Time Evolution of Quantum Systems 2025: Exercise 1

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Baker-Campbell-Hausdorff formula

- **H.1** In the following, we use the Lie brackets as shorthand notation for the commutator [A, B] := AB BA of two matrices A, B. Furthermore, we use Landau's big-O-notation. In particular, the Landau symbol $\mathcal{O}(g)$ in O-notation means that a considered function f grows at most as fast as g. In our case, $\mathcal{O}(\varepsilon^3)$ means that we can neglect all terms of orders $\varepsilon^3, \varepsilon^4, \ldots$ because we assume ε to be small.
 - (a) Show the Baker-Campbell-Hausdorff formula (BCH formula)

$$e^{(A+B)\varepsilon} = e^{A\varepsilon}e^{B\varepsilon}e^{-\frac{1}{2}[A,B]\varepsilon^2 + \mathcal{O}(\varepsilon^3)}$$

up to the second order.

(1 P.)

- (b) Now show $e^{A\varepsilon}e^{B\varepsilon} = e^{B\varepsilon}e^{A\varepsilon}e^{[A,B]\varepsilon^2 + \mathcal{O}(\varepsilon^3)}$. This variant is sometimes also referred to as the BCH formula. (1 P.)
- (c) Finally, show the Lie product formula $e^{A+B} = \lim_{n \to \infty} \left(e^{\frac{1}{n}A} e^{\frac{1}{n}B} \right)^n$. (1 P.)
- (d) Suppose that $[A, B] = c\mathbb{1}$ with $c \in \mathbb{C}$. Show that

$$e^A \cdot e^B = e^B \cdot e^A \cdot e^{c1}.$$

(1 P.)

(e) Prove the Campbell identity for a linear operator on Hilbert space,

$$e^{A}Be^{-A} = \sum_{k=0}^{\infty} \frac{1}{k!} [A, B]_{k}$$

where $[A, B]_0 = B$ and $[A, B]_k = [A, [A, B]]_{k-1}$. (2 P.)

Hint: Replace the operator A by εA with $\epsilon \in \mathbb{R}$, and do a Taylor expansion in ϵ .

Analytic Exact Diagonalization

H.2 Consider a particle in a 3-site 1D chain with periodic boundary conditions. The system is described by the following tight-binding Hamiltonian

$$H = -J \sum_{\langle i,j \rangle} (c_i^{\dagger} c_j + c_j^{\dagger} c_i)$$

where c_i^{\dagger} and c_i are creation and annihilation operators at site i, and J is the hopping amplitude. Assume periodic boundary conditions, meaning site 3 connects back to site 1.

(a) Construct the Hamiltonian matrix in the basis $|1\rangle, |2\rangle, |3\rangle$, where $|i\rangle$ represents the particle localized at site i in the chain. (1 P.)

- (b) Find the eigenvalues of the system by diagonalizing H. (1 P.)
- (c) Determine the eigenvectors. (1 P.)
- (d) Calculate the time evolution operator $U(t) = e^{-iHt}$, expressing it terms of the Hamiltonian eigenvectors. (1 P.)

Numerical Exact Diagonalization

H.3 In this exercise, you will write a small exact diagonalization code to solve a fundamental problem of quantum mechanics: the one-dimensional Heisenberg XXZ model. Consider a system of L spin- $\frac{1}{2}$ particles, subjected to nearest-neighbor interactions, the system is described by the following Hamiltonian,

$$H = J \sum_{j=0}^{L-1} \left(S_j^x S_{j+1}^x + S_j^y S_{j+1}^y + S_j^z S_{j+1}^z \right),$$

with

$$S^x = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad S^y = \frac{1}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad S^z = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix},$$

and J being the interaction strength.

- (a) Write down the computational basis states for L=2 ($2^2=4$ states) and L=3 ($2^3=8$ states). (2 P.)
- (b) Construct the Hamiltonian explicitly as a matrix using the basis from part (a), for both L=2 and L=3. (2 P.)
- (c) Compute the eigenvalues and eigenvectors by numerically diagonalizing the Hamiltonian with J=1 (e.g. numpy.linalg.eigh() or an equivalent function in your preferred programming language), for both L=2 and L=3. (4 P.)

Note: This dense matrix approach is computationally inefficient and becomes impractical for larger system sizes.

(d) Implement a function that takes a time t and returns the complete time evolution operator and returns the full time evolution operator matrix $U(t) = e^{-iHt}$ for L = 2, 3 using the previously determined eigenvalues and -vectors. (2 P.)

Felix Strohling

H.1 - Balver - Campbell - Hausdorff formula

$$O(\mathcal{E}_3) \hookrightarrow$$

Enote: easiest to set E=1 and use BCH-formula, all nested commutators vanish;

but we only considered the validity up to second order for small E)

```
def. f(x) = e^{Ax} e^{Bx}
                                   dfw = AeAxe &x + eAxBe = (A + eAxBe-Ax) f(x)
                [B, A'] = BA - A'B = BA - A'- - A'B = [B, A]A'-+ ABA-- - A'B
                                                              = [BA]A-- + A[BA]An-2 + A2BA--2 - AnB
                                                                                                                                                            Since [A, [A,B]] =0
                                                               = (n-1)[B,A]A^-1 + A^-1BA - A^-1BB
                                                               = nAn-1 [B,A]
            [B, e-Ax] = [B, 2 (-Ax)] = 2 (-1/x) [B, A]
                                                                                = \sum_{n=1}^{\infty} \frac{(-1)^n \times^n}{n!} n + \sum_{n=1}^{\infty} \frac{(-1)^{n-1} \times^{n-1}}{(n-1)!} + \sum_{n=1}^{\infty} \frac{(-1)^n \times^{n-1}}{(n-1)!} + \sum_{n
 → [B, e-Ax] = Be-Ax - e-Ax B = -xeAx [B,A]
       => eAXBeAX = B-X[B,A] = B + X[A,B]
       \Rightarrow \frac{d}{dx} \{(x) = (A + B + x [A, B]) \} (x)
     \Rightarrow \int \{X\} = e^{Ax+Bx} + \frac{1}{2}[ABJx^2] = (A+B) \times e^{\frac{1}{2}[ABJx^2]} = (CABJ) = (CA
     x=1: \(x) = e^A e^B = e^A+B e^{\frac{1}{2}[A,B]}
  125 = [B, A] woll
e) Campbell identity eABe-A = $\frac{1}{4} LABJ, w LABJ = LA, LABJ_{u-1}, LABJ = B
 f(ε) = e<sup>Aε</sup> Bε<sup>-Aε</sup> = Σ f(h)(o) εh
  (E) = AeAEBe-AE + eAEB(-A)e-AE = eAE[A,B]e-AE → (M)(O) = [A,B]
 (min) (E) = dr (eAE [A,B]e-AE) = eAE [A, [A, ..., [A,B]...]e-AE = eAE[A,B], e-AE
> f(E) = Z f(n)(0) Eh > f(N) = Z 1/2 [A,B],
```

4.2- Analytic Exact Diagonalization

Tight-birding Hamiltonian: H = - 3 \(\sum_{\circ_{ij}}\) (c; \(\circ_{j} + c; \circ_{i})\)

Consider a 3-site 10 chain w/ periodic boundary conditions

a) Hamiltonian matrix in basis INO, 127, 130}

H = - 3 { c1 c2 + c2 c1 + c1 c3 + c3 c1 + c2 c3 + c3 c2 }

B det (M-E1) = -E(E2-1)-1. (-E-1)+1. (1+E)

$$= -E^3 + E + 2 + 2E = -E^3 + 3E + 2 = 0$$

$$T. M = E_{\lambda} \vec{v}_{\lambda} = -\vec{v}_{\lambda} \Rightarrow \vec{v}_{\lambda} = \begin{pmatrix} \lambda \\ 0 \\ -\lambda \end{pmatrix}$$

$$T \cdot M\vec{v}_2 = E_2\vec{v}_2 = -\vec{v}_2 \Rightarrow \vec{v}_2 = \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix}$$

$$\boxed{\square} \quad \mathcal{H}_{\mathcal{I}_{\mathbf{z}}} = \mathcal{L}_{\mathbf{z}} \mathcal{I}_{\mathbf{z}} = \mathcal{I}_{\mathbf{z}} \Rightarrow \mathcal{I}_{\mathbf{z}} = \begin{pmatrix} \lambda \\ 1 \end{pmatrix}$$

$$\Rightarrow (I(t) = e^{-itt} = \sum_{u} \frac{(it+t)^{h}}{u!} = \sum_{u} \frac{(it+t)^{h}}{u!} (T+t) = Te^{-itt} dicet T - 1$$

$$\frac{1}{2} = \frac{1}{2} = \frac{1}$$

$$= \frac{1}{3} \begin{pmatrix} 2(e^{idt} - e^{2idt}) & e^{2idt} - e^{-idt} \\ -e^{-idt} - 2e^{2idt} & e^{2idt} + e^{2idt} \\ -e^{-idt} - 2e^{2idt} & e^{2idt} - e^{-idt} \end{pmatrix}$$

$$= \frac{1}{3} e^{-idt} \begin{pmatrix} 2(1 - e^{3idt}) & e^{3idt} - 1 \\ -1 - 2e^{3idt} & e^{3idt} - 1 \end{pmatrix}$$

$$= \frac{1}{3} e^{-idt} \begin{pmatrix} 2(1 - e^{3idt}) & e^{3idt} - 1 \\ -1 - 2e^{3idt} & e^{3idt} - 1 \end{pmatrix}$$

$$= \frac{1}{3} e^{-idt} \begin{pmatrix} 2(1 - e^{3idt}) & e^{3idt} - 1 \\ -1 - 2e^{3idt} & e^{3idt} - 1 \end{pmatrix}$$

H.3 - Numerical Exade Digganalization

One-dimensional Heisenberg XXZ model, L spin-12 particles, NN interactions

$$H = 2 \sum_{r=0}^{j=0} (S_{x}^{j} S_{x}^{j+1} + S_{x}^{j} S_{x}^{j+1} + S_{x}^{j} S_{x}^{j+1})$$

$$(S_i^{\times} \equiv S_{\times_{i}})$$

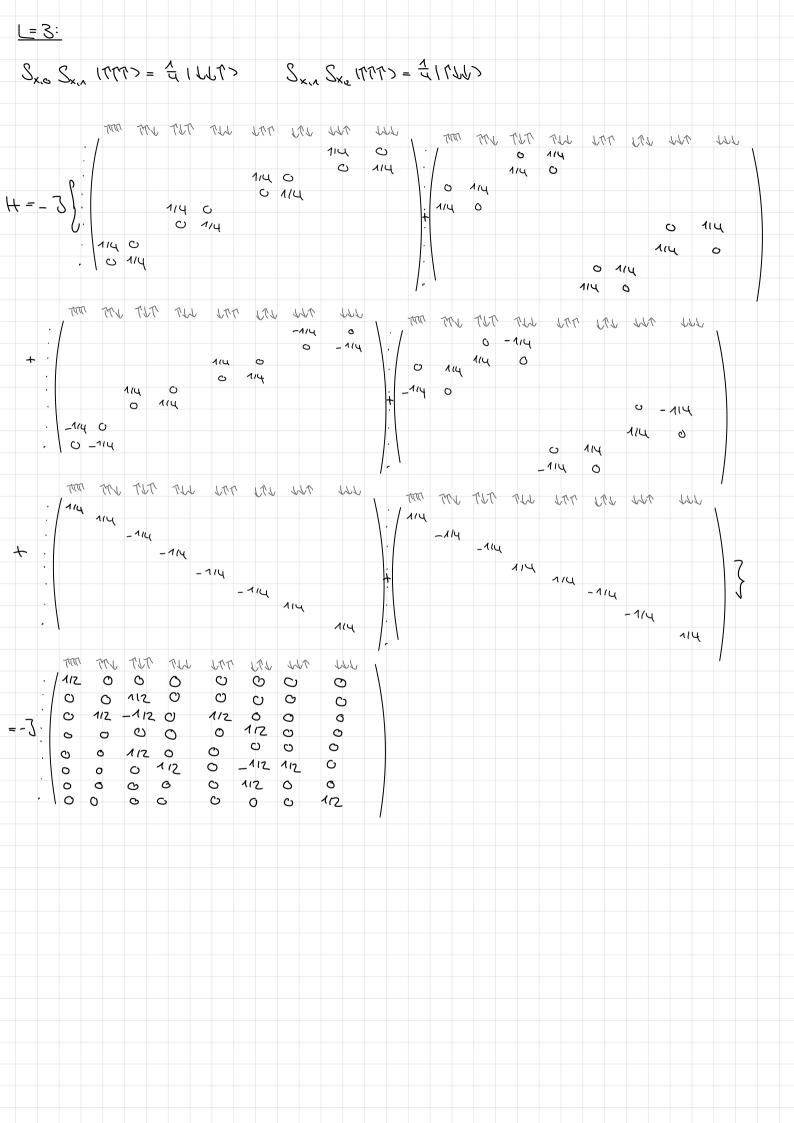
C=3: 17777, (1717), (1717), (1171), (1717), (1717)

6)
$$S_{x}(1) = \frac{1}{2}(1)$$
, $S_{x}(1) = \frac{1}{2}(1)$

$$\underline{\Gamma} = S \cdot S_{x,0} \cdot S_{x,x} \cdot 177 = \frac{1}{4} \cdot 111 \cdot S_{x,0} \cdot S_{x,x} \cdot 117 = \frac{1}{4} \cdot 171 \cdot S_{x,x} \cdot$$

$$S_{x,o}$$
 $S_{x,n}$ $|TU\rangle = \frac{1}{4} |UT\rangle$ $S_{x,o}$ $S_{x,n} |UD\rangle = \frac{1}{4} |TT\rangle$

$$= -\frac{3}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 2 & 0 \\ 0 & 2 & -1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$



TEQS_Exercise_01

April 10, 2025

[1]: import numpy as np

```
1 H.3 a)
[13]: \# L = 2
      H_2 = -1/4 * np.matrix([
          [1, 0, 0, 0],
          [0, -1, 2, 0],
          [0, 2, -1, 0],
          [0, 0, 0, 1]
      ])
      eigenvals2, eigenvecs2 = np.linalg.eig(H_2)
      print("Eigenvalues:", eigenvals2, '\n')
      print("Eigenvectors:\n", eigenvecs2, '\n')
     Eigenvalues: [ 0.75 -0.25 -0.25 -0.25]
     Eigenvectors:
      [[ 0.
                     0.
                                  1.
                                              0.
                                                        ]
      [-0.70710678 0.70710678
                                                       ]
                                 0.
                                             0.
      [ 0.70710678  0.70710678  0.
                                             0.
                                                       ]
      [ 0.
                    0.
                                             1.
                                                       ]]
                                 0.
[14]: \# L = 3
      H_3 = -1/2 * np.matrix([
          [1, 0, 0, 0, 0, 0, 0, 0],
          [0, 0, 1, 0, 0, 0, 0, 0],
          [0, 1, -1, 0, 1, 0, 0, 0],
          [0, 0, 0, 0, 0, 1, 0, 0],
          [0, 0, 1, 0, 0, 0, 0, 0],
          [0, 0, 0, 1, 0, -1, 1, 0],
          [0, 0, 0, 0, 0, 1, 0, 0],
```

```
[0, 0, 0, 0, 0, 0, 0, 1]
      ])
      eigenvals3, eigenvecs3 = np.linalg.eig(H_3)
      print("Eigenvalues:", np.round(eigenvals3, 5), '\n')
      print("Eigenvectors:\n", np.round(eigenvecs3, 5), '\n')
     Eigenvalues: [ 1. -0. -0.5 1. -0.5 0. -0.5 -0.5]
     Eigenvectors:
      [[ 0.
                           0.
                                                                         0.
                                                                                ]
                  0.
                                     0.
                                              0.
                                                       0.
                                                                1.
      Γ0.
                 0.
                          0.
                                  -0.40825 -0.57735 0.70711
                                                               0.
                                                                        0.
                                                                                ]
      [ 0.
                                   0.8165 -0.57735 -0.
                                                                        0.
                                                                               ]
                 0.
                          0.
                                                               0.
      [-0.40825 0.70711 0.57735 0.
                                                                                ]
                                             0.
                                                      0.
                                                               0.
                                                                        0.
                                                                               ]
      [ 0.
                 0.
                          0.
                                   -0.40825 -0.57735 -0.70711
                                                               0.
                                                                        0.
                                                                               ]
      [ 0.8165
                          0.57735 0.
                                             0.
                                                      0.
                                                               0.
                                                                        0.
                 0.
      [-0.40825 -0.70711 0.57735 0.
                                                                        0.
                                                                               ]
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                                                      0.
                                                               0.
      ΓО.
                                                                               11
                 0.
                          0.
                                    0.
                                             0.
                                                      0.
                                                               0.
                                                                        1.
     2 H.3 b)
[19]: # transformation matrix P
      P = [np.matrix(eigenvecs2), np.matrix(eigenvecs3)]
      H_diag = [np.diag(eigenvals2), np.diag(eigenvals3)]
      # time evolution operator
      def U(t, L):
          PL = P[L\%2]
          HL = H_diag[L\%2]
          return PL * np.exp(-1j * HL * t) * PL.T
[22]: # test call
      np.round(U(1, 2), 4)
[22]: matrix([[ 0.9689+0.2474j, 0.
                                       +0.j
                                               , 1.4142+0.j
                1.
                      +0.j
                              ],
              ΓО.
                              , -0.1497-0.2171j, 0.1186+0.4645j,
                      +0.j
                0.
                      +0.j
                              ],
              [ 1.4142+0.j
                              , 0.1186+0.4645j, 1.8503-0.2171j,
                1.4142+0.j
```

```
[ 1. +0.j , 0. +0.j , 1.4142+0.j , 0.9689+0.2474j]])
```

```
Calculations (ignore)
 K.H
a) = = 1 + AE + 2A2 E2 + ...
III. e = 4+ BE + 282 82+ ...
 45 e BE = 1 + BE + 2B2E2 + AE + ABE + 2A2E2 + ...
           = 11 + AE + BE + 2A2E2 + 2B2E2 + ABE2
 V = - 1 [AB] E2 +
 42 e BE e - 1/2 [A,B]E2 = (1/4 + AE + BE + 1/2 A2E2 + 1/2 B2E2 + ABE)(1/4 - 1/2 [A,B]E2)
                       = 11 - 2[ABJE2 + AE + BE + 12A2E2 + 13B2E2 + ABE2
```