

NEXMD with Sander's QM/MM

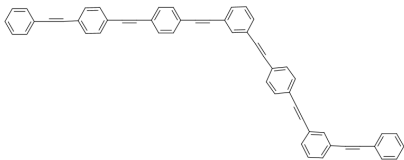
NonAdiabatic Dynamics Using AMBER

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Conjugate Organic Molecules



- They tend to be very good at energy transfer.
- $2 \rightarrow 3 \rightarrow 4$
- Excited state dynamics can speed or slow this transition.

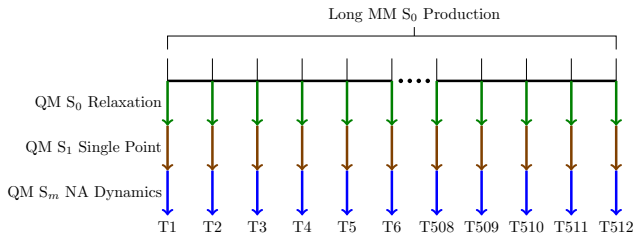
Question

How can we simulate this?

Possible Solutions

- Run Classical Ground State Dynamics then perform Excited State Energy Calculations on the snapshots.
 - ▶ Excited State dynamics could be different.
 - ▶ In the molecules I studied for example, the excited state had a much more planer dynamics than ground state.
- Run dynamics using the average forces of all the relevant excited states (Ehrenfest).
 - ▶ Often the average force of two states is not a very good approximation of either.
- Run a swarm of trajectories that are allowed to switch between different states (FSSH).
 - ▶ The proportion of the population that are on a certain state represent the probability of the state.
 - ▶ This process can be very computationally expensive.

General Outline of Dynamics



Example Times

- Perform a long MM trajectory.
- Take snapshot every 1 PS, run QM Ground State for relaxation.
- Excite the System.
- Allow each of these trajectories relax independently.
 - ▶ Sum the population in each state to get state probabilities over time.
 - ▶ Analyze absorption and emission properties by looking at frequencies and oscillator strengths.

Long MM Trajectory

Sampling

The goal of our analysis is to derive the probability of states of a solute in solution. There are many possible conformations for a molecule. If you then include a solvent, the number of snapshots needed to accurately sample the phase space increased even more. A long run is necessary for decent sampling.

Guidelines

- 1 ps time spacing between MM snapshots at 300K
- 500 samples for 50 atoms in vacuum

Ground State Calculations

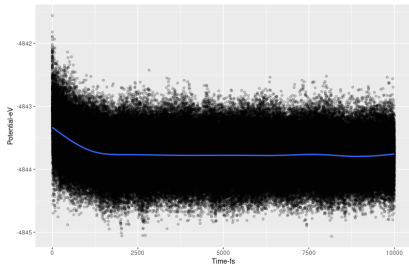
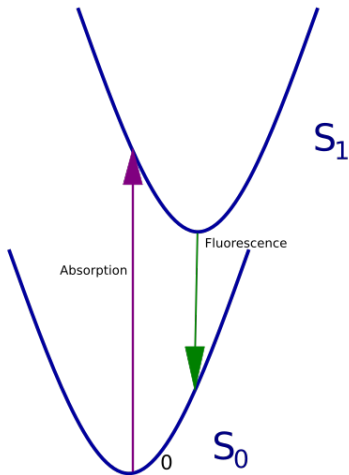


Figure: QM S_0 relaxation of PPV3-NO₂ in methanol.

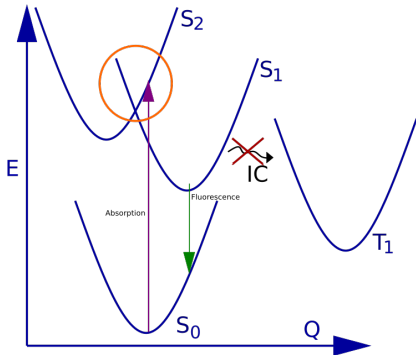
- The algorithm used to decide whether to jump is very sensitive to energies.
- QM/MM energies are not going to be the same as MM and some time is needed to properly relax.

Excite The System



- We use laser excitations with a wave shape similar to that done in experiments.
- Due to quantum uncertainty these states are also chosen based on the Franck-Condon principle.
- Energy is not conserved. (We're shooting a laser at it.)
- Velocities and coordinates are maintained in this step

Let the System Relax



- The molecule will start at a very high excited state.
- When near a crossing of the energies, we have a very strong coupling, the molecule can then hop between the states.
- When there's a switch the velocities of the atoms are adjusted to conserve energy.
- Phosphorescence is not allowed.

Analyze The Results

Population Decay

We can plot the proportion of trajectories in each states as a proxy to quantum probabilities of the states. This provides a detailed view of state decay rates.

Excited State Energy Relaxation

The molecule is suddenly placed outside of equilibrium. Modeling the excited state relaxation provides some insight into how this energy is dissipated to the thermostat, solvent, both.

Evolution of the Electronic Wave Function

The transition density matrix, needed for the calculation of the hopping probabilities can be repurposed to follow the physical movement of the excited electron density during decay.

Change in Vibrational Modes

Certain vibrational mode will be heavily coupled to adiabatic state transitions. Some of these could be very fast (Bonds) and others slower (Torsional).

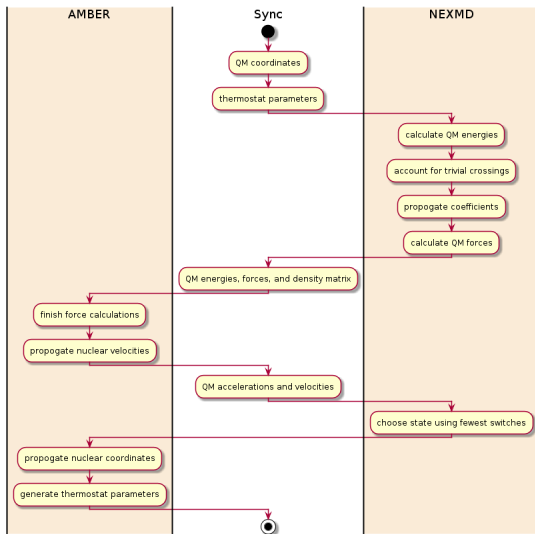
Why add QM/MM?

NEXMD can perform the above calculations already

However

- The choice to switch states is sensitive to energies.
- (Solvatochromic) Shifts in emission and absorption energies are known to occur in solvents.
- Applications such as sensors, fluorescent tags, L.E.D's, etc, will include solvent.
- NEXMD uses SQM as a backend, which is part of AMBERTOOLS.
 - ▶ This should make NEXMD trivial to use a library in AMBER
 - ▶ Unfortunately it didn't.

How We Did It



The Test Systems

polyphenylene vinylene oligomers

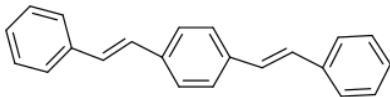


Figure: PPV3

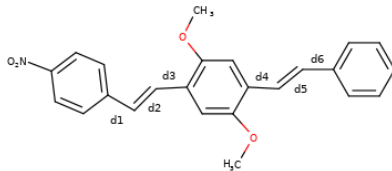
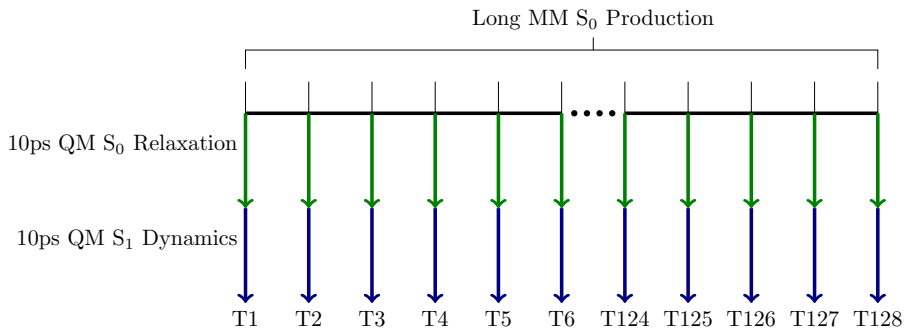


Figure: PPV3-NO₂

- These systems are small.
- These systems are well studied.
- The Hamiltonian is known to provide reasonable results.

PPV3-NO₂ Solvatochromic Shift



Solvent	Dielectric (debye)
CCL ₄	2.24
CH ₃ OH	32.7
CHCl ₃	4.81

Temperature: 300K
Langevin Constant: 2ps⁻¹

PPV3-NO₂ Solvatochromic Shift

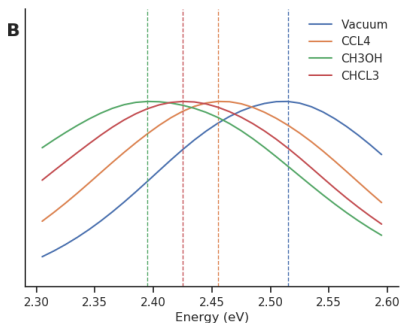


Figure: Emission shift found using explicit solvent.

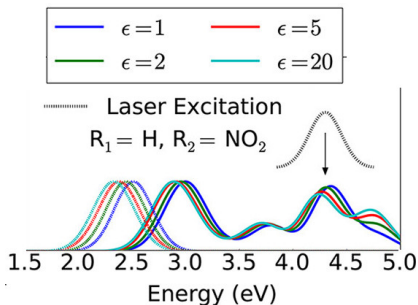
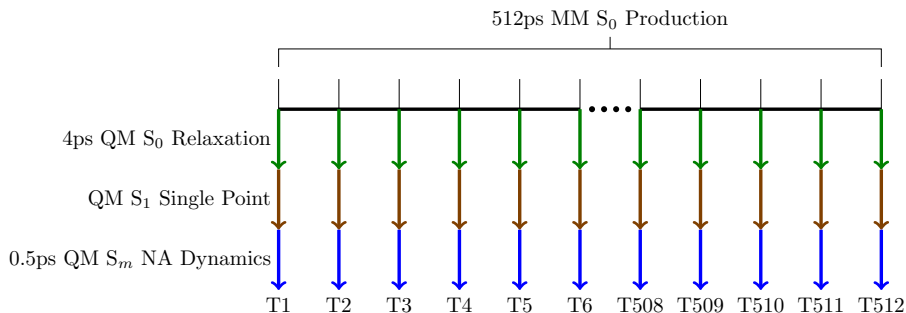


Figure: Emission shift found using implicit solvent. Dotted lines are the emission spectra. ¹

¹Andrew E Sifain et al. "Photoexcited nonadiabatic dynamics of solvated push-pull π -conjugated oligomers with the NEXMD software". In: *Journal of chemical theory and computation* 14.8 (2018), pp. 3955–3966.

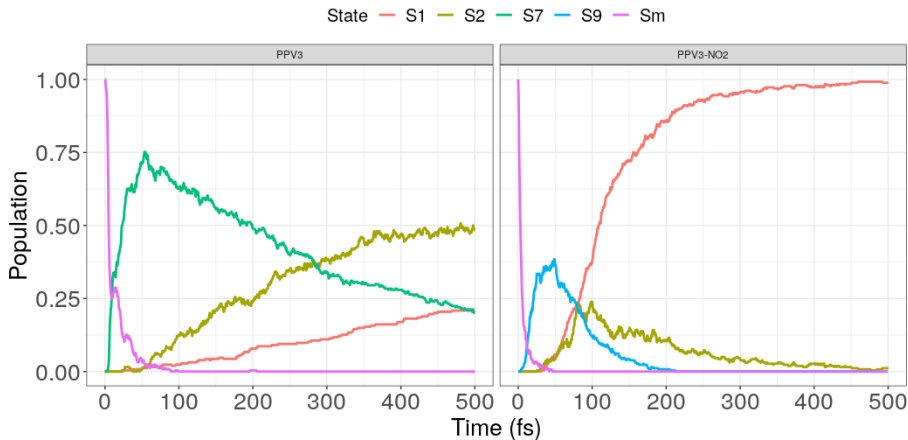
NonAdiabatic Calculation



Temperature: 300K

Langevin Constant: 2ps^{-1}

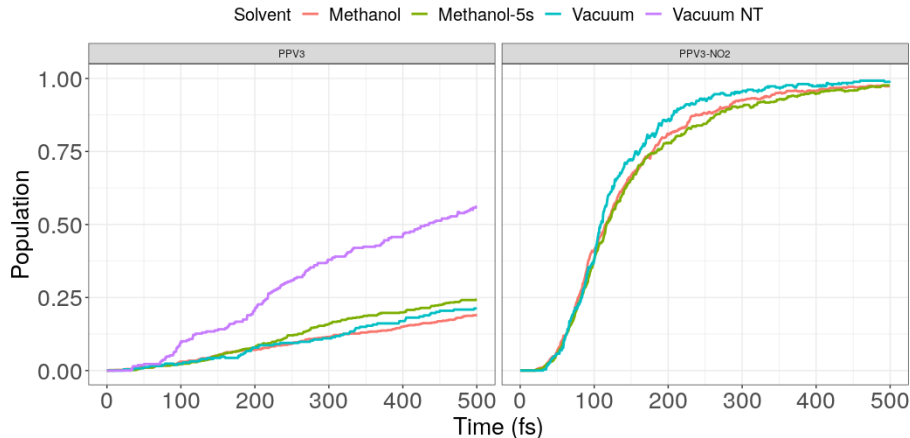
NonAdiabatic Calculation



Populations for the excited states with extended lifespans for PPV₃ and PPV3-NO₂. S_m represents the initial state.

The addition of the CH₃ groups and NO₂ dramatically accelerates decay.

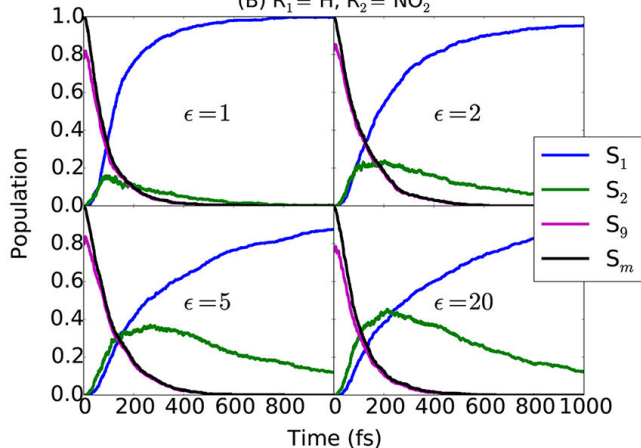
NonAdiabatic Calculation



Population of the first excited state (S_1) of PPV3 and PPV3-NO₂ in various solvents obtained from the fraction of trajectories in each state. NT represents the simulation that did not account for trivial crossings.

NonAdiabatic Calculation

(B) $R_1 = \text{H}$, $R_2 = \text{NO}_2$



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Closest thing to compare the previous results to. An important difference is the initial excited state and the langevin constant. What concerns me is the significant solvent related $S_2 \rightarrow S_1$ decay rate. Currently testing to this.

²Sifain et al., "Photoexcited nonadiabatic dynamics of solvated push-pull

Thank You!



Sifain, Andrew E et al. “Photoexcited nonadiabatic dynamics of solvated push–pull π -conjugated oligomers with the NEXMD software”. In: *Journal of chemical theory and computation* 14.8 (2018), pp. 3955–3966.