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## Chapter 1

# Electromagnetic waves

## 1.1 Electromagnetic theory

### 1.1.1 Macroscopic Maxwell equations

**Theorem 1.1.1** (Gauss' law). *Let  $\omega$  be a solid and  $\rho(\mathbf{r}) : \Omega \rightarrow \mathbb{R}$  its charge density distribution. Then,*

$$\oint_{\partial\Omega} \langle \mathbf{E}, d\mathbf{s} \rangle_I = \frac{1}{\epsilon} \int_{\Omega} \rho(\mathbf{r}) dv, \quad \operatorname{div} \mathbf{E} = \frac{\rho(\mathbf{r})}{\epsilon}. \quad (1.1)$$

**Definition 1.1.1.**

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P}. \quad (1.2)$$

**Theorem 1.1.2** (Gauss' law in media). *Let  $\Omega$  be a dipolar solid and  $\rho_f(\mathbf{r}) : \Omega \rightarrow \mathbb{R}$  its free charge density distribution. Then,*

$$\oint_{\partial\Omega} \langle \mathbf{D}, d\mathbf{s} \rangle_I = \int_{\Omega} \rho_f(\mathbf{r}) dv, \quad \operatorname{div} \mathbf{D} = \rho_f(\mathbf{r}). \quad (1.3)$$

**Theorem 1.1.3** (Gauss' law for magnetism). *Let  $\Omega$  be a solid. Then,*

$$\oint_{\partial\Omega} \langle \mathbf{B}, d\mathbf{s} \rangle_I = 0, \quad \operatorname{div} \mathbf{B} = 0. \quad (1.4)$$

**Theorem 1.1.4** (Faraday's law for magnetism). *Let  $\Sigma$  be a surface formed by a circuit  $\Gamma = \partial\Sigma$ . Then,*

$$\oint_{\Gamma} \langle \mathbf{E}, d\mathbf{r} \rangle_I = -\frac{d}{dt} \int_{\Sigma} \langle \mathbf{B}, d\mathbf{s} \rangle_I, \quad \operatorname{curl} \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}. \quad (1.5)$$

**Theorem 1.1.5** (Ampère's law). *Let  $\Sigma$  be a surface formed by a circuit  $\Gamma = \partial\Sigma$ . Then,*

$$\oint_{\Gamma} \langle \mathbf{B}, d\mathbf{r} \rangle_I = \mu_0 \int_{\Sigma} \langle \mathbf{J}, d\mathbf{s} \rangle_I, \quad \operatorname{curl} \mathbf{B} = \mu_0 \mathbf{J}. \quad (1.6)$$

**Theorem 1.1.6** (Ampère's law for variables fields). *Let  $\Sigma$  be a surface formed by a circuit  $\Gamma = \partial\Sigma$ . Then,*

$$\operatorname{curl} \mathbf{B} = \mu_0 \left[ \mathbf{J} + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right] \quad (1.7)$$

**Definition 1.1.2.**

$$\mathbf{H} = \frac{1}{\mu_0} \mathbf{B} - \mathbf{M}. \quad (1.8)$$

**Theorem 1.1.7** (Ampère's law in media). *Let  $\Sigma$  be a surface formed by a circuit  $\Gamma = \partial\Sigma$  and  $\mathbf{J}_f : \Sigma \rightarrow \mathbb{R}^3$  its free current distribution. Then,*

$$\operatorname{curl} \mathbf{H} = \mathbf{J}_f + \frac{\partial \mathbf{D}}{\partial t}. \quad (1.9)$$

**Proposition 1.1.8.** *If  $\mathbf{P}, \mathbf{M}$  and  $\mathbf{J}_f$  are linear, then the equations are linear and superposition principle holds.*

### 1.1.2 Boundary conditions

### 1.1.3 Lorentz force

### 1.1.4 Material response

**Proposition 1.1.9.** *Let  $\Omega$  be an homogeneous, isotropus and linear solid with  $\rho_f(\mathbf{r}), \sigma_f(\mathbf{r}) = 0$ . Then,*

$$\operatorname{div} \mathbf{E} = 0, \quad \operatorname{div} \mathbf{B} = 0, \quad \operatorname{curl} \mathbf{E} = \frac{\partial \mathbf{B}}{\partial t}, \quad \operatorname{curl} \mathbf{B} = \mu\epsilon \frac{\partial \mathbf{E}}{\partial t}. \quad (1.10)$$

### 1.1.5 Energetic relations

**Proposition 1.1.10.** *Magnetic energy density distribution*

$$\eta_m = \frac{\mu}{2} \|\mathbf{H}\|^2 \approx \frac{\mu_0}{2} \|\mathbf{H}\|^2. \quad (1.11)$$

**Proposition 1.1.11.** *Electric energy density distribution for non-absorbing media*

$$\eta_e = \frac{1}{2} \langle \mathbf{E}, \epsilon_0 \mathbf{E} + \mathbf{P} \rangle_I = \frac{\epsilon}{2} \|\mathbf{E}\|^2. \quad (1.12)$$

**Definition 1.1.3.** Poynting vector

$$\mathbf{S} := \mathbf{E} \times \mathbf{H} = \frac{1}{\mu} \mathbf{E} \times \mathbf{B}. \quad (1.13)$$

## 1.2 Wave equation in dielectric media

**Proposition 1.2.1.** *Let  $\Omega$  be an homogeneous, isotropus and linear solid with  $\rho_f(\mathbf{r}) = 0$ . Then,*

$$\operatorname{div} \mathbf{E} = 0, \quad \operatorname{div} \mathbf{B} = 0, \quad \operatorname{curl} \mathbf{E} = \frac{\partial \mathbf{B}}{\partial t}, \quad \operatorname{curl} \mathbf{B} = \mu \sigma \mathbf{E} + \mu \epsilon \frac{\partial \mathbf{E}}{\partial t}. \quad (1.14)$$

**Proposition 1.2.2.** *Let  $\Omega$  be an homogeneous, isotropus and linear solid with  $\rho_f(\mathbf{r}) = 0$ . Then,*

$$\nabla^2 \mathbf{E} = \mu \sigma \frac{\partial \mathbf{E}}{\partial t} + \mu \epsilon \frac{\partial^2 \mathbf{E}}{\partial t^2}, \quad (1.15)$$

and if  $\sigma(\mathbf{r}) = 0$ , then

$$\nabla^2 \mathbf{E} = \mu \epsilon \frac{\partial^2 \mathbf{E}}{\partial t^2}. \quad (1.16)$$

**Proposition 1.2.3.** *Let*



## Chapter 2

### Extra

## 2.1 Physical constants

**Definition 2.1.1.** Physical constants

$$q_e = 1,602 \times 10^{-19} \text{ C},$$

$$h = 6,62 \times 10^{-34} \text{ kg},$$

$$\epsilon_0 = 8,8542 \times 10^{-12} \text{ C}^2 \text{ N m}^2,$$

$$\mu_0 = 4\pi \times 10^{-7} \text{ N}^2 \text{ s}^2 \text{ C}^{-2}$$