

Analysis of Selected Models for Inelastic Electron Scattering in the KATRIN Gaseous Tritium Source

Master's thesis of

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Introduction

To our current knowledge, neutrinos are at the same time the most elusive and most abundant massive particles in the Universe. Their detection is challenging and pushes experiments at the edge of technical frontiers. Yet, their understanding might shed light on long-standing open questions of modern physics: What is dark matter? Do Majorana particles exist? Why does matter predominate over antimatter? What happened in the earliest stages of our Universe just after the Big Bang? And in all these regards, the yet unknown mass of neutrinos is a key physics parameter.

The KArlsruhe TRItium Neutrino (KATRIN) experiment aims to measure the effective mass of the electron antineutrino with an unprecedented sensitivity of 200 meV (90 % C.L.) based on tritium- β decay. In order to provide this outstanding sensitivity KATRIN features i. a. a gaseous tritium source. Its special characteristics must be well controlled and understood. This thesis focuses on selected effects stemming from electrons scattering in said gaseous source. The effects were included in a high level analysis, meaning their impact on KATRIN's neutrino mass sensitivity was studied where possible.

Outline

[Do not forget to update this.] Chapter 1 is a brief introduction to neutrino physics ToDo with special emphasis on the neutrino mass.

Chapter 2 focuses on the setup of the KATRIN experiment.

Chapter 3 introduces a mathematical model of a KATRIN neutrino mass measurement that can be used in parameter inference.

Chapter 4 integrates the mathematical model into a statistical framework for a high level analysis, especially for neutrino mass inference.

Chapter 5 investigates a refinement of the mathematical model with regard to electrons scattering off gas molecules within KATRIN's gaseous source. In particulars, the dependence of the inelastic scattering cross section on the energy of the incident electrons is studied.

Chapter 6 focuses on a preliminary model for the energy loss of electrons scattering inelastically off deuterium molecules. The model was established by a dedicated subgroup of the KATRIN collaboration based on KARIN data from October 2018. Its relation to KATRIN's sensitivity on the neutrino mass is investigated.

Chapter 7 summarizes the results, draws conclusions and offers an outlook.

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1. Neutrino Physics

This chapter is an introduction to neutrino physics. The primary aim is to give an experimentally-rooted definition of a neutrino. Therefore, in section 1.1 selected experimental milestones are outlined that led to today’s description of a neutrino within the established Standard Model of Particle Physics (SM). In section 1.2 follows an outline of the SM and how it relates to the neutrino. Special attention is paid to the neutrino mass: First, in section 1.3, an extension of the SM that allows for neutrino masses is summarized. Second, in section 1.4, the phenomenon of neutrino flavor oscillations is introduced. Corresponding experiments proved that neutrinos have mass. Third, in section 1.5, experiments for an absolute neutrino mass measurement are presented because as such they relate particularly to the KATRIN experiment.

1.1. Neutrinos until the 1960s

Albeit the neutrino as a hypothetical new particle was not postulated until 1930, its rich scientific history might be seen as already heralded during the preceding 35 years.

In 1895, Becquerel reported results on experiments with phosphorescent substances, especially uranium salts, on photographic plates [Bec95]. These experiments mark the discovery of radioactivity and triggered manifold subsequent investigations.

In 1899, Rutherford published a classification of radioactive decays into α and β types according to their penetration strength [Rut99].

In 1900, Becquerel determined the mass-charge ratio of β -decay particles [Bec00] and identified them as the electron previously described by Thomson [Tho97].

In 1914, Chadwick measured a continuous electron energy spectrum in the β decay of lead-214 and bismuth-214 [Cha14].

In 1927, Ellis and Wooster conducted a calorimetric measurement of the β -decay energy of radium and demonstrated that the continuity of the β spectrum was intrinsic to the decay as opposed to be caused by secondary effects [EW27].

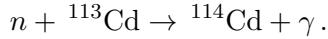
In 1930, a β decay was thought of as a two-body decay ${}^zA \rightarrow {}^{z+1}B + e^-$. Assuming conservation of energy and momentum, in a two-body-decay the momenta of the daughter particles B and e^- are solely determined by their masses and the “energy content”, as Bohr put it, of the parent particle A. According to Bohr, there was no reason to believe that different nuclei of the same element A should have a different energy content in a β decay. Hence, the continuous nature of the β spectrum could not be explained [Boh32]. As a possible solution Pauli suggested the β decay to be a three-body decay and postulated an electrically neutral particle that carries part of the decay energy [Pau30].

In 1934, Fermi developed a quantitative theory of β decay that could describe the preceding experimental results [Fer34]. It comprises a four-fermion contact interaction respectively, a

three-body-decay model. It was the first description of the so-called “weak interaction”. Furthermore, Fermi coined the term “neutrino” for the particle postulated by Pauli. Fermi’s theory inspired the idea to use the so-called “inverse β decay” or “neutrino capture” to detect neutrinos, which in today’s nomenclature is written as



In 1956, Cowan and Reines published results of a corresponding experiment. It was conducted using the high neutrino flux of the nuclear reactor of the Savannah River Plant. The neutrinos originating in the reactor passed a tank of water and cadmium chloride triggering the above process. The emerging neutron was captured by the cadmium which emitted a photon in a 3 MeV to 11 MeV range



The emerging positron annihilated with an electron which produced two photons of 0.5 MeV each. A coincidence measurement of the corresponding photons enabled discriminating signal and background events. Based on their results Cowan and Reines reported the discovery of the free neutrino [Cow+56].

In the same year, 1956, Lee and Yang published an article on parity conservation. Parity conservation implies that a mirrored physical process behaves the same as its non-mirrored counterpart. Here, mirroring means a change of sign of the position vector in the applied physical laws. Lee and Yang pointed out that parity conservation might be violated in weak interactions and suggested several probing methods [LY56].

In 1957 Wu et al. conducted an experiment which employed one of the corresponding probing methods based on β decay. The parity operation respectively “the mirroring” corresponded a change of the magnetic field orientation in the experiment. The results showed that parity is violated in weak interactions [Wu+57].

In 1958, Goldhaber et al. measured the helicity H of the neutrino. Helicity is defined as $H = \hat{\sigma} \cdot \hat{p}$, where $\hat{\sigma}$ is the spin unit vector and \hat{p} is the momentum unit vector (here: of the neutrino). The experiment found $H = -1$ which corresponds to maximum parity violation. In other words, only left-handed neutrinos and right-handed antineutrinos participate in weak interactions [Gol+58].

In 1962, Danby et al. reported on a second type of neutrinos. A beam of pions generated at the Alternating Gradient Synchrotron in Brookhaven decayed according to $\pi^\pm \rightarrow \mu^\pm + \bar{\nu}$. The emerging neutrinos penetrated a 13.5-meter iron shield wall and their interactions were detected in a 10-t aluminum spark chamber. The observed interactions were path-like as opposed to shower-like, which implied the production of muons as opposed to electrons. This was marked as the discovery of the muon neutrino [Dan+62].

The attempts to uniformly describe the manifold discoveries in the field of particle physics in a combined theory converged over the course of the second half of the 20th century into what is known today as the Standard Model of Particle Physics.

1.2. Neutrinos in the Standard Model of Particle Physics

This section introduces the basic concepts of the Standard Model of Particle Physics (SM) in a condensed manner. It aims at giving a description of relevant particle properties in section 1.2.1 and relating them to neutrinos in section 1.2.2.

The SM is a gauge quantum field theory exhibiting the gauge symmetry $SU(3) \times SU(2) \times U(1)$. As such it can be formulated using the principle of least action and a Lagrangian density

\mathcal{L} depending on fields and their derivatives [Zee03]. Albeit it can not account for all known subatomic phenomena, within its known boundaries, the SM is a well-tested and established theory, which is evident by e.g. the extensive review of particle properties of the Particle Data Group [Tan+18].

1.2.1. General Particle Properties

The gap between fields and particles can be bridged as follows: If “[i]n region 1 in spacetime there exists a source that sends out a ‘disturbance in the field’, which is later absorbed by a sink in region 2 in spacetime[,] experimentalists choose to call this a particle” [Zee03]. Intrinsic particle properties can be derived from the relation of their associated fields to the Lagrangian density. E.g., a particle’s mass is encoded by the Yukawa coupling of its field to the higgs doublet through the higgs mechanism and spontaneous symmetry breaking [Hig64]. A further intrinsic property is a particle’s spin, that takes half-integer values in units of \hbar for fermions or integer values for bosons. A particle’s flavor is its eigenstate with respect to the weak interaction, which is described by the $SU(2) \times U(1)$ subgroup (Glashow-Weinberg-Salam model [Gla61; Wei67; Sal68]). According to Noether’s theorem, each symmetry conserves an associated charge [Noe18]. In the case of the $SU(2) \times U(1)$ symmetry, the associated charges are called isospin $\mathbf{T} = (T_1, T_2, T_3)^\top$ and hypercharge Y . A derivative of these charges is the electric charge $Q = T_3 + \frac{1}{2}Y$ [Sch19]. In that sense, each particle has an associated antiparticle that carries the opposite electric charge. As mentioned in the historical overview (section 1.1), a theory consistent with experiment must violate parity. Such theories are called chiral [Zee03]. The SM is a chiral theory and thus its fermion fields can be decomposed in left- and right handed components [Sch19].

Figure 1.1 depicts the particles of the SM along with their selected properties mass and electric charge. It also shows a further categorization among the fermions into quarks and leptons.

1.2.2. Neutrino Properties

With reference to the particle properties listed in the previous section 1.2.1, a neutrino can be described as follows: A neutrino carries a spin of $1/2\hbar$. Thus, it is a fermion. It is categorized as a lepton. It has an electric charge of 0. And there are only left-handed neutrinos and right-handed antineutrinos [Sch19].

The mass of a neutrino will be discussed separately within the following chapters.

Neutrino come in three flavors, typically denoted as ν_e , ν_μ and ν_τ . Some additional remarks about the neutrino flavors can be made: First, the historical overview (section 1.1) mentions the discovery of the electron and muon flavor, but it was not until 2001 that the tau neutrino was discovered by the DONUT collaboration [Kod+01]. Second, a precision measurement of the width of the Z^0 -boson resonance $\Gamma_Z = (2.68 \pm 0.15) \text{ GeV}$ at the Large Electron Positron Collider (LEP) in the 1990s yielded a number of active light neutrino flavors consistent with three. In this context, “light” refers to a neutrino mass smaller than half the mass $M_Z = (91.174 \pm 0.070) \text{ GeV}$ of the Z^0 -boson [Acc+98]. (It should be noted that this refers to active neutrino flavors. Sterile neutrinos as an extension of the SM are not ruled out [OW08].)

1.3. Mechanisms to generate Neutrino Masses

Section 1.4 lists experiments which proof that neutrinos have mass. However, in the SM as described in section 1.2 the neutrino masses are assumed to be zero. Nevertheless, an extension of the SM is possible. The corresponding mass terms are introduced in section 1.3.1. The neutrino-mass formalism entails neutrino flavor mixing as described in section 1.3.2.

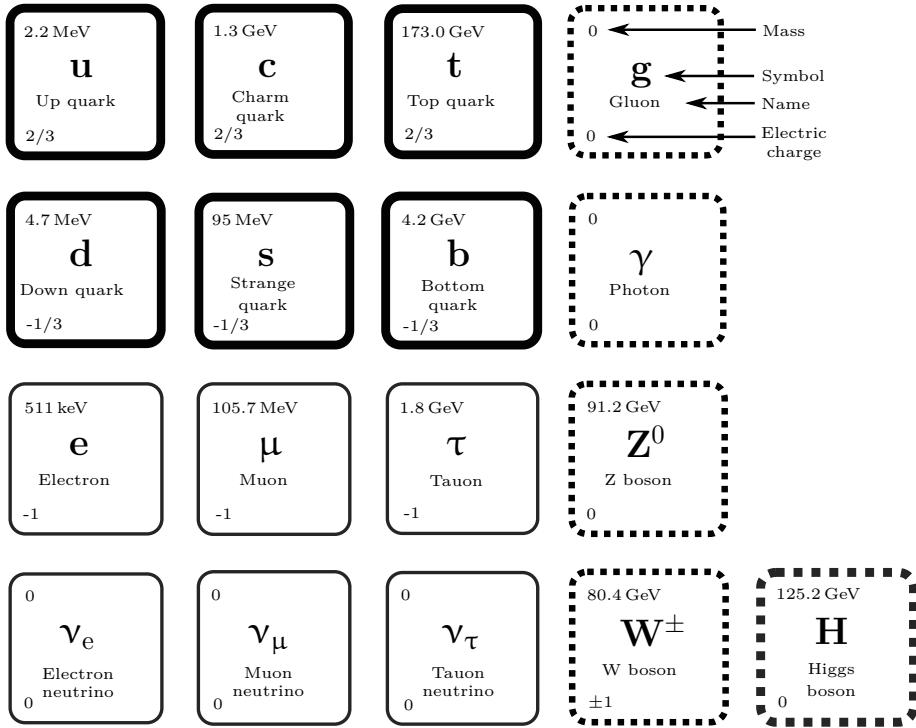


Figure 1.1.: The Standard Model of Particle Physics. The diagram illustrates possible categorizations of particles within the SM. The fermions are framed with continuous and the bosons with dotted lines. Among the fermions the quark sector is marked by a thick frame and the lepton sector by a thin one. The first three columns show the fermions; and the fourth and the fifth the bosons. While the bosons in the fourth column carry a spin of 1, the higgs boson in the fifth column marked with a thicker frame carries a spin of 0. Also shown are the particle masses in natural units and their electric charge in units of the absolute electron charge. All quantities along with uncertainties can be found in the Review of Particle Physics [Tan+18]. (Illustration adapted from [Sei19].)

1.3.1. Neutrino Mass Terms

For a theory to account for neutrino masses, its Lagrangian density must exhibit corresponding mass terms. According to [Zub11] the formalism can be summarized: The form of a mass term is given by the Dirac equation, which is produced by applying the principle of least action to a suitable Lagrangian density \mathcal{L} . The mass terms have to be quadratic in the fermion fields ψ and must leave the Lagrangian density hermitian. Furthermore, a field ψ must have a left- and right-handed component in order for the mass terms not to vanish. Two possible term forms are named after Dirac and Majorana. Whether one or a mixture of both forms corresponds to the neutrino's reality is an open question.

Dirac Masses

A Dirac mass term with mass m_D split in its chiral components (Weyl spinors) $\psi_{L,R}$ has the form [Zub11]

$$\mathcal{L}_D = -m_D \bar{\psi} \psi = -m_D (\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L) . \quad (1.1)$$

Applying this to neutrinos requires both a left- and a right-handed Dirac neutrino. Right-handed neutrinos have not yet been observed. If they exist, they do not interact weakly and hence are called sterile.

Majorana Masses

For Majorana mass terms the CP-conjugate ψ^C of a fermion spinor ψ is used. It should be noted that if ψ is left-handed, ψ^C is right-handed and vice versa. Then, a Majorana field ϕ can be defined and a corresponding mass term \mathcal{L}_M with a mass m_M be constructed [Zub11]:

$$\phi = \psi + \psi^C \quad \mathcal{L}_M = -\frac{1}{2}m_M\bar{\phi}\phi. \quad (1.2)$$

As $\phi^C = \phi$, the described Majorana particle is its own antiparticle, which due to charge conservation is only possible for neutral particles, such as a neutrino.

1.3.2. Neutrino Mixing

If neutrinos have mass, their mass eigenstates $|\nu_i\rangle$ ($i \in \{1, 2, 3\}$) of the free Hamiltonian need not be identical to their flavor eigenstates $|\nu_\alpha\rangle$ ($\alpha \in \{e, \mu, \tau\}$) of the weak interaction [Zub11]. In case they differ, there must be a basis-change matrix. Such a matrix was introduced by Maki, Nakagawa and Sakata in order to explain the so-called neutrino oscillations (see section 1.4) predicted by Pontecorvo [Pon58; MNS62]. Therefore, the matrix U for a basis change is called Pontecorvo-Maki-Nakagawa-Sakata matrix (PMNS matrix)

$$|\nu_\alpha\rangle = \sum_i U_{\alpha i} |\nu_i\rangle. \quad (1.3)$$

As a complex unitary matrix U can be expressed by six parameters. A possible choice are three angles $\theta_{12}, \theta_{23}, \theta_{13} \in [0, 2\pi)$, a phase $\delta \in [0, 2\pi)$ and two Majorana phases $\alpha, \beta \in [0, 2\pi)$:

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta_{23} & \sin \theta_{23} \\ 0 & -\sin \theta_{23} & \cos \theta_{23} \end{pmatrix} \begin{pmatrix} \cos \theta_{13} & 0 & \sin \theta_{13} e^{-i\delta} \\ 0 & 1 & 0 \\ -\sin \theta_{13} e^{i\delta} & 0 & \cos \theta_{13} \end{pmatrix} \begin{pmatrix} \cos \theta_{12} & \sin \theta_{12} & 0 \\ -\sin \theta_{12} & \cos \theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha} & 0 \\ 0 & 0 & e^{i\beta} \end{pmatrix}. \quad (1.4)$$

These parameters are called neutrino mixing parameters. It should be noted that δ is also called “CP-violating” phase. Here, P stands for parity conjugation as it was explained in section 1.1; and C for electric charge conjugation that follows the same idea with a sign change of the electric charge instead of the position vector. Why $\delta \neq 0$ implies CP -violation is shown in section 1.4.1.

One of the consequences of neutrino mixing, namely neutrino oscillations, will be explained in the following section 1.4.

1.4. Neutrino Oscillations

The term “neutrino oscillations” refers to the neutrino’s change of flavor after passing a certain propagation distance. In other words, neutrinos might be detected in another flavor than the one they originated in. Section 1.4.1 introduces a demonstrative formalism that aims at showing the link between oscillations and the masses of neutrinos. Neutrino oscillations also depend on the neutrino mixing parameters introduced in section 1.3.2. The accessibility of these mixing parameters and neutrino masses via neutrino oscillation experiments will be evaluated in section 1.4.4. Neutrino oscillation experiments are manifold. As an exemplary case study the so-called “solar neutrino problem” is discussed in section 1.4.3. Finally, the experimental results on neutrino oscillations will be summarized in section 1.4.4.

1.4.1. Relation to Neutrino Masses

According to [Zub11] a formula demonstrating neutrino oscillations can be derived: Using the PMNS matrix U from equation (1.4) the evolution of a neutrino's flavor eigenstate on a one-dimensional path starting at position $x = 0$ at time $t = 0$ with momentum p_i and energy E_i of its mass eigenstates $|\nu_i\rangle$ is

$$|\nu_\alpha(x, t)\rangle = \sum_i U_{\alpha i} e^{-i(E_i t - p_i x)} |\nu_i\rangle. \quad (1.5)$$

This leads to the transition amplitudes

$$A(\alpha \rightarrow \beta)(t) = \langle \nu_\beta | \nu_\alpha(x) \rangle = \sum_i U_{\beta i}^* U_{\alpha i} e^{-i(E_i t - p_i x)t} \quad (1.6a)$$

$$A(\bar{\alpha} \rightarrow \bar{\beta})(t) = \langle \bar{\nu}_\beta | \bar{\nu}_\alpha(x) \rangle = \sum_i U_{\beta i} U_{\alpha i}^* e^{-i(E_i t - p_i x)t}. \quad (1.6b)$$

It should be noted that if $U \neq U^*$, equation (1.6) implies CP -violation. In reference to section 1.3.1 it holds $U \neq U^* \Leftrightarrow \delta \neq 0$, justifying that δ is called CP -violating phase.

The following assumptions allow for a simple and demonstrative form of the transition probability:

- The neutrinos are relativistic:
 - Their momentum equals approximately their energy which is by far larger than their mass $p_i \approx E_i \gg m_i$. This also implies that the energy can be expanded in the mass-momentum-ratio m_i/p_i .
 - They travel the distance $x = L = ct$ at the speed of light c .
- All neutrino generations have approximately the same momentum $E \approx p \approx p_i$.
- The CP -violating phase vanishes: $\delta = 0$. (This assumption is not necessary, but simplifies the expression for the transition probability. See [Zub11] for $\delta \neq 0$.)

Then, the transition probability from one flavor α to another β in dependence of the neutrino masses and mixing parameters is

$$\begin{aligned} P(\alpha \rightarrow \beta)(L) &= |\langle \nu_\beta | \nu_\alpha(L) \rangle|^2 \\ &= \delta_{\alpha\beta} - 4 \sum_i \sum_{j>i} U_{\alpha i} U_{\alpha j} U_{\beta i} U_{\beta j} \sin^2 \left(\frac{(m_i^2 - m_j^2)}{4} \frac{L}{E} \right). \end{aligned} \quad (1.7)$$

Equation (1.7) shows oscillatory behavior if two conditions are fulfilled. First, the mass eigenvalue of at least two mass eigenstates must differ. Second, the product of the corresponding PMNS matrix elements must not be zero. In other words, neutrino oscillations require at least one neutrino to have mass and to undergo flavor mixing. Furthermore, neutrino oscillation experiments are sensitive to the difference of squared masses

$$\Delta m_{ij}^2 = |m_i^2 - m_j^2|, \quad (1.8)$$

which only yields two independent observables for three masses. Thus, these experiments cannot be used to determine the absolute mass scale of neutrinos.

1.4.2. Experimental Considerations

According to equation (1.7), the ratio L/E determines the sensitivity of an experiment to the oscillation parameters given by the PMNS matrix U (mixing parameters) and Δm_{ij}^2 (mass ordering). L can be tuned by placing the detector in a suitable distance from a Earth-based neutrino source. E can either be tuned by using e. g. particle accelerators as source or it varies naturally, for instance if the source exhibits an energy spectrum like e. g. the Sun.

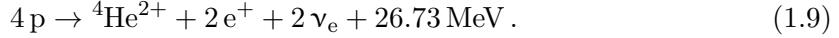
Furthermore, two detection channels can be distinguished. If, on one hand, an experiment is only sensitive to the neutrino flavor which is emitted by the neutrino source, it can detect a weakening of the neutrino flux. This is referred to as disappearance channel. If, on the other side, an experiment is sensitive to a different neutrino flavor than the one which the neutrino source emits, this is referred to as appearance channel [Zub11].

There are four major classes of neutrino sources that can be used to measure the mixing parameters and the mass ordering. They are listed in table 1.1. For each class multiple experiments exist [Zub11]. Not all experiments will be discussed here. Instead, the following section 1.4.3 discusses the so-called “Solar Neutrino Problem” as an exemplary case study on solar neutrino experiments.

1.4.3. Experimental Case Study: The Solar Neutrino Problem

The term “Solar Neutrino Problem” refers to the mismatch of the total neutrino flux arriving on Earth predicted by the standard solar model (SSM) and measured by, first, the Homestake and, later, other experiments. In this section the problem is briefly described and its resolution developed.

At the end of the 1930s, Bethe, von Weizsäcker and Critchfield showed that there are two main fusion cycles in the Sun, the so-called CNO and pp cycle. The latter is the primary source of solar neutrinos [Wei38; BC38; Bet39]. Its multi-step reaction can be summarized as [Zub11]



Its initial reaction and the one with the broadest neutrino energy spectrum (from below 0.1 to ~ 11 MeV [Tan+18]) are



It should be noted that only electron neutrinos are produced in the Sun. Starting from the 1970s the solar electron neutrino flux was measured; the first time by the Homestake

Table 1.1.: Neutrino source classes for neutrino oscillation experiments. Listed are different neutrino source classes and which neutrino flavors they emit [Zub11]. Furthermore, the oscillation parameters they are mainly sensitive to are tabulated along with an example experiment.

source class	flavors	sensitive to	example experiment
nuclear power plants	$\bar{\nu}_e$	$\sin \theta_{13}$	Double Chooz [Abe+16]
accelerators	$\nu_e, \nu_\mu, \bar{\nu}_e, \bar{\nu}_\mu$	$\sin \theta_{12}, \sin \theta_{23}, \Delta m_{12}^2, \Delta m_{23}^2$	MiniBooNE [Agu+09]
atmosphere	$\nu_e, \nu_\mu, \bar{\nu}_e, \bar{\nu}_\mu$	$\sin \theta_{23}, \Delta m_{23}^2$	Super-KamiokaNDE [Fuk+98]
the Sun	ν_e	$\sin \theta_{12}, \Delta m_{21}^2$	SNO [Aha+13]

experiment using the inverse beta decay of ^{37}Cl . It could detect electron neutrinos with an energy threshold of 813 keV. The measured flux was one third of the prediction by the SSM [Cle+98; BPB01]. This is marked as the beginning of the solar neutrino problem. The experiments GALLEX/GNO and SAGE confirmed the results, where the latter could detect electron neutrinos with an energy threshold of 233 keV [Kir98; Alt+05; Abd+09]. The low energy threshold is of importance because the neutrinos emitted by the initial reaction of the pp cycle, equation (1.10), exhibit the highest flux, but at the same time an energy spectrum that ends at approximately 400 keV [Bah96].

Starting from 1999 the SNO experiment measured the neutrino flux of all flavors. It used 1000 t of heavy water D_2O to detect electron neutrinos via charged currents as well as all flavors via neutral currents and elastic neutrino-electron scattering. In order to fully explain the flux data, the so-called Mikheyev-Smirnov-Wolfenstein effect (MSW effect) had to be respected [Wol78; MS86]: Electron neutrinos can undergo charged current interactions with surrounding electrons in a coherent forward scattering process, which alters the flavor transition amplitude. This effect is only significant in areas of high electron densities, such as the Sun. Taking these matter-mediated oscillations into account, the measured flux of all flavors of the ^8B neutrinos, equation (1.10), was in accordance with the electron neutrino flux predicted by the SSM [Aha+13]. Thus, the solar neutrino problem was resolved after more than three decades.

1.4.4. Summary of Experimental Results

This section summarizes the results obtained from neutrino oscillation experiments. The neutrino oscillation parameters consist of the squared neutrino mass differences, equation (1.8), and the mixing parameters of the PMNS matrix, equation (1.4). It should be noted first, that the MSW resonance of solar neutrinos requires $m_1 < m_2$, which still allows for two possible mass orderings [Zub11]:

1. normal ordering $m_1 < m_2 < m_3$ and
2. inverted ordering $m_3 < m_1 < m_2$.

For these two cases a combination of recent experimental results for the neutrino oscillation parameters is given in table 1.2 and illustrated in figure 1.2. All in all, neutrino oscillations are experimentally verified and provide unequivocal that neutrinos have mass.

1.5. Absolute Neutrino Mass Measurements

The absolute masses of neutrinos remain unknown and, as shown in section 1.4.1, neutrino oscillations are only sensitive to squared mass differences. This section presents methods to probe the absolute neutrino mass. The corresponding measurements fall into one of three categories [OW08]:

- observational cosmology (section 1.5.1),
- search for neutrinoless double β decay (section 1.5.2) or
- kinematic measurements of weak decays such as β decay and electron capture (section 1.5.3).

1.5.1. Observational Cosmology

In the early Universe, neutral particles such as light neutrinos could escape from areas of high mass density to areas of low mass density. As they carry away mass, the larger the neutrino mass, the stronger is the suppression of density fluctuations on small scales.



Figure 1.2.: Neutrino mixing parameters and mass ordering. The chart shows how the mass eigenstates ν_i are composed of the flavor eigenstates ν_α in the normal and inverted mass ordering. The composition depends on the phase δ . The mixing is shown for the two extreme cases $\delta = 0^\circ$ (baseline) and $\delta = 180^\circ$ (topline). (Adapted from [Sei19]. Numerical values can be found in [Est+19].)

In a mathematical formulation the power spectrum of the density contrast is examined and a higher neutrino mass would lead to a stronger suppression of small scale structures. Corresponding data are obtained, for example, by the Sloan Digital Sky Survey (SDSS). This survey records the sky's optical and infrared spectrum via telescope in order to map the large scale distribution of galaxies and galaxy clusters [Dor+04]. Furthermore, the temperature anisotropies in the cosmic microwave background (CMB) encode information on the Universe's structure. The latest and most precise data are recorded by the Planck satellite [Agh+18]. Under the assumption that all mass states contribute with the same number density cosmological observations are to first order only sensitive to the sum of all neutrino masses $\sum_i m_i$. A combination of the above data sets yields [Yèc+17]

$$\sum_i m_i < 0.14 \text{ eV} \quad (95\% \text{ C.L.}) .$$

Table 1.2.: Overview of neutrino oscillation parameters. The table lists the observables and their best fit values along with the 1σ -uncertainty range for normal and inverted ordering. (From [Est+19].)

observable	normal ordering	inverted ordering
$\sin \theta_{13}$	$0.310^{+0.013}_{-0.012}$	$0.310^{+0.013}_{-0.012}$
$\sin \theta_{23}$	$0.580^{+0.017}_{-0.021}$	$0.584^{+0.016}_{-0.020}$
$\sin \theta_{13}$	$0.02241^{+0.00065}_{-0.00065}$	$0.02264^{+0.00066}_{-0.00066}$
$\delta / {}^\circ$	215^{+40}_{-29}	248^{+27}_{-29}
$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	$7.39^{+0.21}_{-0.20}$	$7.39^{+0.21}_{-0.20}$
$\frac{\Delta m_{31}^2}{10^{-3} \text{ eV}^2}$	$+2.525^{+0.033}_{-0.032}$	
$\frac{\Delta m_{32}^2}{10^{-3} \text{ eV}^2}$		$-2.512^{+0.034}_{-0.032}$

1.5.2. Search for Neutrinoless Double- β Decay

Double- β decay ($2\nu\beta\beta$) is described as a nucleus of element $X(Z, A)$ with Z protons and $A - Z$ neutrons that decays to a daughter isotope $Y(Z + 2, A)$ via two simultaneous β decays

$$X(Z, A) \rightarrow Y(Z + 2, A) + 2e^- + 2\bar{\nu}_e. \quad (1.11)$$

If the neutrino is its own antiparticle, respectively of Majorana type, the neutrino emitted in the first decay can be absorbed in the second decay resulting in a neutrinoless double decay ($0\nu\beta\beta$). This would require the neutrino to have mass. Such a decay would manifest itself in a peak in the β spectrum two neutrino masses above the endpoint of the continuum [Zub11]. It should be noted that this would violate lepton number conservation. The half-life of such a decay encodes the Majorana mass of the electron neutrino as a coherent sum of all neutrino masses using the PMNS matrix U from equation (1.4)

$$m_{\beta\beta}^2 = \left| \sum_i U_{ei}^2 m_i \right|^2. \quad (1.12)$$

As U contains two unknown Majorana phases partial cancellation might occur. Hence, it is difficult to compare $m_{\beta\beta}$ to masses obtained by other methods. The two most stringent upper limits on $m_{\beta\beta}$ are listed in table 1.3.

1.5.3. Kinematic Measurements of Weak Decays

Several laboratory experiments as well as the supernova event 1987A have provided upper limits of absolute neutrino masses from the analysis of kinematics of weak interactions involving neutrinos or neutrino time-of-flight considerations. Such experiments can not

Table 1.3.: Constraining the neutrino mass by double- β decay experiments. Listed are the two most stringent limits on $m_{\beta\beta}$, equation 1.12, given as ranges.

experiment	isotope	90 % C.L. upper limit on $m_{\beta\beta}$ (eV)
GERDA [Ago+18]	^{76}Ge	0.12–0.26
KamLAND-Zen [Gan+16]	^{136}Xe	0.05–0.16



Figure 1.3.: Feynman graph of neutrinoless double- β decay. The graph depicts the simultaneous transformation of two neutrons into two protons where the down quarks transform into up quarks, whilst two electrons and two neutrinos are produced. The two emitted neutrinos annihilate in a Majorana transition.

Table 1.4.: Neutrino mass constraints from kinematic measurements. The table lists upper limits on the absolute neutrino masses for different neutrino flavors.

flavor	measurement basis	upper limit	reference
ν_e	neutrinos from Supernova 1987A	5.7 eV (95 % credible interval)	[LL02]
ν_μ	muon decay	17 keV (90 % C.L.)	[Ass+96]
ν_τ	tau decay	18.2 MeV (95 % C.L.)	[Bar+98]
$\bar{\nu}_e$	tritium- β decay	2 eV (95 % C.L.)	[Kra+05; Ase+11; Tan+18]

resolve the mass splitting between the squared mass eigenvalues. Therefore, the corresponding observable is a weighted sum of the N neutrino eigenmasses where the weights are the elements of the PMNS matrix from equation (1.4) [OW08]

$$m_{\nu_\alpha}^2 = \sum_i^N |U_{\alpha i}|^2 m_i^2. \quad (1.13)$$

In the scope of this thesis, the measurement of the mass of the electron antineutrino via β^- -decay kinematics is of special interest. Hence, this subject is examined more closely within this section. For completeness, aside from the upper limit on the mass of the electron antineutrino, table 1.4 also lists upper limits for other neutrino flavors obtained by kinematic measurements.

Neutrino Masses from β -decay Kinematics

In β^- decay

$$X(Z, A) \rightarrow Y(Z + 1, A) + e^- + \bar{\nu}_e \quad (1.14)$$

part of the released surplus energy generates the neutrino's mass. This leaves a signature in the β spectrum. (Also see section 3.1.) In a neutrino mass experiment four criteria are important for a suitable β emitter [OW08]:

- The β emitter should have an energy spectrum with a relatively low endpoint, because in the uncertainty on the neutrino mass enter the input uncertainties of neutrino energy and momentum scaled up with the endpoint energy.
- The β emitter must have a sufficiently high activity to provide statistically relevant count rates for quantities that can be handled in the laboratory.
- The β decay should be super-allowed in order for the nuclear matrix element of the decay process to be energy independent.

- The β -emitter molecule should be as simple as possible to allow for a theoretical treatment of its decay kinematics such as the final state of the decay-daughter molecule.

Tritium is an ideal candidate with respect to these criteria [OW08]. The corresponding measurement principle will be explained more closely in the following chapters about the KATRIN experiment. However, KATRIN has several predecessor experiments. The most recent two experiments based on tritium- β decay in Mainz and Troitsk obtained a combined upper limit on the electron antineutrino mass of [Kra+05; Ase+11; Tan+18]

$$m_{\bar{\nu}_e} < 2 \text{ eV} \quad (95\% \text{ C.L.}) .$$

It should be noted that KATRIN aims for a precision that is better by one order of magnitude.

2. The KATRIN Experiment

The Karlsruhe Tritium Neutrino (KATRIN) experiment performs a kinematic measurement of the tritium- β spectrum in order to determine the effective mass of the electron antineutrino (from here forth labeled m_ν) as defined by equation (1.13). In case no neutrino mass signal is observed, KATRIN aims to set an upper limit of

$$m_\nu < 200 \text{ meV} \quad (90\% \text{ C.L.}),$$

which is one order of magnitude more constraining than the one set by its predecessor experiments. KATRIN recorded the first β spectrum in a commissioning run in May 2018 and started neutrino mass measurements in March 2019.

This chapter provides an overview of the KATRIN apparatus. However, given KATRIN's complexity, it can by no means be exhaustive and for a comprehensive treatment the reader is referred to the KATRIN Design Report [KAT05] supplemented by an up-to-date hardware overview that is in the making at the time of writing this thesis¹.

2.1. Overview

The KATRIN experiment comprises a 70-m-long beam line depicted in figure 2.1. It can be divided into two sections:

1. The **source and transport section (STS)** comprises i. a. the gaseous tritium source where the tritium decays and the β electrons are magnetically guided along the beam line. Furthermore, the gas flow from the tritium source to the exit of the STS is reduced by at least 14 orders of magnitude.
2. In the **spectrometer and detector section (SDS)** the β electrons are filtered according to their kinetic energy and finally counted at the detector.

A central concept of the KATRIN setup is the magnetic flux tube. The β electrons must be guided from their point of origin to the detector. Therefore, a magnetic field is created by superconducting solenoids surrounding the beam line in the STS as well as coils around, and superconducting solenoids before and after the spectrometer tank in the SDS. The field lines are intersperse the beam line over the range of the whole experiment. The volume that is mapped onto the detector by this mechanism is called the flux tube. Within the flux tube, charged particles perform cyclotron motions around the field lines and are adiabatically guided from the STS to the detector. Adiabaticity is guaranteed by avoiding strongly varying field strengths on short distances.

A further central concept of KATRIN is the windowless source. As β electrons must not loose energy before energy analysis takes place, the SDS is windowlessly connected to the STS. However, the spectrometer must be kept practically free of any tritium flow

¹K. Altenmüller et al. (KATRIN collaboration), in prep.

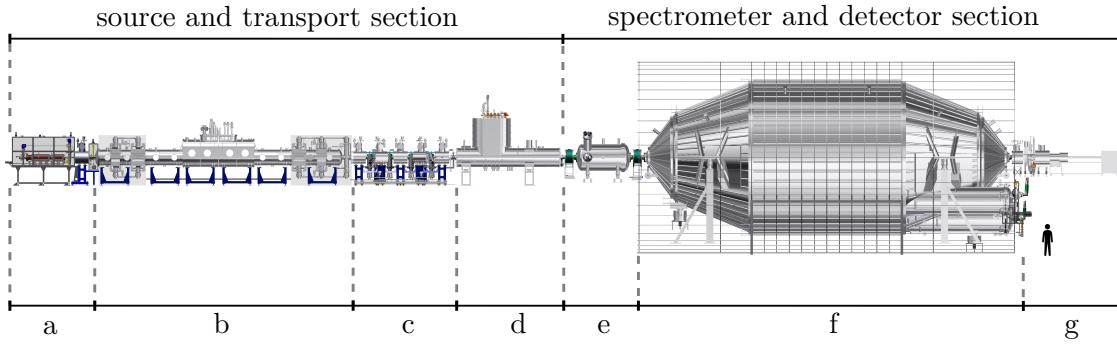


Figure 2.1.: The KATRIN beamline. Shown are the main hardware components:

- a) rear section (see section 2.3)
- b) windowless gaseous tritium source (WGTS) (see section 2.2)
- c) differential pumping section (DPS) (see section 2.4)
- d) cryogenic pumping section (CPS) (see section 2.5)
- e) pre spectrometer (see section 2.6)
- f) main spectrometer (see section 2.6)
- g) detector (see section 2.7)

for safety reasons and to keep the strict background requirements. Therefore, pumping systems reduce the gas inlet pressure of $1.8 \text{ mbar}\ell/\text{s}$ to the tritium partial pressure of below $10^{-14} \text{ mbar}\ell/\text{s}$ of the spectrometer.

The following sections step through the various components along the KATRIN beam line describing their functionality and purpose.

2.2. Windowless Gaseous Tritium Source

The WGTS is a 16-m-long, 1.5-m-wide and 4-m-high cryostat. It is depicted in figure 2.2 and a detailed description can e. g. be found in [Gro+08; Bab+12]. In the following the major features of the WGTS are reviewed:

Tritium purity: The molecular tritium (T_2) is injected in the middle of the 10-m beam tube of 90 mm diameter, where it decays. The design gas column density is $\rho d = 5 \times 10^{17} \text{ molecules}/\text{cm}^2$ with an isotropic tritium purity of $\epsilon_T = 95\%$ [KAT05] or better. At the front and rear of the WGTS, the gas is extracted from the beam tube by turbo molecular pumps in designated differential pumping sections called DPS-1-R (rear) and DPS-1-F (front). The extracted gas is re-injected in the center of the beam tube. The respective circulation system is called the inner loop [PSB15]. The tritium purity ϵ_T must be kept stable on a 0.1 % level [KAT05]. Therefore, a permeator is installed that separates impurities (like e. g. helium) and ejects them into the exhaust loop of the Tritium Laboratory Karlsruhe (TLK). Furthermore, the isotopic composition of the gas is monitored by a designated laser Raman system (LARA) [Sch13].

Injection pressure: The design injection pressure of the tritium gas is $1.8 \text{ mbar}\ell/\text{s}$. It must be kept stable at the 0.1 % level. This is achieved via a pressure- and temperature-controlled buffer vessel within the inner loop [PSB15].

Magnetic field: In order to adiabatically guide the β electrons to the spectrometer section the WGTS is pervaded by a magnetic field parallel to its beam tube axis of up to 5.6 T. It is created by 7 superconducting coils, that surround the beam tube. These magnets are kept at a temperature of 4.2 K by liquid helium [Are+18b].

Temperature: On the one hand, thermal motion smears the energy spectrum of the β electrons (Doppler effect). On the other hand, at low temperatures the gas molecules



Figure 2.2.: The windowless gaseous tritium source (WGTS). The hull and a sketch of the beam tube are shown. Indicated are the 8 turbo molecular pumps (TMP), the 7 magnets, the design temperatures for tritium operation, the maximum magnetic field strengths and a gradient within the beam tube depicting the decreasing gas density from the center to the sides. (Adapted from [Har15].)

start to cluster. $T = 30\text{ K}$ is chosen as a compromise and established by a two-phase neon cooling system. For calibration purposes, it is also possible to operate the WGTS with krypton-83m instead of tritium. This requires a beam tube temperature of $T = 100\text{ K}$ in order for the krypton not to freeze. In this operational mode the neon has to be exchanged for argon [KAT05].

2.3. Rear Section

The rear section (RS) terminates the beam line in the upstream direction and houses monitoring, calibration and control devices. It is depicted in figure 2.3 and a detailed description can e.g. be found in [Bab14]. In the following the major features of the RS are reviewed:

Electron gun: The RS houses an electron gun in order to measure the response function of the experiment (see section 3.2) via a electron source with a well-defined energy resolution of $\sim 0.2\text{ eV}$ and angular resolution of $\sim 4^\circ$. The electrons are guided towards and through the rear wall by a designated electromagnetic guidance system. Furthermore, their flight path can be adjusted by dipole magnets mounted in the WGTS which enable a scanning of the full flux tube [Bab14].

Rear wall and plasma control: The so-called rear wall is a gold-coated stainless-steel disc with a diameter of 6 inches that terminates the beam tube. It has a hole in the center to let electrons from the electron gun pass through. Its main purpose is the control of plasma effects: Space charges, respectively a plasma, forms within the WGTS due to the tritium decay. Therefore, β electrons may start at different potentials which adds uncertainty to the measured β spectrum. Simulations show that the plasma can be influenced by the rear wall potential which can be controlled via a voltage supply in the range of $\pm 10\text{ V}$. Moreover, a UV light illumination of the rear wall can provide a large flux of electrons extracted via the photoelectric effect that can compensate space charges. Therefore, a homogeneous work function of the rear wall with fluctuations less than 20 meV is required [Kuc+18; Kuc16].

Activity monitoring: A super conducting coil designed to create a magnetic field of 4.7 T in the RS ensures that the magnetic flux tube terminates at the rear wall. Hence, per

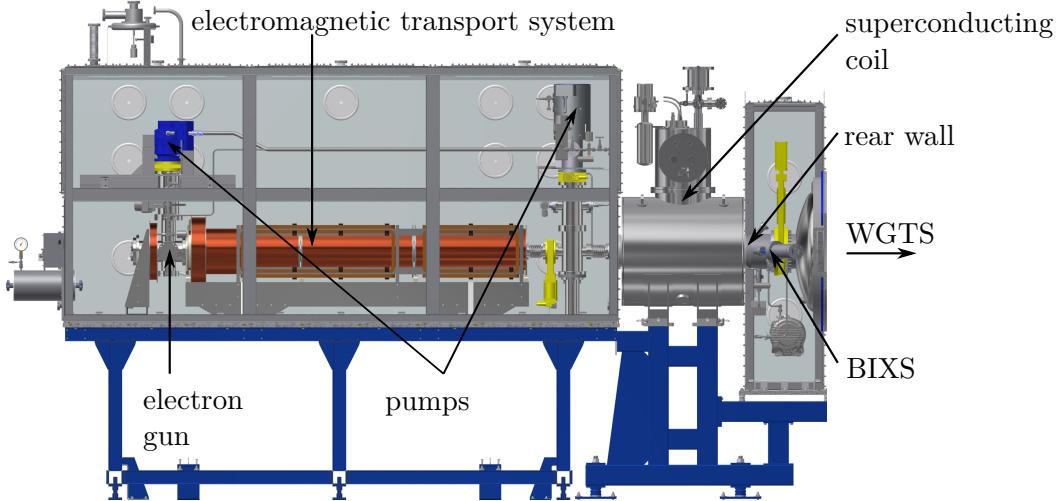


Figure 2.3.: The rear section (RS) terminates the KATRIN beam line and houses several monitoring and calibration devices that are described in the main text. (Adapted from [Sei19].)

design of the magnetic guidance, β electrons either arrive at the detector or hit the rear wall. On hitting the rear wall they emit bremsstrahlung. Two dedicated beta-induced X-ray spectroscopy (BIXS) systems measure the corresponding X-ray spectrum to determine the source strength respectively the gas column density [Röl15].

2.4. Differential Pumping Section

The differential pumping section (DPS) is composed of five elements. It is depicted in figure 2.4 and a detailed description can e. g. be found in [Kos12]. For orientation, in this section, the elements are labeled 1 to 5 from WGTS to cryogenic pumping section (CPS). In the following the major features of the differential pumping section (DPS) are reviewed:

Reduction of tritium flow: The five beam tube elements of the differential pumping section (DPS) form a 20° angle to each other and are arranged in a chicane. β electrons are magnetically guided along the chicane by a magnetic field of up to 5.5 T created by five superconducting solenoids. By contrast, the neutral gas molecules scatter off the walls. This reduces the molecular beaming effect and enhances the pumping probability [ZJH12]. Four turbo molecular pumps mounted between the beam tube elements then reduce the gas flow by approximately five orders of magnitude and feed the gas into the so-called outer loop where it is reprocessed [Kos12].

Ion blocking: In the WGTS, ions such as HeT^+ , T_2^+ , T_3^+ , T_5^+ can form. If not blocked, reach the spectrometer section together with the β electrons and are even accelerated by the retarding voltage (see section 2.6). This would eventually lead to an increased background rate. A potential barrier created by two ring electrodes in element 5 and the pump port between DPS and CPS set to +100 V avoids such a scenario. The positive ions are deflected, and dipole electrodes in the elements 1 to 4 make them drift out of the flux tube. They hit the wall and get neutralized [Klein19].

Ion monitoring: Downstream of the blocking electrodes the remaining ion flux is measured by a Fourier transform ion cyclotron resonance device (FT-ICR) [Ubi+09].

2.5. Cryogenic Pumping Section

The cryogenic pumping section (CPS) is an approximately 7-m-long cryostat. It is depicted in figure 2.5 and a detailed description can e. g. be found in [Jan15]. For orientation, in



Figure 2.4.: The differential pumping section (DPS) reduces the gas flow by five orders of magnitude and blocks tritium ions. Its five elements, each with a separate magnet (M1-M5), and connected by pump ports, are shown. The outer loop connects the turbo molecular pumps to the infrastructure of the TLK. (Adapted from [Sei19].)

this section, its seven elements are labeled 1 to 7 from WGTS to CPS. In the following the major features of the CPS are reviewed:

Reduction of tritium flow: The CPS consists of seven beam tube elements of which the first five are arranged in a similar manner as the beam tube elements of the DPS, in a chicane forming 15° angles. Charged particles are guided along the chicane by a magnetic field of up to 5.6 T created by seven superconducting coils. Neutral molecules hit the walls that are covered by a frozen argon layer cooled down to 3 K in order to cold-trap particles. These low temperatures are achieved via a liquid helium cooling and a heat shield of liquid neon. After the accumulation of about 1 Ci of tritium the argon frost layer has to be renewed. In order to achieve this, the beam tube is warmed up and the argon is pumped off along with the accumulated tritium. Tests and simulations show a reduction of the tritium flow by approximately 10 orders of magnitude [Jan15; Röt19].

The forward beam monitor (FBM): The FBM can be moved horizontally into the pump port between beam tube element 6 and 7 of the CPS with a 2-dimensional spatial resolution of 0.1 mm. Two *pin*-diodes measure the β-electron flux and thus the stability of the gas column density in the WGTS. Furthermore, the forward beam monitor (FBM) equips a temperature and a hall sensor. A second detector board holding a Faraday cup for ion measurements is also available [Klein19]. More information about the FBM can e.g. be found in [Ell+17; Ell19].

The condensed ^{83m}Kr source (CKrS): The condensed ^{83m}Kr source (CKrS) is a sub mono-layer of ^{83m}Kr on a pyrolytic graphite substrate with a diameter of 2 cm. It can be lowered in the pump port of the CPS and moved in a 2-dimensional plane perpendicular to the beam line. This enables the spatial scanning of the properties of the spectrometer using quasi-monoenergetic conversion electron lines of ^{83m}Kr [Bau14; Dyb19; Are+18a].

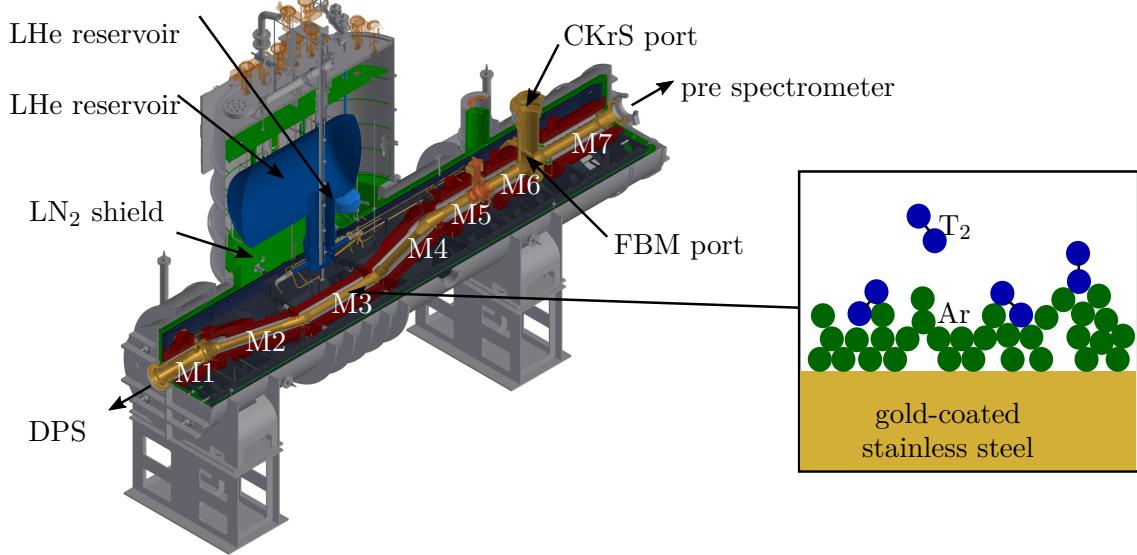


Figure 2.5.: The cryogenic pumping section (CPS) is the coldest part of the KATRIN experiment. It consists of 7 elements, labeled from 1 to 7 from the DPS to the pre spectrometer. Elements 2 to 5 are covered by a frozen argon layer at 3K in order to cold-trap tritium molecules. The low temperatures are established using liquid helium (LHe) and an insulation of liquid nitrogen (LN₂). Each element is enclosed by a super conducting coil (M1 to M7) for magnetic guidance of the β electrons. For the FBM and CKrS the reader is referred to the main text. (Adapted from [Sei19].)

2.6. Pre and Main Spectrometer

The pre and main spectrometer are vacuum vessels designed to filter passing electrons according to their kinetic energy. The pre spectrometer has a length of 3.4 m and a diameter of 1.7 m. Details on its design can e. g. be found in [Val09; Frä10]. The main spectrometer has a length of 23 m and a diameter of 10 m. Details on its design can e. g. be found in [Val09; Val04]. The functionality and purpose of spectrometer related aspects are reviewed in this section. The content is divided into two parts: Section 2.6.1 explains the so-called MAC-E filter principle and section 2.6.2 list several measures to keep the strict KATRIN background requirements.

2.6.1. MAC-E-filter principle

The pre and main spectrometer are based on the principle of the so-called magnetic adiabatic collimation with electrostatic filtering (MAC-E) [BPT80]. It enables the filtering of electrons according to their kinetic energy. Figure 2.6 sketches the MAC-E filter of KATRIN. The following paragraphs outline the basic concepts and their experimental implementation. As the principle is of key importance for the KATRIN experiment, it is additionally treated in a mathematical way in the subsequent section 3.2.3.

Electrostatic filtering: A retarding voltage barrier is applied along the beam axis within the spectrometer, reaching its maximum U at the so-called analyzing plane in the center and dropping off towards the source and the detector. The retarding voltage barrier deflects electrons with kinetic energies below eU . For higher energies it depends on the electron's angle with respect to the beam line axis whether it can pass the spectrometer or not, see equation (3.13).

Magnetic collimation: The electric field gradient of the retarding voltage barrier is parallel to the beam line, but β electrons are emitted in an arbitrary angle with respect



Figure 2.6.: Scheme of the KATRIN main spectrometer and the magnetic adiabatic collimation with electrostatic filtering (MAC-E) filter principle. The KATRIN design magnetic field settings are $B_{PS2} = 4.5$ T, $B_S = 3.6$ T, $B_{\max} = 6.0$ T, $B_D = 3.6$ T, $B_A \approx 3 \times 10^{-4}$ T. \vec{E} denotes the magnetic field regulated by the retarding potential U that reaches its maximum $U_a = U$ at the analyzing plane. (Adapted from [Sei19].)

to the magnetic field lines. In order to analyze their full kinetic energy, they have to be collimated. This is achieved by a magnetic field gradient that drops from $B_S = 3.6$ T in the STS to $B_A \approx 3 \times 10^{-4}$ T in the analyzing plane. In the following a plausibility argument for the momentum collimation due to the field gradient, according to [KAT05], is given: Electrons entering the spectrometer vessel perform cyclotron motions around the magnetic field lines. Their total kinetic energy E_S is split into a longitudinal component E_{\parallel} along the beam axis and a transverse component E_{\perp}

$$E_S = E_{\parallel} + E_{\perp}. \quad (2.1)$$

In the non-relativistic and adiabatic approximation the transverse component can be expressed by the magnetic field strength B and the electron's magnetic moment μ respectively its charge $q = e$, its mass m_e and angular momentum L [Jac75]

$$E_{\perp} = -\mu B = \frac{e}{2m_e} LB. \quad (2.2)$$

Adiabaticity conserves angular momentum L and the total energy of the electron E_S along its trajectory. Hence, when the magnetic field strength B decreases to $B_A = B_{\min}$ in the analyzing plane, the transverse component of the electron's energy E_{\perp} decreases likewise and transforms to longitudinal energy E_{\parallel} .

Magnetic Bottle effect: As the source is placed in a lower magnetic field B_S compared to the maximum field strength along the beam line B_{\max} at the detector side, β electrons traveling downstream to the detector are subject to the magnetic bottle effect [KAT05].

They get reflected and travel upstream to the rear wall if their starting angle θ_S with respect to the beam line axis surpasses θ_{\max} with

$$\sin \theta_{\max} = \sqrt{\frac{B_S}{B_{\max}}}. \quad (2.3)$$

For the KATRIN design values $B_{\max} = 6$ T and $B_S = 3.6$ T one obtains $\theta_{\max} \approx 51^\circ$. A cutting angle θ_{\max} is beneficial because the greater the emission angle of a β electron the larger the distance it travels in the WGTS and the more it is subject to energy losses such as scattering or synchrotron radiation [KAT05].

MAC-E-filter width: Electrons with a kinetic energy below qU cannot pass the spectrometer. Electrons with a kinetic energy above $qU + \Delta E$ do pass the spectrometer. Here, ΔE denotes the filter width [KAT05]

$$\Delta E = \frac{B_A}{B_{\max}} E. \quad (2.4)$$

Electrons with an energy between eU and $eU + \Delta E$ pass the potential barrier only with a certain probability. A quantitative description of this so-called transmission probability is given in the subsequent section 3.2.3. However, it can already be deduced, that a larger ΔE adds a greater uncertainty to the measurement and thus it should be kept as low as possible. ΔE depends on the maximum magnetic field strength along the beam line $B_{\max} = 6$ T, the kinetic energy of β electrons $E \approx 18.6$ keV, and the field in the magnetic field in the analyzing plane $B_A \approx 3 \times 10^{-4}$ T. Hence, its KATRIN design value is $\Delta E \approx 0.93$ eV.

Dimensions of the KATRIN main spectrometer: This paragraph outlines, why the diameter of the KATRIN main spectrometer is 10 m, while the one of its predecessor experiment in Mainz was only 1 m [Kra+05]. KATRIN's envisaged sensitivity requires a relative MAC-E-filter width of at least $\Delta E/E = 1/20000$, which directly corresponds the ratio of the magnetic fields B_{\max}/B_A (see equation (2.4)). For a smaller ΔE , B_A should be chosen as low as possible. However, the lower B_A , the wider the flux tube that must be governed by the spectrometer vessel. Also, the magnetic field must decrease at a sufficiently slow rate from the spectrometer's entrance to the analyzing plane in order to guarantee adiabaticity, which requires a certain spectrometer length. Dimensions that meet the demands and are feasible for the main spectrometer were found to be a radius of 10 m and a length of 23 m [KAT05; Val04].

Now, that the requirements on the magnetic and electrostatic fields are outlined, the following two paragraphs review their technical implementation:

Magnetic field: The main spectrometer is surrounded by a system of coils that shapes the MAC-E filter's magnetic field. Upstream, there is the PS2 magnet ($B_{\max} = 4.5$ T); downstream the pinch ($B_P = B_{\max} = 6.0$ T) as well as the detector magnet ($B_D = 3.6$ T), which are superconducting solenoids. The field is fine-tuned by a system of air coils around the spectrometer hull: There is the Earth magnetic field compensation system (EMCS) with 26 current loops parallel to the beam line axis. Furthermore, there is the low-field correction system (LFCS) with 14 air coils perpendicular to the beam line axis. The combined system constrains the electrons' flux tube to the spectrometer vessel and compensates the Earth's magnetic field as well as effects from ferromagnetic materials in the spectrometer's surroundings [Erh+18]. Additionally, a vertical and radial magnetic measuring system (VMMS and RMMS) are installed outside the spectrometer vessel. The field inside the spectrometer vessel is assessed via samples of these measuring systems combined with simulations [Let+18].

Electrostatic field: A high-voltage system establishes the MAC-E filter's retarding potential. The fluctuation of the retarding voltages must have a standard deviation smaller

than 60 mV for the envisaged sensitivity on the neutrino mass [KAT05]. The antenna-like beam line setup is sensitive to electromagnetic fluctuations of any source, which is why an active post-regulation system is deployed. It monitors the retarding potential and regulates it with the required precision. For the monitoring the monitor spectrometer and a voltage divider are deployed. For details on the later systems the reader is referred to [Thü+09; Erh+14; Zbo11].

2.6.2. Background Mitigation Strategies

The KATRIN sensitivity goal requires a background rate of less than 10 mcps [KAT05]. Several background-related aspects with respect to the spectrometer tanks are:

Vacuum: The spectrometers are operated at a pressure on the order of 1×10^{-11} mbar. This prevents electron scattering on residual gas and minimizes background effects by ionization. Correspondingly, turbo molecular and getter pumps are installed at three pump ports of the spectrometer vessels. Furthermore, the spectrometers can be baked out at up to 350 °C [Are+16].

Wire electrodes: The inner walls of the spectrometer vessels are lined by wire electrodes. Their potential is at a few hundred volts more negative than the spectrometer hull reflecting electrons coming from the vessel walls. Such electrons may be induced by cosmic rays [Val09].

Ion blocking: Analogously to the ones in the CPS (section 2.5), three blocking electrodes are installed; one between the CPS and the pre spectrometer, one between the pre and main spectrometer; and one between the main spectrometer and the detector [Klein19].

Tandem setup: β electrons may scatter on residual gas. This can either directly lead to secondary electrons or create positive ions that travel down the beam line. The positive ions in turn may again yield secondary electrons through scattering. The more β electrons enter the main spectrometer, the higher is the probability to create secondary electrons. In order to reduce the flux of β electrons into the main spectrometer, the retarding potential of the pre spectrometer is set to a few hundred volts more positive than the one of the main spectrometer. On the one hand this is a countermeasure against background events, but on the other hand, charged particles can be trapped between the two spectrometers due to the electromagnetic setup (Penning trap). A sudden discharge may harm the hardware, especially the detector. Therefore, it is possible to sweep a charged wire through the volume in order to collect the trapped particles and avoid this “Penning discharges” [Val09].

2.7. Detector Section

The detector section terminates the beam line in downstream direction. It can be separated from the spectrometer section by closing a gate valve. The detector section is depicted in figure 2.7 and a detailed description can e.g. be found in [Ams+15]. In the following the major features of the detector section are reviewed:

Focal plane detector (FPD): The FPD counts the β electrons that pass the spectrometer section. It is a *pin*-silicon detector with a sensitive area of 9 cm diameter. It is subdivided in 148 pixels of the same area arranged in 12 rings of 12 pixels each and the so called bull’s eye of 4 pixels in the center. This arrangement allows later correction for radial electrical, magnetic and gas dynamical inhomogeneities in the beam line [Ams+15].

Shield and veto system: The radiation shield of the FPD system consists of two nested cylindrical shells: an outer lead shell of 3 cm that reduces photon background and an inner copper shell of 1.27 cm that blocks X-rays originating from the outer lead shell. The shield is surrounded by a veto system to tag incoming muons. Such a system is necessary to keep the strict background requirements [Ams+15].

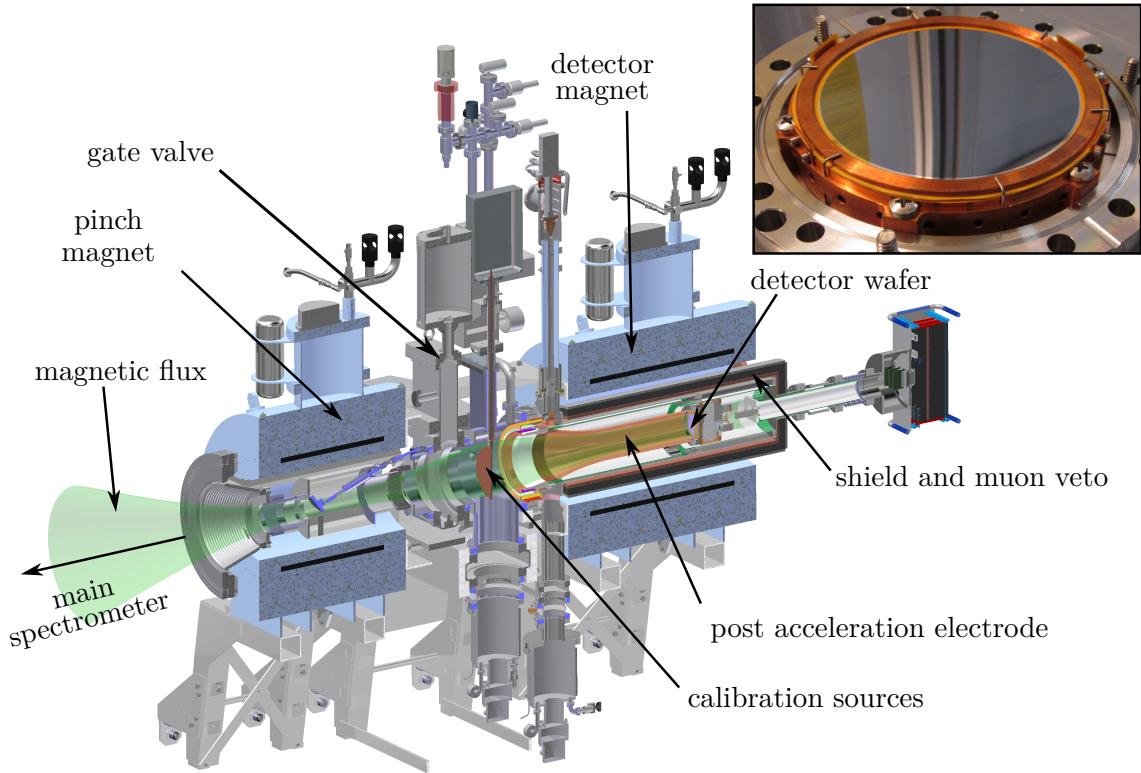


Figure 2.7.: The detector section terminates the KATRIN beam line. Among other instruments it houses the focal plane detector (FPD) for β electrons with the detector wafer at its core. For an explanation of the other components the reader is referred to the main text. (Adapted from [Sei19].)

Calibration: Photoelectron sources can be lowered in the line of sight of the detector. The corresponding photocurrent can be measured with the precision ultra-low current integrating normalization electrometer for low-level analysis (PULCINELLA) system. A comparison of PULCINELLA and the FPD yields the FPD's detection efficiency. It was determined to be $\epsilon_{\text{det}} = 95 \pm 1.8 \pm 2.2\%$ [Ams+15].

Detector magnet: The detector magnet ($B_D = 3.6$ T) allows to form the flux tube near the detector independently of the main spectrometer magnetic field setting. It especially allows its mapping on the the detector [Ams+15].

Post-acceleration electrode: The post-acceleration potential shifts the electrons arriving from the main spectrometer to a more favorable energy region. This increases the detector efficiency and, additionally, β electrons can be distinguished from noise originating in the detector by an energy region of interest cut. An appropriate setting was found to be ~ 10 keV [Ams+15].

3. Mathematical Formalism of a KATRIN Measurement

For neutrino mass inference from data or simulation a mathematical model of a KATRIN neutrino mass measurement is required. As parameter inference is of importance within the scope of this thesis, a mathematical formalism describing a KATRIN neutrino mass measurement is outlined within this chapter. Therefore, an expression for the β -decay rate of a tritium molecule is given in section 3.1. Section 3.2 describes the KATRIN response function, respectively the mathematical modeling of the KATRIN apparatus. Sections 3.3 combines the concepts of the two preceding sections into the β -electron rate at the KATRIN detector. Section 3.4 incorporates the detector efficiency, the background rate and the measurement time in order to translate the rates into an expression for the electron counts measured by the KATRIN detector for a fixed retarding potential. Section 3.6 introduces the concept of a measurement time distribution over retarding potentials. Finally, section 3.7 shows a full simulated KATRIN neutrino mass measurement.

3.1. Differential Tritium- β -Decay Spectrum

This section presents a quantitative expression for the β -decay rate of a tritium molecule in dependence on the kinetic energy of the emitted β electron (differential rate). First the whole mathematical description is denoted, then its components are explained.

Using Fermi theory and Fermi's golden rule the decay rate of a tritium molecule is [Kle+19; OW08]

$$\frac{d\Gamma(E_S)}{dE_S} = \frac{G_F^2 |V_{ud}|^2}{2\pi^3} |M_{\text{nuc}}|^2 \cdot F(Z, E_S) \cdot p(E_S + m_e) \cdot \sum_f P_f \cdot \epsilon_f \cdot \sqrt{\epsilon_f^2 - m_\nu^2} \cdot \Theta(\epsilon_f - m_\nu). \quad (3.1)$$

Its constituents are the kinetic electron energy E_S ; the effective electron-antineutrino mass m_ν defined via the PMNS matrix U (see equation 1.4),

$$m_\nu^2 = |U_{ei}|^2 m_i^2; \quad (3.2)$$

the Fermi constant G_F ; the up-down-quark-coupling given by the Cabibbo angle θ_C [Tan+18]

$$V_{ud} = \cos \theta_C = 0.97425 \pm 0.00022; \quad (3.3)$$

and the nuclear transition matrix element [Tan+18]

$$|M_{\text{nuc}}|^2 = g_V^2 + 3g_A^2 \quad \text{with } g_V = 1 \quad \text{and} \quad g_A/g_V = -1.2646 \pm 0.0035 \quad (3.4)$$

which is independent of the electron's kinetic energy as the decay is super-allowed and given by the vector g_V and axial vector g_A coupling.



Figure 3.1.: Tritium- β spectrum for a vanishing and non-vanishing neutrino mass. The plot shows the differential rate as described by equation (3.1) for a vanishing and non-vanishing neutrino mass. The inset zooms into the endpoint region where a non-vanishing mass causes a shift and a distortion of the spectrum. (Calculated with [Soft-SSC].)

Furthermore, the Fermi function $F(Z, E_S)$ accounts for the Coulomb interaction between the outgoing electron and the daughter nucleus with atomic charge $Z = 2$, which in its relativistic version can be approximated as [Sim81]

$$F(Z, E_S) \approx \frac{2\pi\eta}{1 - \exp 2\pi\eta} \cdot R, \quad (3.5)$$

with Sommerfeld parameter $\eta = \alpha Z/\beta$, fine structure constant α , relativistic velocity β and a relativistic correction factor $R = 1.002037 - 0.001427\beta$.

The phase-space factor of the outgoing electron with momentum p and mass m_e is given by the factor $p(E_S + m_e)$.

The phase space factor of the emitted neutrino depends on multiple quantities: First, there is the β -spectrum endpoint of molecular tritium $E_0 = (18\,574.00 \pm 0.07)$ eV [Mye+15; OW08]. Second, there is the final state energy of the molecule V_f . The excited energy state f is caused by vibration, rotation or electronic excitation of the decaying molecule. A review on tritium molecular final states and tabulated values can e.g. be found in [BPR15] and references therein. The probability that the molecule is in a final state of energy V_f after the decay is denoted by P_f . Then the energy of the neutrino reads

$$\epsilon_f = E_0 - E - V_f. \quad (3.6)$$

Third, there is the neutrino's momentum $\sqrt{\epsilon_f^2 - m_\nu^2}$. Then, the complete phase space factor of the neutrino is a sum over all possible molecular final states labeled f .

Lastly, the Heavyside step function Θ ensures a positive kinetic energy of the neutrino.

The differential rate is depicted in figure 3.1 for a vanishing and non-vanishing effective electron-antineutrino mass. The difference of the two β spectra forms the foundation for neutrino mass inference at KATRIN.

3.2. Response Function

The aim of this chapter is an introduction to the mathematical formalism for the electron rate at the KATRIN detector. The previous section 3.1 gives an expression for the differential β -electron rate. The next step is the inclusion of the characteristics of the KATRIN experimental setup. This can be accomplished by denoting the KATRIN response function. In the outlined formalism, it reflects the probability of an electron emitted in the WGTS to reach the KATRIN detector [Gro15] (the detector efficiency is treated separately.)

First, central concepts and the nomenclature are presented in section 3.2.1. Then, components of the response function are introduced:

- The gas dynamics within the STS needs to be simulated. See section 3.2.2.
- The characteristics of the KATRIN spectrometer can be summarized in the transmission function. See section 3.2.3.
- The passage of electrons through the WGTS is influenced by scattering from gas molecules. The probability for such scattering is discussed in section 3.2.4. Furthermore, the amount of energy an electron loses when scattering is considered in section 3.2.5.

Finally, the described components will be assembled to the KATRIN response function in section 3.2.6.

3.2.1. Concepts and Nomenclature

Before the formalism for the KATRIN response function is developed, this section introduces naming conventions and useful concepts.

Coordinate System

This chapter focuses on a one-dimensional description of the KATRIN response function. The position along the beam line is denoted with z . The origin of the coordinates system is the center of the WGTS as already chosen in previous works, e. g. [Gro15; Kle14]. In this sense, the rear and the front of the WGTS of length d have the coordinates $\mp d/2$.

Pitch Angle

Within this chapter the angle between an electron's direction of motion and the beam line axis, the so-called pitch angle, is denoted by θ .

Parameter Indices

Whether an electron reaches the KATRIN detector depends i. a. on its parameters when originating in the WGTS. Within this chapter these starting parameters are denoted with a lower index S. The three decisive starting parameters are the following:

1. The starting kinetic energy E_S as discussed within the description of the differential rate in equation (3.1).
2. The starting position z_S within the WGTS.
3. The starting pitch angle θ_S within the WGTS.

Parameters that denote quantities in the analyzing plane (see section 2.6) are denoted with a lower index A.

Probabilistic Treatment of the Starting Pitch Angle

It should be noted, that the three listed starting parameters are not known for a single β electron, which suggests a probabilistic treatment. Within the scope of this thesis this is of importance with respect to the starting pitch angle. Therefore, the concept is explained in the following:

Given the distribution $\omega(\theta_S)$ of starting pitch angles, the mean value of any function $g(\theta_S)$ depending on a fixed starting pitch angle θ_S can be calculated within an interval $[0, \theta_{\max}]$ by applying the definition of the mean value

$$\langle g(\theta_S) \rangle = \frac{\int_0^{\theta_{\max}} \omega(\theta_S) g(\theta_S) d\theta_S}{\int_0^{\theta_{\max}} \omega(\theta_S) d\theta_S}. \quad (3.7)$$

An isotropic β -electron emission by a tritium molecule into the unit sphere, meaning all combinations of spherical emission angles $(\varphi, \vartheta = \theta_S)$ are equally likely, yields as distribution for the starting pitch angles [KAT05]

$$\omega(\theta_S) = \sin \theta_S \quad (3.8)$$

with normalization

$$\int_0^{\theta_{\max}} \omega(\theta_S) d\theta_S = \frac{1}{1 - \cos \theta_{\max}}. \quad (3.9)$$

Within this chapter θ_{\max} denotes the maximum acceptance angle due to the magnetic bottle effect as explained in section 2.6.1 with a design value of $\theta_{\max} \approx 51^\circ$ [KAT05]. This calculation of the mean value is applied multiple times within this chapter and in the scope of this thesis.

Experimental Settings

As the response function models the characteristics of the KATRIN apparatus it naturally depends on the experimental settings. The quantities used within this chapter are listed in the following:

- the magnetic field B_S at the place of origin of a β electron within the WGTS;
- the magnetic field B_A within the analyzing plane;
- the maximum magnetic field B_{\max} along the beam line axis;
- the retarding voltage U and the retarding energy qU ;
- the starting potential U_S of a β electron within the WGTS.

For the detailed meaning of these parameters and their KATRIN design values, see section 2.6. It should be noted, that none of these quantities are constant, but they exhibit a spatial, especially a radial, dependency [KAT05]. For ease of notation, the spatial dependency is left implicit within this chapter.

3.2.2. Gas Dynamics

The gas dynamics within the STS has to be simulated. This topic is not treated in detail here. The reader is referred to [Höt12]. In short, in a one-dimensional description the output of such a gas dynamic simulation is the gas molecule density $\rho(z)$. For nominal settings, averaging $\rho(z)$ along the beam line axis and multiplication by the length d of the WGTS yields the design column density $\rho d = 5 \times 10^{17} \text{ cm}^{-2}$ [KAT05].



Figure 3.2.: The KATRIN transmission function as described by equation (3.13). It denotes the probability for an electron with a kinetic energy E to pass through the spectrometer set to a retarding potential of U . The probabilistic treatment of the starting pitch angles of electrons leads to the MAC-E-filter width ΔE with the nominal value of 0.93 eV^2 [KAT05]. (Calculated with [Soft-SSC].)

3.2.3. Transmission Function

The transmission function denotes the probability of an electron to pass the MAC-E filter. It can be characterized by the transmission energy [Gro15]

$$E_{\text{tr}}(qU, E, \theta_S) = \frac{q(U - U_S)}{1 - \sin^2 \theta_S \frac{B_A}{B_S} \frac{\gamma(E) + 1}{\gamma_A + 1}}. \quad (3.10)$$

where $\gamma(E)$ and γ_A denote the relativistic Lorentz factor of the β electrons with energy E and in the analyzing plane. As the electrons are slowed down substantially by the retarding potential in the spectrometer, it holds $\gamma_A \approx 1$. In the following, for ease of notation, also $U_S = 0$ and $\gamma_S = 1$ is assumed.

Electrons pass the MAC-E filter if their energy E when arriving at the spectrometer surpasses the transmission energy E_{tr} , equation (3.10). This condition can be resolved for the starting pitch angle [Gro15]

$$\begin{aligned} E &> E_{\text{tr}}(qU, E, \theta_S) \\ \Leftrightarrow \theta_S &< \theta_{\text{tr}}(E, qU) := \arcsin \left(\sqrt{\frac{E - qU}{E} \frac{B_A}{B_S}} \right). \end{aligned} \quad (3.11)$$

Using equation (3.11), the transmission function depending on the starting pitch angle and the starting energy of electrons can be formulated as a step function

$$\mathcal{T}(E, qU, \theta_S) = \begin{cases} 1 & \text{if } \theta_S < \theta_{\text{tr}}(E, qU) \\ 0 & \text{otherwise} \end{cases}. \quad (3.12)$$

Calculating the mean value of this step function with respect to the probabilistic distributed starting pitch angles of β electrons as described in section 3.2.1 yields the KATRIN transmission function [KAT05]

$$T(E, qU) = \langle \mathcal{T}(E, qU, \theta_S) \rangle = \begin{cases} 0 & \text{if } E < qU \\ \frac{1 - \sqrt{1 - \frac{E - qU}{E} \frac{B_S}{B_A}}}{1 - \sqrt{1 - \frac{\Delta E}{E} \frac{B_S}{B_A}}} & \text{if } qU < E < qU + \Delta E , \\ 1 & \text{if } qU + \Delta E < E \end{cases} \quad (3.13)$$

where

$$\Delta E = E \cdot B_A / B_{\max} \quad (3.14)$$

is the MAC-E-filter width as explained in section 2.6. The transmission function is depicted in figure 3.2 for the KATRIN design values.

3.2.4. Probability of Electron Scattering within the WGTS

This section derives an expression for the probability P_l of an electron to scatter l times within the WGTS.

The electron moves on a spiral track due to its cyclotron motion in the magnetic field in the WGTS. Therefore, when traveling an infinitesimal distance dz in z -direction, it travels a total distance of

$$ds = \frac{1}{\cos \theta_S} dz . \quad (3.15)$$

Remarkably, this expression is independent of the electron energy and the magnetic field strength in the WGTS. The effective column density can then be expressed as a line integral along the electrons path φ over the gas density $\rho(z)$ from the starting position of the electron to the point where it leaves the WGTS

$$\lambda(z_S, \theta_S) = \int_{\varphi} \rho(\vec{r}) ds = \frac{1}{\cos \theta_S} \int_{z_S}^{d/2} \rho(z) dz . \quad (3.16)$$

The expected scattering count then is the product of the effective column density $\lambda(z_S, \theta_S)$ and the scattering cross section σ [Gro15]

$$\mu(z_S, \theta_S) = \lambda(z_S, \theta_S) \sigma . \quad (3.17)$$

Using $\mu(z_S, \theta_S)$, the probability for l -fold scattering can be expressed as a Poisson distribution [Gro15]

$$P_l(z_S, \theta_S) = \frac{\mu(z_S, \theta_S)^l}{l!} e^{-\mu(z_S, \theta_S)} . \quad (3.18)$$

The mean value with respect to the starting positions and the starting pitch angles can be calculated [Gro15]

$$\bar{P}_l = \frac{1}{d} \int_{-d/2}^{d/2} \frac{1}{1 - \cos \theta_{\max}} \int_0^{\theta_{\max}} \sin \theta_S P_l(z_S, \theta_S) d\theta_S dz_S . \quad (3.19)$$

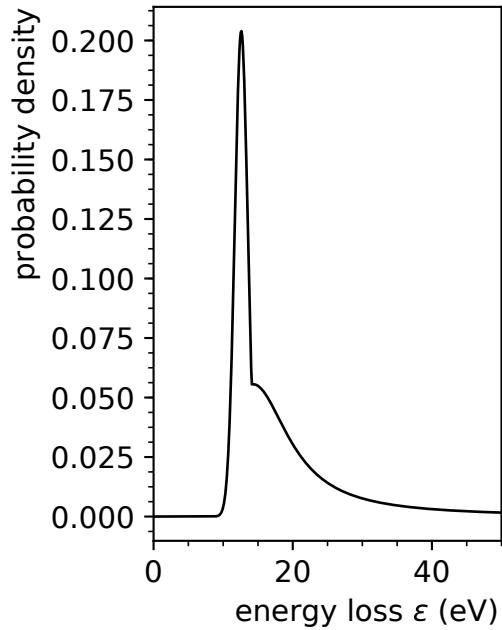
Table 3.1 lists the numerical evaluation of these averaged scattering probabilities. With these probabilities at hand, the next step is the derivation of the energy an electron loses when scattering.

Table 3.1.: Probability for electron scattering within the WGTS averaged over starting positions and pitch angles. Listed are the evaluations of equation (3.19) for the following input parameters: A scattering cross section of $\sigma = 3.456 \times 10^{-22} \text{ m}^2$ [KAT05], a constant gas column density $\rho d = 5 \times 10^{17} \text{ cm}^{-2}$, a WGTS length of $d = 10.0820 \text{ m}$ and a maximum acceptance angle of $\theta_{\max} = 50.7685^\circ$. The same values can be found in [Gro15; Kle14].

scattering count l	scattering probability according to equation (3.19)
0	41.33 %
1	29.27 %
2	16.73 %
3	7.91 %
4	3.18 %

Figure 3.3.: The energy loss probability density due to electron scattering in the WGTS given by equation (3.21) and determined at the Troitsk experiment [Ase+00]. The table below lists the corresponding parameters, where ϵ_1 was fixed and ϵ_c was chosen to make the piece wise defined function continuous.

parameter	value
A_1	0.204 ± 0.001
A_2	0.0556 ± 0.0003
ω_1	$(1.85 \pm 0.02) \text{ eV}$
ω_2	$(12.5 \pm 0.1) \text{ eV}$
ϵ_1	12.6 eV
ϵ_2	$(14.30 \pm 0.02) \text{ eV}$
ϵ_c	14.09 eV



3.2.5. Energy Loss of Electrons due to Scattering

This section describes the “energy loss function” $f_l(\epsilon)$. It denotes the probability density for an electron to loose an energy ϵ when scattering l times. Only the case of inelastic scattering is treated here. For an additional treatment of elastic scattering, which is less likely by one order of magnitude, the reader is referred to [Kle+19].

The energy loss function for no scattering is the Dirac delta function [Kle+19]

$$f_0(\epsilon) = \delta(\epsilon). \quad (3.20)$$

A phenomenological description for 1-fold scattering of electrons from hydrogen isotopologues was derived from data at the Troitsk experiment [Ase+00; Abd+17]

$$f_1(\epsilon) = \begin{cases} A_1 e^{-2\left(\frac{\epsilon-\epsilon_1}{\omega_1}\right)^2} & \text{if } \epsilon < \epsilon_c \\ A_2 \frac{\omega_2^2}{\omega_2^2 + 4(\epsilon - \epsilon_2)^2} & \text{if } \epsilon \geq \epsilon_c \end{cases} \quad (3.21)$$

Figure 3.3 depicts this energy loss function. It should be noted that in the scope of this thesis a more recent, preliminary energy loss model derived from a dedicated subgroup of the KATRIN collaboration is investigated in chapter 6.

For multiple scattering the above function f_1 has to be convoluted with itself and the energy loss function becomes [Kle+19]

$$f_l(\epsilon) = \bigotimes_{i=0}^l f_1(\epsilon) \quad (3.22)$$

where \otimes denotes the convolution

$$(f \otimes f)(\epsilon) = \int_{-\infty}^{\infty} f(\epsilon - \epsilon') f(\epsilon') d\epsilon'. \quad (3.23)$$

3.2.6. Assembly of the KATRIN Response Function

This section gives an expression for the KATRIN response function. It should be noted that the chosen notation differs slightly from those used in the works [Gro15; Kle+19], that this derivation is largely based on. The latter make approximations of the transmission properties and introduce more involved concepts where the approximations do not hold. Here, this approach is inverted: First, the involved concepts are applied and the approximations are introduced in a second step. However, the final results reconcile.

The KATRIN response function in dependence on the starting position and pitch angle of an electron reads

$$\mathcal{R}(E_S, qU, z_S, \theta_S) = \sum_l \int_{-\infty}^{\infty} \mathcal{T}(E_S - \epsilon, qU, \theta_S) \cdot P_l(z_S, \theta_S) \cdot f_l(\epsilon) d\epsilon \quad (3.24)$$

$$= \sum_l \int_0^{\infty} \mathcal{T}(E_S - \epsilon, qU, \theta_S) \cdot P_l(z_S, \theta_S) \cdot f_l(\epsilon) d\epsilon. \quad (3.25)$$

where the integral goes over the energy losses, the sum goes over the scattering count, \mathcal{T} denotes the non-averaged transmission function (3.13), P_l the non-averaged scattering probabilities (3.18) and f_l the energy loss function (3.22). The cut of the lower integral limit is caused by the vanishing energy loss function ($f_l(\epsilon) = 0$ if $\epsilon < 0$). In words, the transmission function is smeared using the energy loss function as a smearing kernel and then a weighted sum is formed over generations of l -fold scattered electrons where the weight is the probability to scatter l times.

The mean value of equation (3.24) with respect to the starting pitch angle can be calculated as described in section 3.2.1. Also the corresponding integral is swapped with the integral over the energy loss and the sum over the scattering count

$$\begin{aligned} R(E_S, qU, z_S) &= \langle \mathcal{R}(E_S, qU, z_S, \theta_S) \rangle \\ &= \sum_l \int_0^{\infty} \int_0^{\theta_{\max}} \frac{\sin \theta_S \cdot \mathcal{T}(E_S - \epsilon, qU, \theta_S) \cdot P_l(z_S, \theta_S)}{1 - \cos \theta_{\max}} d\theta_S \cdot f_l(\epsilon) d\epsilon. \end{aligned} \quad (3.26)$$

This expression can be reformulated to have the same form as the non-averaged response function (3.24). This means, the product “transmission function times scattering probability times energy loss function” can be reestablished, which also reconciles the notation with the expression given in [Gro15]. Therefore, the factor $1 = \bar{P}_l / \bar{P}_l$ with the averaged scattering probabilities from equation (3.19) is introduced into equation (3.26) and the “detailed transmission function” T_l^* is defined

$$\begin{aligned} R(E_S, qU, z_S) &= \sum_l \int_0^{\infty} \underbrace{\int_0^{\theta_{\max}} \frac{\sin \theta_S \cdot \mathcal{T}(E_S - \epsilon, qU, \theta_S) \cdot P_l(z_S, \theta_S)}{(1 - \cos \theta_{\max}) \cdot \bar{P}_l} d\theta_S}_{T_l^*(E_S - \epsilon, qU, z_S)} \cdot \bar{P}_l \cdot f_l(\epsilon) d\epsilon \\ &= \sum_l \int_0^{E_S - qU} T_l^*(E_S - \epsilon, qU, z_S) \cdot \bar{P}_l \cdot f_l(\epsilon) d\epsilon, \end{aligned} \quad (3.27)$$

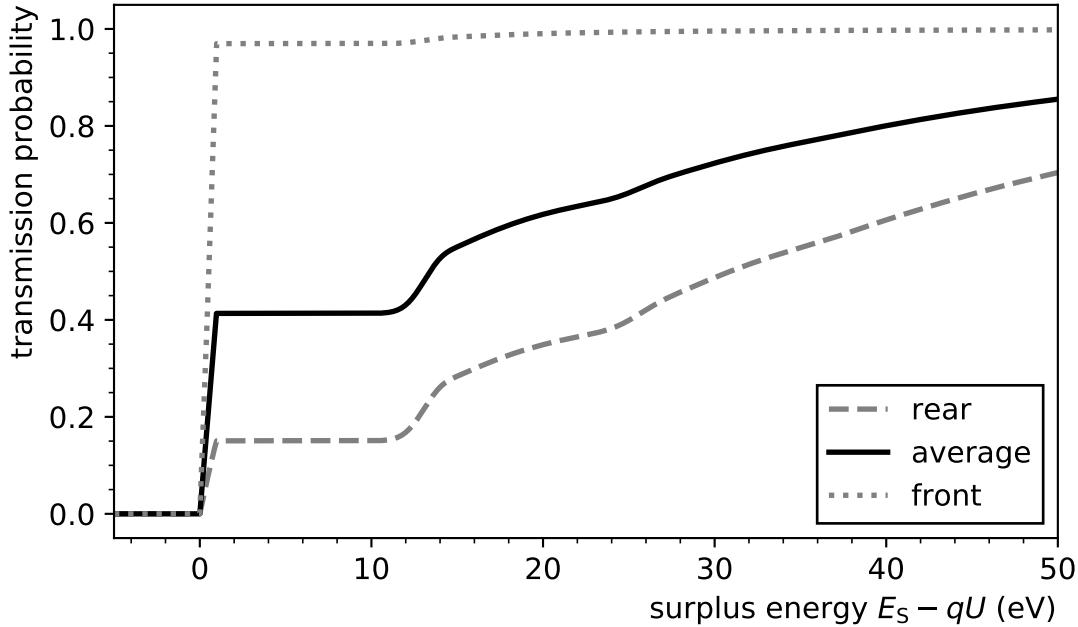


Figure 3.4.: The KATRIN response function at a retarding energy of $qU = 18\,545$ V. It is depicted for three cases: for electrons starting ~ 9 mm from the rear and front of the WGTS and averaged over all starting positions. For a description of the manifold features, the reader is referred to the main text. (Calculated with [Soft-SSC].)

where the cut on the upper integral limit from ∞ to $E_S - qU$ is justified below.

The non-averaged transmission function \mathcal{T} , equation (3.11), within T_l^* is a step function with respect to the starting pitch angle θ_S of an electron. This cuts the upper integral limit from θ_{\max} to θ_{tr} when integrating over θ_S . Furthermore, in analogy to the KATRIN transmission function from equation 3.13, a distinction of cases avoids imaginary square roots. One obtains the detailed transmission function as given in [Gro15; Kle+19]

$$T_l^*(E, qU, z_S) = \begin{cases} 0 & \text{if } E < qU \\ \int_0^{\theta_{\text{tr}}(E, qU)} \frac{\sin \theta_S \cdot P_l(z_S, \theta_S)}{(1 - \cos \theta_{\max}) \cdot \bar{P}_l} d\theta_S & \text{if } qU < E < qU + \Delta E , \\ 1 & \text{if } qU + \Delta E < E \end{cases} \quad (3.28)$$

where θ_{tr} denotes the transmission-pitch angle (3.11) and ΔE the MAC-E-filter width (3.14). As T_l^* vanishes for $E < qU$ the upper integral limit over energy losses in the response function 3.27 can be cut to $E_S - qU$. Furthermore, it was found, that for $l > 3$ scatterings the detailed transmission T_l^* function can be exchanged for the KATRIN transmission function (3.13) without making a significant error [Gro15].

Summary

In equation 3.27 the KATRIN response function was derived, which reconciles with the expressions given in [Gro15; Kle+19]

$$R(E_S, qU, z_S) = \sum_l \int_0^{E_S - qU} T_l^*(E_S - \epsilon, qU, z_S) \cdot \bar{P}_l \cdot f_l(\epsilon) d\epsilon , \quad (3.29)$$

where T_l^* is the detailed transmission function, equation (3.28), \bar{P}_l are the averaged scattering probabilities, equation (3.19) and f_l is the energy loss function, equation (3.22).

The response function denotes the probability of an electron starting with an energy E_S at a position z_S to overcome the retarding energy qU and reach the detector.

Figure 3.4 shows the response function for two different starting positions of electrons as well as averaged over all starting positions. It exhibits many features: For unscattered electrons the response function resembles the transmission function. This causes the steep rise at $E_S - qU = 0$ eV within the interval of the MAC-E-filter width $\Delta E \approx 0.93$ eV. As the transmission probability is weighted by the probability for no scattering, the plateaus resemble the corresponding probabilities (equation 3.18 averaged over starting pitch angles): $\sim 12\%$ (rear), 41.33% (average, see table 3.1) $\sim 98\%$ (front). The discontinuity in the first derivative of the energy loss function at $\epsilon_c = 14.09$ eV causes kinks. As the energy loss function has an onset at $\epsilon_0 \approx 10$ eV, the corresponding kinks are at approximately $n \cdot \epsilon_0 + \epsilon_c$ ($n \in \{0, 1, \dots\}$) and increasingly smoothed for higher n . (Also see figure 3.3 for the energy loss function.) The response function (3.29) can be understood as a weighted sum of smeared transmission functions approximately shifted by the onset of the energy loss function. Electrons starting from the front of the WGTS are unlikely to scatter, which is why the response function almost resembles the transmission function. Electrons starting from the rear are likely to scatter. Thus, the corresponding response function shows the features of multiple scatterings. For multiple scatterings, the sharp edges of the transmission function are smoothed by the energy loss, which is why only one sharp edge and one plateau is apparent.

3.3. Integral Rate

This section gives an expression for the integral β -electron rate at the KATRIN detector.

As already mentioned, the response function (3.29) depends on the starting position of the electrons. To account for this, the WGTS can be thought of being divided into n slices of width $w = d/n$ and an averaged response function for the j th ($j \in \{0, 1, \dots, n-1\}$) slice can be given

$$R(E_S, qU, z_S) \rightarrow R_j(E_S, qU) = \int_{-d/2+jw}^{-d/2+(j+1)w} R(E_S, qU, z_S) dz_S. \quad (3.30)$$

As can be seen from equations (3.29) and (3.28), this averaging integral can be propagated to the scattering probabilities in the enumerator of the detailed transmission function.

The integral rate then reads [Kle+19]

$$\Gamma(qU) = \frac{1}{2} \sum_{j=0}^n N_{j,T} \cdot \int_{qU}^{E_0} \left(\frac{d\Gamma(E_S)}{dE_S} \right) \cdot R_j(E_S, qU) dE_S. \quad (3.31)$$

Here, the integral goes over all starting energies that enable electrons to overcome the retarding potential. The transmission probability vanishes for starting energies smaller than qU and the differential rate vanishes for energies above the β -spectrum endpoint E_0 (see equation 3.1), which yields the two integral limits. The sum goes over all slices of the WGTS. $N_{j,T}$ is the number of tritium nuclei in the j th slice of the WGTS. And the factor $1/2$ accounts for the fact, that on average only half the β electrons are emitted towards the detector.

3.4. Detector Counts

This section gives an expression for the electron counts measured by the KATRIN detector.

Therefore, the detector efficiency $\epsilon_{\text{det}} \in [0, 1]$ has to be taken into account. (For a description, its determination and value see section 2.7.) Furthermore, the background rate Γ_{bg} with a

nominal value of 10 cps [KAT05] has to be considered. Also a relative rate factor $A_S = 1$ between the background and the β -electron rate is introduced as it can be used in fitting procedures (see section 4.3.) Assuming a measurement time of $t(qU)$ attributed to a retarding energy qU , the detector counts are [Kle14]

$$N(qU) = t(qU) \cdot \epsilon_{\text{det}} \cdot (A_S \cdot \Gamma(qU) + \Gamma_{\text{bg}}) , \quad (3.32)$$

where $\Gamma(qU)$ denotes the integral rate (3.31).

3.5. Model Amendments

The outlined formalism that lead to the expression for the detector counts (3.32) forms a scaffold for the mathematical formalism that describes a KATRIN measurement. Modifications of isolated terms can incorporate further effects. For a comprehensive list, the reader is referred to [Kle+19]. Selected examples are listed below:

- **Doppler effect:** Gas flow and temperature move the tritium molecules and hence, smear the kinetic energy distribution of β electrons (see section 2.2). This can be modeled by convolving the differential rate (3.1) with a Maxwellian distribution or by applying corrections to the final energy states of the decaying molecules.
- **Plasma potential:** Space charges, respectively a plasma, forms within the WGTS due to the tritium decay (see section 2.2). β electrons may originate at higher/lower potentials due to space charges. This can be modeled by an adaption of the starting potential U_S in the transmission energy (3.10).
- **3-dimensional description:** This chapter focuses on a 1-dimensional formalism. However, as noted, input parameters such as the magnetic fields are not solely z -dependent, which requires a 3-dimensional approach and an incorporation of the segmentation of the detector. This can be accomplished by calculating the detector counts (3.32) separately for each detector pixel exploiting that the magnetic flux tube maps specific volumes of the WGTS onto specific areas of the analyzing plane and pixels.

3.6. Measurement Time Distribution

KATRIN measures electron counts as described in equation (3.32) at a set of retarding energies $\{qU_i\}$. How much measurement time $t(qU_i)$ is attributed to a certain retarding energy is specified in a measurement time distribution (MTD). The MTD influences the experiment's sensitivity to the neutrino mass. An optimal MTD balances the following aspects:

1. Some measurement time has to be attributed to retarding energies beyond the endpoint of the integral tritium- β spectrum to determine the background rate. The optimal duration depends on the background rate, but can generally take up a sizable fraction (of order 30 percent) of the overall measurement time. [KAT05; Kle14].
2. Near its endpoint, the shape of the integral tritium- β spectrum depends most strongly on the neutrino mass. Hence, most measurement time should be attributed to this region [KAT05; Kle14].
3. Retarding voltage bins deeper into the spectrum increase the count rate and hence, lower the statistical uncertainty due to Poisson statistics. These measurements mainly determine the endpoint from extrapolating the slope of the integral tritium- β spectrum [KAT05; Kle14].

4. The theoretical description of the integral tritium- β spectrum is optimized for the endpoint region. E.g. the molecular final states for β -electron energies 40 eV below the endpoint would need further investigation [Dos+06]. Hence, deeper scans introduce modeling uncertainties. However, it is expected that continuous modeling efforts decrease these uncertainties as needed.

The KATRIN Design Report [KAT05] suggests five 3-year-long MTDs for different measurement ranges $[E_0 - \alpha \text{ eV}, E_0 + 5 \text{ eV}]$ with $\alpha \in \{20, 25, 30, 40, 50\}$ and the conclusion that $\alpha = 30$ yields the best sensitivity to the neutrino mass.

As an energy-dependent effect is investigated within this thesis, it should be noted that scans beyond the 50 eV range have already been performed and may also be performed again in the future. E.g. searches for sterile neutrinos at the keV-scale would require deeper scans [Mer+19]. On top of that, within several measurement campaigns deeper scans were conducted: The First Tritium (FT) commissioning campaign successfully proved the apparatus functioning. The corresponding MTD covered a range starting at $\sim E_0 - 1.6 \text{ keV}$. The KATRIN neutrino mass measurement campaign 1 (KNM1) is being evaluated during the writing of this thesis. It set out to establish an unprecedented limit on the neutrino mass by β -decay measurements. Its MTD starts at $\sim E_0 - 90 \text{ eV}$, but the analysis range for neutrino mass inference remains still to be determined.

3.7. A Simulated KATRIN Neutrino Mass Measurement

In summary, a KATRIN measurement yields a set of electron counts $\{N(qU_i)\}$, equation (3.32), distributed over retarding voltage bins $\{qU_i\}$, where the counts fluctuate statistically [KAT05]. A possible model for the fluctuations is a Poissonian distribution [Kle14]. Figure 3.5 shows a KATRIN measurement for an MTD starting at $E_0 - 30 \text{ eV}$ and a total measurement time of three years. The distortion of the measured integral rate by a non-vanishing neutrino mass can be seen approximately 4 eV below the endpoint E_0 . This distortion can be used to infer the squared electron antineutrino mass from a KATRIN neutrino mass measurement. Chapter 4 presents corresponding statistical methods.

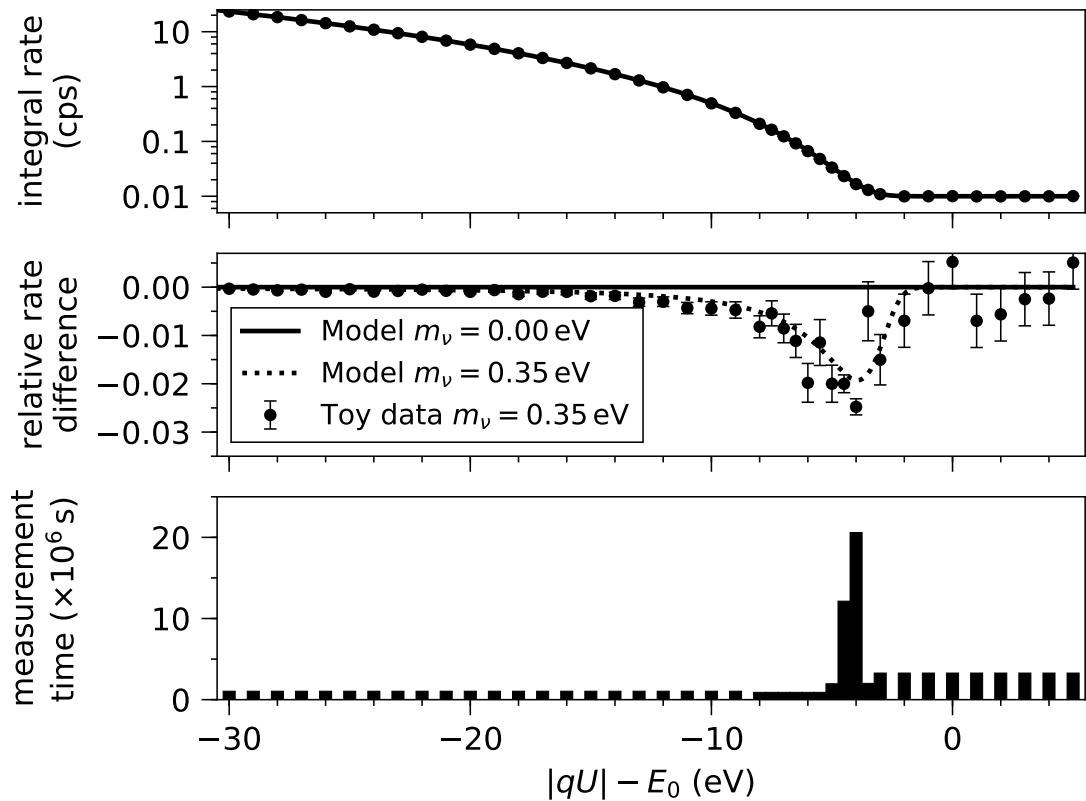


Figure 3.5.: Simulated KATRIN measurement for a non-vanishing neutrino mass. The total measurement time is three years. The top panel shows the measured integral rate Γ in dependence of the retarding energy. The center panel shows the relative rate difference for a non-vanishing neutrino mass $\Gamma(m_\nu = 0.35 \text{ eV})/\Gamma(m_\nu = 0 \text{ eV}) - 1$. The difference is $\sim 2\%$ at a retarding energy approximately 4 eV below the endpoint (simulated as $E_0 = 18575 \text{ eV}$). The bottom panel shows the MTD where most measurement time is attributed to the most sensitive region. This is also reflected by the uncertainty bars of the toy data. (Adapted from [Sei19].)

4. Statistical Methods and Neutrino Mass Inference at KATRIN

The best estimator for the neutrino mass m_ν alongside with an uncertainty or an upper limit will be retrieved by comparing the output of the KATRIN measurement with theoretical predictions within the process of parameter inference. This chapter reviews a selection of statistical approaches suitable in relation to the KATRIN experiment.

Section 4.1 outlines the principle of the maximum likelihood estimator (MLE). Section 4.2 relates the principle of the MLE to a KATRIN measurement and neutrino mass inference. Section 4.3 introduces the formalism of a nominal neutrino mass fit at KATRIN. Section 4.4 reviews the concept of uncertainty intervals and how confidence intervals can be extracted from the likelihood. Section 4.5 introduces the software framework that was used within this thesis. Section 4.6 relates the principle of uncertainty to neutrino mass inference and explains the origin of the often quoted 200 meV (90 % C.L.) KATRIN sensitivity.

4.1. Maximum Likelihood Estimation

The likelihood is the probability of a measurement outcome given a hypothesis. A hypothesis depending on a parameter vector $\boldsymbol{\theta}$ is called a composite hypothesis. A measurement outcome can be quantified by a vector of observed values \mathbf{x} . The probability P of \mathbf{x} given a hypothesis in dependence of $\boldsymbol{\theta}$ is called the likelihood function [Tan+18]

$$L(\boldsymbol{\theta}) = P(\mathbf{x} \mid \boldsymbol{\theta}). \quad (4.1)$$

If p denotes the probability for one observed value x_i in \mathbf{x} , then the likelihood function can be written as a product [Tan+18]

$$L(\boldsymbol{\theta}) = \prod_i p(x_i \mid \boldsymbol{\theta}). \quad (4.2)$$

The parameter vector $\hat{\boldsymbol{\theta}}$ that maximizes the likelihood function is called the maximum likelihood estimator (MLE) for the true values of $\boldsymbol{\theta}$.

4.2. The Likelihood of a KATRIN Measurement

The MLE-method can be applied to a KATRIN measurement as follows: The data vector is given by a set of n electron counts $\{N_{\text{obs},i}\}$ measured at different retarding potentials $\{qU_i\}$. The hypothesis is that these counts follow a Poisson distribution with predicted expected electron counts $\{N_{\text{theo},i}\}$ as per equation (3.32) [Kle14]. For sufficiently high counts (> 25 [Kle+19]) the Poisson distribution can be approximated by a Gaussian distribution $\mathcal{N}(x, \mu, \sigma)$ with mean $\mu = N_{\text{theo},i}(\boldsymbol{\theta})$ and standard deviation $\sigma = \sqrt{N_{\text{obs},i}}$. The likelihood function then reads [Kle14]

$$L(\boldsymbol{\theta}) = \prod_i^n \mathcal{N}\left(x = N_{\text{obs},i}, \mu = N_{\text{theo},i}(\boldsymbol{\theta}), \sigma = \sqrt{N_{\text{obs},i}}\right). \quad (4.3)$$

Commonly, instead of maximizing the likelihood function, its negative logarithm is minimized and a factor 2 is introduced [Tan+18]. This yields

$$-2 \ln L(\boldsymbol{\theta}) = \chi^2(\boldsymbol{\theta}) = \sum_i^n \left(\frac{N_{\text{obs},i} - N_{\text{theo},i}(\boldsymbol{\theta})}{\sqrt{N_{\text{obs},i}}} \right)^2 + \text{constants}. \quad (4.4)$$

The minimization of equation (4.4) yields the MLE estimator $\hat{\boldsymbol{\theta}}$ for $\boldsymbol{\theta}$.

Equation (4.4) is a sum of n standard normal distributed random variables. Hence, evaluated at the MLE, this chi-square expression $\chi^2(\hat{\boldsymbol{\theta}})$ follows the Pearson's chi-square statistic with $n - \dim \boldsymbol{\theta}$ degrees of freedom. Accordingly, the value $\chi^2(\hat{\boldsymbol{\theta}})$ is a measure for the goodness-of-fit [Tan+18]. In conclusion, equation (4.4) can be used for neutrino mass inference via the maximum likelihood method.

4.3. A Nominal KATRIN Neutrino-Mass Fit

In regard to a KATRIN neutrino mass measurement, the parameter of interest in the parameter vector $\boldsymbol{\theta}$ is the squared neutrino mass m_ν^2 . Furthermore, $\boldsymbol{\theta}$ typically comprises the endpoint of the tritium- β spectrum E_0 , equation (3.6), an overall normalization factor for the β -electron counts A_S and the background rate Γ_{bg} [Kle14; KAT05], that are treated as nuisance parameters. For the latter two see equation (3.32). Hence, in order to infer the neutrino mass, the four-dimensional likelihood has to be minimized. Following this procedure with simulated data enables the determination of KATRIN's sensitivity (see subsequent section 4.6).

4.4. Uncertainty Intervals

The presented maximum likelihood method (section 4.1) provides point estimates $\hat{\boldsymbol{\theta}}$. However, additional information can be provided by interval estimates. There are two main approaches to statistical inference, which may be called Bayesian and frequentist [Tan+18]. They differ in their interpretation of probability, which becomes especially evident by the interval estimates associated with the two approaches: Credible and confidence intervals. Both interval types can be given with reference to quantiles of the Gaussian distribution. E.g. the 1- and 2- σ levels define 68 % and 95 % intervals. The following two sections 4.4.1 and 4.4.2 explain the matter in more detail.

4.4.1. Bayesian Credible Intervals

The likelihood $L(\mathbf{x} | \boldsymbol{\theta})$ is a probability distribution for the data \mathbf{x} given the parameters $\boldsymbol{\theta}$. The likelihood can be transformed into a probability density for the parameters $\boldsymbol{\theta}$ by multiplication with a prior distribution $\pi(\boldsymbol{\theta})$ and normalization to one using Bayes theorem. One obtains the posterior distribution [Tan+18]

$$P(\boldsymbol{\theta} | \mathbf{x}) = \frac{L(\mathbf{x} | \boldsymbol{\theta})\pi(\boldsymbol{\theta})}{\int L(\mathbf{x} | \boldsymbol{\theta}')\pi(\boldsymbol{\theta}') d\boldsymbol{\theta}'}. \quad (4.5)$$

Credible regions, in which the true parameters lie with a certain probability can be extracted. When $\boldsymbol{\theta}$ is one dimensional a credible region is also called a credible interval.

4.4.2. Frequentist Confidence Intervals

This section first gives definitions for the terms “confidence interval”, “coverage probability” and “confidence level”. It is explained, how confidence intervals can be extracted from a likelihood on the basis of these definitions. This approach was applied within the scope of

this thesis in order to study the impact of model uncertainties on KATRIN’s sensitivity in chapter 6.

In frequentist statistics, probability is interpreted as the frequency of the outcome of a repeatable experiment. The boundary of a confidence region is given by a function of the data. There is some freedom of choice for the corresponding function. It should be noted that in this sense, the term confidence region is somewhat “unqualified” [Tan+18]. But it obtains a deeper meaning in combination with a coverage probability. First, it should be noted that the boundary of the confidence region would fluctuate if one were to repeat the experiment many times. One would obtain an ensemble of confidence regions. The coverage probability α refers to the fraction of regions in such an ensemble that contains the true parameter values $\boldsymbol{\theta}_T$ [Tan+18]. If an ensemble of confidence regions covers the true parameter values $\boldsymbol{\theta}_T$ at least a fraction of α times, the confidence interval is understood to have a confidence level of α [Tan+18]. When $\boldsymbol{\theta}$ is one dimensional a corresponding confidence region is called a confidence interval.

In a practical context, a prescription is required on how to construct confidence intervals of a eligible confidence level. The Neyman construction [Ney37] or the unified approach by Feldman and Cousins [FC98] are such prescriptions.

A further method of constructing confidence intervals is to consider a test (see hypothesis testing in [Tan+18]) of the hypothesis that the parameter values $\boldsymbol{\theta}$ have the true values $\boldsymbol{\theta}_T$ [Tan+18]. In this construction the choice of test to be used is free. One possibility is a test statistic based on the likelihood ratio between the MLE $\hat{\boldsymbol{\theta}}$ and $\boldsymbol{\theta}$ [Tan+18]

$$\lambda(\boldsymbol{\theta}) = \frac{L(\boldsymbol{\theta})}{L(\hat{\boldsymbol{\theta}})}. \quad (4.6)$$

In the case of a construction via a hypothesis test, all parameter values $\boldsymbol{\theta}$ are excluded from the confidence interval of level α that are rejected by the test with a significance of α [Tan+18].

If the likelihood follows the form of a multivariate Gaussian distribution in $\boldsymbol{\theta}$, then the above test statistic (4.6) can be evaluated and the hyper surface defined by

$$\ln L(\boldsymbol{\theta}) = \ln L(\hat{\boldsymbol{\theta}}) - \frac{s^2}{2} \quad (4.7)$$

encloses a $s\sigma$ confidence region for $\boldsymbol{\theta}$ [Tan+18]. (Here, $s\sigma$ denotes the corresponding quantile of a Gaussian distribution.)

It should be noted that the extraction of confidence intervals from the KATRIN likelihood requires its extrapolation to nonphysical negative squared neutrino masses [Kle14] - a complication that can be avoided when using Bayesian methods.

In conclusion, in this section a constructive approach has been presented, that enables the extraction of confidence regions for the parameters of a KATRIN measurement from the KATRIN likelihood. An extension of this formalism (namely, the profile likelihood method) finds application in the scope of this thesis in chapter 6.

4.5. The KaFit and SSC Software Frameworks

With respect to neutrino mass inference at KATRIN two formalisms have been presented: the model for a KATRIN neutrino mass measurement in chapter 3 and a statistical framework for parameter inference within the current chapter 4.

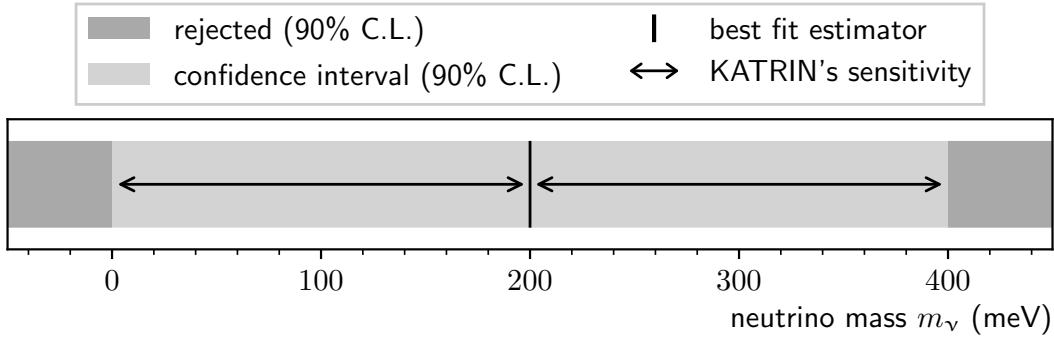


Figure 4.1.: Illustration of KATRIN’s sensitivity to the neutrino mass. The graph illustrates KATRIN’s sensitivity of 200 meV [KAT05] as per equation (4.8). It shows a hypothetical measurements of to the neutrino mass with a symmetric confidence interval (90 % C.L.) centrally located around the best fit estimator. If such a classical confidence interval is constructed, KATRIN can reject the null hypothesis of a vanishing neutrino mass at 90 % C.L. if it estimates a neutrino mass of at least 200 meV.

The two formalisms are implemented within two modules of the “KATRIN Analysis and Simulations Package” (KASPER) [Soft-KAS]:

1. The **source and spectrum calculation (SSC)** [Soft-SSC] module implements the formulas for the differential and integrated spectrum calculations. Therefore, it follows the formulas given in chapter 3. Additionally, it also includes aspects beyond the given description, such as the gas dynamics within the WGTS [Höt12; Gro15; Kle+19; Käf12].
2. The **KaFit** [Soft-KF] module translates the β spectrum calculated by SSC into expected detector counts. Furthermore, KaFit implements several statistic tools tailored to the KATRIN experiment. One of which is the extraction of confidence intervals according to the profile likelihood method (see section 4.4.2), which is of importance in the scope of this thesis. The actual minimization and profiling are done by the interfaced MINUIT2 and MINOS package from the ROOT¹ [Ant+09] analysis framework [Kle14].

Both packages were used and extended to allow for the analysis done within the scope of this thesis as will be explained in chapters 5 and 6.

4.6. KATRIN’s Sensitivity to the Electron Antineutrino Mass

This section explains the origin of the often quoted KATRIN design sensitivity to the neutrino mass of 200 meV (90 % C.L.). First, a definition of the sensitivity is given in section 4.6.1. Then, section 4.6.2 and 4.6.3 show how this definition was applied by different works to deduce KATRIN’s sensitivity.

4.6.1. Definition and Construction of KATRIN’s Sensitivity

KATRIN’s sensitivity can be understood as half the width of a symmetric and central confidence interval (90 % C.L.) for the neutrino mass obtained from a KATRIN neutrino mass measurement. As such it can be constructed from a $1-\sigma$ uncertainty $\sigma(m_\nu^2)$ on the squared neutrino mass [KAT05]

$$S_{m_\nu}(90\%) = \sqrt{1.645 \cdot \sigma(m_\nu^2)}, \quad (4.8)$$

¹<http://root.cern.ch/>

Table 4.1.: KATRIN’s sensitivity to the neutrino mass from ensemble tests.
The table lists KATRIN’s sensitivity S_{m_ν} (90 %) as defined by equation (4.8). Several works reevaluated the statistical uncertainty according to experimental and theoretical progress. For each reevaluation a systematic uncertainty of $\sigma_{\text{sys}}(m_\nu^2) = 0.017 \text{ eV}^2$ was assumed. A value derived from an ensemble test is a random variable. Corresponding uncertainty intervals are reprinted where originally stated.

$\sigma_{\text{stat}}(m_\nu^2)$ (eV ²)	S_{m_ν} (90 %) (meV)	comment	reference
0.018	200	design value	[KAT05]
0.0165 ± 0.0001	198	updated β -spectrum calculation	[Höt12]
0.0162 ± 0.0001	197	further updated β -spectrum calculation	[Kle14]
0.01490	193	optimized MTD	[Kle14]

where the factor 1.645 translates the 68.3 % interval into a 90 % interval.

In other words, KATRIN’s sensitivity to the neutrino mass can be understood as the minimal neutrino mass that has to be inferred from a KATRIN neutrino mass measurement to exclude the null hypothesis of a vanishing neutrino mass [Kle14] when constructing a symmetric and central confidence interval (90 % C.L.). Figure 4.1 illustrates this statement.

For a more comprehensive picture, where not only a symmetric and central confidence interval is considered, but also the unified approach according to Feldmann and Cousins as well as Bayesian statistics, the reader is referred to [Kle+19].

4.6.2. Sensitivity from Simulated Ensembles

In the KATRIN Design Report the sensitivity to the neutrino mass was evaluated using ensemble tests. An ensemble of many KATRIN measurements was simulated (see section 3.7 on how such a simulation can be conducted) with a true neutrino mass of 0 eV. From each simulated measurement the squared neutrino mass was inferred in a standard KATRIN four-parameter fit (see section 4.3). Probability was interpreted as the frequency of an outcome of such a fit. The central 1- σ interval of the obtained ensemble of squared neutrino masses was taken as the statistical uncertainty on the squared neutrino mass [KAT05]

$$\sigma_{\text{stat}}^{\text{TDR}}(m_\nu^2) = 0.018 \text{ eV}^2 \quad (4.9)$$

A systematic uncertainty was estimated to be approximately 0.01 eV. Due to the early stage of the experiment the systematic uncertainty was conservatively enlarged to a systematic budget at approximately the same scale as the statistical uncertainty [KAT05]

$$\sigma_{\text{sys}}^{\text{TDR}}(m_\nu^2) = 0.017 \text{ eV}^2 \quad (4.10)$$

Adding the statistic and systematic uncertainty quadratically and applying the definition (4.8) yields KATRIN’s design sensitivity [KAT05]

$$S_{m_\nu}^{\text{TDR}}(90 \%) = \sqrt{1.645 \cdot \sqrt{\sigma_{\text{stat}}^{\text{TDR}}(m_\nu^2)^2 + \sigma_{\text{sys}}^{\text{TDR}}(m_\nu^2)^2}} \approx 200 \text{ meV}. \quad (4.11)$$

The corresponding investigations were redone in the scope of several works. Table 4.1 lists selected results.

4.6.3. Sensitivity from the Profile Likelihood Method

In the scope of this thesis, the sensitivity on the neutrino mass as obtained by the profile likelihood method is of importance. (See chapter 6 for a description and application of the profile likelihood method.) In this regard, previous results from [Kle14] based on the

profile likelihood method are shortly reviewed here. Two uncertainties were obtained for two different MTDs in a KATRIN standard 4-parameter fit (see section 4.3). The first MTD had been specially optimized with regard to KATRIN’s sensitivity and the resulting statistical uncertainty is $\sigma_{\text{stat}}(m_\nu^2) = 0.014\,94 \text{ eV}^2$. For the second result the nominal MTD from the KATRIN Design Report was used as introduced in section 3.6. The corresponding profile likelihood was plotted and $\sigma_{\text{stat}}(m_\nu^2)$ can be extracted to be between 0.0155 eV^2 – 0.0165 eV^2 . Both results are in agreement on the 10^{-3} level with the results from ensemble tests (see [Kle14] in table 4.1). This is an indicator for the general validity of the profile likelihood method in the context of a KATRIN measurement. A more detailed discussion on the matter is given in chapter 6.

5. The Energy-Dependence of Inelastic Electron Scattering within the KATRIN WGTS

[Clean up!] A β electron's probability to scatter when traveling through the STS depends on the scattering cross section σ (??). The cross section σ in turn depends on the kinetic energy E_{kin} of the electron: $\sigma \equiv \sigma(E_{\text{kin}})$. This dependence has been neglected in the modeling of the electron counts in section ??.

To Do

5.1. Cross Section of Tritium- β Electrons Scattering from Hydrogen Isotopologues

In a KATRIN measurement there exists a minimum retarding energy qU_{\min} and only electrons with a kinetic energy greater than qU_{\min} can reach the detector. A scattering cross section averaged over the kinetic energy of the electrons reaching the detector was assumed

$$\sigma = \sigma_{\text{avg}} = \frac{1}{\Delta E_{\text{kin}}} \int_{qU_{\min}}^{qU_{\min} + \Delta E_{\text{kin}}} \sigma(E_{\text{kin}}) dE_{\text{kin}} \quad \text{with} \quad \Delta E_{\text{kin}} = E_0^{\text{eff}} - qU_{\min}. \quad (5.1)$$

Here, the “effective endpoint” E_0^{eff} denotes the highest kinetic energy of β electrons reaching the detector. It does not necessarily match the endpoint of the tritium β spectrum as experiment specific effects might shift it. $E_0^{\text{eff}} = 18.575 \text{ eV}$ is assumed. Instead of an average cross section an energy dependent formula can be used. Incorporating the energy dependence makes the model more complicated and slower to compute; neglecting it may lead to wrong results. More light will be shed on both aspects within this chapter.

The total scattering cross section is a sum of the cross section for elastic and inelastic scattering.

$$\sigma(E_{\text{kin}}) = \sigma_{\text{el}} + \sigma_{\text{inel}}(E_{\text{kin}}). \quad (5.2)$$

Only the energy dependence of the inelastic scattering cross section will be considered. This is justified for now as according to [KAT05] the cross section for elastic scattering is about an order of magnitude smaller than the one for inelastic scattering. An expression for the inelastic cross section for electrons scattering from hydrogen molecules can be found in [Liu73]. Two expressions are given, one for relativistic incident particles and one for non-relativistic incident particles. For the maximum relativistic β factor of β electrons from tritium decay one finds

$$\beta(E_{\text{kin}}, m) = \sqrt{1 - \frac{1}{(\frac{E_{\text{kin}}}{m} + 1)^2}} \quad (5.3)$$

$$\beta_{\max} = \beta(E_0^{\text{eff}}) \approx 18.6 \text{ keV}, m_e \approx 511 \text{ keV} \approx 0.26$$

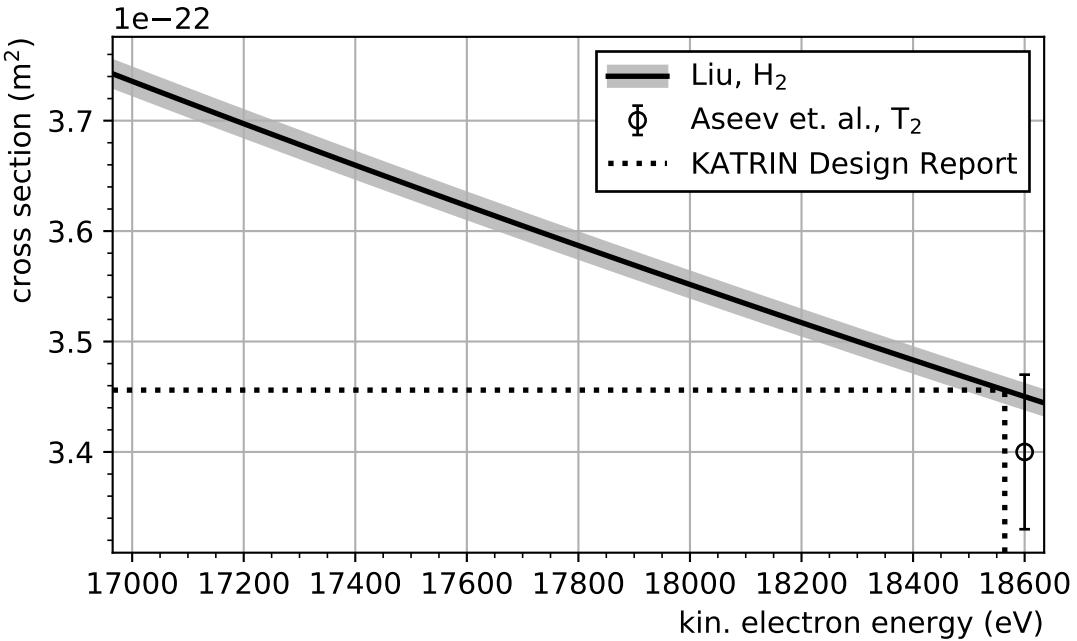


Figure 5.1.: Inelastic cross section for non-relativistic electrons scattering from molecular hydrogen isotopologues. Shown is the theoretical calculation by Liu along with its uncertainty [Liu73], the measurement by Aseev et. al. at the Troitsk experiment [Ase+00] and the value stated in the KATRIN design report [KAT05]. (The values are also listed in table 5.1.) The depicted energy interval matches the measurement interval of the First Tritium (FT) campaign. **FIX et al.**

Traveling at approximately a forth of the speed of light, the β electrons are assumed to be non-relativistic. Then, the given expression for the energy dependent cross section is

$$\sigma(E_{\text{kin}}) = (4\pi a_0^2) \cdot \left(\frac{E_{\text{kin}}}{R} \right)^{-1} \cdot \left[C_1 \cdot \ln \left(\frac{E_{\text{kin}}}{R} \right) + C_2 \right] \quad (5.4)$$

with the Bohr radius a_0 , the Rydberg energy R and two constants given as

$$C_1 = 1.5487 \quad \text{and} \quad C_2 = 2.2212 \pm 0.0434. \quad (5.5)$$

Note that in other works $C_2 = 2.4036$ [Liu87] and $C_2 = 1.53$ [Ger75] are given. The value in (5.5) from [Liu73] was chosen to enable comparability with the KATRIN Design Report [KAT05].

Furthermore, the total inelastic scattering cross section was measured at the Troitsk experiment and the KATRIN Design Report states a reference value. The values are listed in table 5.1. The reference value matches the theoretical calculation using (5.4) at a kinetic energy of $E_{\text{kin}} \approx 18564.4$ eV which would be the center of a $\Delta E_{\text{kin}} = 20$ eV KATRIN measurement interval. Note that in the KATRIN reference measurement interval $[E_0^{\text{eff}} - 30 \text{ eV}, E_0^{\text{eff}}]$ the cross section varies $\sim 0.14\%$. This variation is below its theoretical uncertainty given by (5.5). In the measurement interval of the FT campaign it varies $\sim 8\%$.

5.2. Motivation and Significance for the KATRIN Experiment

The motivation to implement an energy dependent cross section into the source and spectrum calculation (SSC) framework is two-fold:

Table 5.1.: Inelastic cross section for electrons scattering off molecular hydrogen isotopologues for different incident energies. Listed are important values with reference to KATRIN. Except for the measured value, they are obtained using (5.4). The values are given relative to an assumed endpoint of the measured spectrum $E_0^{\text{eff}} = 18\,575 \text{ eV}$.

kin. energy	cross section (10^{-22} m)	Note
$E_0^{\text{eff}} - 1600 \text{ eV}$	3.740	largest range in FT campaign
$E_0^{\text{eff}} - 90 \text{ eV}$	3.469	largest range in KNM1 campaign
$E_0^{\text{eff}} - 30 \text{ eV}$	3.459	KATRIN reference measurement interval [KAT05]
$E_0^{\text{eff}} - 10 \text{ eV}$	3.456	KATRIN reference value [KAT05]
E_0^{eff}	3.454	endpoint of tritium β spectrum
18 600 eV	3.40 ± 0.07	measured at the Troitsk experiment [Ase+00]

Deep scans: According to [Gro15] for a neutrino mass determination with the precision goals of KATRIN, the inelastic scattering cross section has to be known with an upper uncertainty of $0.0055 \times 10^{-22} \text{ m}$ (0.16 %). This requirement might be surpassed by neglecting that the scattering cross section is not constant, but varies with energy. According to table 5.1 the requirement is fulfilled in a measurement interval of $[E_0^{\text{eff}} - 30 \text{ eV}, E_0^{\text{eff}}]$. According to (5.4) it is violated if the lower bound of the measurement interval is below 18 531 eV. Scans deeper into the tritium β spectrum increase the count rates and hence, improve the statistic uncertainty on the neutrino mass. Also, for the search of sterile neutrinos deeper scans are necessary. On top of that, deeper scans have already been performed e.g. in the FT campaign and help to establish an even better understanding of the KATRIN apparatus. In these cases the energy dependence is not negligible.

Comparability: A possible cross-check for software is its comparison to other software that was developed independently. SSC is part of a data fitting suit. Other fitting software exists within the KATRIN collaboration that uses an energy dependent scattering cross section, e.g. FITRIUM [Kar+19]. Furthermore, SSC is commonly cross-checked against a Monte-Carlo particle tracking software called KASSIOPEIA [KATRINCOL2019] which also implements an energy dependent scattering cross section.

5.3. An Energy-Dependent Scattering Model within the KATRIN Formalism

5.3.1. Model Description

Table 5.2.: The scattering probabilities averaged over all starting positions and starting angles in the WGTS. Both the values from a Monte Carlo (MC) simulation and the values according to (??) are given. The later ones can also be found in [Gro15; Kle14] and in the code of the SSC framework [KATRINCOL2019].

scattering count	MC particle tracking [Gro15]	equation (??)
0	0.415 ± 0.002	0.41334
1	0.292 ± 0.002	0.29266
2	0.166 ± 0.001	0.16733
3	0.079 ± 0.001	0.07913
4	0.031 ± 0.001	0.03178

The energy-dependence of the scattering cross section enters into the calculation of the scattering probabilities (??). In the derivation the dependence on the starting energy E_S

was neglected. Instead an average starting energy and hence, an average scattering cross section $\sigma(18\,564.4\,\text{eV}) = 3.456 \times 10^{-22}\,\text{m}^2$ was assumed. For this case table 5.2 lists the scattering probabilities averaged over all starting positions and pitch angles in the WGTS. In the energy dependent case the scattering cross section has to be adapted to the starting energy of the β electrons. In other words, the scattering probabilities follow a Poisson distribution averaged over starting positions and pitch angles of β electrons. Analogously to the derivation in section ?? they read:

$$\mu(E_S, z_S, \theta_S) = \frac{\sigma(E_S)}{\cos \theta_S} \int_{z_S}^{L/2} \rho(z) dz \quad (5.6a)$$

$$P_l(E_S, z_S, \theta_S) = \text{Poisson}(\mu(E_S, z_S, \theta_S), l) \quad (5.6b)$$

$$\bar{P}_l(E_S) = \frac{1}{L \cdot (1 - \cos(\theta_{\max}))} \int_{-L/2}^{L/2} \int_0^{\theta_{\max}} \sin(\theta_S) \text{Poisson}(\mu(E_S, z_S, \theta_S), l) d\theta_S dz_S. \quad (5.6c)$$

Here, $\bar{P}_l(E_S)$ in the final equation (5.6c) denotes the probability of l -fold scattering for a β electron with a starting energy E_S averaged over all starting positions and pitch angles. This model will be denoted Poisson model. It is expected to be accurate for the probability of no scattering $\bar{P}_0(E_S)$.

But, depending on the required accuracy, for 1 or more scatterings this Poisson model does not necessarily hold. A scattering electron loses energy. Due to a lower energy the scattering cross section increases and the electron becomes more likely to scatter again. In other words, the probability of individual scattering processes are no longer independent. This violates one of the preconditions to model the scattering probabilities via a Poisson distribution. Another model is suggested. It assumes a constant gas density ρ , a fixed energy loss of ϵ per scattering and a source length of L :

$$\mu(E, \theta_S) = \frac{\sigma(E)\rho L}{\cos \theta_S} \quad (5.7a)$$

$$p_0(E, \theta_S, n) = \left(1 - \frac{\mu(E, \theta_S)}{N}\right)^n \quad (5.7b)$$

$$p_l(E, \theta_S, n) = \sum_{k=l}^n p_{l-1}(E, \theta_S, k-1) (1 - p_0(E - (l-1)\epsilon, \theta_S, 1)) p_0(E - l\epsilon, \theta_S, n-k) \quad (5.7c)$$

$$\bar{p}_l(E, \theta_S) = \frac{1}{L} \int_0^L p_l(E, \theta_S, \left\lceil N \frac{z_S}{L} \right\rceil) dz_S \quad (5.7d)$$

$$P_l^*(E_S, \theta_S) = \lim_{N \rightarrow \infty} \bar{p}_l(E_S, \theta_S) \quad (5.7e)$$

$$\bar{P}_l^*(E_S) = \frac{1}{1 - \cos \theta_{\max}} \int_0^{\theta_{\max}} P_l^*(E_S, \theta_S) d\theta_S. \quad (5.7f)$$

Here, $\bar{P}_l^*(E_S)$ in the final equation (5.7f) denotes the probability of l -fold scattering for a β electron with a starting energy E_S averaged over all starting positions and pitch angles assuming a fixed energy loss ϵ per scattering. This model will be denoted extended model. A descriptive derivation is given in appendix A.1.

The extended model is evaluated numerically as it includes one limit and two integrals. A balance between the numerical accuracy and the evaluation run time had to be found. Details can be found in the appendix A.2. The probability for 1-fold scattering could be calculated with an accuracy on the 10^{-5} level and for 2-fold scattering with an accuracy on the 10^{-3} level.

The suggested extended model is only valid for a fixed energy loss per scattering ϵ . The value $\epsilon = 14.1$ eV was chosen as it is the most probable energy loss for electrons travelling through tritium gas according to [Ase+00]. A more accurate description would use the energy loss probability density. This has not yet been considered.

Figure 5.2 shows the Poisson model along with the suggested extended model. The results are discussed in the following paragraphs.

5.3.2. Model Discussion

Model compatibility

For 1-fold scattering the difference between the Poisson and the extended model is on the 10^{-4} level. For 2-fold scattering the accuracy of the numerical evaluation of the extended model is not yet sufficient to distinguish it from the Poisson model. Table 5.2 lists the scattering probabilities for an energy independent Poisson model and a reference cross section $\sigma(18\,564.4\text{ eV}) = 3.456 \times 10^{-22}\text{ m}$ given in the KATRIN Design Report [KAT05]. It also lists the outcome of a Monte Carlo simulation taken from [Gro15]. As expected the energy dependent Poisson model recovers the energy independent model exactly at the corresponding energy of 18 564.4 eV. Furthermore, within the endpoint region of the tritium β spectrum all models match the Monte Carlo simulation within its uncertainty.

Trend of energy-dependence

For a decreasing starting energy of β electrons the probability for no and 1-fold scattering also decreases while the probability for 2-fold scattering increases. This is expected for the following reasons. For a starting energy of $E_S = 18\,564.4\text{ eV}$ the expected amount of scatterings μ (5.6a) averaged over all starting positions z_S and pitch angles θ_S is

$$\bar{\mu} = \frac{1}{L \cdot (1 - \cos(\theta_{\max}))} \int_{-L/2}^{L/2} \int_0^{\theta_{\max}} \sin(\theta_S) \mu(E_S, z_S, \theta_S) d\theta_S dz_S = 1.077. \quad (5.8)$$

In other words, β electrons with a starting energy near the endpoint of the tritium β spectrum are expected to scatter around $\bar{\mu} = 1.077$ times when travelling through the WGTS. If the starting energy decreases and hence, scattering becomes more likely it becomes more likely to scatter more than $\bar{\mu}$ times and less likely to scatter less than $\bar{\mu}$ times. Thus, the probability for no and 1-fold scattering decreases along with the starting energy of the β electrons. At the same time the probability for more than 1 scatterings increases for lower starting energies. This intuitive reasoning is also reflected by deriving the scattering probabilities (5.6b) for the starting energy where μ denotes the expected scattering count (5.6a)

$$\frac{dP_l(E_S, z_S, \theta_S)}{dE_S} = \underbrace{\text{Poisson}(\mu(E_S, z_S, \theta_S), l)}_{\geq 0} \cdot \underbrace{\left(\frac{l}{\mu(E_S, z_S, \theta_S)} - 1 \right)}_{(*)} \cdot \underbrace{\frac{\int_{z_S}^L \rho(z) dz}{L \cos \theta_S}}_{>0} \cdot \underbrace{\frac{d\sigma(E_S)}{dE_S}}_{<0}. \quad (5.9)$$

The sign of the derivative is determined by the sign of $(*)$ which follows approximately the above reasoning being negative for $l < \bar{\mu}$ and positive for $l > \bar{\mu}$.

5.3.3. Model Implementation

5.3.4. Performance and Accuracy

The suggested extended model (5.7) at its current stage is computationally too expensive to be used in a fitting procedure. Note that a similar model has been derived for a fixed change of the pitch angle θ per scattering in [Gro15]. This model for angular changes is implemented in the SSC software and can be used in fitting procedures. Though, an important difference is that the determination of the count rate requires an integral over the starting energy E_S in (??). Hence, the energy dependent scattering probabilities have to be recomputed in every step of the numerical integration. This is not the case for the model considering the angular changes.

Albeit the energy dependent Poisson model might not fully hold in the case of an energy dependent cross section, figure 5.2 shows that for larger measurement intervals it is more accurate than assuming energy independent scattering probabilities. Hence, it is a reasonable choice to implement the energy dependent Poisson model into the SSC software. This was done. As mentioned the energy dependent scattering probabilities have to be recomputed in every step of the numerical integration in (??) over the β electron energy. The impact on the run time was probed for 5 different MTDs. They differ in the amount and range of retarding potentials. Both aspects should influence the run time. The run time should be approximately linear in the amount of retarding potentials. The run time should get longer the wider the range of the retarding potentials is as the numerically evaluated integral over energies in (??) then stretches over a wider range. Figure 5.3 shows that the run times increase by a factor of approximately 40 – 120 between assuming a constant cross section and an energy dependent one.

5.4. Effect on the Inferred Neutrino Mass

5.4.1. Neutrino Mass Shift within a Standard KATRIN Measurement Interval

- ToDo** **[Numbers can not be reconciled with shifting the cross section. What to do about it?]** In order to determine the influence of the energy dependence of the scattering cross section on the neutrino mass determination the so-called “shift method” was used. A description can e.g. be found in [Sei19]. **[Write own description]** In short, an integrated rate using the energy dependent Poisson model for the scattering probabilities was simulated and fitted to a model using the constant Poisson model. One obtains a deviation of the fitted and simulated neutrino mass, the neutrino mass shift. The procedure was repeated for multiple measurement intervals using the settings from the KATRIN Design Report [KAT05] for a 3-year measurement. A comparison of the statistical uncertainty on the neutrino mass and its shift is shown in figure 5.4 depending on the measurement interval. For a measurement interval of approximately 35 eV the shift and the statistical uncertainty are approximately equal

$$\sigma(m_\nu) \approx \sqrt{0.015\text{eV}} = 122\text{ meV (68 \% C.L.)} . \quad (5.10)$$

5.4.2. Neutrino Mass Shift in an Extended Measurement Interval

In spring 2019 the so-called KNM1 campaign started. The measurement interval was 90 eV. The MTD had been optimized according to the up-to-date status of the KATRIN experiment which did not match the KATRIN Design Report in all aspects. The study on the neutrino mass shift was redone scaling the MTD of the KNM1 campaign to 3 years and using the corresponding KNM1-settings. The neutrino mass shift then would be 297 meV.

5.5. Conclusion and Outlook

The energy dependence of the scattering cross section enters into the calculation of the scattering probabilities. An accurate modelling of the energy dependent scattering probabilities is challenging due to performance reasons, but modelling them according to a Poisson distribution is possible. It was shown that the difference between the Poisson model and a more accurate model for 1-fold scattering is on the 10^{-4} level. The cases for more than 1 scatterings need further investigation. Also a fixed energy loss per scattering was assumed instead of a energy loss probability distribution. Future work might consider these aspects and what influence a more accurate modeling on the scattering probabilities has on the neutrino mass determination.

Given the KATRIN uncertainty budget, when modeling the energy dependent scattering probabilities via a Poisson distribution, the energy dependence of the scattering cross section is not negligible for measurement intervals that extend more than 35 eV below the endpoint of the tritium β spectrum.

Including the energy dependence in the analysis increases the run time of the fitting procedure significantly. Future work might consider to precalculate the scattering probabilities for different fixed energies and use interpolation techniques for energies in-between the fixed ones.

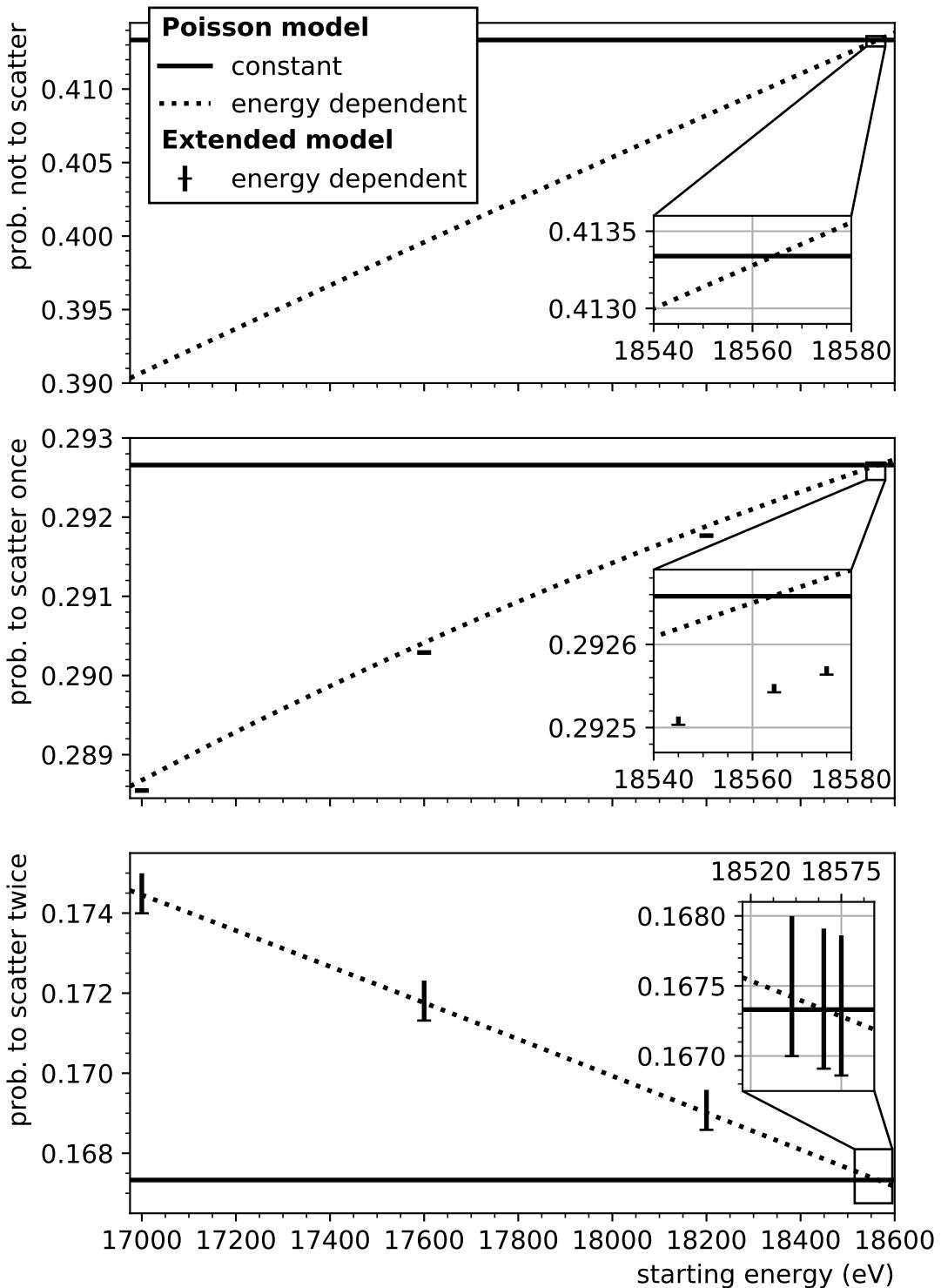


Figure 5.2.: Energy-dependent probability of electron scattering within the WGTS. Shown are from top to bottom the probability for no, 1-fold and 2-fold scattering averaged over all starting positions and starting pitch angles of β electrons. Depending on the required accuracy, for 1-fold and 2-fold scattering the Poisson model does not hold. The markers show an extended model as described in the text. Its numerical evaluation suffers from a one-sided uncertainty (of $\sim 10^{-5}$ for 1-fold and $\sim 10^{-3}$ for 2-fold scattering) depicted as uncertainty bars. The inset shows an energy span around the endpoint of the tritium β spectrum. The full energy range matches the measurement interval of the First Tritium (FT) campaign.

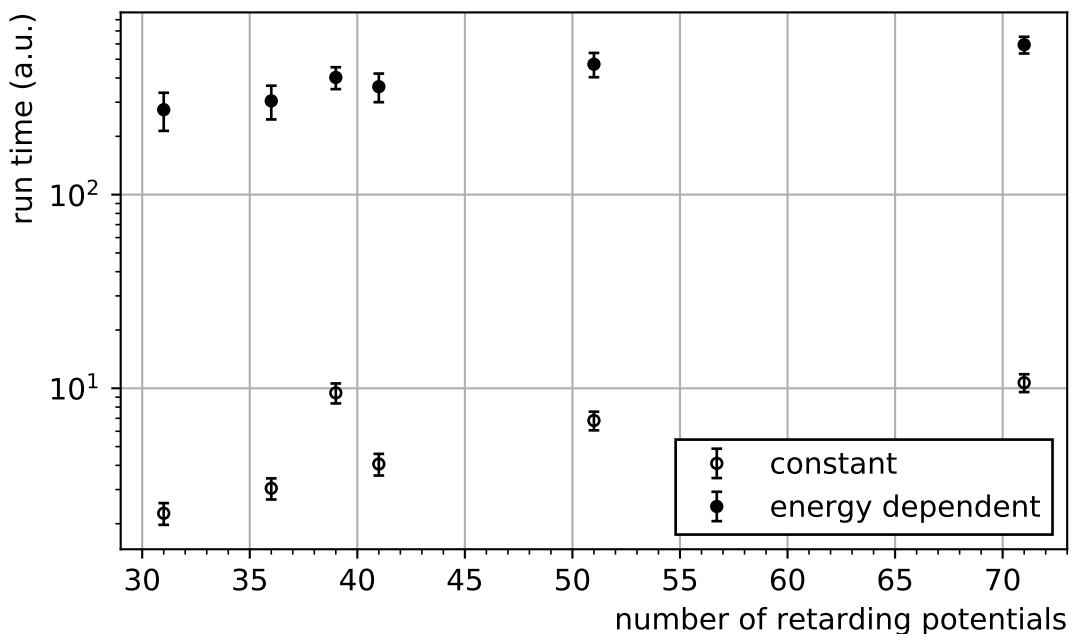


Figure 5.3.: Run-time comparison between using an energy-dependent and -independent inelastic scattering cross section. Shown are the run times for a fit of a simulated model to itself for different MTDs. The MTDs for the ranges 20, 25, 30, 40 and 50 eV are taken from the KATRIN design report [KAT05] and have a rising number of retarding potentials. A further 90 eV-wide MTD of the KNM1 campaign was used. It contains 39 retarding potentials. The corresponding run time lies above the trend line as the width of the measurement interval is wider. For each MTD 30 fits were performed. The run times were clocked for a constant and an energy-dependent cross section.

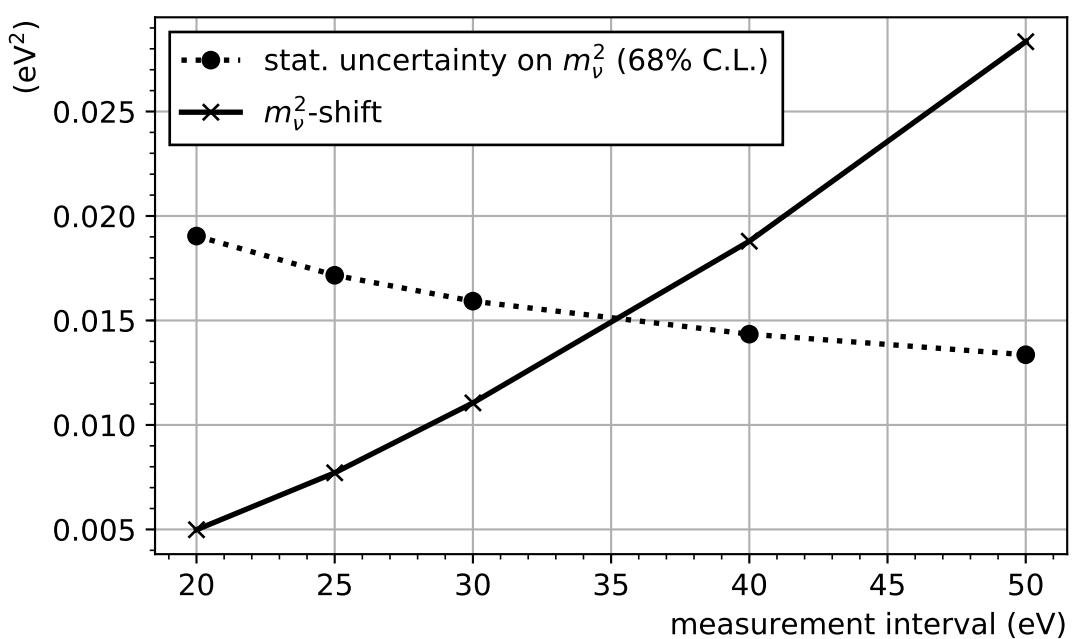


Figure 5.4.: Neutrino mass shift due to an energy dependent inelastic scattering cross section as calculated by SSC for 5 measurement intervals. The configuration for the calculation, especially the measurement time distribution (MTD), follows the KATRIN Design Report [KAT05]. For comparison the statistical uncertainty is plotted as well. It is derived from the profile likelihood. The lines between the 5 markers are linear interpolations.

6. Impact on KATRIN’s Sensitivity by an Energy Loss Model for Inelastically-Scattering Electrons Derived from KATRIN Data

A quantitative accurate description of the scattering processes of β electrons within KATRIN’s gaseous tritium source is of crucial importance for the neutrino-mass-sensitivity goal. In modeling the corresponding effects the energy loss function as described in section 3.2.5 plays an important role. With reference to the described energy loss function, the KATRIN Design Report states, that its precision is not sufficient for the KATRIN-sensitivity goal. It was planned to deduce a sufficiently accurate model from data taken at KATRIN [KAT05]. In that regard, a preliminary model has successfully been established for electrons scattering from deuterium molecules based on data taken in October 2018 by a dedicated subgroup of the KATRIN collaboration. This preliminary energy loss model, from here forth called the “KATRIN model”, has improved uncertainties with respect to the model from literature (section 3.2.5), from here forth called the “Aseev model” after the primary author of the corresponding publication [Ase+00]. Within the scope of this thesis the impact on KATRIN’s sensitivity by the exchange of the Aseev model for the KATRIN model was studied. Therefore, this chapter is structured as follows: Section 6.1 presents the general idea of this study. Section 6.2 outlines the KATRIN model. Section 6.3 discusses the scope of the validity of this study, for example what is expected of the comparison of a model for electrons scattering off deuterium and another model for scattering off tritium. Section 6.4 introduces the necessary statistical tools. Section 6.5 lists and discusses the results. And section 6.6 concludes and offers an outlook.

6.1. Introduction and Motivation

For the conducted sensitivity study KATRIN neutrino mass measurements were simulated assuming a neutrino mass of 0 eV and using the KATRIN model. A confidence interval was deduced using the profile likelihood method as described in a subsequent section 6.4.2. This enables the treatment of the model uncertainties and their correlations as nuisance parameters. The implementation of the required statistical tools into the available software frameworks is of general interest as the approach taken in the scope of this thesis may be applied to other model uncertainties, not only the ones stemming from the energy loss function.

6.2. The Empirical KATRIN Energy Loss Model

The presented KATRIN energy loss model is a phenomenological description fitted to data taken at the KATRIN experiment in October 2018. It was established by a dedicated subgroup of the KATRIN collaboration. The model is outlined in the following with an emphasis on its statistical properties.

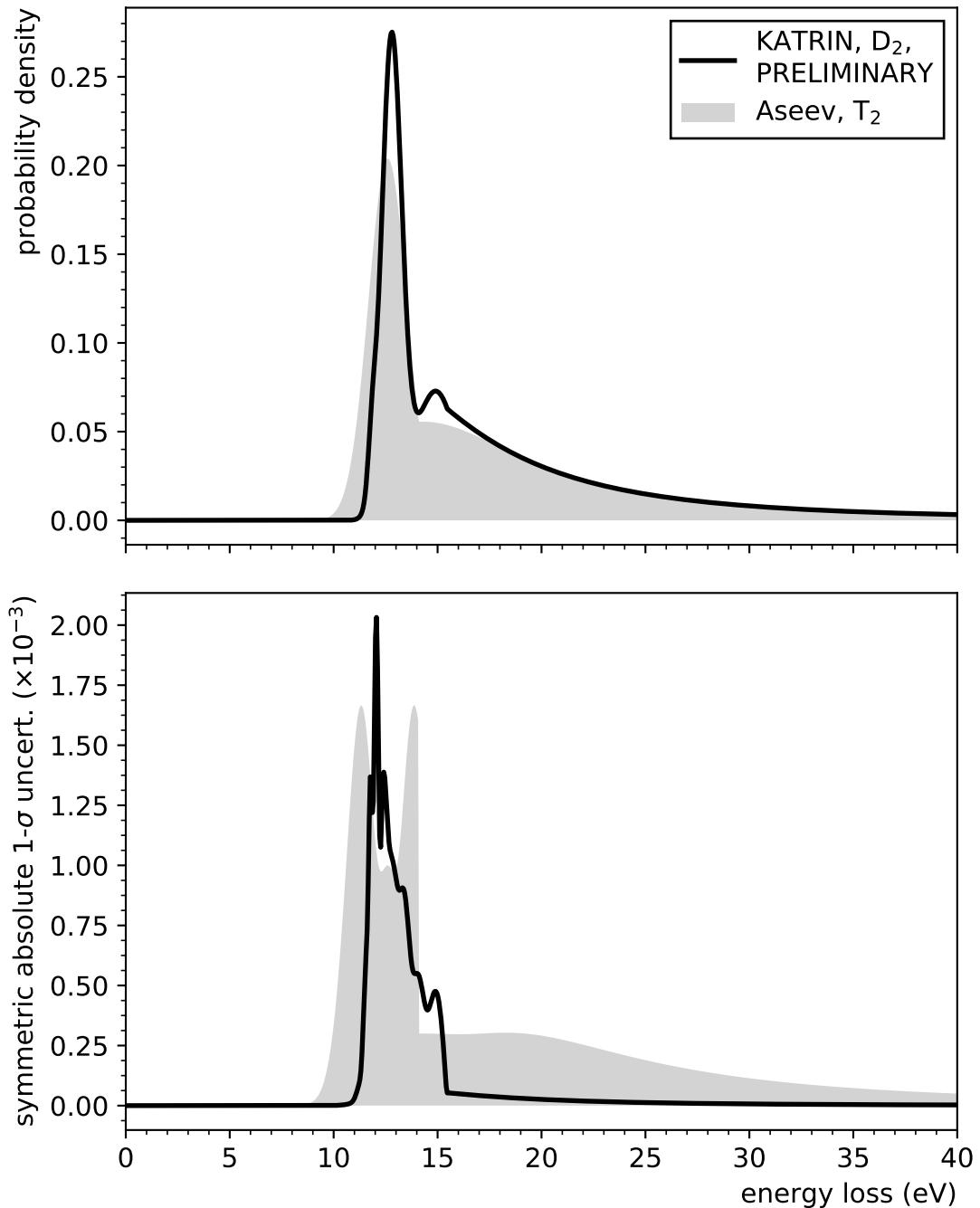


Figure 6.1.: The preliminary model for the energy loss of electrons scattering from deuterium molecules established at KATRIN. The black line shows the energy loss model established by a dedicated subgroup of the KATRIN collaboration for scattering off deuterium molecules (KATRIN model). The corresponding model for scattering off tritium molecules as established by [Aseev+00] (Aseev model) is shown for comparison as a shaded area. (That the latter is plotted as area instead of as a line solely serves readability as the two functions overlap strongly.) The top panel shows the probability densities of the energy loss and the bottom panel shows the corresponding absolute symmetric 1- σ uncertainties. The uncertainties were obtained through uncertainty propagation via derivatives from the uncertainties on the model parameters (see figure 3.3 for the Aseev model and appendix D for the KATRIN model). Correlations are respected for the KATRIN model. However, there are no published correlations for the Aseev model.

6.2.1. Description

Figure 6.1 shows the KATRIN model in comparison to the Aseev model. They are not expected to be fully compatible within their uncertainties as they describe the scattering off two different hydrogen isotopologues. For a comparison of the KATRIN model for deuterium with an energy loss model for deuterium from literature, the reader is referred to [Pre-Rod19]. Furthermore, the parametrization of the KATRIN model comprises a second peak for excitation and molecular dissociation of deuterium. As the model is still in its early stages this work refrains from a detailed physical interpretation and instead focuses on the statistical features of the KATRIN model. For a more detailed physical interpretation the reader is referred to the KATRIN documents [Pre-Rod19; Pre-Han19-1; Pre-Han19-2]. With respect to the statistical properties, the KATRIN model shows an improved uncertainty in the ionization tail and also in large parts of the excitation peak.

6.2.2. Parametrization

In section 3.2.5, the energy loss function $f_1(\epsilon)$ was introduced. It denotes the probability density for an energy ϵ that an electron loses when scattering once. The KATRIN model is such an energy loss function. It can be divided into two parts. The first part is the phenomenological description of the excitation peak region by the sum of three scaled Gaussian distributions \mathcal{N}_i ($i \in \{1, 2, 3\}$). This part can be described by 9 parameters which comprise the scales A_i , means m_i and standard deviations s_i of the three Gaussian distributions [Pre-Han19-2]

$$\boldsymbol{\pi}_{\text{eloss}} = (A_1, m_1, s_1, A_2, m_2, s_2, A_3, m_3, s_3)^T \quad (6.1)$$

$$\frac{d\sigma_{\text{exit}}^{\text{phen}}(\epsilon | \boldsymbol{\pi}_{\text{eloss}})}{d\epsilon} = \sum_{i=1}^3 A_i \cdot \mathcal{N}_i(\epsilon, \mu = m_i, \sigma = s_i). \quad (6.2)$$

The second part, the ionization tail, follows the binary-encounter-dipole (BED) model. For a full formula, the reader is referred to [KR94]. Here, this part will be denoted

$$\frac{d\sigma_{\text{ion}}^{\text{BDE}}(\epsilon)}{d\epsilon}. \quad (6.3)$$

The transition between the two parts is introduced at the ionization energy of deuterium $E_{\text{ion},D_2} = 15.467 \text{ eV}$ [Shi+93]. In order for the transition to be continuous a scaling factor for the ionization tail is introduced

$$c = \left(\frac{d\sigma_{\text{exit}}^{\text{phen}}(E_{\text{ion},D_2} | \boldsymbol{\pi}_{\text{eloss}})}{d\epsilon} \right) \Bigg/ \left(\frac{d\sigma_{\text{ion}}^{\text{BDE}}(E_{\text{ion},D_2})}{d\epsilon} \right). \quad (6.4)$$

Then, the full parametrization for the KATRIN model reads [Pre-Han19-1]

$$f_1^{\text{KATRIN}}(\epsilon | \boldsymbol{\pi}_{\text{eloss}}) = \begin{cases} \frac{d\sigma_{\text{exit}}^{\text{phen}}(\epsilon | \boldsymbol{\pi}_{\text{eloss}})}{d\epsilon} & \text{if } \epsilon \leq E_{\text{ion},D_2} \\ c \cdot \frac{d\sigma_{\text{ion}}^{\text{BDE}}(\epsilon)}{d\epsilon} & \text{if } \epsilon > E_{\text{ion},D_2} \end{cases}. \quad (6.5)$$

6.2.3. Uncertainties

The uncertainties of the KATRIN model as evaluated at the time of writing this thesis can be divided into two sets: The KATRIN model itself comprises nine parameters. However, it was obtained in a 15-parameter fit. The other six fit parameters are correlated with the parameters of the KATRIN model and hence are not necessarily negligible with regard to uncertainties. In the scope of this thesis they were incorporated in the statistical

treatment of the uncertainties. Whether they may be neglected in future studies needs further investigation.

For a detailed description of the full fit of the KATRIN model to the recorded data, the reader is referred to [Pre-Han19-1], but a short description is given in the following. In order to describe the further fit parameters a short overview of the used data sets is required. Four integral spectra were recorded analogously to the described integral β spectrum in chapter 3, but with the electron gun (see section 2.3) as source instead of tritium. For each data set a different deuterium column density (see section 2.2) was set in the WGTS: 0 %, 15 %, 50 %, 100 % of the nominal column density of $\rho d = 5 \times 10^{17}$ molecules/cm². The fit model for the 15 %-measurement was rescaled with respect to the other measurements. The corresponding scaling factor $N_{\text{int},15}$ was a free fit parameter. A further scaling factor N_K for the KATRIN model was introduced in the fit as a free parameter in order for the whole KATRIN model to keep its properties of a probability density and integrate to unity. Additionally, a 5th data set recorded at 15 % column density in a time-of-flight mode [Bon+99] was fitted simultaneously. For each data set a different expected scattering count $\mu_{\text{tov}}, \mu_{\text{int},1}, \mu_{\text{int},2}, \mu_{\text{int},3}$ (see equation 3.17) was fitted, which adds four further parameters (there is no scattering for the 0 % measurement). In summary, the additional parameters in the fit of the KATRIN model are

$$\boldsymbol{\pi}_{\text{eloss+}} = (N_K, \mu_{\text{tov},15}, \mu_{\text{int},15}, \mu_{\text{int},50}, \mu_{\text{int},100}, N_{\text{int},15})^T. \quad (6.6)$$

The best best-fit values and the covariance matrix of the full 15-parameter fit ($\boldsymbol{\pi}_{\text{eloss}}$ and $\boldsymbol{\pi}_{\text{eloss+}}$) can be found in appendix D. The aim of this chapter is to study the impact of the model uncertainties from these 15 parameters on KATRIN's sensitivity to the neutrino mass.

6.3. Scope of Validity

This section lists considerations that should be kept in mind with regard to the study presented in this chapter.

First, it should be noted that the KATRIN model as presented is still preliminary and may be subject to change.

Furthermore, during a neutrino mass measurement, the nominal tritium purity in the WGTS is approximately 95 %. Here, an energy loss model for deuterium is used. Two different, but similar energy loss models are expected for the two different gas species (compare [Abd+17] and [Ase+00]). Within the scope of this thesis, the impact on KATRIN's sensitivity from the reduction of uncertainties on the energy loss model are of primary interest. As long as the KATRIN model for scattering off deuterium molecules is sufficiently similar to the model for scattering off tritium molecules, it is a plausible approach to use it for simulations of KATRIN neutrino mass measurements. (For the similarity, see figure 6.1.) In that regard, the results for KATRIN's sensitivity derived in this chapter are expected to be meaningful.

It should also be noted that the covariances used in this thesis for the parameter sets $\boldsymbol{\pi}_{\text{eloss}}$ and $\boldsymbol{\pi}_{\text{eloss+}}$ (see equations 6.1 and 6.6) were estimated using the MIGRAD algorithm of the ROOT software framework [Pre-Han19-1]. A cross-check by a more reliable estimation of the covariances, for example by means of scanning the likelihood with the MINOS algorithm was not yet done at the time of writing this thesis.

Furthermore, the KATRIN model does not yet incorporate systematic uncertainties. Corresponding efforts for their incorporation are made at the time of writing this thesis. A future version of the KATRIN model might exhibit larger uncertainties than the version used in this thesis.

6.4. Statistical Prerequisites

This section develops the statistical tools used in the scope of this thesis in order to evaluate the impact of the KATRIN model on KATRIN’s sensitivity to the neutrino mass. The methods are described in a general manner (and could be applied to study model uncertainties in general) and then related to the KATRIN model. Section 6.5 presents a concept for the combination of multiple measurements in parameter inference. Section 6.4.2 introduces the profile likelihood method for the treatment of nuisance parameters. And section 6.4.3 outlines how the statistical concepts were implemented into the KaFit (see section 4.5) software framework.

6.4.1. Combination of Commissioning and Neutrino Mass Measurements

If two measurements share a set of parameters $\boldsymbol{\theta}_s$, but have additionally an individual set of parameters $\boldsymbol{\theta}_1$ and $\boldsymbol{\theta}_2$ and different sets of observations a combined likelihood is given by the product of the single likelihoods L_1 and L_2 [Tan+18]

$$\begin{aligned} -2 \ln L(\boldsymbol{\theta}_s, \boldsymbol{\theta}_1, \boldsymbol{\theta}_2) &= -2 \ln L_1(\boldsymbol{\theta}_s, \boldsymbol{\theta}_1) - 2 \ln L_2(\boldsymbol{\theta}_s, \boldsymbol{\theta}_2) \\ &\equiv -2 \ln L_1(\boldsymbol{\theta}_{1,s}) - 2 \ln L_2(\boldsymbol{\theta}_{2,s}). \end{aligned} \quad (6.7)$$

where, for ease of notation, the combined parameter vectors $\boldsymbol{\theta}_{s,1} \equiv (\boldsymbol{\theta}_s, \boldsymbol{\theta}_1)$ and $\boldsymbol{\theta}_{s,2} \equiv (\boldsymbol{\theta}_s, \boldsymbol{\theta}_2)$ are introduced. The first measurement could be sensitive to the neutrino mass whereas the second measurement could have been a calibration or commissioning measurement and be sensitive to parameters of the response function (see equation 3.29). Combining both likelihoods would incorporate the uncertainties on the parameters of the response function in the neutrino mass determination.

For practicality, in this thesis, an approximation is applied: The calibration measurement is seen as evaluated independently and one obtains estimates $\hat{\boldsymbol{\theta}}_{2,s}$, and an estimated covariance matrix $\hat{V}_{2,s}$. These can in turn be used to approximate the likelihood L_2 . A choice that stands to reason for the approximation of L_2 is a multivariate normal distribution \mathcal{N} . For the purpose of parameter inference through the maximum likelihood method $-2 \ln L_2$ needs only to be accurately approximated within the contour, that is needed to extract uncertainty intervals. The choice of a multivariate normal distribution corresponds a symmetric approximation of $-\ln L_2$ around its minimum by a second order function. The combined likelihood then reads

$$\begin{aligned} -2 \ln L(\boldsymbol{\theta}_s, \boldsymbol{\theta}_1, \boldsymbol{\theta}_2) &\approx -2 \ln L'(\boldsymbol{\theta}_s, \boldsymbol{\theta}_1) \\ &= \underbrace{\chi^2(\boldsymbol{\theta}_{1,s})}_{(1)} + \underbrace{(\boldsymbol{\theta}_s - \hat{\boldsymbol{\theta}}_{s,2})^\top \hat{V}_{s,2}^{-1} (\boldsymbol{\theta}_s - \hat{\boldsymbol{\theta}}_{s,2})}_{(2)} + \text{constants} \\ &= \chi^2(\boldsymbol{\theta}_{1,s}) - 2 \ln \mathcal{N}(\boldsymbol{\theta}_{2,s}, \hat{\boldsymbol{\theta}}_{s,2}, \hat{V}_{s,2}^{-1}) + \text{constants} \end{aligned} \quad (6.8)$$

Here, (1) is the chi-square likelihood for a KATRIN neutrino mass measurement (see equation 4.4). And (2) resembles the negative log likelihood of the calibration measurement approximated by a multivariate normal distribution. Terms having a form like (2) are also sometimes called “pull terms” because in the minimization of the likelihood they “pull” the parameters $\boldsymbol{\theta}_s$ towards the corresponding values in $\hat{\boldsymbol{\theta}}_{s,2}$.

The chi-square term (1) is a sum of n standard normal distributed random variables. Hence, as discussed in section 4.2, a likelihood only composed of the chi-square term (1) offers a goodness-of-fit criteria via the Pearson chi-square statistic. Whether the same criteria can be applied to the combined likelihood has to be investigated individually from case to case.

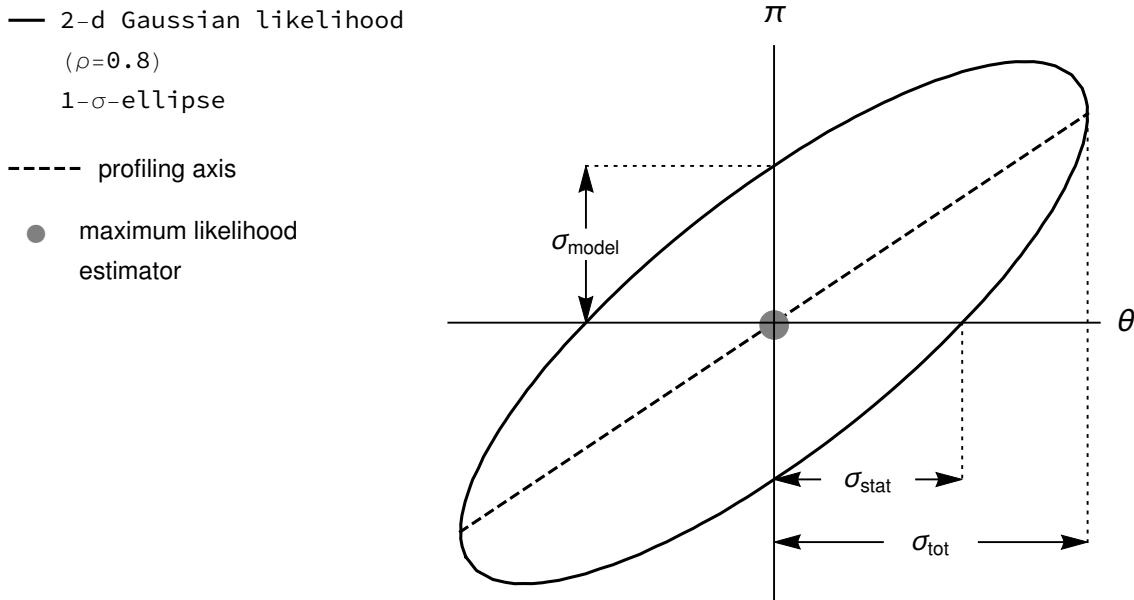


Figure 6.2.:

Application to the KATRIN Energy Loss Model

With regard to the study presented in this chapter, the following identification can be made. θ_1 comprises the parameter of a nominal four-parameter KATRIN neutrino mass fit (see section 4.3). Furthermore, the commissioning/calibration measurement can be identified with the measurement of the KATRIN energy loss model $\hat{\theta}_s = \hat{\pi}_{\text{eloss}}$, $\hat{\theta}_2 = \hat{\pi}_{\text{eloss+}}$ and the corresponding covariance matrix $\hat{V}_{s,2}$. The numerical values for the later three can be found in appendix D.

6.4.2. Nuisance Parameters and the Profile Likelihood Method

Apart from the parameters of interest θ (usually the squared neutrino mass), the KATRIN likelihood depends on further so-called nuisance parameters π . The dimensionality may pose difficulties when deriving a confidence region for the combined parameter set. Furthermore, as indicated by the naming conventions, the dimensions of the nuisance parameters in the confidence region are not of interest. Hence, in order to derive a confidence interval with just one dimension, a test statistic, similar to the one in equation 4.6, but that solely depending on the parameters of interest, has to be found. The following paragraph outlines, how a corresponding test statistic can be constructed using the profile likelihood method.

A corresponding derivation may start with the definition of the profile likelihood: The profile likelihood only depends on the parameters of interest θ and not on the nuisance parameters π . Its values correspond the likelihood values evaluated at θ in the dimensions of the parameters of interest and maximized in the dimensions of the nuisance parameters [Tan+18]

$$L_p(\theta) = L(\theta, \hat{\pi}(\theta)), \quad (6.9)$$

where the double-hat indicates the maximization respectively the profiling. Also, the profile likelihood ratio can be defined [Tan+18]

$$\lambda_p(\theta) = \frac{L_p(\theta)}{L_p(\hat{\theta})}. \quad (6.10)$$

According to Wilks' theorem [Wil38], the distribution of $-2 \ln \lambda_p(\hat{\theta})$, where $\hat{\theta}$ is the MLE (see section 4.1), approaches a chi-square distribution in the limit of a large data sample

size, independently of the values of the nuisance parameters π [Tan+18]. Hence, the profile likelihood ratio offers a test statistic, from which a confidence interval for the parameters of interest can be derived. However, it should be noted that whether the limit of a large sample size is given has to be investigated individually from case to case.

6.4.3. Extension of the KaFit Software Framework

The likelihood $L(\boldsymbol{\theta})$ can be multiplied by a function $g(\boldsymbol{\theta})$

$$-2 \ln L'(\boldsymbol{\theta}) = -2 \ln L(\boldsymbol{\theta}) - 2 \ln g(\boldsymbol{\theta}). \quad (6.11)$$

This procedure may have different interpretations and usage scenarios. E.g. a comparison with (??) shows, if g is a prior probability distribution, L' becomes a non-normalized posterior distribution that can be used in a Bayesian analysis. A further interpretation is given in section 6.4.3.

KaFit allowed to choose g in 6.11 as a product of one-dimensional Gaussian distributions. Within this thesis the software was extended to allow products of other functions. Three function types were explicitly made available through a configuration file.

1. A reimplementation of a one-dimensional Gaussian distribution: The reimplementation was necessary to conveniently enable the combination of function types.
2. A multivariate Gaussian distribution: This enables the treatment of uncertainties quantified by calibration or monitor measurements as described in section 6.4.3. It can also be used as a prior distribution in a Bayesian analysis. Particularly, correlations can be respected.
3. A one-dimensional probability density, that is constant in the square root of a parameter, if it is positive and 0 otherwise:

$$g(\theta) = \begin{cases} 0 & \text{if } \theta \leq 0 \\ \text{constant} \cdot \frac{1}{\sqrt{\theta}} & \text{if } \theta > 0 \end{cases}. \quad (6.12)$$

This can be used as a uniform prior on the neutrino mass ($\theta = m_\nu^2$). Formerly, it was only possible to use a uniform prior on the squared neutrino mass. A derivation of the form of g can be found in appendix B.

An example on how to configure KaFit using the new feature is given in appendix [\[Add ToDo appendix\]](#).

6.5. A Sensitivity Study using the Recent Preliminary KATRIN Energy Loss Model

6.6. Conclusion and Outlook

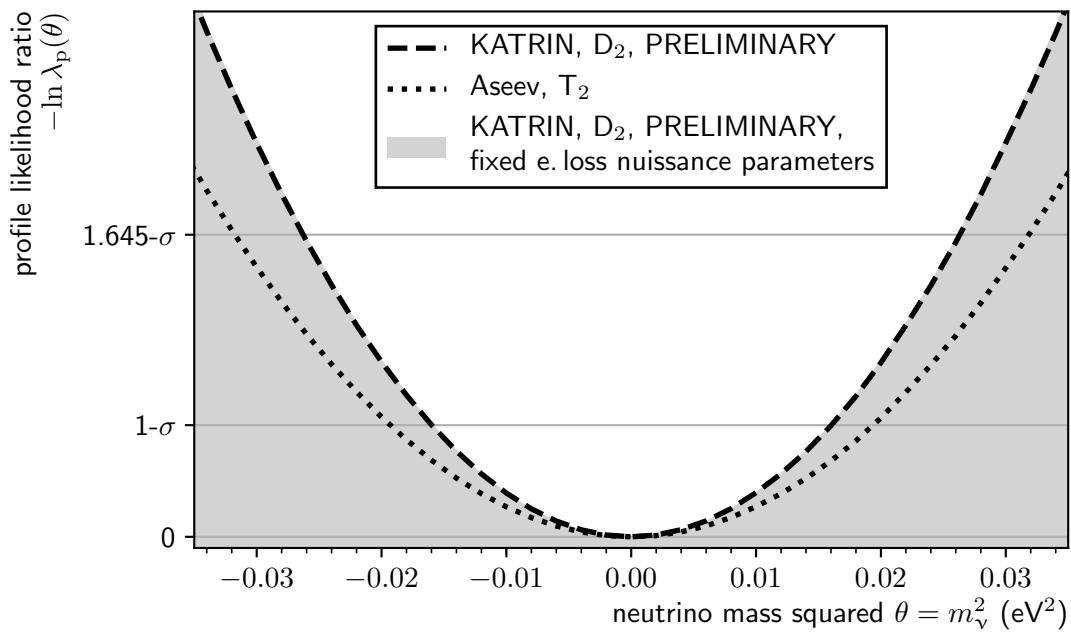


Figure 6.3.:

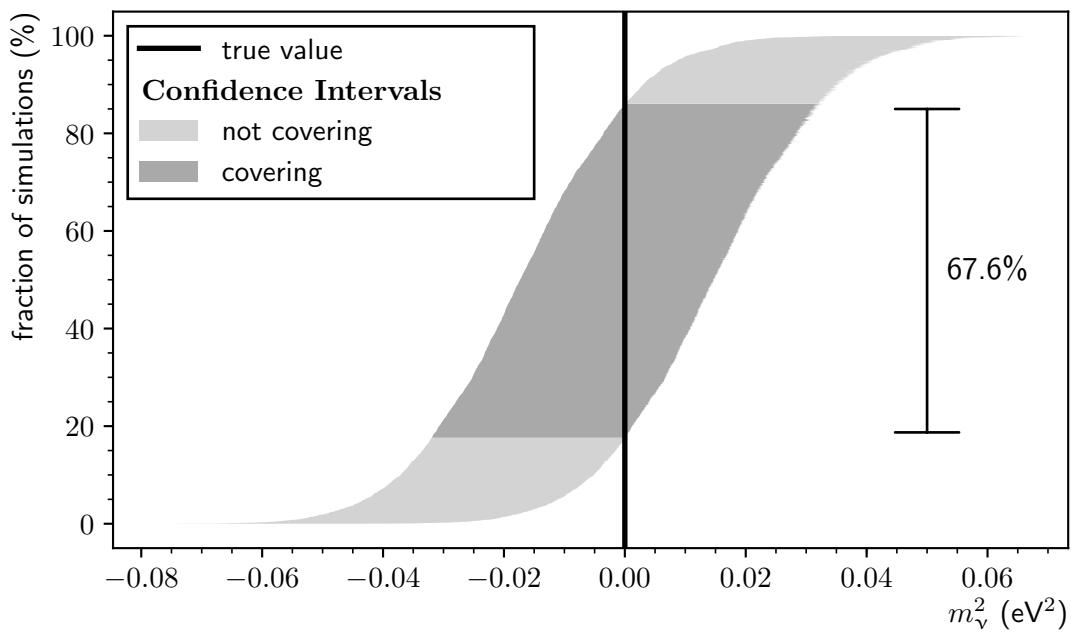


Figure 6.4.:

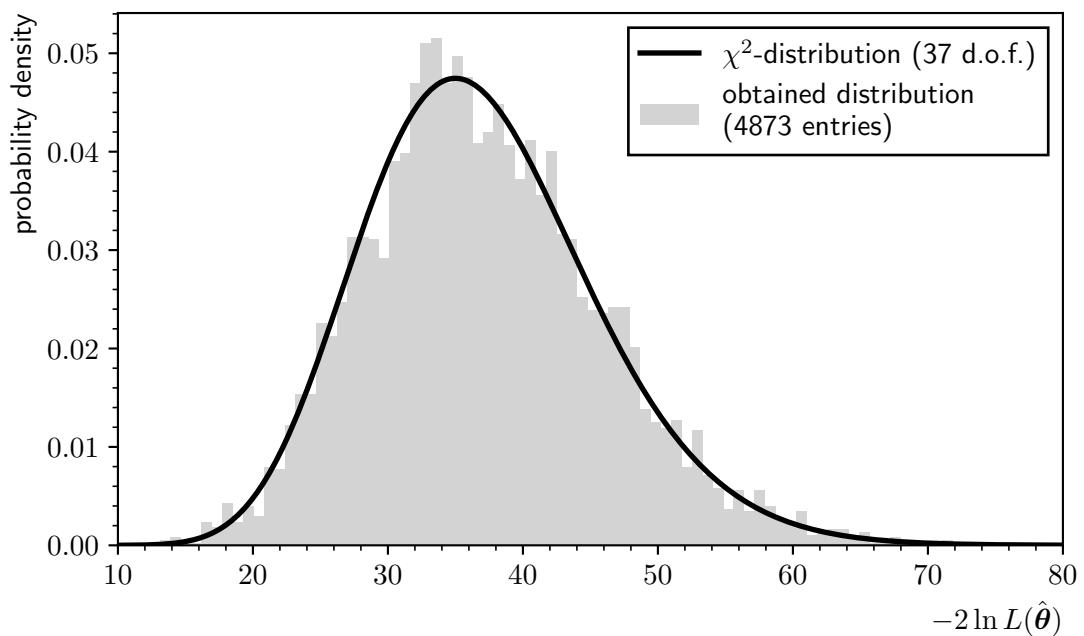


Figure 6.5.:

7. Conclusions and Outlook

Appendix

A. Energy-Dependence of the Scattering Probabilities

[Better subcaption, introduction]

ToDo

A.1. Modeling

A model for energy dependent scattering probabilities is derived. The expected amount of scatterings for a β electron when traveling through the whole WGTS volume of length L filled with a gas of constant density ρ is

$$\mu(E, \theta_S) = \frac{\sigma(E)\rho L}{\cos \theta_S}, \quad (\text{A.1})$$

where E denotes the electron's kinetic energy; θ_S the starting pitch angle; and $\sigma(E)$ the energy dependent scattering cross section.

The volume of the WGTS is divided into N slices of equal width $w = L/N$. N is chosen sufficiently large that the probability for a β electron to scatter twice within one slice is essentially zero. Then, for large N the probability to scatter within one slice is $\mu(E, \theta_S)/N$. The probability not to scatter within $n < N$ slices is

$$p_0(E, \theta_S, n) = \left(1 - \frac{\mu(E, \theta_S)}{N}\right)^n. \quad (\text{A.2})$$

Using the well known limit for the Euler constant, one obtains for $n = N$ and $N \rightarrow \infty$ that p_0 is a Poisson distribution with expectation μ evaluated at 0

$$\lim_{N \rightarrow \infty} p_0(E, \theta_S, N) = \lim_{N \rightarrow \infty} \left(1 - \frac{\mu(E, \theta_S)}{N}\right)^N = e^{-\mu(E, \theta_S)}. \quad (\text{A.3})$$

Assuming a constant energy loss per scattering of ϵ the probability to scatter l times within $n < N$ slices can be expressed recursively

$$p_l(E, \theta_S, n) = \underbrace{\sum_{k=l}^n}_{(4)} \underbrace{p_{l-1}(E, \theta_S, k-1)}_{(1)} \underbrace{(1 - p_0(E - (l-1)\epsilon, \theta_S, 1))}_{(2)} \underbrace{p_0(E - l\epsilon, \theta_S, n-k)}_{(3)}. \quad (\text{A.4})$$

The terms have the following meaning:

- (1) Probability to scatter $l-1$ times within $k-1$ slices with a kinetic energy of E .
- (2) Probability to scatter once within the k th slice with a kinetic energy of $E - (l-1)\epsilon$.
- (3) Probability not to scatter within the remaining $N-k$ slices.
- (4) Sum over all slices k . The sum starts at l because the probability to scatter $l-1$ times within less than $k = l-1$ slices (term (1)) is 0.

The probability to scatter l times can be averaged over all starting positions. The averaging sum can be further approximated by an integral as this helps cutting down on run time in a numerical evaluation

$$\bar{p}_l(E, \theta_S) = \frac{1}{N} \sum_{n_S=1}^N p_l(E, \theta_S, n_S) \approx \frac{1}{L} \int_0^L p_l(E, \theta_S, \left\lceil N \frac{z_S}{L} \right\rceil) dz_S. \quad (\text{A.5})$$

Then the limit $N \rightarrow \infty$ is applied

$$P_l^*(E_S, \theta_S) = \lim_{N \rightarrow \infty} \bar{p}_l(E_S, \theta_S) \quad (\text{A.6})$$

$P_l^*(E_S, \theta_S)$ denotes the probability for a β electron to scatter l times when traveling through the whole WGTS with a starting energy E_S and pitch angle θ_S averaged over all starting positions. Finally, this expression has to be averaged over all starting pitch angles in order to obtain the energy dependent scattering probabilities

$$\bar{P}_l^*(E_S) = \frac{1}{1 - \cos \theta_{\max}} \int_0^{\theta_{\max}} \sin \theta_S P_l^*(E_S, \theta_S) d\theta_S. \quad (\text{A.7})$$

A.2. Numerical Evaluation

The energy dependent probability to scatter once \bar{P}_1^* in (A.7) was evaluated numerically. In (A.6) taking the limit $N \rightarrow \infty$ was replaced by choosing a large N . The averaging integral over the starting positions (A.5) and pitch angles (A.7) was computed using Gaussian quadrature.

Figure A.1 shows the result in dependence of N and the order of the Gaussian quadrature. Both should be chosen as low as possible to cut down on run time but sufficiently high for the required accuracy. As a cross check, for an energy loss of $\epsilon = 0$ eV per scattering, the suggested extended model (A.7) must recover the Poisson model (see (5.6)). For $N = 10^6$ and using Gaussian quadrature of order 5, they differ approximately 3×10^{-6} . Furthermore, the numerical calculation converges from below, which can be interpreted as a one-sided numerical inaccuracy. The calculations for $\epsilon = 14.1$ eV are shown in the lower row of figure A.1. They also show convergence on the 10^{-5} level. Conclusively, the results make it plausible to assume a one-sided numerical inaccuracy on the 10^{-5} level for $N = 10^5$ and using Gaussian quadrature of order 5 for the integrals.

The corresponding run time to compute \bar{P}_1^* is in $\mathcal{O}(N)$ as it requires a sum over all N slices. Note that the extended model is defined recursively and therefore, the run time for l -fold scattering is in $\mathcal{O}(N^l)$. Hence, computing the probability for 2-fold scattering would take 500000 times as long as for 1-fold scattering for the same 10^{-5} accuracy. This was found not to be feasible. For 2-fold scattering $N = 5000$ and also Gaussian quadrature of order 5 was chosen for a numerical accuracy on the 10^{-3} level.

B. Neutrino Mass Prior

A KATRIN analysis derives the squared neutrino mass m_ν^2 . Hence, in a Bayesian analysis a prior on m_ν^2 is required. A uniform prior on the neutrino mass m_ν is a possible choice. The following lines derive a prior on m_ν^2 that resemble a uniform prior on m_ν . Let $f(m_\nu) = C = \text{constant}$ be the prior on m_ν and $g(m_\nu^2)$ be the prior on m_ν^2 . Starting from conservation of probability one derives

$$\begin{aligned} f(m_\nu) dm_\nu &= g(m_\nu^2) dm_\nu^2 \\ \Rightarrow g(m_\nu^2) &= f(m_\nu) \left(\frac{dm_\nu^2}{dm_\nu} \right)^{-1} \\ \Rightarrow g(m_\nu^2) &= C \frac{1}{2\sqrt{m_\nu^2}}. \end{aligned}$$

C. Likelihood Extensions

Within this thesis the implementation of the likelihood in the KaFit framework was extended as described in section ???. It is possible to multiply the likelihood by different function types. Example excerpts from KaFit-XML-configurations are given below. **[Double]** is used as a placeholder for a number; and **[Index*]** for a parameter index. E.g. the neutrino mass squared has the parameter index 0. A new **Penalty**-tag was introduced as sub-tag of the **LoglikelihoodKatrin**-tag:

```
<LoglikelihoodKatrin
  Name="myKatrinLogL" PDF="Gauss" RunSource="myRunGen"
  SpectrumSimulator="mySpecSim">
  <Penalty>
    <!-- Penalty Type -->
  </Penalty>
</LoglikelihoodKatrin>
```

<!-- Penalty Type --> can be substituted by one ore more of the following tags:

Multivariate Gaussian Distribution

The attribute **Mean** specifies the mean of a parameter; **Std** the standard deviation; and one ore more **Correlation**-sub-tags the correlations between the parameters specified by the **Parameter**-tags.

```
<MultivarNorm>
  <Parameter Index="[Index1]" Mean="[Double]" Std="[Double]" />
  <Parameter Index="[Index2]" Mean="[Double]" Std="[Double]" />
  <!-- ... -->
  <Correlation Index1="[Index1]" Index2="[Index2]" Value="[Double]" />
  <!-- ... -->
</MultivarNorm>
```

One-Dimensional Gaussian Distribution

Analogously to the multivariate, a one-dimensional Gaussian distribution can be used:

```
<Gaussian ParamIndex="[Index]" Mean="[Double]" Std="[Double]" />
```

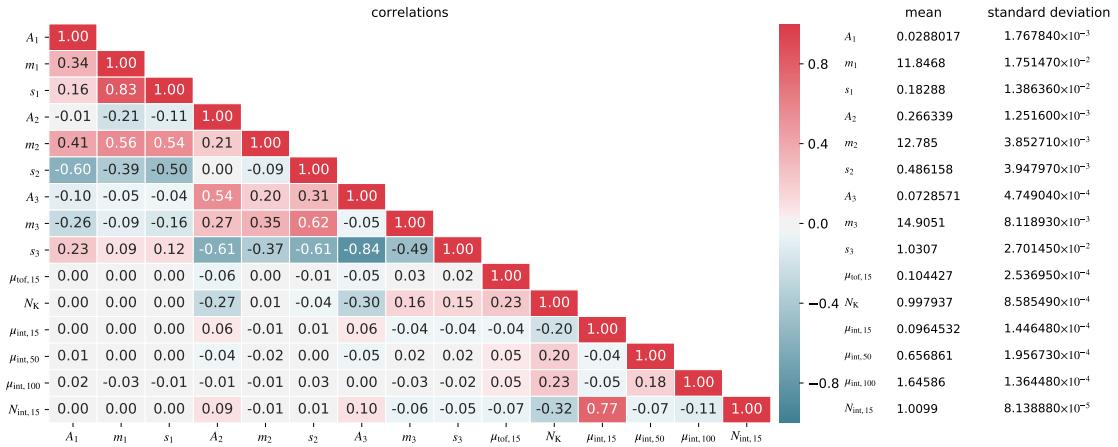
Uniform Neutrino Mass Prior

A constant prior on the neutrino can be set via the following tag:

```
<ConstInSqrt ParamIndex="0" />
```

D. Parameter Values of the Preliminary KATRIN Energy Loss Model

The best fit values, standard deviations and correlations of the fit parameters for the KATRIN model for the energy loss of electrons scattering off deuterium molecules as used in this thesis are listed below [Pre-Han19-1]. The parameter names follow section 6.2.



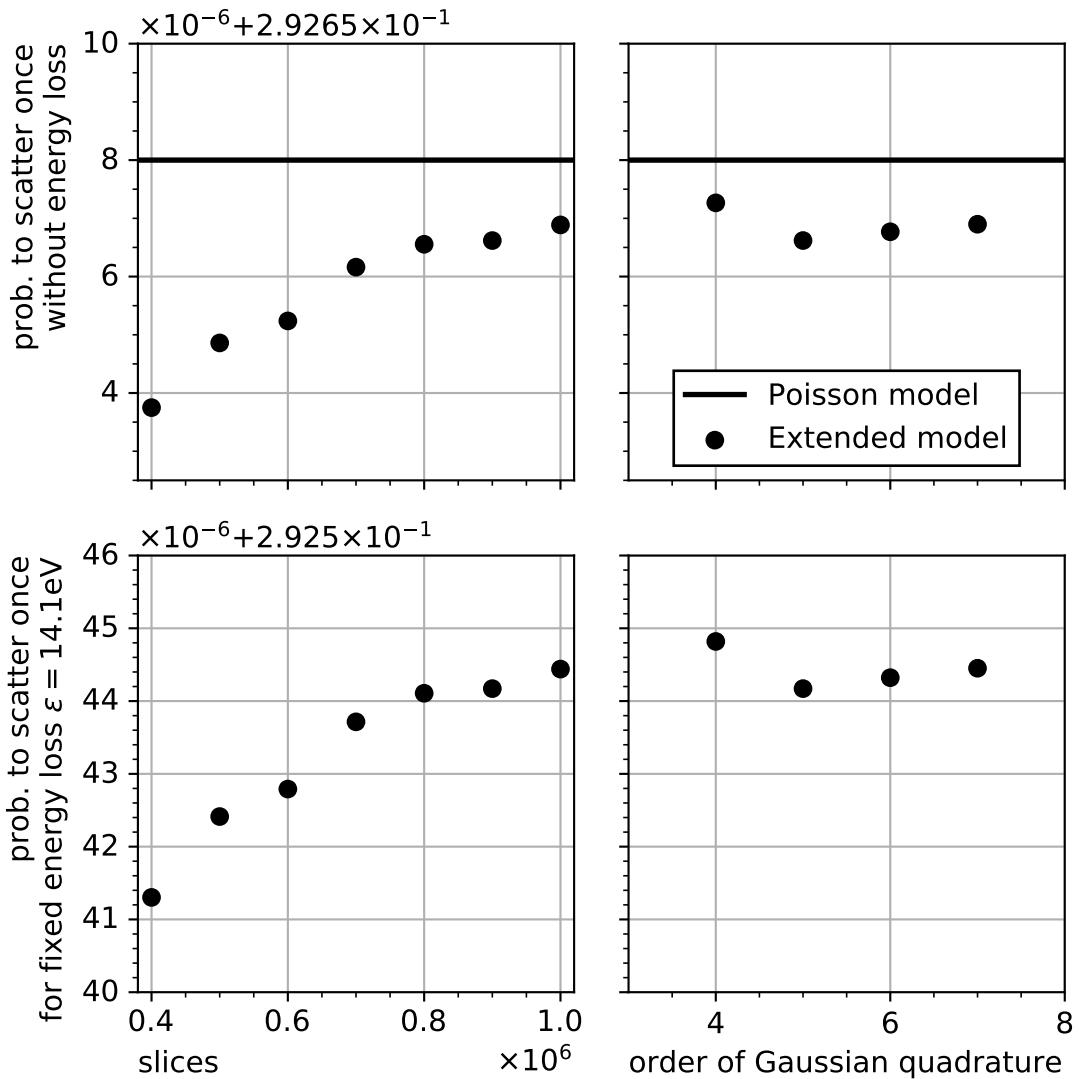


Figure A.1.: Numerical accuracy of the extended model for the probability of electron scattering within the WGTS. The extended model is given by (A.6). The left column shows the dependence on the number of slices of the WGTS. The right column shows the order of Gaussian quadrature that was used to evaluate the two integrals in the model. For the left column the order of Gaussian quadrature was fixed to 5. For the right column the number of slices was fixed to 5×10^5 . The upper row shows the extended model for no energy loss per scattering (markers). Its exact solution is given by the Poisson model (see (5.6)). The lower row shows the model for an energy loss of 14.1 eV per scattering. The results make it plausible to assume a one-sided numerical inaccuracy on the 10^{-5} level for $N = 10^5$ and using Gaussian quadrature of order 5.

Acronyms

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