{-1}^0 \dots = ? \quad \int_{-2}^0 \dots = ?$$

Stirling formula

We want to estimate how fast $N!$ goes to infinity as $N \rightarrow \infty$. We will apply the Laplace's method to the gamma function:

$$(25) \quad \Gamma(\lambda) = \int_0^{\infty} x^{\lambda-1} e^{-x} dx \quad \lambda > 0$$

Exercise: Show that $\Gamma(\lambda+1) = \lambda \Gamma(\lambda)$ hence $\lambda! = \Gamma(\lambda+1)$ which generalize the factorial to complex numbers.

Let's consider $\Gamma(\lambda+1) = \int_0^{\infty} x^{\lambda} e^{-x} dx$. If we write this as $\int_0^{\infty} e^{-x} e^{\lambda \ln x} dx$ we cannot apply Laplace's method.

It is more beneficial to consider the max of the function $f(x) = -x + \lambda \ln x$ and set $g(x) = 1$. The max occurs at $x = \lambda$ which suggests a change of vars: $x = \lambda t$ (so the max is now fixed w.r.t. t). We get

$$\Gamma(\lambda+1) = \int_0^{\infty} x^{\lambda} e^{-x} dx = \lambda^{\lambda+1} \int_0^{\infty} t^{\lambda} e^{-\lambda t} dt$$

\uparrow
 $x = \lambda t$

$t^{\lambda} e^{-\lambda t} = e^{\lambda(\ln t - t)}$ $h(t) = \ln t - t$, $h'(1) = 0$, $h''(1) = -1$.
In this way we can apply eq. (22) with $g(t) = 1$ ($t_0 = 1$).

$$(26) \quad \Gamma(\lambda+1) = \lambda! \simeq \lambda^{\lambda+1} e^{-\lambda} \sqrt{\frac{2\pi}{\lambda}} \quad \text{as } \lambda \rightarrow \infty.$$

Exercise: Show that at leading order

$$\int_0^{\infty} e^{-\lambda t} e^{-\frac{1}{t}} dt \cong \frac{\sqrt{\pi} e^{-2\sqrt{\lambda}}}{\lambda^{3/4}} \quad \text{as } \lambda \rightarrow \infty$$

Example:

Let's consider the class of integrals:

$$I_m(x) = \int_0^{\infty} t^m e^{-\frac{t^2}{2} - \frac{x}{t}} dt \quad x > 0.$$

We calculate $I_m(x)$ for large x and fixed m .

$\frac{t^2}{2} + \frac{x}{t}$ has a movable min at $\frac{d}{dt} \left(\frac{t^2}{2} + \frac{x}{t} \right) = 0$, $\bar{t} = x^{1/3}$
So we introduce the new variable τ :

$$t = x^{1/3} \tau$$

So

$$I_m(x) = x^{\frac{m+1}{3}} \int_0^{\infty} \tau^m e^{-x^{2/3} \left(\frac{\tau^2}{2} + \frac{1}{\tau} \right)} d\tau$$

The min of the exponent occurs at $\tau = 1$ (interior)
so we can apply eq. (22):

$$I_m(x) = x^{\frac{m+1}{3}} e^{-\frac{3}{2} x^{2/3}} \sqrt{\frac{2\pi}{x^{2/3} \cdot 3}}$$

$$= x^{m/3} e^{-\frac{3}{2} x^{2/3}} \sqrt{\frac{2\pi}{3}}$$

as $x \rightarrow \infty$.
leading order, fixed m

Exercises:

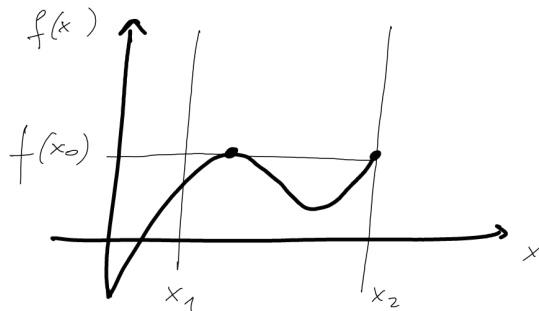
$$\int_{-2}^0 e^t e^{\lambda(3t^2+2t^3)} dt \approx e^{\lambda-1} \sqrt{\frac{\pi}{3\lambda}}$$

$$\int_0^1 e^t e^{\lambda(3t^2+2t^3)} dt \approx \frac{e^{5\lambda+1}}{12\lambda}$$

$$\int_0^1 \sqrt{1+t} e^{\lambda(2t-t^2)} dt \approx e^{\lambda} \sqrt{\frac{\pi}{2\lambda}}$$

$$\int_{-1}^2 e^{\lambda(t^3-1)} (1+t^2) dt = \frac{5e^{6\lambda}}{11\lambda}$$

Question: Let's assume that $f(x)$ behaves like in the fig:



Where does the leading order come from?