

FORMULARIO DI MEMS & MICROSENSORS

the SPRING-MASS-DAMPER SYSTEM

described by:

$$m\ddot{x} + b\dot{x} + kx = F_{ext}$$

↓ L (eaplace)

$$\begin{aligned} T_{XF}(s) &= \frac{x(s)}{F_{ext}(s)} \quad \begin{matrix} \text{TRANSFER FUNCTION} \\ \text{DISPLACEMENT} \\ \text{FORCE} \end{matrix} \\ &= \frac{s/m}{(s^2 + s \frac{\omega_0}{Q} + \omega_0^2)} \\ &= \frac{1/m}{(s^2 + \frac{b}{m}s + \frac{k}{m})} \end{aligned}$$

• for $\omega \ll \omega_0$:

$$|T_{XF}(j\omega)| = \frac{1/m}{\omega^2} = \frac{1}{k}$$

• for $\omega = \omega_0$:

$$|T_{XF}(j\omega)| = \frac{Q/m}{\omega^2} = \frac{Q}{k}$$

• for $\omega \gg \omega_0$:

$$|T_{XF}(j\omega)| = \frac{1/m}{\omega^2}$$

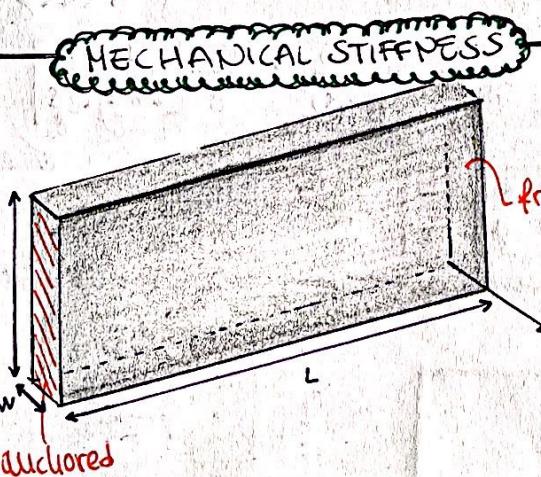
$$\omega_0 = \sqrt{k/m} \quad \begin{matrix} \text{RESONANCE} \\ \text{FREQUENCY} \end{matrix}$$

$$Q = \omega_0 \frac{m}{b} \quad \begin{matrix} \text{QUALITY} \\ \text{FACTOR} \end{matrix}$$

$$\tau = Q / (\pi f_0) \quad \begin{matrix} \text{TIME-CONSTANT} \end{matrix}$$

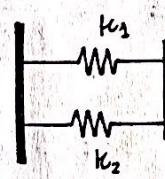
since $b \propto \sqrt{T}$

$$\hookrightarrow Q \propto 1/\sqrt{T}$$

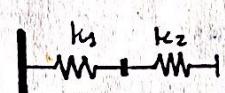


$$k_x = \frac{E}{4} \cdot \frac{h}{l} \left(\frac{w}{l} \right)^3$$

FREE END
+
ANCHORED END
IP MOTION STIFFNESS
YOUNG'S MODULUS



$$k_{TOT \text{ PAR}} = \sum_i k_i \quad \begin{matrix} \text{SPRINGS} \\ \text{IN PARALLEL} \end{matrix}$$



$$\frac{1}{k_{TOT \text{ SER}}} = \sum_i \frac{1}{k_i} \quad \begin{matrix} \text{SPRINGS IN} \\ \text{SERIES} \end{matrix}$$

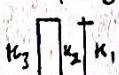
• SPRINGS USED ARE OFTEN GUIDED-END:

as two free springs in series

$$k_{spring} = E h \left(\frac{w}{l} \right)^3$$

GUIDED-
END
MECHANICAL
STIFFNESS
SPRING

• IF A PARAMETER IS NOT ALLOWED WE CAN FOLD SPRINGS:
(to get lower ω).



$$k_{folded} = \frac{1}{N_{FOLD}} k_{spring} \quad \begin{matrix} \text{FOLDED} \\ \text{SPRING} \\ \text{STIFFNESS} \end{matrix}$$

TORSIONAL MECHANICAL STIFFNESS and SPRING-MASS-DAMPER SYSTEM

We define:

- **SHEAR MODULUS G :**

ratio between shear stress and shear strain

- **TORQUE \vec{M} :** $\vec{e} \times \vec{F}$

- **TORSIONAL STIFFNESS k_θ :**

$$k_\theta = G \frac{h}{3} \cdot \frac{W^3}{L} \quad [\text{N} \cdot \text{m}]$$

- **MOMENT OF INERTIA I :** it's a moment of mass with respect to distance from an axis

$$I = \int_M r^2 dm$$

DAMPING COEFFICIENT b

$$b \propto V_{\text{max}} \cdot n_{\text{max}} = \frac{P}{R_B T} \sqrt{k_B T} \propto \sqrt{T}$$

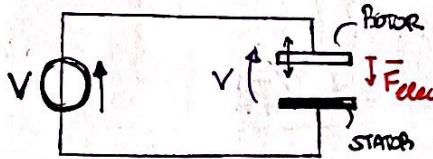
$$b = b_{\text{area}} \cdot 2 \cdot L_{\text{pp}} \cdot H_{\text{pp}} = (\text{const.} \cdot P_0) 2 L_{\text{pp}} H_{\text{pp}}$$

DAMPING COEFF. HEIGHT OF PROCESS LENGTH OF P.P.

HEMS MASS

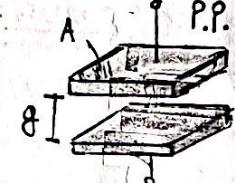
simply exploit the MATERIAL DENSITY ρ and volume dimensions

ELECTROSTATIC FORCES and EQUIVALENT ELECTROSTATIC STIFFNESS



$$F_{\text{elec}} = \frac{V^2}{2} \cdot \frac{\partial C}{\partial x}$$

ELECTROSTATIC FORCE (always directed to inside of capacitor)



- **GAP VARYING CAPACITOR (parallel plate)**

$$C_{\text{pp}} = \epsilon_0 \frac{AN}{g \pm x}$$

$$\frac{\partial C}{\partial x} = \epsilon_0 \frac{AN}{(g-x)^2} = \frac{C_0}{8} \left[\frac{1}{1 + \left(\frac{x}{g}\right)^2} - \frac{1}{g} \right]$$

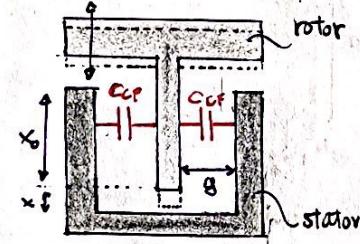
usually $x \ll g$

- **AREA VARYING CAPACITOR (cub-finger)**

$$C_{\text{cf}} = 2 \epsilon_0 N \frac{H(x_0 \pm x)}{g}$$

number geometry

$$\frac{\partial C}{\partial x} = 2 \epsilon_0 \frac{HN}{g}$$



obtained from F_{elec}

$$k_{\text{elec}} = -2V^2 \frac{C_0}{g^2}$$

EQUIVALENT ELECTROSTATIC STIFFNESS (only PP)

IT'S USEFUL TO MOVE THE SIGNAL AWAY FROM FLICKER NOISE and FROM DC OFFSET

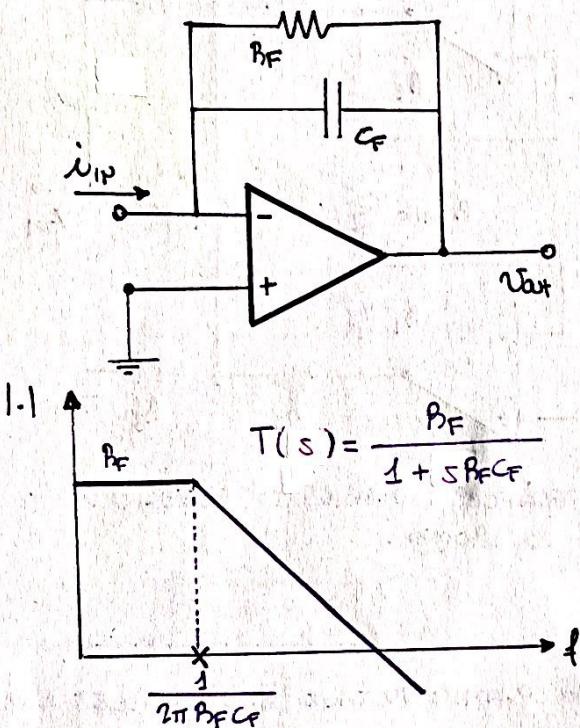
↳ the CAPACITIVE variation is multiplied by a SINUSOIDAL wave at f_{mod}

→ WE ARE MODULATING THE ROTOR VOLTAGE, THAT MODULATES THE CURRENT BUT DOES NOT CHANGE THE TRANSFER FUNCTION OF THE ELECTRONIC VOICE TO THE OUTPUT

After DEMODULATION and LPF, the SIGNAL is moved back to BASE-BAND and LOW FREQUENCY NOISE is SHIFTED AT HIGH FREQ. and then CUT-OFF.

COMMON ELECTRONICS BLOCKS

CHARGE AMPLIFIER and TRANSIMPEDANCE AMPLIFIER



IT'S THE GENERIC TOPOLOGY OF THE FRONT-END STAGE THAT READS OUT THE MOTIONAL CURRENT FLOWING THROUGH THE SENSE ELECTRODE OF A MEMS STRUCTURE

ITS POLE IS TO READOUT THE MOTIONAL CURRENT AND TRANSLATE IT INTO A VOLTAGE

- **TIA:** in a TRANSIMPEDANCE AMPLIFIER, the feedback is DOMINATED by the RESISTANCE. Indeed the pole MUST BE at least a DECADE AFTER the operating system frequency

$$\omega_F = \frac{1}{R_F C_F} \gg \omega_0$$

$$|T_{TIA}(s)| \Big|_{\omega \approx \omega_0} \approx R_F$$

- **CA:** in a CHARGE AMPLIFIER, the feedback is DOMINATED by the CAPACITANCE. Indeed the pole should be at least a decade before the working freq.

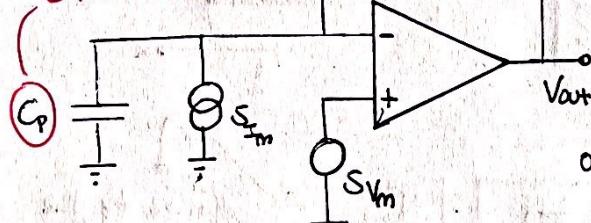
$$\omega_F = \frac{1}{R_F C_F} \ll \omega_0$$

$$|T_{CA}(s)| \Big|_{\omega \approx \omega_0} = \frac{1}{s C_F}$$

- CA can be preferred because the SNR is dependent on $\sqrt{R_F}$ but not on capacitance. Indeed we prefer the highest R_F which happen with low pole. We can make feedback resistance noise negligible
- in addition with TIA we cannot independently choose SNR and the GAIN as they both depend on R_F

TIA and CA - NOISE

PARASITIC CAPACITANCE



for noise contributions:

$$\sqrt{S_{m,tot}} = \sqrt{\sum_i S_{m,i}} \left[\frac{V^2}{Hz} \right]$$

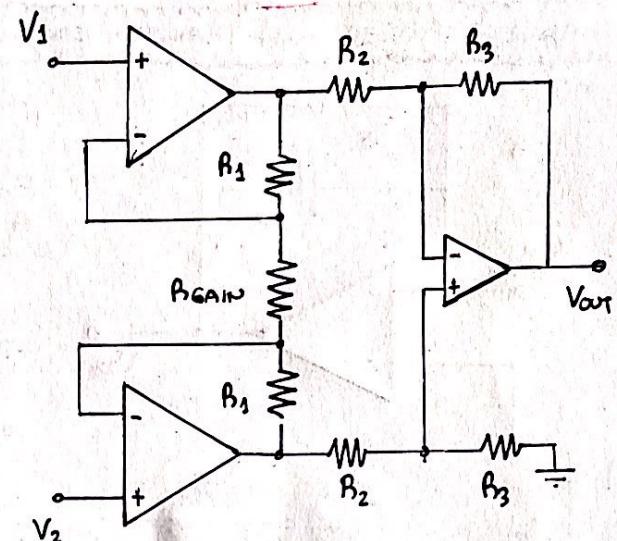
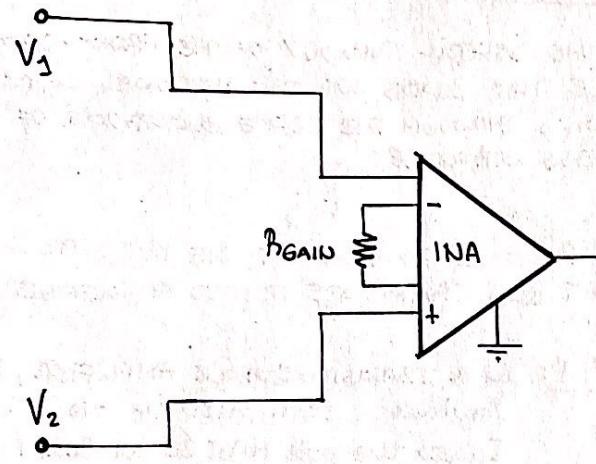
$$S_{Vm}^{out} = S_{Vm} \left(1 + \frac{R_F || \frac{1}{s C_F}}{s C_P} \right)^2 \underset{W R_F C_F \gg 1}{\approx} S_{Vm} \left(1 + \frac{C_P}{C_F} \right)^2 [V^2 / Hz]$$

(by doing $\sqrt{S_{Vm}^{out}} / \text{SENSITIVITY}$ we find the noise in the quantity to be measured)

$$S_{Im}^{out} = S_{Im} \left(\frac{R_F}{1 + s C_F R_F} \right)^2 \underset{W R_F C_F \gg 1}{\approx} S_{Im} \left(\frac{1}{s C_F} \right)^2 \left[\frac{V^2}{Hz} \right]$$

$$\frac{S_{Vm}^{out}}{R_F m} = \frac{4 k_B T}{R_F} \left(\frac{R_F}{1 + s C_F R_F} \right)^2 \underset{W R_F C_F \gg 1}{\approx} \frac{4 k_B T}{R_F} \left(\frac{1}{s C_F} \right)^2 \left[\frac{V^2}{Hz} \right]$$

4 INA - INSTRUMENTATION AMPLIFIER



IT HAS LOW OFFSET AND NOISE, HIGH INPUT IMPEDANCE (THANKS TO BUFFERS).
IT AMPLIFIES A MEASURED ΔV_{out} AND REFER IT TO GROUND.

INA NOISE

INPUT REFERRED INA NOISE IS DUE TO THE TWO OP-AMPS VOLTAGE NOISE SOURCES AND THE THERMAL NOISE OF GAIN RESISTANCE

$$\frac{S_m}{V}_{\text{INA}} = 2 \frac{S_m}{V}_{\text{OP-AMP}} + \frac{S_m}{V}_{\text{RINA}}$$

due:

$$\frac{S_m}{V}_{\text{RINA}} = 4 f_B k_B T R_{\text{GAIN}}$$

GAIN RESISTANCE NOISE

OP-AMP VOLTAGE NOISE

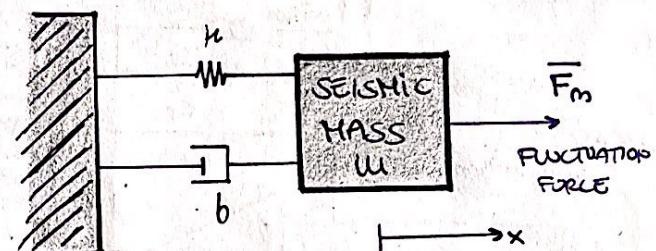
$$\frac{S_m}{V}_{\text{OP-AMP}} = 2 \cdot \frac{4 f_B k T}{g_{\mu\mu}}$$

$= 2/3$
HOSPEL
in SATURATION

two transistors

MECHANICAL NOISE

ANY DISSIPATIVE MECHANISM THAT RESULTS IN A MECHANICAL DAMPING MUST BE BALANCED BY A FLUCTUATION FORCE TO OBTAIN MACROSCOPIC ENERGY BALANCE. IT'S THE FORCE THAT PASS BETWEEN THE MEMS AND THE ENVIRONMENT.



described by
a
NOISE SPECTRUM

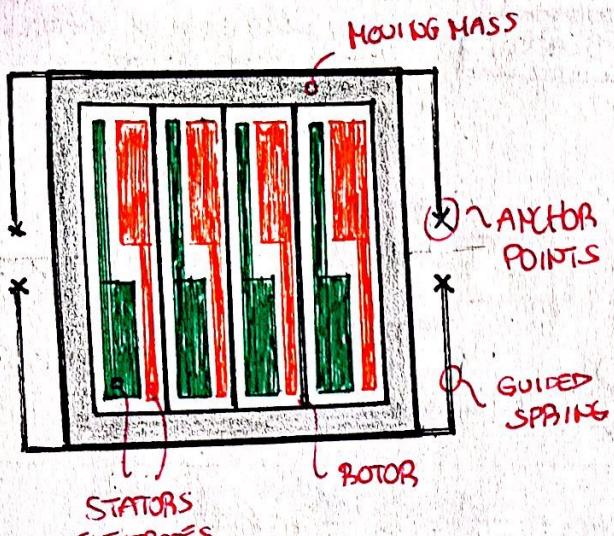
$$S_{F_m} = 4 f_B k_B T b$$

MEMS NOISE POWER SPECTRAL
DENSITY

$$\left[\frac{N^2}{Hz} \right]$$

divide by the
sensitivity²
to obtain
noise in
the quantity
to be measured

ACCELEROMETER



Parallel Plates Accelerometers

IT'S A SPRING-MASS-DAMPER SYSTEM described by:

$$m\ddot{x} + b\dot{x} + K_{TOT}x = m\ddot{a}_{ext}$$

where

$$K_{TOT} = K_{mech} + K_{elec} \quad (<0)$$

an ACCELEROMETER can be

- PARALLEL PLATE / COMB-FINGER
- IN-PLANE MOTION / OUT-OF-PLANE MOTION

it usually works at LOW FREQUENCIES ($\omega \ll \omega_0$), indeed

$$|T_{xp}(j\omega)| \approx 1/K_{TOT}$$

THE MECHANICAL SENSITIVITY IS:

$$S_{mech} = \frac{dC}{da} = (\text{?}) \frac{dx}{da} \cdot \frac{dc}{dx}$$

MECHANICAL SENSITIVITY

DISPLACEMENT TO CAPACITIVE GAIN

DIFFERENTIAL GAIN

ACCELERATION TO DISPLACEMENT GAIN

DISPLACEMENT v ACCELERATION

$\frac{dx}{da}$ term is given by the spring-mass-damper system description. Usually the ACCELEROMETER working region is $\omega \ll \omega_0$ then

$$\left| \frac{X(s)}{F(s)} \right| = \left| \frac{X(s)}{m\ddot{a}(s)} \right| \approx \frac{1}{K_{TOT}}$$

$$\frac{dx}{da} = \frac{m}{K_{TOT}} = \frac{1}{\omega_0^2}$$

$$S_{mech} = 2 \cdot \frac{C_0}{g} \cdot \frac{1}{\omega_0^2}$$

PARALLEL-PLATE

$$S_{mech} = 2 \cdot \frac{C_0}{\omega_0} \cdot \frac{1}{\omega_0^2}$$

COMB-FINGER

CAPACITIVE VARIATION

$\frac{dC}{dx}$ term depends on the type of ACCELEROMETER

• PARALLEL PLATE:

$$= \frac{C_0}{g}$$

• COMB-FINGER:

$$= \frac{C_0}{\omega_0}$$

NOTE THAT COMB-FINGER:

- IS good BECAUSE THE ELECTROSTATIC FORCE IS NOT FUNCTION OF DISPLACEMENT (NO POLE-IN-PATH) and

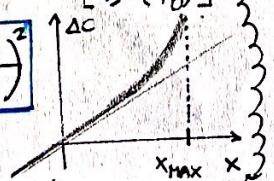
NO NON-LINEARITY IN CAPACITIVE READOUT
NO TRADE-OFF SENSITIVITY v FSR

- IS bad BECAUSE WITH SAME PARAMETERS THE NUMBER OF FINGERS WHICH CAP FIT IS LOWER

the real capacitance variation is not $2 \frac{C_0}{g} x$ but $C_0 \left[\frac{2x/g}{1-(x/g)^2} \right]$

$$\epsilon_{\text{err}}(x) = \frac{\Delta C_{\text{real}} - \Delta C_{\text{lin}}}{\Delta C_{\text{real}}} \rightarrow \epsilon_{\text{err}}^{\text{MAX}} = 100 \left(\frac{x_{\text{MAX}}}{g} \right)^2$$

PP LINEARITY ERROR



The non-linearity is:

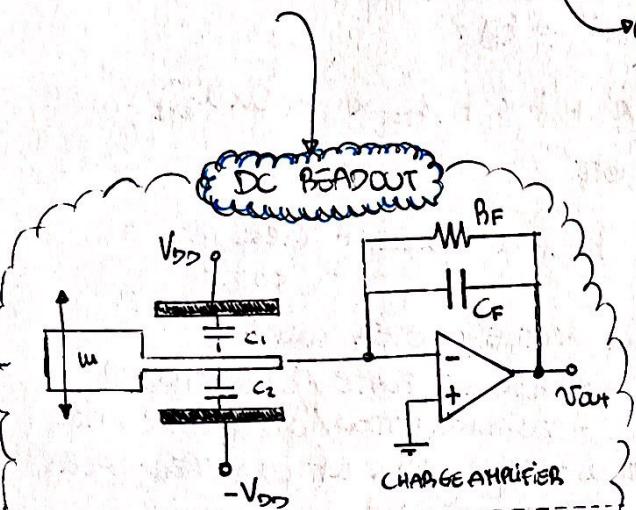
$$\epsilon_{\text{err}}(x) = \frac{\Delta C_{\text{real}} - \Delta C_{\text{lin}}}{\Delta C_{\text{real}}} \rightarrow \epsilon_{\text{err}}^{\text{MAX}} = 100 \left(\frac{x_{\text{MAX}}}{g} \right)^2$$



6

READOUT CONFIGS

THERE ARE TWO CONFIGS



$$i_{C_i} = \frac{dQ_i}{dt} = C_i \frac{dV_i}{dt} + V_i \frac{dC_i}{dt} = V_i \frac{dC_i}{dt}$$

$$V_{out} = -i \frac{1}{SCF} = -[V_{DD} \pm C(s)] \frac{1}{SCF}$$

$$\frac{dV_{out}}{dC} = \frac{V_{DD}}{C_F}$$

CAPACITIVE VARIATION TO VOLTAGE GAIN (DC)

CANNOT MEASURE DC ACCELERATIONS.
(WORKING REGION, AFTER RFFC)
POLE; WE NEED C.A.

$$S_{TOT} = \frac{dV_{out}}{da} = \frac{dV_{out}}{dC} \text{ such asy if } \text{ MODULATION APPROACH}$$

$$S_{TOT}^{PP} = \frac{2}{2} \frac{V_{DD}}{C_F} \cdot \frac{C_0}{2} \cdot \frac{1}{\omega_0^2}$$

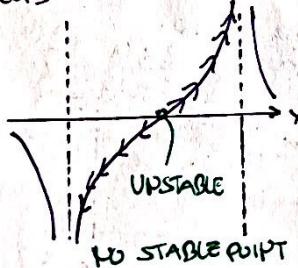
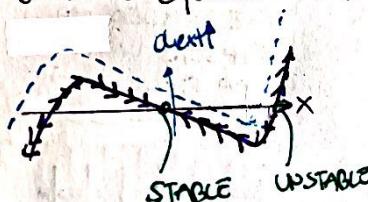
$$S_{TOT}^{CF} = \frac{2}{2} \frac{V_{DD}}{C_F} \cdot \frac{C_0}{2} \cdot \frac{1}{\omega_0^2}$$

OVERALL SENSITIVITY

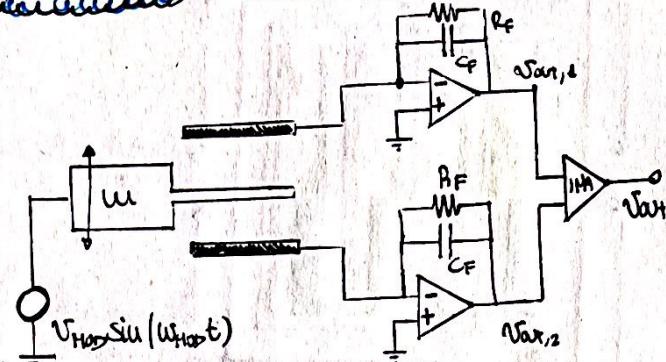
ONLY IF WE HAVE DIFFERENT READOUT

PULL-IN AND STABILITY

we can analyze the system total force
wrt its equilibrium points



IT CONSIST IN A HIGH-FREQUENCY MODULATION OF THE SUSPENDED MASS, WITH EACH OF STATORS KEPT IN VIRTUAL GROUND



$$i_{C_i} = \frac{dQ_i}{dt} = C \frac{dV}{dt} + V \frac{dC}{dt}$$

by assuming $C = C_0 + C_a [\cos(\omega_a t)]$
(ω_a : acceleration freq.)

$$\frac{C \frac{dV}{dt}}{V \frac{dC}{dt}} \approx \frac{\omega_{HOI}}{\omega_a}$$

if $\omega_{HOI} \gg \omega_a$
 $V \frac{dC}{dt}$ is negl.

$$V_{out,1} = -\frac{1}{C_F} \int C \frac{dV}{dt} dt = -\frac{C_0 \pm \Delta C}{C_F} V_{DD} \sin(\omega_a t)$$

$$\frac{dV_{out}}{dC} = \frac{V_{HOI}}{C_F} \sin(\omega_a t)$$

CAPACITIVE TO VOLTAGE GAIN (AC)

CAN MEASURE DC ACCELERATIONS

- MOVE AWAY FROM FILTER NOISE AND DC OFFSET

- NEED TO DEMODULATE AND LPF TO CUT-OFF NOISE SHIFTED AT HIGH FREQ.

- $V_{DD} = V_{HOI}$

IT COULD HAPPEN THAT THERE ARE NO STABLE POINT AND THE MOTOR PLATES SNAP UP TO THE STATOR ONES (PULL-IN CONDITION).

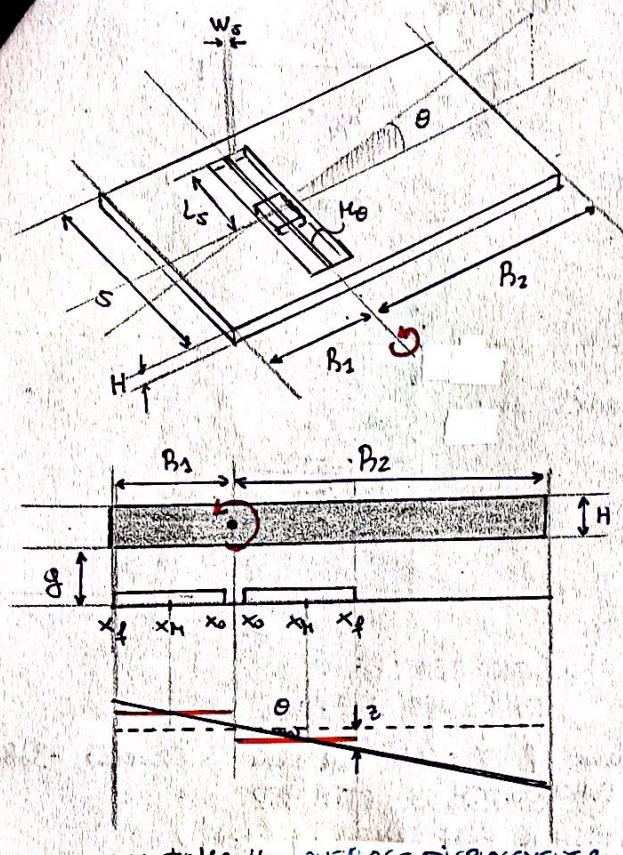
WE NEED:

$$V_{DD} < V_{PULL-IN} = \sqrt{\frac{g^3 k}{2 \epsilon_0 A N}}$$

NO PULL-IN CONDITION

$$|k_{mech}| > |k_{electr}|$$

1-OF-PLANE ACCELEROMETER



We take the AVERAGE DISPLACEMENT z for CAPACITIVE VARIATIONS (gap) in x_M

$$C_i = \frac{\epsilon_0 (x_f - x_o) \cdot \epsilon_{pp}}{g \pm z}$$

$$z = x_M \tan(\theta) \approx x_M \theta$$

$$S_{\text{mech}} = 2d \frac{C_0}{g} \cdot x_M \cdot \frac{1}{\omega_0^2} \frac{\partial u}{I}$$

MECHANICAL SENSITIVITY

$$\theta = \frac{1}{\omega_0^2} \frac{H}{I} = \frac{1}{\omega_0^2} \frac{\partial x(u)}{I}$$

comes from

$$K_{\text{elec}} = -2d \frac{C_0}{g} V_{DD}^2 x_M^2$$

$$K_{\text{mech}} = 2G \frac{H}{3} \frac{W_s^3}{E_s}$$

STIFFNESS K_{TOT}

the MOMENT OF INERTIA IS

$$I = I_1 + I_2 = \int_{u_1} r^2 du + \int_{u_2} r^2 du$$

$$du = SHP dt$$

$$I = \frac{\beta_1^2 u_1 + \beta_2^2 u_2}{3}$$

MOMENT OF INERTIA

MECHANICAL OFFSET

DUE TO PROCESS INEQUALITIES IN FABRICATION PROCESS

- GAPS OF TWO DIFFERENTIAL PP
- RESIDUAL MECHANICAL STRESS IN WAFER BENDING

IT'S CALLED 260-ZERO G OUTPUT

$$260(T) = \frac{x_{os}(T)}{\frac{dx}{da_{ext}}}$$

THE SYSTEM IS DESCRIBED BY:

$$I \ddot{\theta} + b \dot{\theta} + k_{\text{tot}} \theta = M_{\text{ext}} \text{ and } |T_{\text{ext}}| \approx \frac{1}{k_{\text{tot}}} \frac{1}{\omega_0^2}$$

$$\Delta C_{\text{real}} = \epsilon_0 \frac{A_{\text{PP}}}{g-z} - \epsilon_0 \frac{A_{\text{PP}}}{g+z} = C_0 \frac{2 z/g}{1 - (z/g)^2} \alpha$$

$$\Delta C_{\text{lin}} = \Delta C_{\text{real}} \Big|_{z \ll g} \approx 2 C_0 \frac{z}{g} \alpha$$

$$\epsilon_{\text{eli}} \% = 100 \left(\frac{z_{\text{FSR}}}{g} \right)^2 = 100 \left(\frac{x_M \cdot \theta_{\text{FSR}}}{g} \right)^2$$

LINEARITY ERROR

$$M_{\text{FSR}} = M_1 + M_2 = \bar{P}_1 \times \bar{F}_1 + \bar{P}_2 \times \bar{F}_2 = - \frac{P_1}{2} u_1 F_{\text{FSR}} g + \frac{P_2}{2} u_2 F_{\text{FSR}} g$$

force applied at well

$$\theta_{\text{FSR}} = \frac{M_{\text{FSR}}}{k} = \frac{1}{\omega_0^2} \cdot \frac{M_{\text{FSR}}}{I}$$

$$\omega_0 = \sqrt{\frac{k}{I}}$$

$$\Delta C_{\text{FSR}} = 2d \frac{C_0}{g} z_{\text{FSR}}$$

HOLES TRANSDUCTION COEFF.

$$M_{\text{elec}} = 2F_{\text{elec}} x_M = 2d \frac{C_0}{g^2} V_{DD}^2 z x_M \\ = 2d \frac{C_0}{g^2} V_{DD}^2 \theta x_M^2$$

$$K_{\text{elec}} = -2d \frac{C_0}{g^2} V_{DD}^2 x_M^2$$

$$K_{\text{mech}} = 2G \frac{H}{3} \frac{W_s^3}{E_s}$$

2x (two springs)

IT ADDS NON-LINEARITY

$$\Delta C_{\text{diff}} = C_0 \frac{x}{g} \left[\frac{-1}{1 + \frac{x}{g-x_0}} + \frac{1}{1 - \frac{x}{(g+x_0)}} \right]$$

To correct it we can

- apply WAFER-LEVEL CALIBRATION
- apply IN-MOTION CALIBRATION

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NEAD

it's the **NOISE EQUIVALENT ACCELERATION DENSITY**

$$\text{NEAD} = \sqrt{S_{Am}} = \frac{\sqrt{S_{fm}}}{m} \left[\frac{m/s^2}{\sqrt{Hz}} \right]$$

$$= \sqrt{\frac{4 k_B T b}{m}} = \sqrt{\frac{4 k_B T_{iso}}{m Q}}$$

for overall contribution we need bandwidth BW

$$\bar{a}_m = \sqrt{S_{Am} \cdot BW}$$

$$= \text{NEAD} \cdot \sqrt{BW}$$

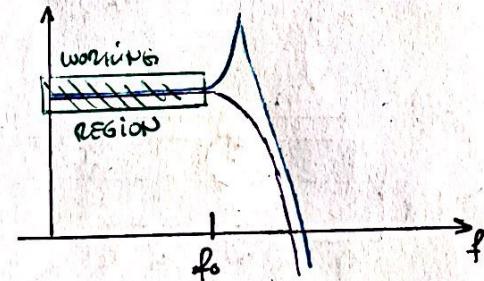
CHOICE OF QUALITY FACTOR Q

the **QUALITY FACTOR Q** is a key factor; it appears in:

- NEAD (HIGH Q, low noise)
- TIME RESPONSE
- SPECTRUM TRANSFER FUNCTION

IF MEMS NOISE DOMINATES \rightarrow HIGH Q

IF ELECTRONIC NOISE DOMINATES \rightarrow $Q \approx 0.5$



\rightarrow REQUIRED LPF

TO LOWER THE PEAK

\rightarrow because if $Q < 0.5$ the gain is lowered -3dB and BANDWIDTH IS OPTIMIZED
(if HIGH Q feed region is less!)

TRADE-OFFS

about SENSITIVITY:

$$S_{tot} = 2 \frac{V_{dd}}{C_F} \cdot \frac{C_o}{g} \cdot \frac{m}{(K - 2V_{dd} \frac{C_o}{g^2})}$$

- $g \downarrow$: $S_{tot} \uparrow$ PULL-IN RISK \uparrow

- $m \uparrow$: $S_{tot} \uparrow$ AREA \uparrow

BANDWIDTH \downarrow

PULL-IN RISK \uparrow

BANDWIDTH \downarrow

PULL-IN RISK \uparrow

POWER CONSUMPTION \uparrow

AREA, m, $g \uparrow$

- $k_{tot} \downarrow$: $S_{tot} \uparrow$

- $V_{dd} \uparrow$: $S_{tot} \uparrow$

- H: $S_{tot} \uparrow$

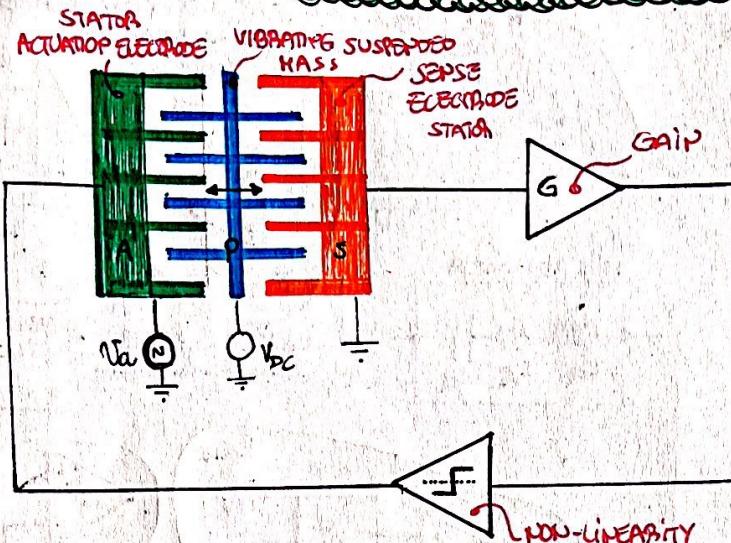
about BANDWIDTH:

NO PULL-IN $k_{tot} \uparrow$, $f_0 \uparrow$ SENSITIVITY \downarrow

$\alpha_{MAX} \uparrow$, $W_0^2 \uparrow$ SENSITIVITY \downarrow (because $\alpha_{MAX} = \alpha_{MAX} W_0^2$)

BANDWIDTH-NOISE (about Q)

RESONATORS and OSCILLATORS



PARAMETERS

N_{CF} : number of rotator fingers

g : gap

L_{ov} : finger overlap at rest

m : mass

k : stiffness

b : damping coeff.

H : process height

A : $H \cdot L_{ov}$ area

ELECTROSTATIC FORCES

BY CHOOSING

$$V_p = V_{dc}$$

$$V_A = V_a \sin(\omega t)$$

$$V_s = 0V$$

$$|F_{elec}| = \frac{V^2}{2} \cdot \frac{\partial C}{\partial x}; \text{ for COMB FINGER } \frac{\partial C}{\partial x} = \epsilon_0 \frac{N_{CF} H}{8} = \frac{C_0}{x}$$

$$\begin{aligned} |F_{elec, tot}| &= |F_{elec, 1} - F_{elec, 2}| \\ &= \left| \frac{1}{2} \cdot \frac{\epsilon_0 N_{CF} H}{8} \left[(V_{dc} - V_a \sin(\omega t))^2 - V_{dc}^2 \right] \right| \\ &= \frac{1}{2} \epsilon_0 \frac{N_{CF} H}{8} \left[2 V_{dc} V_a \sin(\omega t) + V_a^2 \sin^2(\omega t) \right] \\ &\approx \frac{\epsilon_0 N_{CF} H}{8} V_{dc} V_a \sin(\omega t) \quad \text{for } V_a \ll V_{dc} \end{aligned}$$

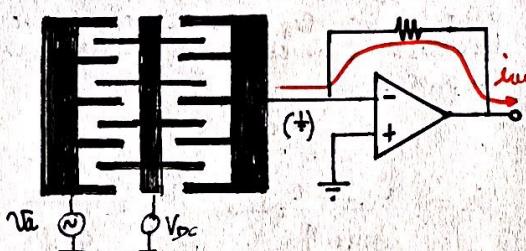
ACTUATION

TRANSDUCTION FACTOR

$$M_A = \frac{F_{elec}(s)}{V_a(s)} = V_{dc} \cdot \frac{dC_A}{dx}$$

SENSE PORT READOUT

IN ORDER TO READOUT ROTOR DISPLACEMENT AND CALIBRATE THE LOOP GAIN WE NEED TO SENSE THE CURRENT, i_{ss} , CAUSED BY CAPACITANCE VARIATION



$$i_{ss} = V_{dc} \cdot \frac{dc_s}{dt} = V_{dc} \frac{dc_s}{dx} \cdot \frac{dx}{dt} = V_{dc} \frac{dc_s}{dx} \dot{x}$$

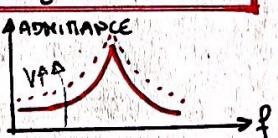
$$M_S = \frac{i_{ss}(s)}{s x(s)} = V_{dc} \frac{dc_s}{dx} \quad \text{ELECTROMECHANICAL TRANSDUCTION FACTOR}$$

$$i_{ss}(s) = \frac{1}{V_a(s) \left(L_{eq}s + R_{eq} + \frac{1}{s C_{eq}} \right)}$$

(rate) $\frac{1}{s C_{eq}}$ L_{eq} P_{eq} C_{eq}
 $L_{eq} = m/m^2$ $P_{eq} = b/m^2$ $C_{eq} = m^2/k$

$$L_{eq} = m/m^2 \quad P_{eq} = b/m^2 \quad C_{eq} = m^2/k$$

EQUIVALENT CIRCUIT OF RESONATOR



ELECTRICAL MODEL

CONSIDERING $M_A = M_S$ WE FIND:

$$\frac{x(s)}{F_{elec}(s)} = \frac{1}{m s^2 + b s + k} \quad (\text{ADMITTANCE})$$

$$\left(\frac{i_{ss}(s)}{V_a(s)} = M_S^2 \frac{1/s}{(s^2 + b/m s + k/m)} = M_S^2 s \frac{1/m}{(s^2 + s \omega_0^2 + \omega_0^2)} \right)$$

$$\omega \ll \omega_0 : T_{Ia} = \frac{m^2}{k} s \rightarrow \infty$$

$$\omega = \omega_0 : T_{Ia} = m^2/b \rightarrow \infty$$

$$\omega \gg \omega_0 : T_{Ia} = \frac{m^2}{k} \cdot \frac{1}{s} \rightarrow 0$$

BARKHAUSEN CRITERIA and SUSTAINING CIRCUIT

IN ORDER TO HAVE A STABLE OSCILLATION
THE BARKHAUSEN CRITERIA SHOULD
BE MET :

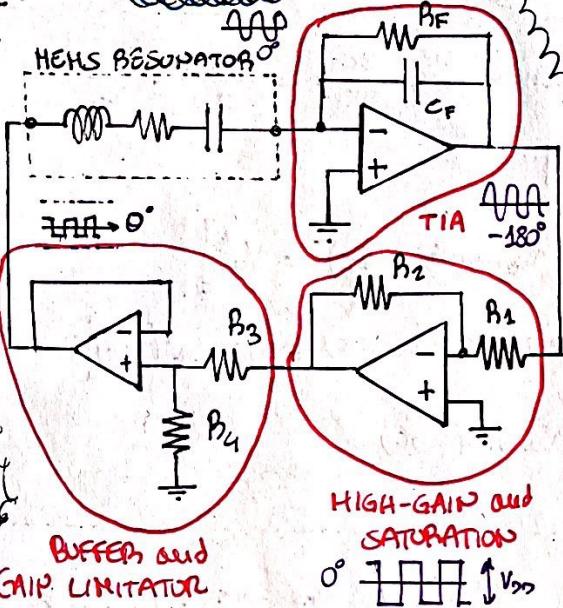
$$|G_{\text{loop}}(j\omega_0)| = 1$$

$$\times [G_{loop}(j\omega_0)] = 360^\circ$$

BABYHAUSEN CRITERIA

TRANSIMPEDANCE AMPLIFIER

HEHS RESPIRATOR



$$\circ \boxed{\text{TIA}} : G_{\text{TIA}}(j\omega) \Big|_{\omega \approx \omega_0} \simeq R_F \quad ($$

R_F dominates over C_F

$$w_f = \frac{1}{\beta e(C)} \rightarrow w_0$$

HIGH GAIN AND SATURATION

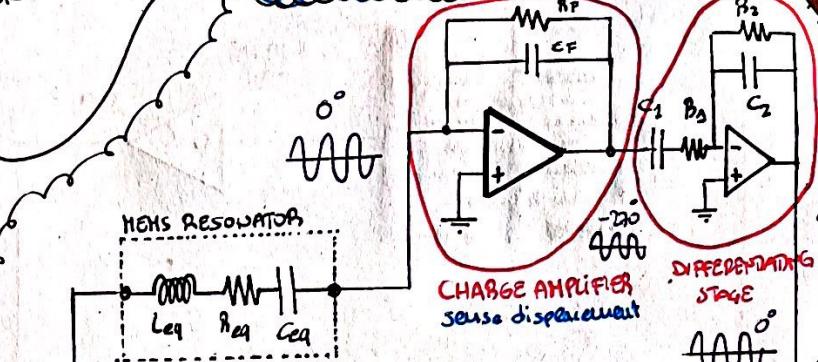
$$G_{HG} = - \frac{B_2}{B_1}$$

OUTPUT SATURATES at $\pm V_{DD}$

• BUFFER and GAIN LIMITATOR

CA is to be preferred than TIA
because of S/N & $\sqrt{P_{RF}}$

CHARGE AMPLIFIER CONFIG.



The diagram illustrates a circuit stage consisting of two main sections: a 'BUFFER + GAIN LIMITATOR' and a 'HARD-LIMITER'. The 'BUFFER + GAIN LIMITATOR' section on the left contains a operational amplifier (op-amp) configured as a voltage follower. Its output is connected to the non-inverting input (+) of a second op-amp, which is labeled 'HARD-LIMITER'. The inverting input (-) of the 'HARD-LIMITER' is grounded. The output of the 'HARD-LIMITER' is shown as a square-wave signal. Two resistors, R_{DG1} and R_{DG2} , are connected between the output of the first op-amp and the inputs of the 'HARD-LIMITER' op-amp.

- **CHARGE AMPLIFIER**: sense the displacement

$$G_{CA}(j\omega) = \frac{1}{j\omega C_F} \quad \Phi = -180^\circ - 90^\circ$$

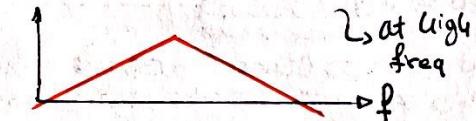
$\approx -270^\circ$

add $\Phi = 360^\circ - \tan^{-1}(B_0/f_m)$
 - **DIFFERENTIATING STAGE**: adjust the phase by adding -90° (total reaches -360°)

$$CA(j\omega) \Big|_{\omega \approx \omega_0} = \frac{1}{j\omega C_F}$$

Differentiating Stage: adjust the phase by adding -90° (totae reads -360°)

$$G_{\text{DF}}(s) = \frac{-sC_1R_2}{(s+sC_1R_1)(s+sC_2R_2)}$$



- **HARD-CLIPPER**: High voltage, saturates at $\pm V_{op}$. It's output cannot be used to directly drive the resonator because it would violate small signal hypothesis.

$$\frac{V_a}{g} \ll V_{DC}$$

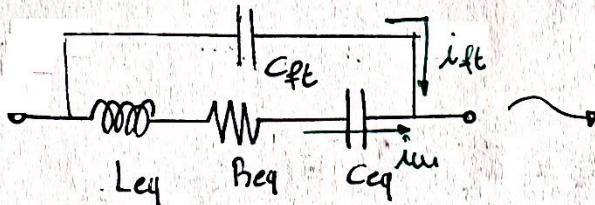
$$N_a = \frac{x_{a \text{ max}}}{Q_{\text{max}} M} \rightarrow N_{a_{sq}} = \frac{N_a}{4/\pi}$$

- ## • BUFFER and GAIN LIMITATOR

$$G_{DG} = \frac{V_{DD}}{V_{DD} - V_{GS}} = \frac{R_{DG2}}{R_{DG1} + R_{DG2}}$$

$$G_{loop}(j\omega_0) = \frac{1}{P_{eq}} G_{CA} G_{diff} G_{TL} G_{DG}$$

EFFECTS OF FEEDTHROUGH CAPACITANCE



TRANSFER FUNCTION WITH FEEDTHROUGH CAPACITANCE

$$\frac{i(s)}{V_{in}(s)} = \frac{1}{sL_{eq} + B_{eq} + \frac{1}{sC_{eq}}} + sC_{ft}$$

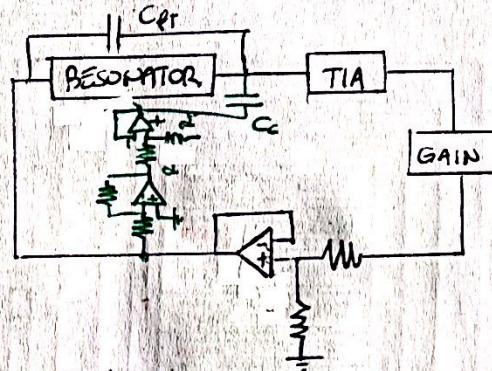
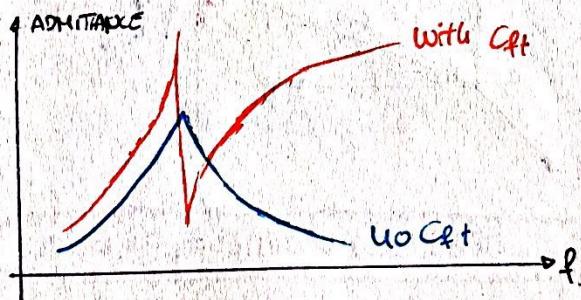
→ we want to lower the B_{eq} right after the resonance peak

- we can solve it with a 2nd pole near after the peak we can avoid other freq. to be sustained
- introduce an opposite capacitance to the feedthrough contribution using a suitably capacitive gain excited by a current in antiphase to feedthrough signal

$$i_C = C_C \frac{dV_{in}}{dt} \left(-\frac{B_T}{B_T + B} \right)$$

$$C_{ft} = C_C \left(\frac{A_T}{B_T + B} \right)$$

try to cancel with i_{CC}
lift



TUNING FREQUENCY

WE TRY TO EXPLOIT ELECTROSTATIC SOFTENING TO DOWNTUNE WITH PARALLEL PLATES

N_{pp} : number of pp

C_{pp} : overall capacitance

g_{pp} : gap

V_{tun} : drive tuning

$$\Delta \omega = \sqrt{\frac{k_{TOT}}{m}} - \sqrt{k - 2(V_{dc} - V_{tun})^2 \frac{C_{pp}}{g_{pp}^2}}$$

$$= \sqrt{\frac{|k_{act}-k_{el}|}{m}} - \sqrt{\frac{k_{act} - |k_{act} + k_{tun}|}{m}}$$

$$k_{tun} = 2V_{tun}^2 \frac{C_{pp,tun}}{g_{pp}^2}$$

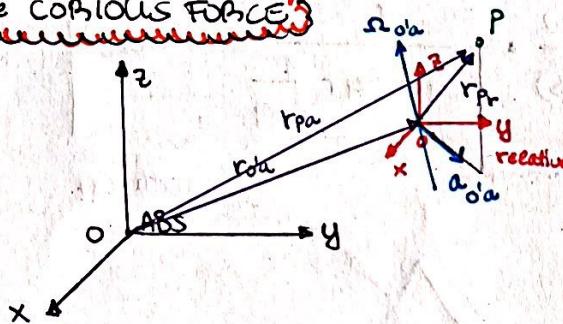
FREQ. SHIFT

TRADE-OFFS

- an HIGH Q could be undesirable as electric gain could be lower
- an HIGH BIAS VOLTAGE on BOTOB would increase η but also increases power consumption
- we need low displacements to avoid non-linearities

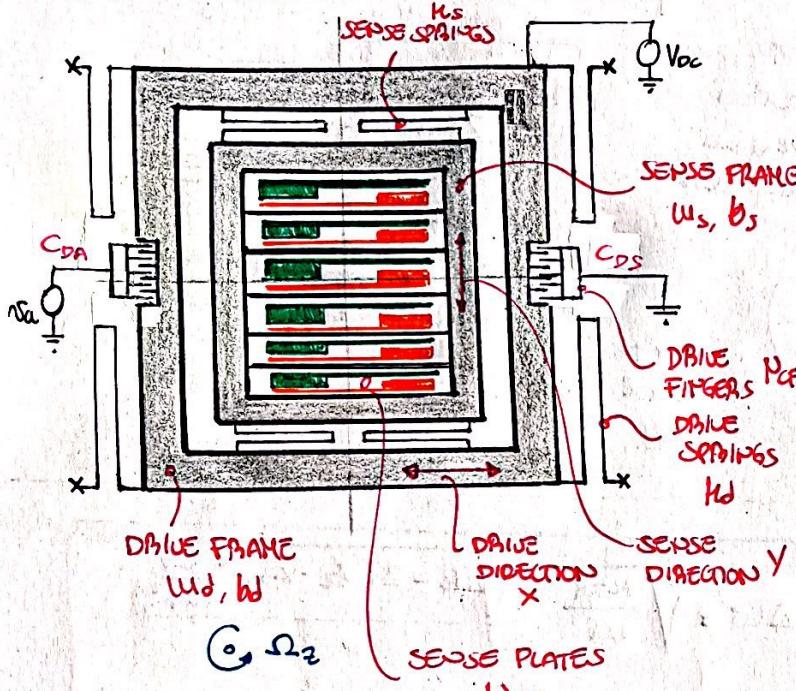
GYROSCOPES

the Coriolis Force



$$\vec{a}_{Pa} = \vec{a}_{Pr} + \vec{a}'_{da} + \cancel{(\Omega_{da} \times \vec{v}_{Pr})} + \cancel{\Omega_{da} \times (\vec{a}'_{da} + \vec{v}_{Pr})} + 2(\vec{\omega}_{da} \times \vec{v}_{Pr}) \text{ CORIOLIS ACCELERATION}$$

ARCHITECTURE and WORKING PRINCIPLE



there are two motion directions of interest

- **DRIVE MODE**: along which the device oscillates creating a desired \dot{x}_d
- **SENSE MODE**: direction along which device moves only under angular rates

it's a combination of
 COMB DRIVEN + PARALLEL PLATE
 RESONATOR ACCELEROMETER

MODE MATCHED GYROSCOPE

$$\left. \begin{aligned} C_{da} &= \frac{2\epsilon_0 H (L_{av} - x) N_{cf}}{8} \\ C_{dd} &= \frac{2\epsilon_0 H (L_{av} + x) N_{cf}}{8} \end{aligned} \right\}$$

$$|F_{eexc}| \approx \frac{\epsilon_0 H N_{cf}}{8} [2V_{dc} \omega_{av} \sin(\omega t)]$$

$$Q = \frac{\omega}{\omega_0} \quad \Delta \omega_{BW} = \frac{\omega_0}{2Q}$$

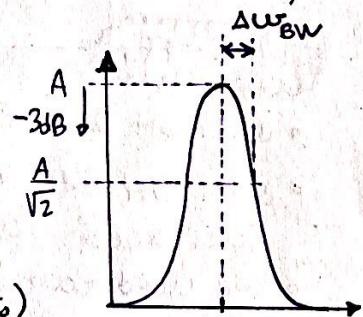
$$x_{D,0} = F_{eexc,0} \quad \frac{Q_D}{k_D} = \frac{Q_D}{k_D} \cdot \frac{\epsilon_0 H N_{cf}}{8} 2V_{dc} \omega_a$$

$$\omega_D = \dot{x}_D \rightarrow \omega_{D,0} = \omega_0 x_{D,0}$$

$$F_{Corr} = M_s \cdot a_{Corr} = -2M_s \omega_D \Omega \quad (\text{modulated at } \omega_0)$$

$$\rightarrow y_{S,D} = F_{Corr,0} \cdot \frac{Q_S}{k_S} = 2M_s x_{D,0} \omega_0 \frac{Q_S}{k_S} \Omega \Rightarrow$$

$$\frac{y_{S,D}}{\Omega} = \frac{x_{D,0}}{b_s / 2M_s} = \frac{x_{D,0}}{\Delta \omega_{BW}}$$



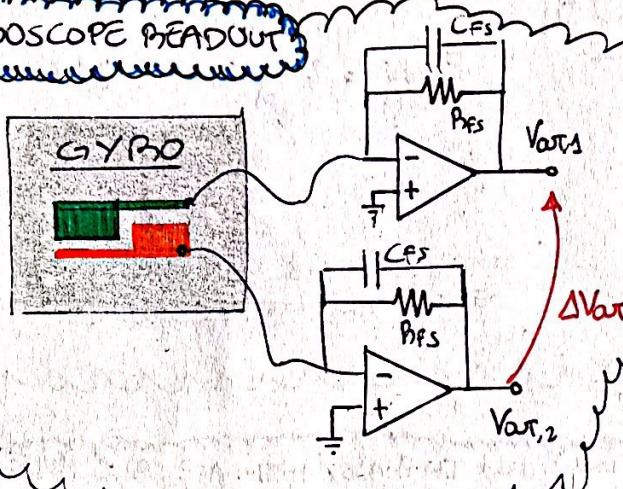
We need to find $\frac{\partial C_s}{\partial y}$:

WITH PARALLEL PLATES SENSING:

$$C_s = \frac{\epsilon_0 A_{\text{plate}}}{g \pm y} \rightarrow \frac{\partial C_s}{\partial y} \approx \frac{C_s}{g} \quad (y \ll g_s)$$

$$S_{\text{mech}} = 2 \frac{dC}{dy} \cdot \frac{dy}{d\omega} = 2 \cdot \frac{C_s}{g} \cdot \frac{x_{D,0}}{\Delta \omega_{\text{BW}}} \quad \text{MECHANICAL SENSITIVITY}$$

GYROSCOPE READOUT



$$S = \frac{\Delta V_{\text{out},0}}{1 \Omega} = \frac{V_{DC}}{CFS} \cdot S_{\text{mech}}$$

OVERALL SENSITIVITY

DRIVE MODE SYSTEM OVERVIEW

IT MAKES THE DRIVE FRAME OSCILLATE AND SETS $x_{D,0}$ WHICH APPEARS IN S_{mech}

$x_{D,0}$ SHOULD BE STABLE

HOW ABOUT TEMPERATURE CHANGES?

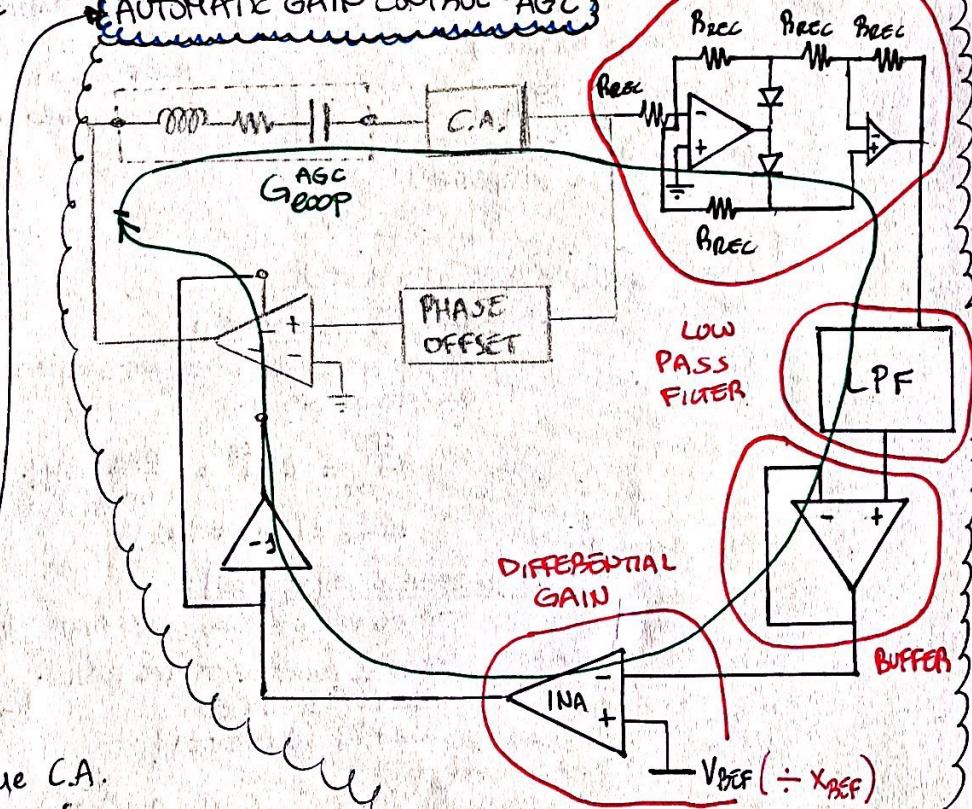
$$Q \propto 1/\sqrt{T}$$

$$\frac{\Delta Q}{Q} = -\frac{1}{2} \frac{\Delta T}{T}$$

We need a circuit to drive correctly $x_{D,0}$ to a reference value

If $x_0 > x_{0,\text{ref}}$ the CA output voltage will be higher. It will be compared with V_{REF} ($> V_{\text{REF}}$) and the voltage biasing the comparator will be lowered, reducing x_0 until...

AUTOMATIC GAIN CONTROL - AGC



DIFFERENTIAL GAIN

IN A

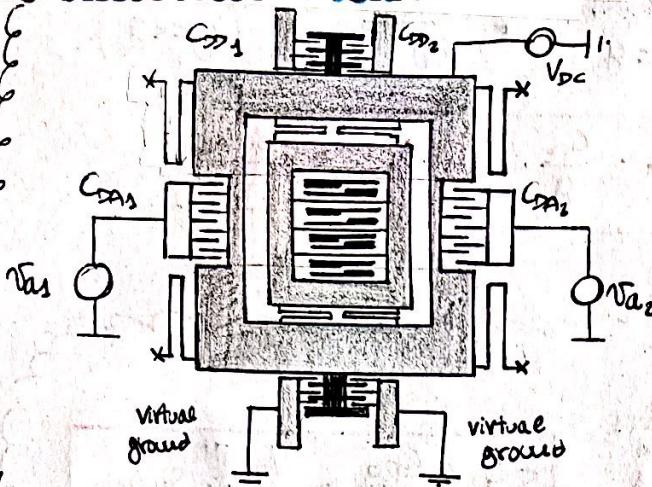
LOW PASS FILTER

LPF

BUFFER

$$x_0 \approx x_{0,\text{ref}}$$

PUSH-PULL ACTUATION for DUAL MODE



$$\ddot{V}_{da1} = \omega_a \sin(\omega_d t)$$

$$\ddot{V}_{da2} = -\omega_a \sin(\omega_d t)$$

THERE ARE TWO DRIVE SIGNALS AND
TWO PAIR OF ELECTRODES FOR
ACTUATION (C_{DA}) AND DETECTION (C_{DA})
SUPPOSED TO BE EQUAL IN PAIRS

BY CALCULATING THE TOTAL F_{elec}
DRIVE DETECTION TERMS GET CANCELED
AND WE HAVE:

$$F_{elec} = 2M_{DA} \omega_a \sin(\omega_d t)$$

TRANSMISSION COEFF.
OF ONE DRIVE-ACTUATION
ELECTRODE

IN THIS WAY THE 2nd TERM IS
DELETED AND NOT PEGLECTED!

EFFECTS OF EXTERNAL ACCELERATIONS and IN-PHASE/ANTIPHASE DUAL MODE

$$T_{YF} \Big|_{w \rightarrow w_0} = \frac{1/M}{w_0^2}$$

$$\hookrightarrow Y_s = \frac{a}{w_0^2} \gg y_0$$

(due to rotation)

IT'S IMPORTANT TO
NOT HAVE FREQS OF GYROS
FALLING INTO DISTURBANCES
RANGE (SOUND OR VIBRATION)

TO SAVE THIS WE CAN (COUPLING WITH TUNING FORK)

- USE TWO GYROS WITH OPPOSITE DUAL
DIRECTION TO HAVE
 - CORIOLIS FORCES IN OPPOSITE DIRECTION
 - ACCELERATIONS IN THE SAME
- IN ORDER TO CANCEL AS COMMON
MODE SIGNAL THE COMMON ACCELERATIONS
(ALSO DOUBLED SENSITIVITY but DOUBLED AREA)

→ THERE ARE TWO MODES

I) IN-PHASE DRIVE MODE:

TUNING FORK SPRINGS DOES NOT BEND AND DON'T CONTRIBUTE
ONLY DRIVE-SPRINGS

$$\left. \begin{aligned} k_d^{IP} &= 2k_d^P = \frac{n_{par}}{h_{fold}} EH \left(\frac{w_d}{l_d} \right)^3 \\ w_{tot} &= (w_d + w_s)_2 \end{aligned} \right\} f_d^{IP} = \frac{1}{2\pi} \sqrt{\frac{\frac{3P}{2}}{w_{tot}}}$$

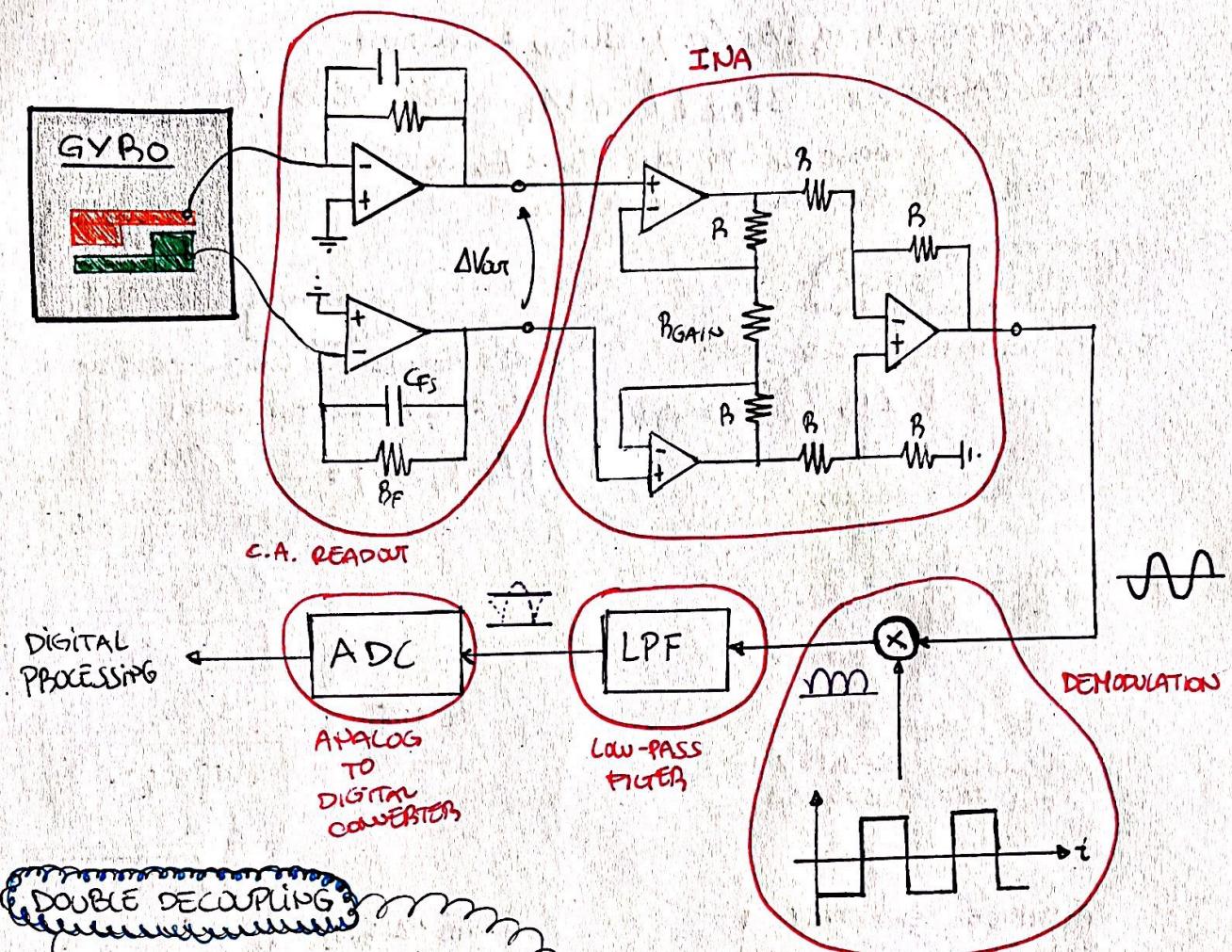
II) ANTI-PHASE DRIVE MODE:

TUNING FORK SPRINGS TAKES PART
CENTRAL POINT OF TUNING FORK
VIRTUALLY ANCHORED
Used in sensitivity!

$$k_{fp} = \frac{n_{par}}{h_{fold}} EH \left(\frac{w_{fp}}{l_{fp}} \right)^3$$

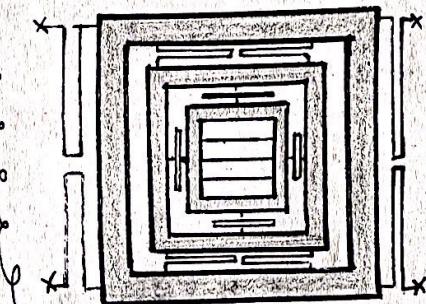
$$k_{tot} = k_{fp} + k_{IP}$$

EPSE MODE SYSTEM OVERVIEW



SENSE MODE CAN SEE SMALL CAPITATIVE CHANGES BECAUSE OF FRINGE EFFECT

FURTHER DECOUPLING



- DRIVE MODE NOISE NEGIGIBLE
- SENSE MODE

$$S_{fm} = S_{fm} \cdot \left(\frac{Q_s}{k_s} \right)^2 \left[\frac{\text{m}^2}{\text{Hz}} \right]$$

$$S_{\omega_m} = S_{fm} \cdot \left(\frac{\Delta \omega_{bm}}{\omega_{go}} \right)^2 \left[\frac{(\text{rad/s})^2}{\text{Hz}} \right]$$

$$NEPD = \sqrt{S_{\omega_m}} \quad \begin{matrix} \text{NOISE} \\ \text{EQUIVALENT} \\ \text{RATE} \\ \text{DENSITY} \end{matrix}$$

in dps by multiplying $\times 180/\pi$

• FEEDBACK RESISTANCE: $\sqrt{S_{\omega_m, \beta_F}} = \left(\sqrt{2 \cdot \frac{4 \pi k T}{\beta_F} \cdot \left(\frac{1}{\omega_{CFS}} \right)^2} \right) / \frac{\Delta V_{out}}{\Omega} \cdot \left(\frac{180}{\pi} \right)$

• OP-AMP: $\sqrt{S_{\omega_m, CA}} = \left(\sqrt{2 S_{m, op}} \left(1 + \frac{C_F}{C_{FS}} \right) \right) / \frac{\Delta V_{out}}{\Omega} \left(\frac{180}{\pi} \right)$

MODE-MATCHED GYRO TRADE-OFFS

- GAIN/BANDWIDTH $(\Delta V_{in}/\omega) \cdot \Delta \omega_{BW} = \text{constant}$
- NOISE/BANDWIDTH $\Delta f_{noise} \propto b_s \propto 1/Q$ while $\text{NEPD} \propto b_s$
- USEFUL TO HAVE LARGE QD TO HAVE $V_a \ll V_{dc}$ in DRIVE MODE
- NOT HIGH Qs FOR SENSE MODE (trade-off above)

MODE-SPLIT GYROSCOPE

ISSUES WITH MODE-MATCHED GYROS:

- usually $f_s \neq f_d$ (for non-uniformity)
 - above tradeoffs
 - sensitivity changes for T changes
- $$\frac{dQ_s}{Q_s} = \frac{dbs}{bs} = \frac{1}{2} \frac{\delta T}{T}$$

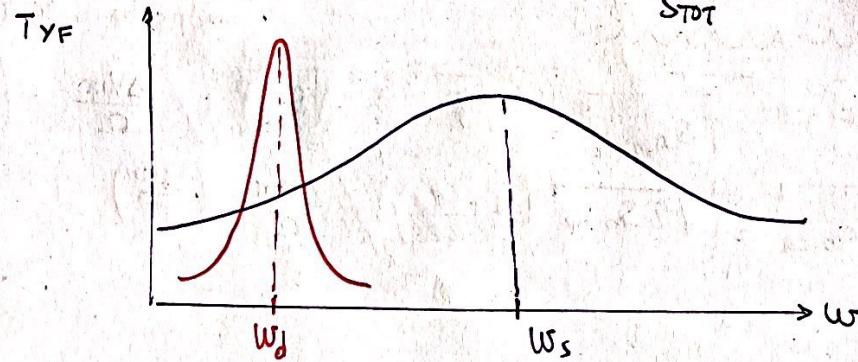
WE WANT A GYRO THAT WORKS UNIFORMLY IN DIFFERENT CONDITIONS

SACRIFICE SENSITIVITY FOR BANDWIDTH

$$y_{s,o} = \frac{x_{d,o}}{\omega} \cdot \Delta \omega_{BW}$$

MODE-SPLIT SENSITIVITY

- 😊 • DOES NOT DEPEND ANYMORE ON b AND Q , THEN WE CAN LOWER b TO MINIMIZE HED
- MUCH MORE TOLERANT TO T CHANGES BECAUSE OF T CHANGES
- NO BW-SENSITIVITY TRADE-OFF



Mode Split
Let's moderate the Coriolis-force at $\omega_0 < \omega_s$ (before the peake)

$$\begin{aligned} |T_{YF}(j\omega_0)| &= \frac{1/\omega_s}{(-\omega_0^2 + j \frac{\omega_0 \omega_s}{Q_s} + \omega_s^2)} \\ &= \frac{1/\omega_s}{\sqrt{(\omega_s^2 - \omega_0^2)^2 + (\frac{\omega_0 \omega_s}{Q_s})^2}} \\ \omega_{HS} \ll \omega_s &\approx \frac{1/\omega_s}{\sqrt{4\omega_s^2 \Delta \omega_{HS}^2 + \frac{\omega_s^4}{Q_s^2}}} \\ &= \frac{1/\omega_s}{2\omega_s \sqrt{\Delta \omega_s^2 + \Delta \omega_{BW}^2}} \\ \Delta \omega_{HS} \gg \Delta \omega_{BW} &\approx \frac{1}{k_s} \cdot \frac{\omega_s}{2\Delta \omega_s} = \frac{1}{k_s} Q_{eff} \end{aligned}$$

$$Q_{eff} = \frac{\omega_s}{2\Delta \omega_{HS}}$$

EFFECTIVE
QUALITY
FACTOR

- 😢 • GAIN REDUCTION
• NEED HIGH PERFORMANCE ELECTRONIC

$$\frac{dS_{TOT}}{S_{TOT}} = \alpha \Delta T \text{ with } \alpha = -30 \frac{\mu\text{V}}{\text{K}}$$

MAGNETOMETER 3

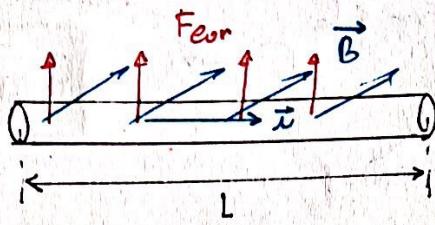
The LORENTZ FORCE

$$\vec{F}_{\text{ext}} = q(\vec{v} \times \vec{B})$$

LORENTZ FORCE

$$\vec{F}_{\text{ext}} = i(L) \times \vec{B}$$

CURRENT



$$F_{\text{ext}} \approx 1 \text{ mN}$$

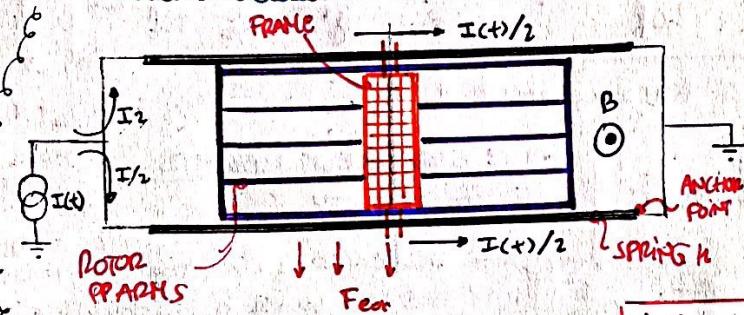
with

$$L = 1000 \mu\text{m}$$

$$i = 0.2 \text{ mA}$$

ARCHITECTURE

Z-AXIS MAGNETOMETER - SIMPLE ARCH



$$\text{if } \Delta C_{\text{pp}} = 2 \frac{C_0}{g} \times \sim$$

$$\frac{\Delta C}{B} = \frac{C_0}{g} \cdot \frac{Q}{k} i L = S_{\text{mech}}$$

MECHANICAL
SENSITIVITY

$$\hookrightarrow S_{\text{tot}} = \frac{V_{\text{dd}}}{C_F} \cdot S_{\text{mech}} \quad (\text{with charge amplifier and INA readout})$$

AGAIN, LIKE FOR GYROS THE INFORMATION ABOUT B IS IN THE LORENTZ FORCE PROPORTIONAL TO ω_0 OF CURRENT GEN.

MAXIMUM BANDWIDTH

$$\Delta \omega_{\text{BW}} = \frac{\omega_0}{2Q} = \frac{b}{2\pi}$$

(-3dB off resonance peak)

RESONANT MODE NOISE

$$\sqrt{S_{\text{fm}}} = \sqrt{4k_B T b}$$

$$\text{NEMD} = \sqrt{S_{\text{Bm}}} = \left[\left(\sqrt{S_{\text{fm}}} \cdot \frac{Q}{k} \right) / \left(\frac{b}{B} \right) \right]$$

$$= \frac{4}{iL} \sqrt{4k_B T b}$$

RESONANT OPERATION MODE:

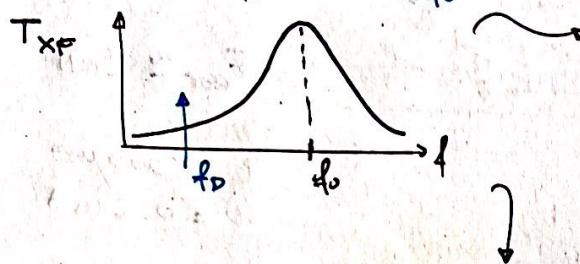
TRADE-OFFS

o again, BW vs NEMD

OFF-PERSONANCE OPERATION MODE

MULTI-LOOP

TO SOLVE THE NOISE-GAIN-ON TRADE-OFF LET'S CHOOSE A DRIVE FREQ. f_D OF FEED $< f_0$



WE SHOULD FIX SENSITIVITY AND FORMULA WITH

$$Q_{eff} = \frac{Q}{2\Delta f} = \frac{\omega_0}{2\Delta\omega} \quad \text{EFFECTIVE QUALITY FACTOR}$$

FREQ OF MAGNETOMETER SENSING MODE

TO FURTHER IMPROVE THE LOST GAIN DUE TO $Q_{eff} < Q$ WE CAN REUSE THE CURRENT IN LOOPS AND BOOST FEED BY NEOP.

$$S_{\text{mech}} = \frac{\Delta C}{B} = 2 \frac{C_0}{g} \cdot \frac{Q_{eff}}{2k_{1/2}} iL N_{\text{loop}} = \frac{C_0}{g k_{1/2}} iL \frac{f_0}{2\Delta f} N_{\text{loop}}$$

$$S = \frac{\Delta V_{out}}{B} = \frac{V_{DD}}{C_F} \cdot \frac{C_0}{g} \cdot \frac{iL}{k_{1/2}} \cdot \frac{f_0}{2\Delta f} N_{\text{loop}} = \frac{V_{DD}}{C_F} \cdot \frac{\epsilon_A N_{\text{loop}}}{g^2} \cdot \frac{iL}{k_{1/2}} \cdot \frac{f_0}{2\Delta f} N_{\text{loop}}$$

OVERALL SENSITIVITY

$\hookrightarrow S \cdot \text{BW} = \text{constant}!$

$$k_{1/2} = k_{beam} + k_{V2TF}$$

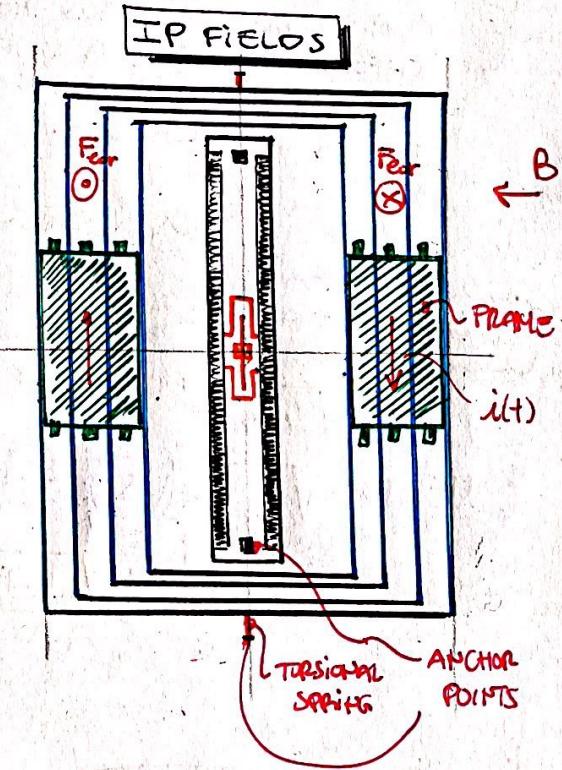
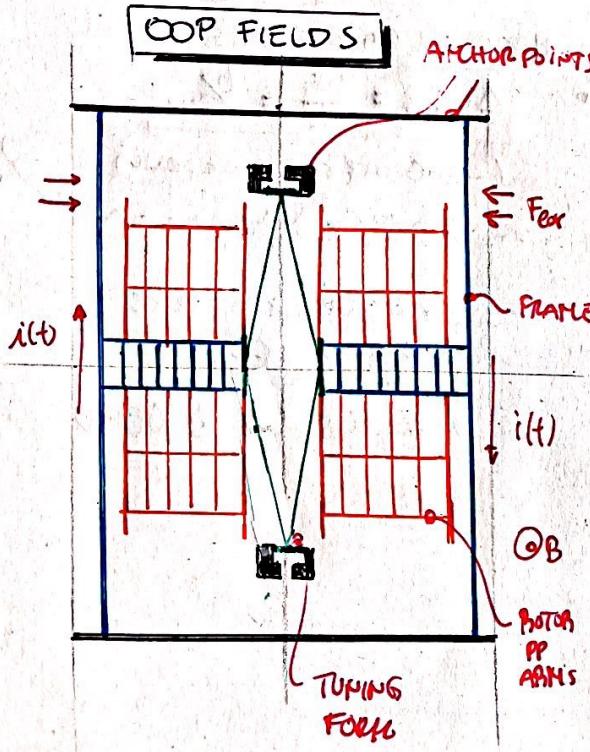
TUNING FORK STRUCTURE (HALF DEVICE)

$$\text{NEMD} = \frac{4}{N_{\text{loop}} iL} \sqrt{k_B T_b}$$

NOISE EQUIVALENT MAGNETIC FIELD DENSITY

- DECREASED $\propto 1/N_{\text{loop}}$ (also ELECTRONIC)
- WE CAN LOWER b WITHOUT LIMITATIONS
- AREA DECREASED

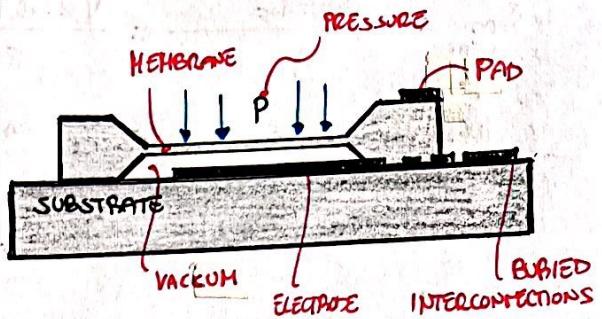
DIFFERENTIAL ARCHITECTURES



PRESSURE SENSOR

SENSITIVITY

$$S = \frac{\Delta V_{out}}{P} = \frac{V_{DD}}{C_F} \cdot \frac{C_0}{g} \cdot \frac{\delta y}{\delta P} = \frac{V_{DD} \cdot C_0 \cdot A}{C_F \cdot g \cdot 2k}$$



same parameters:

$$P_{SPL} = 20 \log_{10} \frac{P}{P_{REF}}$$

$$P_{REF} = 20 \mu\text{Pa} \quad (\text{THRESHOLD OF HUMAN HEARING})$$

$$k_{M} = 16 E \cdot \frac{\pi h^3}{r^2} \quad (\text{STIFFNESS})$$

$$m = \pi r^2 h p \quad (\text{MASS})$$

$$b = b_{\text{area}} \pi r^2 \quad (\text{DAMPING})$$

THE INPUT SIGNAL OF A MICROPHONE IS SOUND PRESSURE P AS AN AC SOUND WAVE

$$p = P_0 \sin(2\pi f_a t)$$

ACOUSTIC FREQUENCY RANGE
20Hz - 20kHz

THE VIBRATING MEMBRANE CAN BE APPROXIMATED AS A 3D MASS-SPRING-DAMPER SYSTEM, GIVEN A UNIFORMLY DISTRIBUTED FORCE APPLIED ONTO THE MEMBRANE

$$\omega_R = \sqrt{\frac{k_M}{m}}$$

$$y = F / k_M$$

THE REQUIREMENTS ARE

$$P_{a,HIN} = P_{REF} \cdot 10 \left(\frac{P_{a,HIN,SPL}}{20} \right) = 893 \mu\text{Pa}$$

$$P_{a,MAX} = P_{REF} \cdot 10 \left(\frac{P_{a,MAX,SPL}}{20} \right) = 28.3 \text{ Pa}$$

the FORCE APPLIED is

$$F = P A = P \pi r^2$$

NOTE THAT BOTH THE CHAMBERS FACED BY THE MEMBRANE ARE KEPT AT THE ATMOSPHERIC PRESSURE, SO THAT THE DC DISPLACEMENTS INDUCED BY THE ATMOSPHERIC PRESSURE ARE CANCELED

THE MINIMUM DETECTABLE SIGNAL BY THE BANDWIDTH OF OUR SENSOR

$$S_{MDR} = \frac{P_{a,HIN}}{\sqrt{f_{\text{max}} - f_{\text{min}}}} \approx \frac{P_{a,HIN}}{\sqrt{f_{\text{max}}}}$$

THE PRESSURE DIFFERENCE APPLIED TO THE MEMBRANE IS

$$P_{diff} = (P_{atm} + P_{sound}) - P_{atm} = P_{sound}$$

$$S_m = \frac{S_{m,HEAS}}{A} = \frac{\sqrt{4k\theta T b}}{\pi r^2} = \sqrt{\frac{4k\theta T b \text{area}}{\pi r^2}}$$

also showed be $S_{m,HEAS,p} = \frac{S_{m,TORP}}{\Gamma^2}$ for a good balance

WE FIND:

$$\left| \frac{\partial y}{\partial p} \right|_{WCCUR} = \frac{\text{Area}}{\mu m} \quad \text{DISPLACEMENT v PRESSURE}$$

and some linearity error

$$\varepsilon_{lin} = \left(\frac{y}{g} \right)^2 = (\varepsilon_{lin})^2 \quad \text{LINEARITY ERROR}$$

with a CHARGE AMPLIFIER READOUT:

$$T_{TOT}(j\omega) = -\pi r^2 \frac{1}{\mu + j\omega b - \omega^2 m} \cdot \frac{\partial C}{\partial y} j\omega V_{DC} \frac{R_F}{1 + S_C R_F}$$

OVERALL SENSITIVITY WITH TCA

$$T_{TOT}(j\omega) = \frac{V_{out}(j\omega)}{P(j\omega)}$$

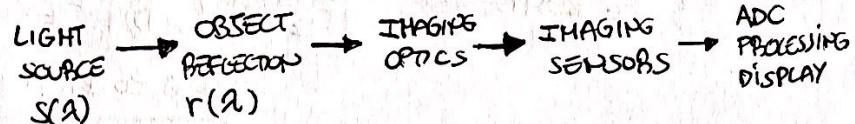
CHOS IMAGING SENSORS

BASICS OF OPTICS

RAYS are not colored!
they are made of PHOTONS
with a certain ENERGY
which corresponds a
WAVELENGTH (and a
certain color perception
from the eye)

$$E_{ph} = \frac{h c}{\lambda}$$

ENERGY
of
PHOTONS



$$B = \int_{400}^{700} s(\lambda) \cdot r(\lambda) \cdot t(\lambda) \cdot b(\lambda) d\lambda$$

↓ SPECTRAL TRANSMISSION
↓ ILLUMINANT
↓ BLUE PIXEL RESPONSE

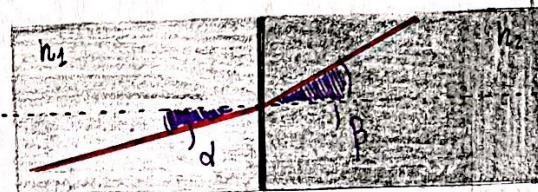
we can approximate as almost
perfect BLACK BODIES RADIATORS

$$\lambda_{peak} \cdot T = 2,9 \cdot 10^{-3} \text{ m} \cdot \text{K}$$

WIEN'S LAW

some approximations ...

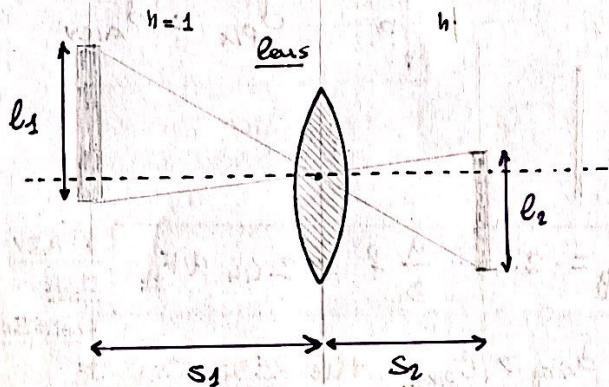
- light is described by RAYS which travel in STRAIGHT LINES until they are deflected by an OPTICAL ELEMENT
- the DEFLECTION is described by the SNELL LAW



$$n_1 \sin(d) = n_2 \sin(\beta)$$

SNELL'S LAW

LENS SYSTEMS and PARAMETERS



$$\frac{1}{s_1} + \frac{1}{s_2} = \frac{1}{f} \quad \text{with} \quad \frac{1}{f} = (n-1) \left[\frac{1}{R_1} - \frac{1}{R_2} \right]$$

LENS MAKER FORMULA

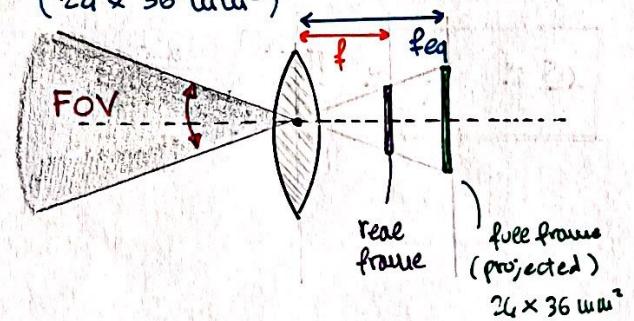
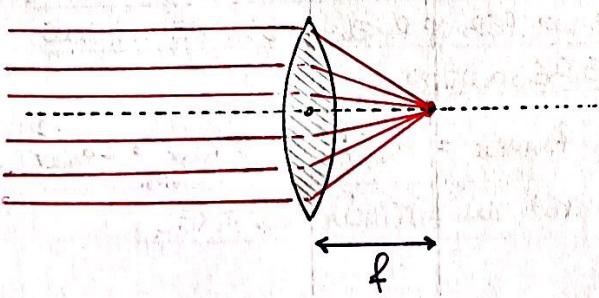
- since $s_1 \gg s_2 \rightarrow f \approx s_2$

FOCAL LENGTH

$f \uparrow$ with POSITIVE CONVEX LENS

$f \uparrow$ with NEGATIVE CONCAVE LENS

- we can calculate the EQUIVALENT FOCAL LENGTH by referring to FULL-FRAME FORMAT ($24 \times 36 \text{ mm}^2$)



We also define :

- MAGNIFICATION FACTOR

$$m = \frac{l_2}{l_1} = \left(\frac{s_2}{s_1} \approx \frac{f}{s_1} \right) = \frac{\text{Area}_{\text{pix}}^{\text{MIN}}}{\text{Area}_{\text{MAX}}^{\text{SCENE}}} \quad \text{MAGNIFICATION FACTOR}$$

- F# NUMBER

$$F\# = f/D \quad \begin{matrix} \text{Diameter} \\ \text{of the lens} \end{matrix}$$

useful to find the

$$\theta = \arctan\left(\frac{D}{2f}\right) = \arctan\left(\frac{1}{2F\#}\right) \quad \begin{matrix} \text{Maximum} \\ \text{aperture} \\ \text{of folded} \\ \text{rays} \end{matrix}$$

ABERRATIONS

CAUSES IMPERFECTIONS IN
REPLICAS OF OBJECTS

WORSENS THE RESOLUTION

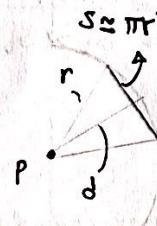
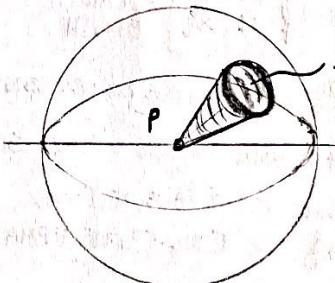
WE DON'T HAVE A
FOCAL POINT

instead we have a
disk of confusion
can be reduced using



ASPHERIC LENS

SOLID ANGLES



$$\Omega = \frac{S}{d^2} \approx \frac{\pi r^2}{d^2}$$

SOLID ANGLE
[sr]

HAPPENS DUE TO :

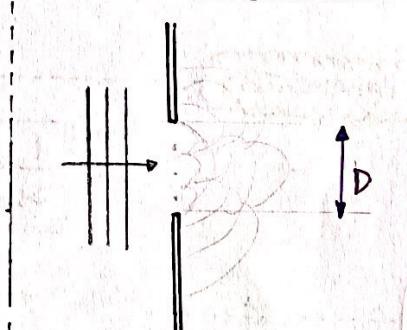
- errors during 1st order approximations
(3rd order spherical aberrations)
- the radiation is not 'monochromatic' and
the refractive index is function of λ
(chromatic aberration)

can be reduced with :

- LOW DISPERSION GLASSES
- ACHROMATIC DOUBLETS

DIFFRACTION

to have a
BALANCED SYSTEM
usually choose
 $d_{\text{pix}} = d_{\text{Airy}}$

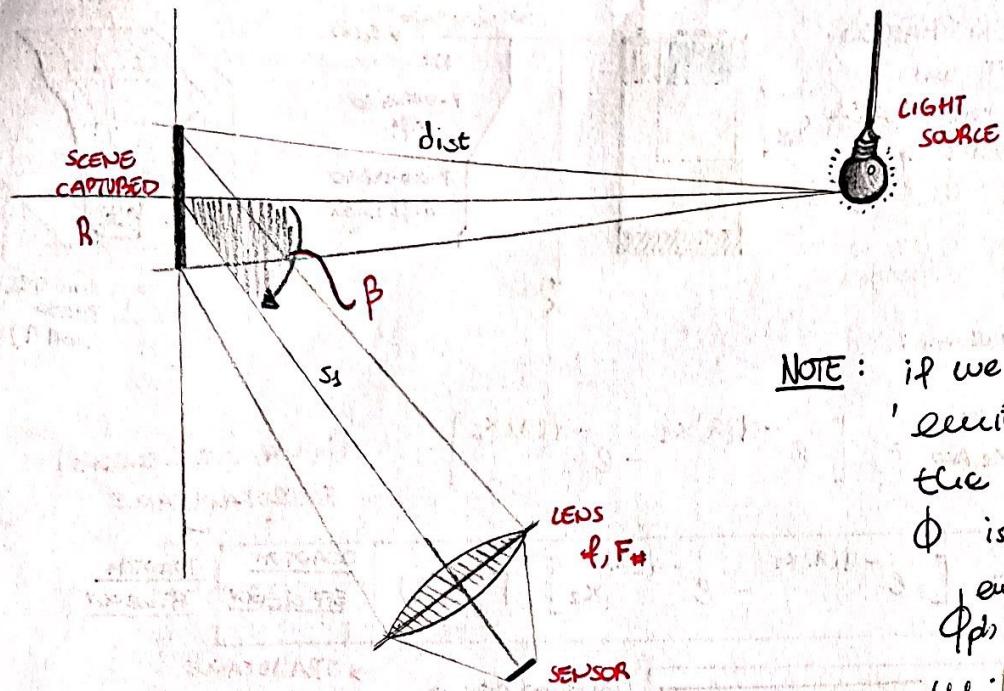


$$d_{\text{Airy}} = 2.44 \cdot \frac{\lambda}{D} f = 2.44 \lambda F\# \quad \begin{matrix} 1^{\text{st}} \text{ AIRY} \\ \text{DISK} \\ \text{DIAMETER} \end{matrix}$$

- if $d_{\text{pix}} > d_{\text{Airy}}$ the RESOLUTION OF THE SENSOR IS LIMITED. INCREASING THE NUMBER OF PIXELS DOES NOT IMPROVE THE RESOLUTION.

- if $f_{\text{SAMPLE}} = l_{\text{pix}}^{-1} < 2f_{\text{MAX}} = 2d_{\text{Airy}}^{-1}$
we need an ANTIALIASING FILTER

OPTICAL POWER and PHOTON FLUX



NOTE: if we have an 'emitted' Φ_{ph} the absorbed Φ is given by

$$\Phi_{ph} \cdot \text{APD} \cdot \frac{1}{M^2}$$

 w.e.: magnification factor

1

CALCULATE THE AREA OF THE SCENE WHICH CORRESPONDS TO ONE PIXEL

$$A_{\text{SCENE}} = \frac{A_{\text{PIXEL}}}{(M)^2}$$

MAGNIFICATION FACTOR

Only photons coming from that area are captured by corresponding pixels

2

STARTING FROM OPTICAL POWER OF THE SOURCE WE EVALUATE THE IMPINGING POWER ON A SCENE PIX

- for an ISOTROPIC LIGHT SOURCE

$$I_{\text{SOURCE}} = \frac{P_{\text{SOURCE}}}{4\pi \text{ sr}} \left[\frac{\text{W}}{\text{sr}} \right]$$

- we calculate the solid angle seen from source to $A_{\text{SCENE PIX}}$:

$$\Omega_{\text{SCENE PIX SOURCE}} = \frac{A_{\text{SCENE}}}{(dist)^2} \left[\text{sr} \right]$$

- the OPTICAL POWER IMPINGING ON SCENE AREA WHICH CORRESPONS TO ONE PIXEL

$$P_{\text{SCENE PIX}}^{(I)} = I_{\text{SOURCE}} \cdot \Omega_{\text{SCENE PIX SOURCE}} \left[\text{W} \right]$$

EVALUATE THE OPTICAL POWER REFLECTED

$$P_{\text{SCENE}}^{(R)} = P_{\text{SCENE PIX}}^{(I)} \cdot R_{\text{SCENE}}$$

SCENE REFLECTANCE

AND THEN THE INTENSITY DEPENDING ON TYPE OF REFLECTING OBJECT:

- ISOTROPIC REFLECTION:

$$I_R^{\text{ISO}} = P_{\text{SCENE}}^{(R)} / 2\pi \text{ sr}$$

ISO

- LAMBERTIAN REFLECTION:

$$I_R^{\text{LAM}} = \frac{P_{\text{SCENE}}^{(R)}}{\pi} \cos(\beta)$$

4 WE EVALUATE THE POWER IMPINGING ON A PIXEL

$$A_{\text{PIXEL}} = \pi \left(\frac{D}{2} \right)^2$$

$$\Omega_{\text{LEPS}} = \frac{A_{\text{PIXEL}}}{S_1^2}$$

$$P_{\text{PIX}} = I_R \cdot \Omega_{\text{LEPS}} \frac{1}{S_1^2}$$

DIVIDE BY ENERGY OF A PHOTON

$$\bar{E}_{ph} = \frac{hc}{\lambda}$$

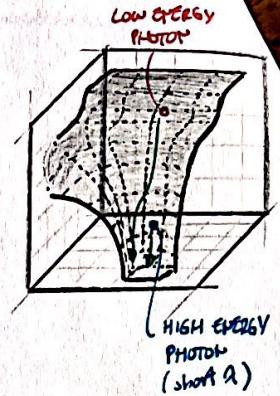
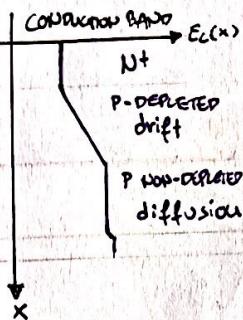
$$N_{ph} = \frac{P_{\text{PIX}}}{\bar{E}_{ph}} \left[\frac{\text{ph}}{\text{s}} \right]$$

$$\Phi_{ph} = N_{ph} \left[\frac{\text{ph}}{\text{W} \cdot \text{s}} \right]$$

PHOTON FLUX

the PHOTODIODE

CHARGE COLLECTION HAPPEN WITH BOTH DRIFT and DIFFUSION into the N^+ WELL WHERE PHOTOCURRENT IS GENERATED



let's define some values:

$$\text{PHOTON ABSORBED BETWEEN } x_1 \text{ AND } x_2 = [e^{-\alpha(\lambda)x_1} - e^{-\alpha(\lambda)x_2}] \quad (\text{in CFA, COLOR FILTER ARRAYS})$$

$$\frac{\text{collected } e^-}{\text{incident ph}} = M(\lambda) = [e^{-\alpha(\lambda)x_1} - e^{-\alpha(\lambda)x_2}] T_{Si}(\lambda) T_{FIL}(\lambda) \quad (\text{QUANTUM EFFICIENCY})$$

$$x_{depe} = \sqrt{\frac{2\varepsilon_0 \varepsilon_{Si} (V_B + V_Bi)}{q N_A}} \quad (\text{DEPLETION REGION LENGTH})$$

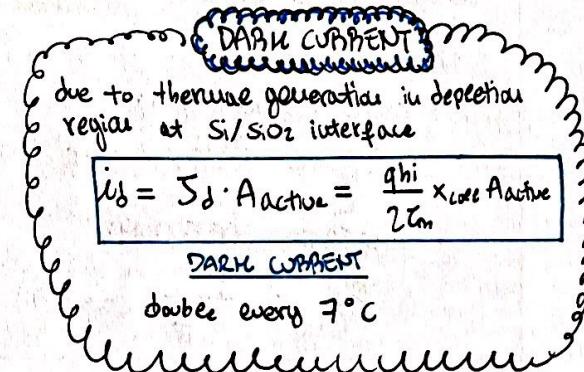
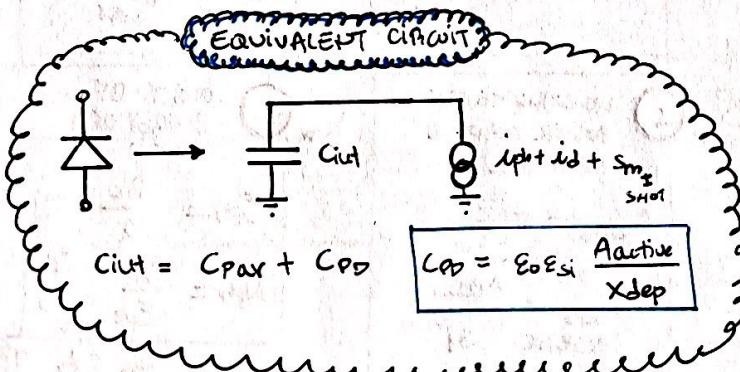
$$R_2(\lambda) = \frac{i_{ph}}{P_{in}} = \frac{q \cdot (\text{collected electrons per second})}{E_{ph} \cdot (\text{incident photons per second})} = \frac{q}{E_{ph}} M(\lambda) \quad (\text{RESPONSIVITY})$$

INPUT OPTICAL POWER

given the power impinging on a pixel P_{pix} or the photon flux Φ_{ph}

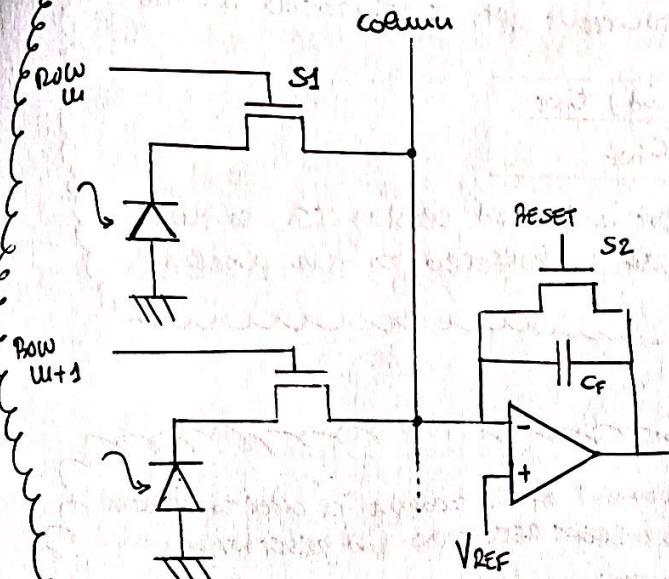
$$i_{ph} = q \cdot M(\bar{\lambda}) \cdot \Phi_{ph} \cdot A_{pixel} = R_2(\bar{\lambda}) P_{pix} \quad (\text{GENERATED PHOTOCURRENT})$$

$$i_{ph} = q M_{pix} T_{fil} \frac{\lambda}{hc} \cdot \frac{P_{source} A_{pix} f_{scen}}{16\pi (\text{dist})^2 F_*^2} \quad (\text{PHOTON FLUX})$$



READOUT TOPOLOGIES

PASSIVE PIXEL SENSOR - PPS



each pixel is formed just by photosensitive element and a selection transistor



- SMALL PIXEL SIZE



- DOES NOT INTRODUCE PIXEL TO PIXEL VARIATIONS
- LIMITED READOUT SPEED
- SPECIALIZED TECHNOLOGY
- HIGH VOLTAGE REQUIRED FOR CHARGE TRANSFER

1. OUTPUT INITIALLY RESET TO V_{REF} BY CLOSING BOTH S_1 AND S_2 . PHOTOCURRENT OF THE DIODE DOES NOT DETERMINE ANY CHANGE IN THE OUTPUT VOLTAGE (AS IT FLOWS THROUGH THE SHORT CIRCUIT DUE TO S_2 FROM THE LOW-IMPEDANCE OUTPUT OF THE OPERATIONAL AMPLIFIER)

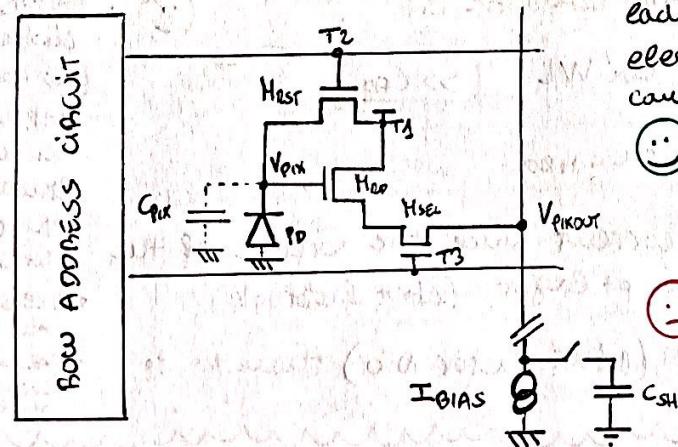
2. DURING INTEGRATION S_1 IS KEPT OPEN TOGETHER WITH THE CAMERA SHUTTER UNTIL READOUT HAPPEN. PHOTOCURRENT AND DARK CURRENTS ARE COLLECTED ON PHOTODIODE CAPACITANCE, WHOSE VOLTAGE DECREASES DURING INTEGRATION TIME

$$V_{PD}(t) = V_{REF} - \Delta V_{PD}(t) \quad \Delta V_{PD}(t) = \frac{(i_{ph} + i_d)t}{C_{int}}$$

3. IN THE READOUT PHASE, S_1 IS CLOSED AGAIN SOON AFTER S_2 IS OPEN. THE VOLTAGE CHANGE ACROSS THE DIODE CAPACITANCE DETERMINES A CURRENT PULSE FLOWING IN THE FEEDBACK LOOP OF THE AMPLIFIER

$$V_{out} = V_{REF} + \Delta V_{PD} = V_{REF} + \frac{(i_{ph} + i_d)t_{int}}{C_F}$$

ACTIVE PIXEL SENSOR - APS: 3T TOPOLOGY



each pixel is formed by the photosensitive element and a small charge to voltage conversion



- HIGH SPEED READOUT
- STANDARD FABRICATION PROCESS
- LOWER VOLTAGE
- ONLY SELECTED PIXELS CONSUME POWER



- INTRODUCE GAIN/OFFSET NONUNIFORMITIES
- SAME NODE FOR PHOTOSENSITIVE ELEMENT AND CONVERSION (GAIN - FF TRADE OFF)
- NONLINEARITY DUE TO C_{DEP} VARIABILITY

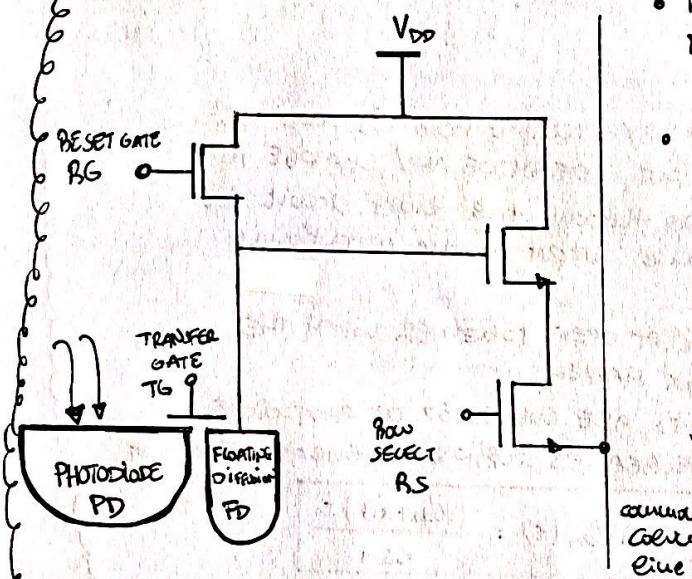
3T topology:

1. the INTEGRATION node is first RESET to a value equal to VDD by closing the RESET switch.
 2. then T2 is open and the photocurrent gets INTEGRATED at the integration capacitance.

$$V_{pix} = V_{DD} - \frac{(V_{ph} + id) t_{int}}{C_{int}}$$

3. finally, the follower transistor is biased closing T3, so that the voltage at the follower gate is buffered to the pixel output

APS - LT TOPOLOGY



- DIFFUSION CAPACITANCE can be made very small since it's no longer responsible of photons collection

 - improvement on 3T topology in order to obtain DARK CURRENT REDUCTION and RESET NOISE REDUCTION
 - IT CONSIST IN A PITTED PHOTODIODE and FOUR TRANSISTORS
 - HTX TRANSFER GATE
 - MBS RESET TRANSISTOR
 - MHD AMPLIFIER TRANSISTOR
 - MSEL SELECT TRANSISTOR
 - WE HAVE ALSO A DOUBLE SAMPLING
WE SUBTRACT TWO MEASUREMENT:
 - one with photocurrent and offset and noise
 - one only with reset and noise (no signal)
 - Overall noise is thus reduced!

$$C_{FD} = \frac{E_0 E_R}{\frac{WL}{X_{def}}}$$

while INPUT CAPACITANCE OF HVD TRANSISTOR IS

$$C_{iu, MHD} = \text{Cox}' WL \quad (\gg C_{FD})$$

$$C_{iu} = C_{PD} + C_{iu, HRG}$$

- Also we have lower dark current since the surface of the pinned photodiode is shielded by a p+ layer (about factor N_{20})
 - also KTC rms value is reduced (about factor N_{20}) thanks to double sampling

- BETTER LINEARITY
(no edger dependent
on photodiode
capacitance
that changes
with signal)
 - PIXELS SHARE
SAME ELECTRONICS
 - BACKSIDE
ILLUMINATION
ALLOWED

SENSITIVITY, SIGNAL, NON-LINEARITIES

with a constant photocurrent, we can evaluate the collected electrons:

$$Q_{ph} = i_{ph} \cdot t_{int} \rightarrow N_{ph} = Q_{ph}/q$$

and the gain:

$$G = \frac{\Delta V_{PD}}{N_{ph}} = \frac{Q_{ph}/q}{C_{PD} \cdot Q_{ph}} = \frac{q}{C_{PD}} \rightarrow \Delta V_{PD} = G \cdot N_{ph}$$

and finally the sensitivity:

$S = \frac{\Delta V_{PD}}{\Phi_{ph}}$	$\left[\frac{V}{\frac{ph}{W^2 S}} \right]$	<u>SENSITIVITY</u>
---------------------------------------	---	--------------------

However in 3T topology, we should take into account that photodiode capacitance varies with voltage:

$$\Delta V_{PD} = \frac{(i_{ph} + i_d) \cdot t_{int}}{C_{PD}}$$

$$V_{PD} = V_{DD} - \Delta V_{PD}$$

$$x_{dep} = \sqrt{\frac{2 \epsilon_0 \epsilon_R (V_{DD} + V_{bi})}{q N_A}}$$

$C_{PD} = \epsilon_s \frac{A}{x_{dep}} = \frac{C_{PD}^0}{\sqrt{1 + \frac{V_{PD}}{V_{bi}}}}$	<u>REAL PHOTODIODE CAPACITANCE</u>
---	------------------------------------

$$C_{PD}^0 = C_{PD} \Big|_{V_{PD}=0}$$

we can also find:

$Q_{ph, max}^{real} = 2 C_0 V_{bi} \left[1 - \sqrt{1 - \frac{V_{DD}}{V_{bi}}} \right]$

MAXIMUM COLLECTABLE CHARGE

and

$E_{lin} = \frac{Q_{ph, max}^{real} - Q_{ph, max}^{lin}}{Q_{ph, max}^{real}}$

with $Q_{ph, max}^{lin} = V_{DD} \cdot C_{PD} \Big|_{V_{PD}=V_{DD}}$

NOISE CONTRIBUTIONS

		DARK	ILLUMINATED	
SEPARATE	TEMPORAL	ELECTRONIC OFFSET NONUNIFORMITY	PHOTODIODE RESPONSE NONUNIFORMITY	at sensor level
		DARK CHARGE TRAP-UP NONUNIFORMITY	HOT PIXELS	
SEPARATE	TEMPORAL	DARK CURRENT SHOT NOISE	PHOTON CURRENT SHOT NOISE	at pixel level
		READ NOISE (AMPLIFIER)	RESET NOISE	
				QUANTIZATION NOISE

TEMPORAL NOISE

We are referring to 3T-topology!

CURRENT SHOT NOISE

A READOUT PERFORMED AS THE APPROACH RESET + INTEGRATION CORRESPONDS TO A GATED-INTEGRATOR SCHEME FOR FINITE TIME t_{int} WITH A BANDWIDTH $1/2t_{int}$

$$S_{i, \text{shot}}^2 = 2q(i_{ph} + i_d) \left[\frac{A^2}{Hz} \right]$$

We can convert it in terms of charge:

$$\sigma_{q, \text{shot}}^2 = \sigma_{i, \text{shot}}^2 \cdot t_{int}^2 = \frac{S_{i, \text{shot}}^2 t_{int}^2}{2t_{int}}$$

$$\sigma_{q, \text{shot}}^2 = q(i_{ph} + i_d)t_{int} \left[C^2 \right]$$

CURRENT SHOT NOISE

QUANTIZATION NOISE

REFERRED TO THE ADC

$$\sigma_{\text{quant}, v}^2 = \frac{LSP}{\sqrt{12}} = \frac{V_{REF}/2^{N_{bit}}}{\sqrt{12}}$$

$$\sigma_{\text{quant}, a}^2 = \frac{V_{REF}/2^{N_{bit}}}{\sqrt{12}} C_{diode}$$

QUANTIZATION NOISE

to choose ADC we can force:

$$\sigma_{\text{quant}} < \sigma_{a, \text{reset}} \text{ KTC}$$

OTHERS NEGIGIBLE NOISE

- 1/f NOISE
- thermal noise of source-follower

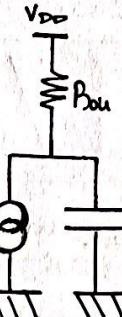
REDUCE TEMPORAL NOISE

- by DOUBLE SAMPLING
- by AVERAGING:

take k consecutive images for a selected exposure time and calculate the average of the series for each pixel

KTC NOISE (RESET NOISE)

WHEN RESET SWITCH OPENS, ONE RANDOM VOLTAGE VALUE WITHIN THE NOISE POWER SPECTRAL DENSITY IS 'SAMPLED' AND FROZEN AS STARTING VOLTAGE FOR THE INTEGRATION. THE BANDWIDTH IS REPRESENTED BY THE $1/(4BC_0)$ BANDWIDTH



$$S_{V_{\text{reset}}}^2 = 4kT R_{\text{out}} \left[\frac{V^2}{Hz} \right]$$

$$\sigma_{V_{\text{reset}}}^2 = \frac{S_{V_{\text{reset}}}^2}{4P_{\text{clock}}} = \frac{kT}{C_{\text{int}}} \left[V^2 \right]$$

$$\sigma_{q_{\text{KTC}}}^2 = kT C_{\text{int}} \left[C^2 \right]$$

KTC NOISE

S/NR - SIGNAL TO NOISE RATIO

$$\text{SNR} = 20 \log_{10} \frac{N_{ph}}{\sqrt{\sigma_{\text{shot}}^2 + \sigma_{a, \text{KTC}}^2 + \sigma_{\text{quant}}^2}}$$

SIGNAL TO NOISE RATIO

MAXIMUM SNR

$$\text{SNR}_{\text{max}} = \frac{N_{ph, \text{max}}}{\sqrt{\sigma_{a, \text{shot}}^2 + \sigma_{a, \text{KTC}}^2}} = \sqrt{FWC} \text{ dB}$$

TRADE-OFF DB-SNR ON t_{int} $t_{int} \uparrow \text{DB} \uparrow \text{SNR}$

DB - DYNAMIC RANGE

saturation of the pixel

$$DB = \frac{\text{MAXIMUM DETECTABLE SIGNAL}}{\text{MINIMUM DETECTABLE SIGNAL}}$$

set by noise

$$N_{ph, \text{min}} \approx \frac{1}{q} \sqrt{\sigma_{q, \text{shot}}^2 + \sigma_{a, \text{reset}}^2}$$

quant. noise
read noise

$$N_{ph, \text{max}} \approx \frac{1}{q} (V_{DD} C_{PD})$$

(otherwise use real
| $Q_{ph, \text{max}}$)

FPN SPATIAL NOISE

OPTICAL RESPONSES NONUNIFORMITIES
CAUSES DIFFERENT PHOTO-SIGNALS
FOR NOMINALLY IDENTICAL ILLUMINATED
PIXELS

→ We used to store in a memory the
correction values for offset and gain
for each pixel

- ↳ UNWANTED GAIN CHANGES FROM
PIXEL TO PIXEL DUE TO
 - SOURCE FOLLOWER V_{GS} VARIATIONS
 - DARK CURRENT

$$\Delta V_{\text{PIXOUT}} = \frac{\Phi(\lambda) \eta(\lambda) q A_{\text{pix}} \text{FF} T_{\text{CFA}}(\lambda) T_{\text{OPT}}(\lambda) t_{\text{int}}}{C_{\text{dep}} + C_g} G_{\text{SF}} + \left(\frac{i_{\text{distr}}}{C_{\text{dep}} + C_g} G_{\text{SF}} + \Delta V_{\text{OFFSET}} \right)$$

PHOTOGENERATED CURRENT TERM

DARK CURRENT TERM

$$\Delta V_{\text{PIXOUT}} = I_{\text{pix}} \cdot G_{\text{pixel}} \left(1 + \frac{\sigma_{\text{PNR}}}{100} \right) + O_{\text{pix}} \left(1 + \frac{\sigma_{\text{DSNU}}}{100} \right)$$

DSNU - DARK SIGNAL NON UNIFORMITIES

IT'S A MEASURE OF THE DISTRIBUTION OF DARK OUTPUT VOLTAGES FOR ALL THE PIXELS
MEASURABLE BY TAKING PHOTOS IN DARK CONDITIONS.

• to model:

- 1) take N consecutive images and remove
temporal noise by averaging

$$\bar{V}_{\text{out}, d_{\text{SNR}}} = \frac{1}{N} \sum_i V_{\text{out}, d_{\text{SNR}} i}$$

- 2) DSNU is the standard deviation of all
the values across the sensor

- 3) % DSNU is found by normalizing to mean value

$$\sigma_d^2 = \frac{1}{M \cdot N} \sum_i (V_{\text{out}, d_{\text{SNR}} i} - \bar{V}_{\text{out}, d_{\text{SNR}}})^2$$

$$\sigma_{\text{DSNU}} = \frac{\sigma_d}{\bar{V}_{\text{out}, d_{\text{SNR}}}} \cdot 100 \rightarrow \sigma_{\text{DSNU}} = i_{\text{distr}} \frac{\sigma_{\text{DSNU}}}{100}$$

PNRU - PHOTO CURRENT NON UNIFORMITIES

IT'S MEASURED AS THE DEVIATION OF
THE COLOR COORDINATES OF A PIXEL
FROM THE AVERAGE DUE TO GRIP FLUCTUATIONS

② DIVIDE THE RESULTS
BY THE AVERAGE
RGB OF THAT COLOR
FOR ALL PIXELS

④ REPEAT FOR MORE COLORS

- ① CAPTURE A COLOR WITH ALL
PIXEL (BY AVERAGING TO REMOVE
TEMPORAL NOISE)

- ③ WE EVALUATE FOR
EACH PIXEL THE
DEVIATION FROM
COLOR VALUE
AVERAGE IS σ_{PNRU}

$$\sigma_{\text{PNRU}} = i_{\text{pht}} t_{\text{int}} \frac{\sigma_{\text{PNRU}}}{100}$$

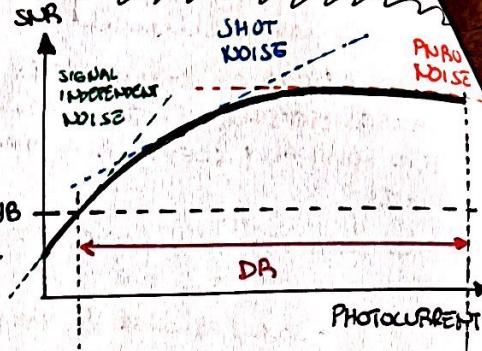
OVERALL NOISE

$$\sigma_{TOT} = \sqrt{\sigma_{READ}^2 + \sigma_{SHOT}^2 + \sigma_{PNR}^2}$$

ALL NOISE SOURCES SIGNAL INDEPENDENT (KTC, QUANTIZATION)

RELATED TO SIGNAL (CURRENT)

ASSOCIATED WITH PIXEL-TO-PIXEL RESPONSE DIFFERENCES



PTC - PHOTON TRANSFER CURVE

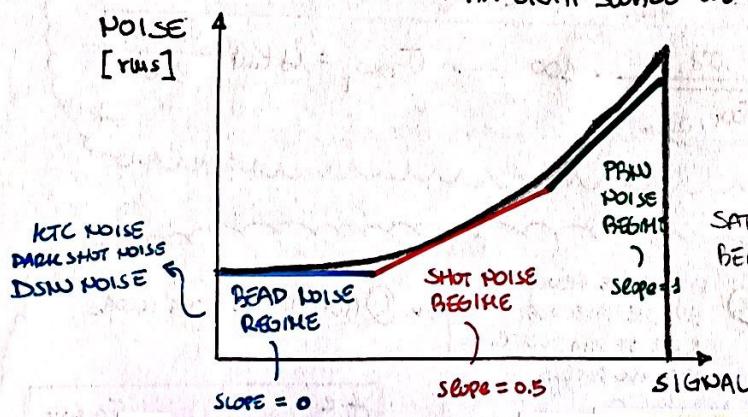
IT'S A TESTING METHODOLOGY EMPLOYED IN THE CHARACTERIZATION, OPTIMIZATION, CALIBRATION AND APPLICATION.

WE CAN EXTRACT QUANTUM EFFICIENCY, NOISE, CHARGE COLLECTION EFFICIENCY, SENSE NODE CAPACITANCE, FULL CHARGE, DYNAMIC RANGE, PIXEL RESPONSIVITY

PTC TREATS A CAMERA SYSTEM AS A BLACK BOX. WE NEED ONLY TO EXPOSE THE CAMERA TO A LIGHT SOURCE AND MEASURE THE SIGNAL AND NOISE OUTPUT RESPONSES

ACTUALLY THE SENSITIVITY OF AN IMAGE SENSOR IS A FUNCTION OF WAVELENGTH. IN THIS CONTEXT A LINEAR GAIN K (instead of an integral law over the visible spectrum) can be defined only if radiation spectrum is monochromatic. IF IT'S NOT WE CAN ASSUME A DOMINANT WAVELENGTH

BY EXPOSING TO A UNIFORM LIGHT SOURCE WE PLOT THE NOISE OVER AVERAGE OUTPUT:



- We can calculate the AVERAGE OUTPUT SIGNAL DN, its standard deviation σ_{DN} . we can plot (for various photon fluxes) average signal on x-axis^b and its standard deviation on y-axis^c [D.N. = DIGITAL NUMBERS]

from PTC curve ($\text{DN} - B$ x-axis, $\sigma_B - \text{DN}_{\text{rns}}$ y-axis) we can extract:

- the GAIN:

$$K = \left[\frac{\sigma_B^2}{B} \right]_{\sigma_B=1} = \frac{1}{\text{DN}_{\text{shot}}}$$

from theory:

$$K = \frac{\text{DN}}{N_{ph}} = \eta \left[\frac{2^{1/2}}{\text{FWC}} \right]_{\text{full-well charge}}$$

DN_{shot} is where $1/2$ shot noise line intercept the 1DN_{rns} noise

- the PNU:

$$\% \text{PNU} = \left[\frac{\sigma_{\text{PNU}}}{\text{DN}} \right]_{\sigma_{\text{PNU}}=1} = \frac{1}{\text{DN}_{\text{PNU}}}$$

DN_{PNU} is where PNU line intercept the 1DN_{rns} line

- the MAXIMUM NUMBER OF ELECTRONS:

$$N_{\text{sat}} = \frac{\text{DN}_{\text{sat}} \eta}{K} \quad [\text{number of } e^-]$$

- the READ NOISE:

$$\sigma_{\text{READ}, N} = \frac{\sigma_{\text{READ, IDN}} \eta}{K}$$

maximum slope = 0 line intercept y-axis [number of e^-]

- the DYNAMIC RANGE DR:

$$DR = \frac{N_{\text{sat}}}{\sigma_{\text{READ}, N}} \quad [\text{dB}]$$