

Newman Janis Algorithm

1 Calculations

The standard kerr metric is given by

$$ds^2 = - \left(1 - \frac{2Mr}{\rho^2} \right) dt^2 - \frac{4Mar \sin^2 \theta}{\rho^2} dt d\phi + \frac{\Sigma}{\rho^2} \sin^2 \theta d\phi^2 + \frac{\rho^2}{\Delta} dr^2 + \rho^2 d\theta^2$$

where,

$$\rho^2 = r^2 + a^2 \cos^2 \theta, \Delta = r^2 - 2Mr + a^2 \text{ and } \Sigma = (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta$$

$$g_{\mu\nu} = \begin{vmatrix} - \left(1 - \frac{2Mr}{\rho^2} \right) & 0 & 0 & - \frac{2Mar \sin^2 \theta}{\rho^2} \\ 0 & \frac{\rho^2}{\Delta} & 0 & 0 \\ 0 & 0 & \rho^2 & 0 \\ - \frac{2Mar}{\sin^2 \theta} \rho^2 & 0 & 0 & \frac{\Sigma}{\rho^2} \sin^2 \theta \end{vmatrix}$$

and the Schwarzschild metric is given by

$$ds^2 = - \left(1 - \frac{2M}{r} \right) dt^2 + \frac{1}{1 - \frac{2M}{r}} dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

Switching to Eddington-Finkelstein coordinate with $x^{0'} = t + r$, we have

$$ds^2 = - \left(1 - \frac{2M}{r} \right) dv^2 + 2dvdr + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

And the metric looks like:

$$g_{\mu\nu} = \begin{vmatrix} - \left(1 - \frac{2M}{r} \right) & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & r^2 & 0 \\ 0 & 0 & 0 & r^2 \sin^2 \theta \end{vmatrix} \quad \text{and} \quad g^{\mu\nu} = \begin{vmatrix} 0 & 1 & 0 & 0 \\ 1 & 1 - \frac{2M}{r} & 0 & 0 \\ 0 & 0 & 1/r^2 & 0 \\ 0 & 0 & 0 & 1/(r^2 \sin^2 \theta) \end{vmatrix}$$

Based on the choice of normalization $l \cdot n = 1$ and $m \cdot \bar{m} = -1$, we have:

$$g^{\mu\nu} = l^\mu n^\nu + n^\mu l^\nu - m^\mu \bar{m}^\nu - \bar{m}^\mu m^\nu$$

In advanced Eddington-Finkelstein coordinate, the null geodesic is directed towards ∂_v , thus let us take that as l^μ and assuming $m_v = m_r = 0$. Thus, the

null tetrad can be calculated from the metric tensor in this coordinate and is given by:

$$l^\mu = (0, 1, 0, 0) \quad (1)$$

$$n^\mu = \left(1, \frac{1}{2} \left(1 - \frac{2M}{r}\right), 0, 0\right) \quad (2)$$

$$m^\mu = \frac{1}{\sqrt{2}r} \left(0, 0, 1, \frac{\iota}{\sin \theta}\right) \quad (3)$$

$$\bar{m}^\mu = \frac{1}{\sqrt{2}r} \left(0, 0, 1, -\frac{\iota}{\sin \theta}\right) \quad (4)$$

Under the Newman-Janis transformation, $r' \rightarrow r + ia \cos \theta$ and $v' \rightarrow v - ia \cos \theta$, we have:

$$\begin{aligned} l'^a &= \frac{\partial x'^a}{\partial x^a} l^a = \frac{\partial x'^a}{\partial r} l^r \\ &= \delta_r^a \end{aligned}$$

$$\begin{aligned} n'^a &= \frac{\partial x'^a}{\partial x^a} n^a = \frac{\partial x'^a}{\partial v} n^v + \frac{\partial x'^a}{\partial r} n^r \\ &= \delta_v^a - \frac{f(r)}{2} \delta_r^a \\ &= \delta_v^a - \frac{1}{2} \left(1 - M \frac{1}{r} - M \frac{1}{\bar{r}}\right) \delta_r^a \\ &= \delta_v^a - \frac{1}{2} \left(1 - \frac{2Mr}{r^2 + a^2 \cos^2 \theta}\right) \delta_r^a \end{aligned}$$

The last two of the tetrads can be calculated as follows:

$$\frac{\partial r'}{\partial \theta} = \frac{\partial}{\partial \theta} (r + ia \cos \theta) = -ia \sin \theta \quad (5)$$

$$\frac{\partial v'}{\partial \theta} = \frac{\partial}{\partial \theta} (v - ia \cos \theta) = ia \sin \theta \quad (6)$$

Now, moving onto the tetrads

$$m^a = \frac{1}{\sqrt{2}r} \left(\delta_\theta^a + \frac{1}{\sin \theta} \delta_\phi^a \right) \quad \text{and} \quad \bar{m}^a = \frac{1}{\sqrt{2}\bar{r}} \left(\delta_\theta^a - \frac{1}{\sin \theta} \delta_\phi^a \right)$$

Transformed tetrads are given as:

$$\begin{aligned}
m'^a &= \frac{\partial x'^a}{\partial x^a} m^a = \frac{\partial x'^a}{\partial x^\theta} m^\theta + \frac{\partial x'^a}{\partial \phi} m^\phi \\
&= \left(\frac{\partial x'^a}{\partial r'} \frac{\partial r'}{\partial \theta} + \frac{\partial x'^a}{\partial v'} \frac{\partial v'}{\partial \theta} + \frac{\partial x'^a}{\partial \theta'} \frac{\partial \theta'}{\partial \theta} \right) m^\theta + \frac{\partial x'^a}{\partial \phi} m^\phi \\
&= (-1a \sin \theta \delta_{r'}^a + 1a \sin \theta \delta_{v'}^a + \delta_\theta^a) m^\theta + \frac{\partial x'^a}{\partial \phi} m^\phi \\
&= \frac{1}{\sqrt{2}r} \left(\delta_\theta^a + (\delta_{v'}^a - \delta_{r'}^a) 1a \sin \theta + \frac{1}{\sin \theta} \delta_\phi^a \right) \quad \left(\text{using } m^\theta = \frac{1}{\sqrt{2}r} \delta_\theta^\theta \right) \\
&= \frac{1}{\sqrt{2}(r' - 1a \cos \theta)} \left(\delta_\theta^a + (\delta_{v'}^a - \delta_{r'}^a) 1a \sin \theta + \frac{1}{\sin \theta} \delta_\phi^a \right)
\end{aligned}$$

Now, \bar{m}^a

$$\begin{aligned}
\bar{m}'^a &= \frac{\partial x'^a}{\partial x^a} \bar{m}^a = \frac{\partial x'^a}{\partial x^\theta} \bar{m}^\theta + \frac{\partial x'^a}{\partial \phi} \bar{m}^\phi \\
&= \left(\frac{\partial x'^a}{\partial \bar{r}'} \frac{\partial \bar{r}'}{\partial \theta} + \frac{\partial x'^a}{\partial \bar{v}'} \frac{\partial \bar{v}'}{\partial \theta} + \frac{\partial x'^a}{\partial \theta'} \frac{\partial \theta'}{\partial \theta} \right) \bar{m}^\theta + \frac{\partial x'^a}{\partial \phi} \bar{m}^\phi \\
&= (1a \sin \theta \delta_{r'}^a - 1a \sin \theta \delta_{v'}^a + \delta_\theta^a) \bar{m}^\theta + \frac{\partial x'^a}{\partial \phi} \bar{m}^\phi \\
&= \frac{1}{\sqrt{2}\bar{r}} \left(\delta_\theta^a - (\delta_{v'}^a - \delta_{r'}^a) 1a \sin \theta - \frac{1}{\sin \theta} \delta_\phi^a \right) \quad \left(\text{using } \bar{m}^\theta = \frac{1}{\sqrt{2}\bar{r}} \delta_\theta^\theta \right) \\
&= \frac{1}{\sqrt{2}(r' + 1a \cos \theta)} \left(\delta_\theta^a - (\delta_{v'}^a - \delta_{r'}^a) 1a \sin \theta - \frac{1}{\sin \theta} \delta_\phi^a \right)
\end{aligned}$$

Thus the relevant metric tensor is given as:

$$\begin{aligned}
g^{vv} &= l^v n^v + n^v l^v - m^v \bar{m}^v - \bar{m}^v m^v \\
&= 0 \times \delta_v^v + \delta_v^v \times 0 - \frac{1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \times \frac{-1a \sin \theta}{\sqrt{2}(r' + 1a \cos \theta)} \\
&\quad - \frac{-1a \sin \theta}{\sqrt{2}(r' + 1a \cos \theta)} \times \frac{1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \\
&= -\frac{a^2 \sin^2 \theta}{r'^2 + a^2 \cos^2 \theta} \\
\\
g^{rr} &= l^r n^r + n^r l^r - m^r \bar{m}^r - \bar{m}^r m^r \\
&= 1 \times -\frac{f(r)}{2} - \frac{f(r)}{2} \times 1 - \frac{-1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \times \frac{1a \sin \theta}{\sqrt{2}(r' + 1a \cos \theta)} \\
&\quad - \frac{1a \sin \theta}{\sqrt{2}(r' + 1a \cos \theta)} \times \frac{-1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)}
\end{aligned}$$

$$\begin{aligned}
&= - \left(1 - \frac{2Mr}{r'^2 + a^2 \cos^2 \theta} \right) - \frac{a^2 \sin^2 \theta}{r'^2 + a^2 \cos^2 \theta} \\
&= - \left(\frac{r^2 + a^2 \cos^2 \theta - 2Mr + a^2 \sin^2 \theta}{r'^2 + a^2 \cos^2 \theta} \right) = - \left(\frac{r^2 - 2Mr + a^2}{r'^2 + a^2 \cos^2 \theta} \right)
\end{aligned}$$

$$\begin{aligned}
g^{\theta\theta} &= l^\theta n^\theta + n^\theta l^\theta - m^\theta \bar{m}^\theta - \bar{m}^\theta m^\theta \\
&= 0 \times 0 + 0 \times 0 - \frac{1}{\sqrt{2}(r' - 1a \cos \theta)} \frac{1}{\sqrt{2}(r' + 1a \cos \theta)} \\
&\quad - \frac{1}{\sqrt{2}(r' + 1a \cos \theta)} \frac{1}{\sqrt{2}(r' - 1a \cos \theta)} \\
&= - \frac{1}{(r'^2 + a^2 \cos^2 \theta)}
\end{aligned}$$

$$\begin{aligned}
g^{\phi\phi} &= l^\phi n^\phi + n^\phi l^\phi - m^\phi \bar{m}^\phi - \bar{m}^\phi m^\phi \\
&= 0 \times 0 + 0 \times 0 - \frac{1}{\sqrt{2}(r' - 1a \cos \theta) \sin \theta} \frac{-1}{\sqrt{2}(r' + 1a \cos \theta) \sin \theta} \\
&\quad - \frac{-1}{\sqrt{2}(r' + 1a \cos \theta) \sin \theta} \frac{1}{\sqrt{2}(r' - 1a \cos \theta) \sin \theta} \\
&= - \frac{1}{(r^2 + a^2 \cos^2 \theta) \sin^2 \theta}
\end{aligned}$$

$$\begin{aligned}
g^{vr} &= l^v n^r + n^v l^r - m^v \bar{m}^r - \bar{m}^v m^r \\
&= 0 \times \frac{-f(r)}{2} + 1 \times 1 - \frac{1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \frac{1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \\
&\quad - \frac{-1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \frac{-1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \\
&= 1 + \frac{a^2 \sin^2 \theta}{r^2 + a^2 \cos^2 \theta}
\end{aligned}$$

$$\begin{aligned}
g^{rv} &= l^r n^v + n^r l^v - m^r \bar{m}^v - \bar{m}^r m^v \\
&= 1 \times 1 + 1 \times 0 - \frac{-1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \frac{-1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \\
&\quad - \frac{1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \frac{1a \sin \theta}{\sqrt{2}(r' - 1a \cos \theta)} \\
&= 1 + \frac{a^2 \sin^2 \theta}{r^2 + a^2 \cos^2 \theta}
\end{aligned}$$

$$\begin{aligned}
g^{v\theta} &= l^v n^\theta + n^v l^\theta - m^v \bar{m}^\theta - \bar{m}^v m^\theta \\
&= 0 \times 0 + 0 \times 0 - \frac{1a \sin \theta}{\sqrt{2}(r - 1a \cos \theta)} \times \frac{1}{\sqrt{2}(r' + 1a \cos \theta)} \\
&\quad - \frac{-1a \sin \theta}{\sqrt{2}(r + 1a \cos \theta)} \times \frac{1}{\sqrt{2}(r' - 1a \cos \theta)} \\
&= 0
\end{aligned}$$

$$g^{v\phi} = l^v n^\phi + n^v l^\phi - m^v \bar{m}^\phi - \bar{m}^v m^\phi$$

$$\begin{aligned}
&= 0 \times 0 + 0 \times 0 - \frac{1a \sin \theta}{\sqrt{2}(r - 1a \cos \theta)} \frac{-1}{\sqrt{2}(r - 1a \cos \theta) \sin \theta} \\
&\quad - \frac{-1a \sin \theta}{\sqrt{2}(r + 1a \cos \theta)} \times \frac{1}{\sqrt{2}(r + 1a \cos \theta) \sin \theta} \\
&= -\frac{a}{r^2 + a^2 \cos^2 \theta} \\
\\
g_{r\phi} &= l^r n^\phi + n^r l^\phi - m^r \bar{m}^\phi - \bar{m}^r m^\phi \\
&= 0 \times 0 + 0 \times 0 - \frac{-1a \sin \theta}{\sqrt{2}(r - 1a \cos \theta)} \frac{-1}{\sqrt{2}(r - 1a \cos \theta) \sin \theta} \\
&\quad - \frac{1a \sin \theta}{\sqrt{2}(r + 1a \cos \theta)} \times \frac{1}{\sqrt{2}(r + 1a \cos \theta) \sin \theta} \\
&= \frac{a}{r^2 + a^2 \cos^2 \theta}
\end{aligned}$$

Therefore the metric tensor takes the following form:

$$g^{\mu\nu} = \begin{vmatrix} -\frac{a^2 \sin^2 \theta}{r'^2 + a^2 \cos^2 \theta} & 1 + \frac{a^2 \sin^2 \theta}{r^2 + a^2 \cos^2 \theta} & 0 & -\frac{a}{r^2 + a^2 \cos^2 \theta} \\ 1 + \frac{a^2 \sin^2 \theta}{r^2 + a^2 \cos^2 \theta} & -\frac{r^2 - 2Mr + a^2}{r'^2 + a^2 \cos^2 \theta} & 0 & \frac{a}{r^2 + a^2 \cos^2 \theta} \\ 0 & 0 & -\frac{1}{(r'^2 + a^2 \cos^2 \theta)} & 0 \\ -\frac{a}{r^2 + a^2 \cos^2 \theta} & \frac{a}{r^2 + a^2 \cos^2 \theta} & 0 & -\frac{1}{(r^2 + a^2 \cos^2 \theta) \sin^2 \theta} \end{vmatrix}$$

The inverse metric is given by

$$g^{\mu\nu} = \begin{vmatrix} 1 - \frac{2Mr}{r^2 + a^2 \cos^2 \theta} & 1 & 0 & \frac{2Mar \sin^2 \theta}{r^2 + a^2 \cos^2 \theta} \\ 1 & 0 & 0 & -a \sin^2 \theta \\ 0 & 0 & -\frac{1}{r^2 + a^2 \cos^2 \theta} & 0 \\ \frac{2Mar \sin^2 \theta}{r^2 + a^2 \cos^2 \theta} & -a \sin^2 \theta & 0 & -\left[(r^2 + a^2) + \frac{2Mar \sin^2 \theta}{r^2 + a^2 \cos^2 \theta}\right] \sin^2 \theta \end{vmatrix}$$

This result matches with Chandrasekhar's Mathematical Theory of Black Hole, pg 307. The Newman Janis Algorithm is Einstein Field Invariant, it doesn't break the diffeomorphism invariance.