

# 1 Harmonic oscillator

**Definition 1.1.** Let  $\mathcal{H}$  be a Hilbert space in a harmonic potential

$$V(\hat{x}) = \frac{\omega^2}{2} \hat{x}^2, \quad \omega^2 = \frac{k}{m}. \quad (1)$$

We define the *creation* and *annihilation operators* as

$$\hat{a}^\dagger := \frac{\alpha}{\sqrt{2}} \left( \hat{x} - \frac{i}{m\omega} \hat{p} \right), \quad (2)$$

$$\hat{a} := \frac{\alpha}{\sqrt{2}} \left( \hat{x} + \frac{i}{m\omega} \hat{p} \right), \quad (3)$$

$$\alpha := \sqrt{\frac{m\omega}{\hbar}}. \quad (4)$$

**Proposition 1.1.** Let  $\mathcal{H}$  be a Hilbert space in a harmonic potential. Then,

$$\langle x | \hat{a}^\dagger = \frac{\alpha}{\sqrt{2}} \left( x - \frac{1}{\alpha^2} \frac{d}{dx} \right), \quad (5)$$

$$\langle x | \hat{a} = \frac{\alpha}{\sqrt{2}} \left( x + \frac{1}{\alpha^2} \frac{d}{dx} \right), \quad (6)$$

$$\alpha = \frac{m\omega}{\hbar}. \quad (7)$$

**Proposition 1.2.** Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. Then,

$$\hat{x} = \frac{1}{\sqrt{2}\alpha} (\hat{a}^\dagger + \hat{a}), \quad \hat{p} = i\hbar \frac{\alpha}{\sqrt{2}} (\hat{a}^\dagger - \hat{a}). \quad (8)$$

**Proposition 1.3.** Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. Then,

1.  $\hat{a}, \hat{a}^\dagger$  are not hermitian.

2.  $[\hat{a}, \hat{a}^\dagger] = \hat{I}$ .

3.  $\hat{H} = \hbar\omega \left( \hat{a}^\dagger \hat{a} + \frac{1}{2} \right)$ .

**Definition 1.2.** Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. We define the *number operator* as

$$\hat{N} := \hat{a}^\dagger \hat{a}. \quad (9)$$

**Proposition 1.4.** Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. Then,

1.  $\hat{H}$  is hermitian.

2.  $[\hat{N}, \hat{a}] = -\hat{a}, [\hat{N}, \hat{a}^\dagger] = \hat{a}^\dagger,$

3.  $\hat{H} = \hbar\omega \left( \hat{N} + \frac{1}{2} \hat{I} \right)$ .

**Proposition 1.5.** Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. Then,  $\hat{H}$  and  $\hat{N}$  have a common basis of eigenvectors, which is countable, and

$$\hat{a}^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle, \quad \hat{a} |n\rangle = \sqrt{n} |n-1\rangle, \quad (10)$$

$$\hat{N} |n\rangle = n |n\rangle, \quad \hat{H} |n\rangle = \hbar\omega \left( n + \frac{1}{2} \right) |n\rangle, \quad (11)$$

$$n \in \mathbb{N}. \quad (12)$$

**Corollary 1.6.** Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. Then,

$$|n\rangle = \frac{1}{\sqrt{n!}} (\hat{a}^\dagger)^n |0\rangle. \quad (13)$$

**Proposition 1.7.** Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. Then, the eigenstates form a non-degenerate basis.

**Definition 1.3** (Fock states). Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. We define the *Fock states* as the states that determine the basis  $(|n\rangle)$  and have a well-defined number of excitations.

**Definition 1.4.** Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. We call the fundamental Fock state the *vacuum*.

**Proposition 1.8.** Let  $\mathcal{H}$  be a Hilbert space with a harmonic potential. Then,  $\hat{a}, \hat{a}^\dagger$  and  $\hat{N}$  have the following matrix representation in the basis  $(|n\rangle)$ .

$$[\hat{N}]_B = \begin{pmatrix} 0 & 0 & 0 & \cdots \\ 0 & 1 & 0 & \cdots \\ 0 & 0 & 2 & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}, \quad (14)$$

$$[\hat{a}]_B = \begin{pmatrix} 0 & \sqrt{1} & 0 & \cdots \\ 0 & 0 & \sqrt{2} & \cdots \\ 0 & 0 & 0 & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}, \quad (15)$$

$$[\hat{a}^\dagger]_B = \begin{pmatrix} 0 & 0 & 0 & \cdots \\ \sqrt{1} & 0 & 0 & \cdots \\ 0 & \sqrt{2} & 0 & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}, \quad (16)$$

or in coefficient representation,

$$[\hat{N}]_{ij} = (i-1)\delta_{ij}, \quad (17)$$

$$[\hat{a}]_{ij} = \sqrt{j-1}\delta_{i,j-1}, \quad (18)$$

$$[\hat{a}^\dagger]_{ij} = \sqrt{i-1}\delta_{i-1,j}. \quad (19)$$