

PH 102, Electromagnetism,

Post Mid Semester

Lecture 11

Electromagnetic Waves in vacuum & matter

D. J. Griffiths: 9.2, 9.3

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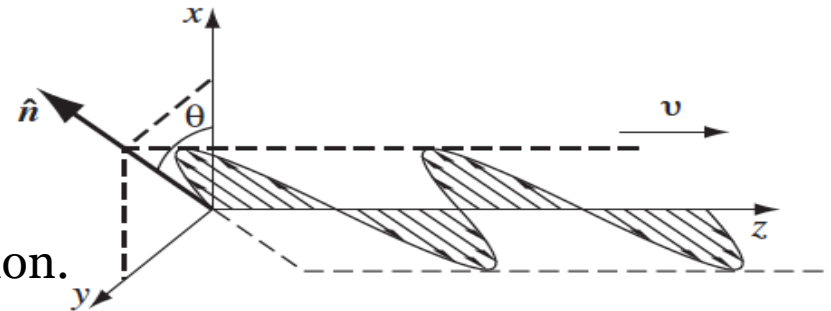
Electromagnetic Waves

Polarization :

Transverse wave: two independent states of polarization (plane of vibration)

along any other direction in the xy plane

$$\tilde{\mathbf{f}}(z, t) = \tilde{A}e^{i(kz - \omega t)} \hat{\mathbf{n}}.$$



Polarization vector $\hat{\mathbf{n}}$ defines the plane of vibration.

Transverse waves : $\hat{\mathbf{n}}$ is perpendicular to the direction of propagation

$$\hat{\mathbf{n}} \cdot \hat{\mathbf{z}} = 0.$$

In terms of the polarization angle θ , $\hat{\mathbf{n}} = \cos \theta \hat{\mathbf{x}} + \sin \theta \hat{\mathbf{y}}$.

Superposition of two waves—one horizontally polarized and the other one vertically:

$$\tilde{\mathbf{f}}(z, t) = (\tilde{A} \cos \theta)e^{i(kz - \omega t)} \hat{\mathbf{x}} + (\tilde{A} \sin \theta)e^{i(kz - \omega t)} \hat{\mathbf{y}}.$$

Electromagnetic Waves

The Wave Equation for \mathbf{E} and \mathbf{B} in vacuum:

For no charge or current, Maxwell's equations: coupled, first-order, pde for \mathbf{E} & \mathbf{B} .

$$(i) \quad \nabla \cdot \mathbf{E} = 0, \quad (iii) \quad \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t},$$

$$(ii) \quad \nabla \cdot \mathbf{B} = 0, \quad (iv) \quad \nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}.$$

and

$$\nabla^2 \mathbf{E} = \mu_0 \epsilon_0 \frac{\partial^2 \mathbf{E}}{\partial t^2}, \quad \nabla^2 \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial^2 \mathbf{B}}{\partial t^2}.$$

Interdependent,
one is the other!

Price for decoupling: 2nd order equations.

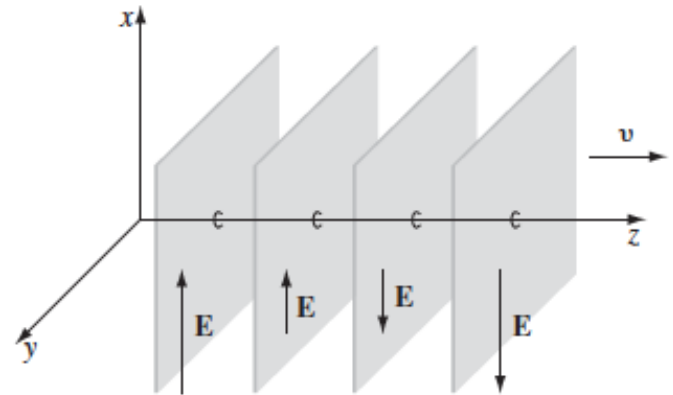
Electromagnetic Waves

Monochromatic Plane Waves :

- Sinusoidal waves of frequency ω : different frequencies in the visible range correspond to different colors, such waves are called ***monochromatic***.
- Waves traveling in the z direction, no x or y dependence: Plane waves
- The fields are ***uniform*** over ***every*** plane perpendicular to the propagation direction.

$$\tilde{\mathbf{E}}(z, t) = \tilde{\mathbf{E}}_0 e^{i(kz - \omega t)}, \quad \tilde{\mathbf{B}}(z, t) = \tilde{\mathbf{B}}_0 e^{i(kz - \omega t)},$$

where $\tilde{\mathbf{E}}_0$ and $\tilde{\mathbf{B}}_0$ are the (complex) amplitudes



Electromagnetic Waves

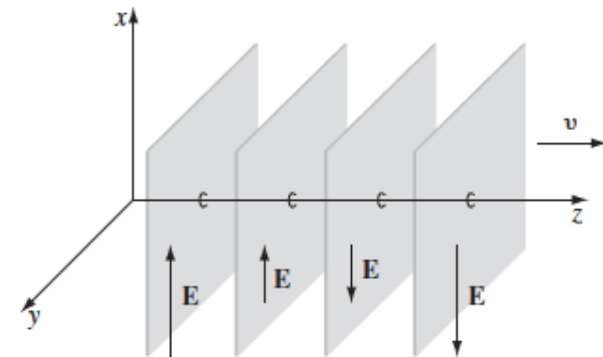
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Electromagnetic Waves

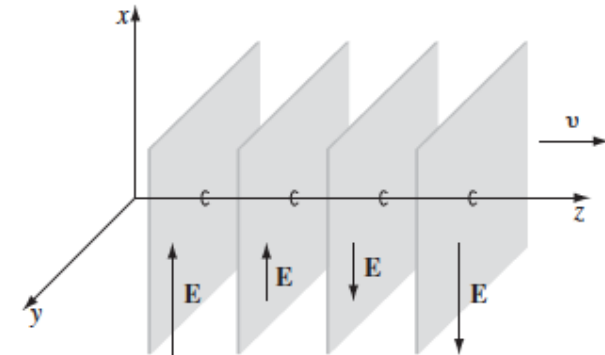
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$$\tilde{\mathbf{E}}(z, t) = \tilde{\mathbf{E}}_0 e^{i(kz - \omega t)}, \quad \tilde{\mathbf{B}}(z, t) = \tilde{\mathbf{B}}_0 e^{i(kz - \omega t)},$$

where $\tilde{\mathbf{E}}_0$ and $\tilde{\mathbf{B}}_0$ are the (complex) amplitudes



Extra constraints on $\tilde{\mathbf{E}}_0$ and $\tilde{\mathbf{B}}_0$ from Maxwell's equations

$$\nabla \cdot \mathbf{E} = 0 \text{ and } \nabla \cdot \mathbf{B} = 0, \quad \longrightarrow \quad (\tilde{E}_0)_z = (\tilde{B}_0)_z = 0. \quad \boxed{\mathbf{E}, \mathbf{B} \perp \mathbf{z}}$$

EM plane waves: No longitudinal component.

EM waves are purely transverse!

Electromagnetic Waves

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$$\tilde{\mathbf{E}}(z, t) = \tilde{\mathbf{E}}_0 e^{i(kz - \omega t)}, \quad \tilde{\mathbf{B}}(z, t) = \tilde{\mathbf{B}}_0 e^{i(kz - \omega t)},$$

where $\tilde{\mathbf{E}}_0$ and $\tilde{\mathbf{B}}_0$ are the (complex) amplitudes

Extra constraints on $\tilde{\mathbf{E}}_0$ and $\tilde{\mathbf{B}}_0$ from Maxwell's equations

Faraday's law,

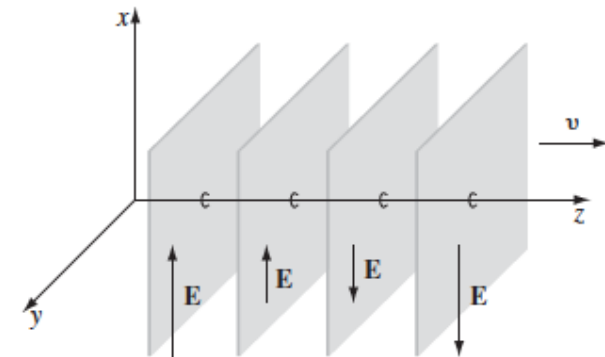
$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \quad \longrightarrow \quad -k(\tilde{E}_0)_y = \omega(\tilde{B}_0)_x, \quad k(\tilde{E}_0)_x = \omega(\tilde{B}_0)_y,$$

$$\tilde{\mathbf{B}}_0 = \frac{k}{\omega} (\hat{\mathbf{z}} \times \tilde{\mathbf{E}}_0).$$

\mathbf{E} & \mathbf{B} in phase & mutually perpendicular.

The real amplitudes,

$$B_0 = \frac{k}{\omega} E_0 = \frac{1}{c} E_0.$$



Electromagnetic Waves

Monochromatic Plane Waves :

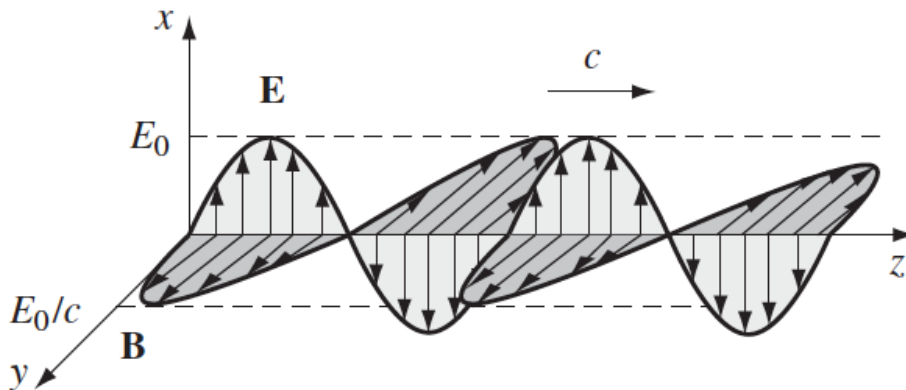
$$\tilde{\mathbf{E}}(z, t) = \tilde{\mathbf{E}}_0 e^{i(kz - \omega t)}, \quad \tilde{\mathbf{B}}(z, t) = \tilde{\mathbf{B}}_0 e^{i(kz - \omega t)},$$

Now, if \mathbf{E} points in the x direction then \mathbf{B} points in the y direction $[\tilde{\mathbf{B}}_0 = \frac{k}{\omega}(\hat{\mathbf{z}} \times \tilde{\mathbf{E}}_0).]$

$$\tilde{\mathbf{E}}(z, t) = \tilde{E}_0 e^{i(kz - \omega t)} \hat{\mathbf{x}}, \quad \tilde{\mathbf{B}}(z, t) = \frac{1}{c} \tilde{E}_0 e^{i(kz - \omega t)} \hat{\mathbf{y}},$$

The real part of the waves functions

$$\mathbf{E}(z, t) = E_0 \cos(kz - \omega t + \delta) \hat{\mathbf{x}}, \quad \mathbf{B}(z, t) = \frac{1}{c} E_0 \cos(kz - \omega t + \delta) \hat{\mathbf{y}}.$$



Monochromatic plane wave!

Note: Direction of E specify polarization direction

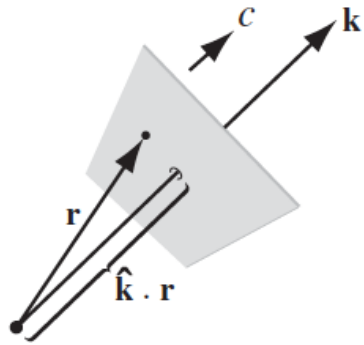
Electromagnetic Waves

Monochromatic Plane Waves :

Generalize to monochromatic plane waves traveling in an arbitrary direction.

(nothing special about the z direction)

Introducing, the wave vector \mathbf{k} (pointing in the direction of propagation)



The scalar product $\mathbf{k} \cdot \mathbf{r}$ is the appropriate generalization of kz so,

$$\tilde{\mathbf{E}}(\mathbf{r}, t) = \tilde{E}_0 e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \hat{\mathbf{n}},$$

$$\tilde{\mathbf{B}}(\mathbf{r}, t) = \frac{1}{c} \tilde{E}_0 e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} (\hat{\mathbf{k}} \times \hat{\mathbf{n}}) = \frac{1}{c} \hat{\mathbf{k}} \times \tilde{\mathbf{E}},$$

$\hat{\mathbf{n}}$ polarization vector
 \mathbf{E} being transverse, $\hat{\mathbf{n}} \cdot \hat{\mathbf{k}} = 0.$

The real \mathbf{E} & \mathbf{B} fields in a monochromatic plane wave with $(\mathbf{k}, \hat{\mathbf{n}})$

$$\mathbf{E}(\mathbf{r}, t) = E_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t + \delta) \hat{\mathbf{n}},$$

$$\mathbf{B}(\mathbf{r}, t) = \frac{1}{c} E_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t + \delta) (\hat{\mathbf{k}} \times \hat{\mathbf{n}}).$$

Electromagnetic Waves

Monochromatic Plane Waves :

Problem 9.9 Write down the (real) electric and magnetic fields for a monochromatic plane wave of amplitude E_0 , frequency ω , and phase angle zero that is (a) traveling in the negative x direction and polarized in the z direction; (b) traveling in the direction from the origin to the point $(1, 1, 1)$, with polarization parallel to the xz plane. In each case, sketch the wave, and give the explicit Cartesian components of \mathbf{k} and $\hat{\mathbf{n}}$.

Answer (a) $\mathbf{k} = -\frac{\omega}{c} \hat{\mathbf{x}}; \hat{\mathbf{n}} = \hat{\mathbf{z}}.$

$$\mathbf{k} \cdot \mathbf{r} = \left(-\frac{\omega}{c} \hat{\mathbf{x}}\right) \cdot (x \hat{\mathbf{x}} + y \hat{\mathbf{y}} + z \hat{\mathbf{z}}) = -\frac{\omega}{c} x;$$

$$\mathbf{k} \times \hat{\mathbf{n}} = -\hat{\mathbf{x}} \times \hat{\mathbf{z}} = \hat{\mathbf{y}}.$$

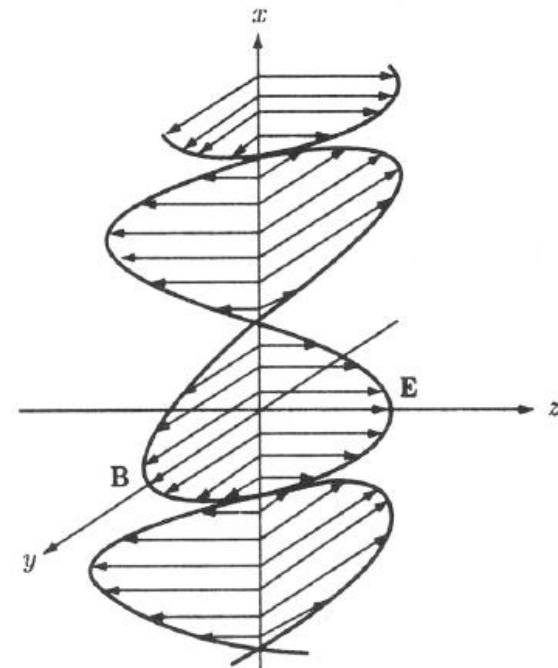
$$\mathbf{E}(\mathbf{r}, t) = E_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t + \delta) \hat{\mathbf{n}},$$

$$\mathbf{B}(\mathbf{r}, t) = \frac{1}{c} E_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t + \delta) (\hat{\mathbf{k}} \times \hat{\mathbf{n}}).$$

Therefore,

$$\mathbf{E}(x, t) = E_0 \cos\left(\frac{\omega}{c} x + \omega t\right) \hat{\mathbf{z}};$$

$$\mathbf{B}(x, t) = \frac{E_0}{c} \cos\left(\frac{\omega}{c} x + \omega t\right) \hat{\mathbf{y}}.$$



Electromagnetic Waves

Monochromatic Plane Waves :

Problem 9.9 Write down the (real) electric and magnetic fields for a monochromatic plane wave of amplitude E_0 , frequency ω , and phase angle zero that is (a) traveling in the negative x direction and polarized in the z direction; (b) traveling in the direction from the origin to the point $(1, 1, 1)$, with polarization parallel to the xz plane. In each case, sketch the wave, and give the explicit Cartesian components of \mathbf{k} and $\hat{\mathbf{n}}$.

Answer (a) $\mathbf{k} = \frac{\omega}{c} \left(\frac{\hat{\mathbf{x}} + \hat{\mathbf{y}} + \hat{\mathbf{z}}}{\sqrt{3}} \right); \hat{\mathbf{n}} = \frac{\hat{\mathbf{x}} - \hat{\mathbf{z}}}{\sqrt{2}}.$

$$\begin{aligned} \mathbf{E}(\mathbf{r}, t) &= E_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t + \delta) \hat{\mathbf{n}}, \\ \mathbf{B}(\mathbf{r}, t) &= \frac{1}{c} E_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t + \delta) (\hat{\mathbf{k}} \times \hat{\mathbf{n}}). \end{aligned}$$

Since $\hat{\mathbf{n}}$ is parallel to the xz plane, it must have the form $\alpha \hat{\mathbf{x}} + \beta \hat{\mathbf{z}}$;

$$\hat{\mathbf{n}} \cdot \mathbf{k} = 0, \beta = -\alpha \text{ \& \text{since it is a unit vector, } } \alpha = 1/\sqrt{2}.$$

$$\mathbf{k} \cdot \mathbf{r} = \frac{\omega}{\sqrt{3}c} (\hat{\mathbf{x}} + \hat{\mathbf{y}} + \hat{\mathbf{z}}) \cdot (x \hat{\mathbf{x}} + y \hat{\mathbf{y}} + z \hat{\mathbf{z}}) = \frac{\omega}{\sqrt{3}c} (x + y + z);$$

$$\hat{\mathbf{k}} \times \hat{\mathbf{n}} = \frac{1}{\sqrt{6}} (-\hat{\mathbf{x}} + 2\hat{\mathbf{y}} - \hat{\mathbf{z}}).$$

$$\mathbf{E}(x, y, z, t) = E_0 \cos \left[\frac{\omega}{\sqrt{3}c} (x + y + z) - \omega t \right] \left(\frac{\hat{\mathbf{x}} - \hat{\mathbf{z}}}{\sqrt{2}} \right);$$

$$\mathbf{B}(x, y, z, t) = \frac{E_0}{c} \cos \left[\frac{\omega}{\sqrt{3}c} (x + y + z) - \omega t \right] \left(\frac{-\hat{\mathbf{x}} + 2\hat{\mathbf{y}} - \hat{\mathbf{z}}}{\sqrt{6}} \right).$$

Sketch the wave!

Electromagnetic Waves

Energy and Momentum in Electromagnetic Waves

Energy per unit volume in electromagnetic fields is, $u = \frac{1}{2} \left(\epsilon_0 E^2 + \frac{1}{\mu_0} B^2 \right)$

For monochromatic plane wave, $B_0 = \frac{k}{\omega} E_0 = \frac{1}{c} E_0$. thus, $B^2 = \frac{1}{c^2} E^2 = \mu_0 \epsilon_0 E^2$,

i.e., the electric and magnetic contributions are equal!!

$$\begin{aligned} \mathbf{E}(z, t) &= E_0 \cos(kz - \omega t + \delta) \hat{\mathbf{x}}, \\ \mathbf{B}(z, t) &= \frac{1}{c} E_0 \cos(kz - \omega t + \delta) \hat{\mathbf{y}}. \end{aligned}$$

The wave carries the energy along with it.

$$u = \epsilon_0 E^2 = \epsilon_0 E_0^2 \cos^2(kz - \omega t + \delta)$$

Electromagnetic Waves

Energy flux density (energy per unit area, per unit time) transported by the fields is given by the Poynting vector,

$$\mathbf{S} = \frac{1}{\mu_0} (\mathbf{E} \times \mathbf{B})$$

$$\begin{aligned}\mathbf{E}(z, t) &= E_0 \cos(kz - \omega t + \delta) \hat{\mathbf{x}}, \\ \mathbf{B}(z, t) &= \frac{1}{c} E_0 \cos(kz - \omega t + \delta) \hat{\mathbf{y}}.\end{aligned}$$

For monochromatic plane waves propagating in the z direction,

$$\mathbf{S} = c\epsilon_0 E_0^2 \cos^2(kz - \omega t + \delta) \hat{\mathbf{z}} = cu \hat{\mathbf{z}}.$$

energy density (u)
times
wave velocity ($c \hat{\mathbf{z}}$)

EM fields also carry momentum,

the momentum density stored in the fields is,

$$\mathbf{g} = \frac{1}{c^2} \mathbf{S}$$

For monochromatic plane waves,
$$\mathbf{g} = \frac{1}{c} \epsilon_0 E_0^2 \cos^2(kz - \omega t + \delta) \hat{\mathbf{z}} = \frac{1}{c} u \hat{\mathbf{z}}.$$

Electromagnetic Waves

- For light, the wavelength is very short ($\sim 10^{-7}$ m) and so is the period ($\sim 10^{-15}$ s).
- Any macroscopic measurement will encompass many cycles.
- Time average: more interesting compare to fluctuating cosine-squared:

$$\mathbf{S} = c\epsilon_0 E_0^2 \cos^2(kz - \omega t + \delta) \hat{\mathbf{z}} = cu \hat{\mathbf{z}}. \quad \& \quad \mathbf{g} = \frac{1}{c^2} \mathbf{S}$$

$$\frac{1}{T} \int_0^T \cos^2(kz - 2\pi t/T + \delta) dt = 1/2.$$

Time average over a complete cycle

$$\langle u \rangle = \frac{1}{2} \epsilon_0 E_0^2,$$

$$\langle \mathbf{S} \rangle = \frac{1}{2} c \epsilon_0 E_0^2 \hat{\mathbf{z}},$$

$$\langle \mathbf{g} \rangle = \frac{1}{2c} \epsilon_0 E_0^2 \hat{\mathbf{z}}.$$

Average power per unit area transported by an EM wave: Intensity

$$I \equiv \langle S \rangle = \frac{1}{2} c \epsilon_0 E_0^2.$$

Electromagnetic Waves

Light (normal incidence) on perfect absorber: Delivers momentum to the surface.

In a time Δt , the momentum transfer is $\Delta p = \langle g \rangle A c \Delta t$,

Radiation pressure (average force per unit area)

$$P = \frac{1}{A} \frac{\Delta p}{\Delta t} = \frac{1}{2} \epsilon_0 E_0^2 = \frac{I}{c}$$

Qualitatively :

- \mathbf{E} drives charges in the \mathbf{x} direction
- \mathbf{B} exerts on them a force $q(\mathbf{v} \times \mathbf{B})$ in the \mathbf{z} direction.
- The net force on all the charges in the surface produces the pressure!

$$\begin{aligned}\mathbf{E}(z, t) &= E_0 \cos(kz - \omega t + \delta) \hat{\mathbf{x}}, \\ \mathbf{B}(z, t) &= \frac{1}{c} E_0 \cos(kz - \omega t + \delta) \hat{\mathbf{y}}.\end{aligned}$$

Electromagnetic Waves

Radiation pressure (average force per unit area)

$$P = \frac{1}{A} \frac{\Delta p}{\Delta t} = \frac{1}{2} \epsilon_0 E_0^2 = \frac{I}{c}$$

Problem 9.10 The intensity of sunlight hitting the earth is about 1300 W/m^2 . If sunlight strikes a perfect absorber, what pressure does it exert? How about a perfect reflector? What fraction of atmospheric pressure does this amount to?

Answer: $P = \frac{I}{c} = \frac{1.3 \times 10^3}{3.0 \times 10^8} = \boxed{4.3 \times 10^{-6} \text{ N/m}^2}.$

For a perfect reflector the pressure is twice as great: $8.6 \times 10^{-6} \text{ N/m}^2$.

(momentum switches direction)

Atmospheric pressure is $1.03 \times 10^5 \text{ N/m}^2$, so the pressure of light on a reflector is

$$(8.6 \times 10^{-6}) / (1.03 \times 10^5) = 8.3 \times 10^{-11} \text{ atmospheres.}$$

Electromagnetic Waves

Electromagnetic Waves in Matter:

Inside matter, no free charge and no free current, (i) $\nabla \cdot \mathbf{D} = 0$, (iii) $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$,

$$(ii) \quad \nabla \cdot \mathbf{B} = 0, \quad (iv) \quad \nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t}.$$

If the medium is linear and homogeneous,

$$\mathbf{D} = \epsilon \mathbf{E}, \quad \mathbf{H} = \frac{1}{\mu} \mathbf{B}, \quad \text{and}$$

$$(i) \quad \nabla \cdot \mathbf{E} = 0, \quad (iii) \quad \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t},$$

$$(ii) \quad \nabla \cdot \mathbf{B} = 0, \quad (iv) \quad \nabla \times \mathbf{B} = \mu \epsilon \frac{\partial \mathbf{E}}{\partial t},$$

Evidently, EM waves propagate through a linear homogeneous medium at a speed,

$$v = \frac{1}{\sqrt{\epsilon \mu}} = \frac{c}{n}, \quad \text{where, index of refraction } n \equiv \sqrt{\frac{\epsilon \mu}{\epsilon_0 \mu_0}}$$

For most materials, Permeability

$$\mu \equiv \mu_0(1 + \chi_m) \approx \mu_0 \quad n \cong \sqrt{\epsilon_r} > 1$$

Light travels slowly through matter! Optics!

Electromagnetic Waves

Electromagnetic Waves in Matter:

The energy density, $u = \frac{1}{2} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right)$

Poynting vector is $\mathbf{S} = \frac{1}{\mu} (\mathbf{E} \times \mathbf{B})$

Monochromatic plane waves: $\omega = kv$, Amplitude of $B \sim 1/v \times$ amplitude of E

The intensity is $I = \frac{1}{2} \epsilon v E_0^2$

Wave passes from one transparent medium to another,

$$(i) \epsilon_1 E_1^\perp = \epsilon_2 E_2^\perp, \quad (iii) \mathbf{E}_1^\parallel = \mathbf{E}_2^\parallel,$$

$$(ii) B_1^\perp = B_2^\perp, \quad (iv) \frac{1}{\mu_1} \mathbf{B}_1^\parallel = \frac{1}{\mu_2} \mathbf{B}_2^\parallel.$$

All previous results
carry over with

$$\begin{aligned} \epsilon_0 &\rightarrow \epsilon, \\ \mu_0 &\rightarrow \mu \\ &\& \\ c &\rightarrow v \end{aligned}$$

Electromagnetic Waves

Electromagnetic Waves in Matter:

Reflection and Transmission at Normal Incidence:

- Suppose, the xy plane forms the boundary between two linear media.
- A plane wave of frequency ω , traveling in the z direction and polarized in the x direction, approaches the interface from the left

Incident Wave:

$$\tilde{\mathbf{E}}_I(z, t) = \tilde{E}_{0_I} e^{i(k_1 z - \omega t)} \hat{\mathbf{x}},$$

$$\tilde{\mathbf{B}}_I(z, t) = \frac{1}{v_1} \tilde{E}_{0_I} e^{i(k_1 z - \omega t)} \hat{\mathbf{y}}.$$

Reflected Wave:

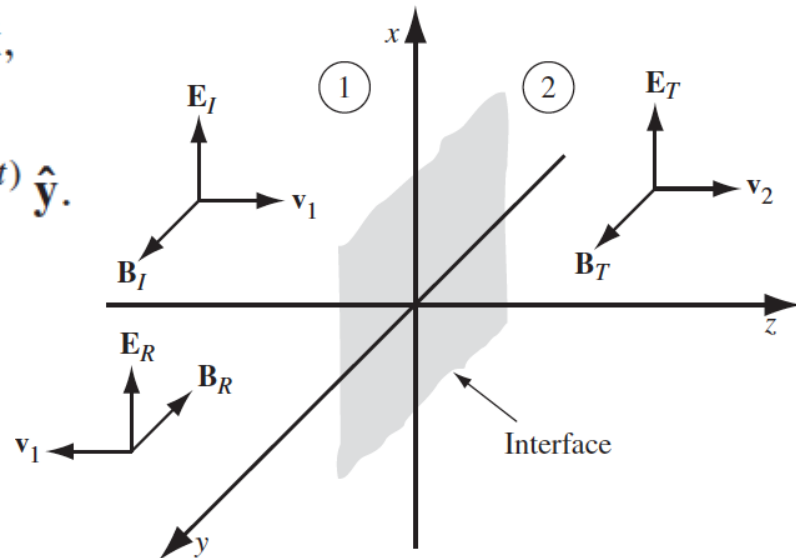
$$\tilde{\mathbf{E}}_R(z, t) = \tilde{E}_{0_R} e^{i(-k_1 z - \omega t)} \hat{\mathbf{x}},$$

$$\tilde{\mathbf{B}}_R(z, t) = -\frac{1}{v_1} \tilde{E}_{0_R} e^{i(-k_1 z - \omega t)} \hat{\mathbf{y}},$$

Transmitted Wave:

$$\tilde{\mathbf{E}}_T(z, t) = \tilde{E}_{0_T} e^{i(k_2 z - \omega t)} \hat{\mathbf{x}},$$

$$\tilde{\mathbf{B}}_T(z, t) = \frac{1}{v_2} \tilde{E}_{0_T} e^{i(k_2 z - \omega t)} \hat{\mathbf{y}},$$



Note the minus sign in $\tilde{\mathbf{B}}_R$

$$\tilde{\mathbf{B}}(\mathbf{r}, t) = \frac{1}{c} \tilde{E}_0 e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} (\hat{\mathbf{k}} \times \hat{\mathbf{n}})$$

Electromagnetic Waves

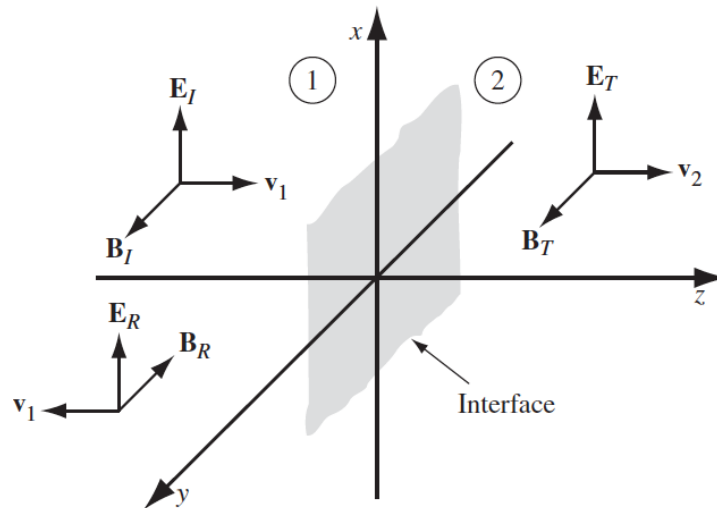
Reflection and Transmission at Normal Incidence:

At $z = 0$, the E & B fields on left (from both I and R) should join the fields on right (from T) in the boundary conditions.

$$(i) \epsilon_1 E_1^\perp = \epsilon_2 E_2^\perp, \quad (iii) \mathbf{E}_1^\parallel = \mathbf{E}_2^\parallel,$$

There are no components perpendicular to the surface, (i) & (ii) trivial.

$$(ii) B_1^\perp = B_2^\perp, \quad (iv) \frac{1}{\mu_1} \mathbf{B}_1^\parallel = \frac{1}{\mu_2} \mathbf{B}_2^\parallel.$$



Electromagnetic Waves

Reflection and Transmission at Normal Incidence:

On the interface,

$$(i) \epsilon_1 E_1^\perp = \epsilon_2 E_2^\perp,$$

$$(ii) B_1^\perp = B_2^\perp,$$

$$(iii) \mathbf{E}_1^\parallel = \mathbf{E}_2^\parallel,$$

$$(iv) \frac{1}{\mu_1} \mathbf{B}_1^\parallel = \frac{1}{\mu_2} \mathbf{B}_2^\parallel.$$

$$(iii) \tilde{E}_{0_I} + \tilde{E}_{0_R} = \tilde{E}_{0_T},$$

$$(iv) \frac{1}{\mu_1} \left(\frac{1}{v_1} \tilde{E}_{0_I} - \frac{1}{v_1} \tilde{E}_{0_R} \right) = \frac{1}{\mu_2} \left(\frac{1}{v_2} \tilde{E}_{0_T} \right)$$

Incident Wave:

$$\tilde{\mathbf{E}}_I(z, t) = \tilde{E}_{0_I} e^{i(k_1 z - \omega t)} \hat{\mathbf{x}},$$

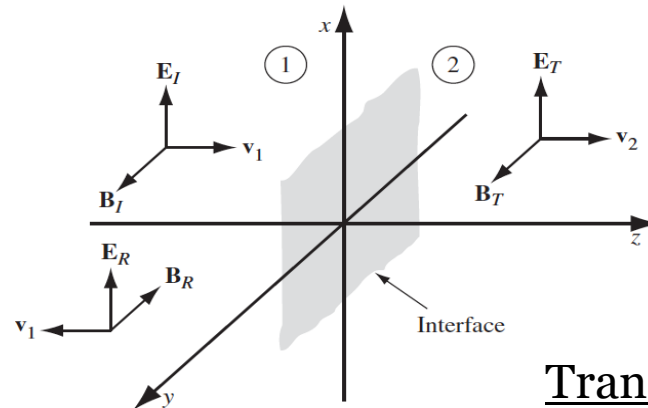
$$\tilde{\mathbf{B}}_I(z, t) = \frac{1}{v_1} \tilde{E}_{0_I} e^{i(k_1 z - \omega t)} \hat{\mathbf{y}}.$$

Reflected Wave:

$$\tilde{\mathbf{E}}_R(z, t) = \tilde{E}_{0_R} e^{i(-k_1 z - \omega t)} \hat{\mathbf{x}},$$

$$\tilde{\mathbf{B}}_R(z, t) = -\frac{1}{v_1} \tilde{E}_{0_R} e^{i(-k_1 z - \omega t)} \hat{\mathbf{y}},$$

Note the minus sign in $\tilde{\mathbf{B}}_R$



Transmitted Wave:

$$\tilde{\mathbf{E}}_T(z, t) = \tilde{E}_{0_T} e^{i(k_2 z - \omega t)} \hat{\mathbf{x}},$$

$$\tilde{\mathbf{B}}_T(z, t) = \frac{1}{v_2} \tilde{E}_{0_T} e^{i(k_2 z - \omega t)} \hat{\mathbf{y}},$$

Electromagnetic Waves

Reflection and Transmission at Normal Incidence:

$$(iii) \quad \tilde{E}_{0_I} + \tilde{E}_{0_R} = \tilde{E}_{0_T},$$

$$(iv) \quad \frac{1}{\mu_1} \left(\frac{1}{v_1} \tilde{E}_{0_I} - \frac{1}{v_1} \tilde{E}_{0_R} \right) = \frac{1}{\mu_2} \left(\frac{1}{v_2} \tilde{E}_{0_T} \right)$$

$$\tilde{E}_{0_I} - \tilde{E}_{0_R} = \beta \tilde{E}_{0_T},$$

where, $\beta \equiv \frac{\mu_1 v_1}{\mu_2 v_2} = \frac{\mu_1 n_2}{\mu_2 n_1}.$

Solving, (iii) & (iv),

$$\tilde{E}_{0_R} = \left(\frac{1 - \beta}{1 + \beta} \right) \tilde{E}_{0_I}, \quad \tilde{E}_{0_T} = \left(\frac{2}{1 + \beta} \right) \tilde{E}_{0_I}.$$

Permeability, $\mu \equiv \mu_0(1 + \chi_m) \approx \mu_0$ $\beta = v_1/v_2,$

$$\tilde{E}_{0_R} = \left(\frac{v_2 - v_1}{v_2 + v_1} \right) \tilde{E}_{0_I}, \quad \tilde{E}_{0_T} = \left(\frac{2v_2}{v_2 + v_1} \right) \tilde{E}_{0_I},$$

In phase for $v_2 > v_1$

In phase

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$$\tilde{E}_{0_R} = \left(\frac{v_2 - v_1}{v_2 + v_1} \right) \tilde{E}_{0_I}, \quad \tilde{E}_{0_T} = \left(\frac{2v_2}{v_2 + v_1} \right) \tilde{E}_{0_I},$$

In terms of the real amplitudes,

$$\begin{aligned} E_{0_R} &= \left| \frac{v_2 - v_1}{v_2 + v_1} \right| E_{0_I}, & E_{0_T} &= \left(\frac{2v_2}{v_2 + v_1} \right) E_{0_I}, \\ &= \left| \frac{n_1 - n_2}{n_1 + n_2} \right| E_{0_I}, & &= \left(\frac{2n_1}{n_1 + n_2} \right) E_{0_I}. \end{aligned} \quad \begin{array}{l} \text{(in terms of} \\ \text{refraction indices)} \end{array}$$

Let us see, what fraction of the incident energy is reflected / transmitted?

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Intensity (average power per unit area), $I = \frac{1}{2} \epsilon v E_0^2.$

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$$\begin{aligned} E_{0R} &= \left| \frac{v_2 - v_1}{v_2 + v_1} \right| E_{0I}, & E_{0T} &= \left(\frac{2v_2}{v_2 + v_1} \right) E_{0I}, \\ &= \left| \frac{n_1 - n_2}{n_1 + n_2} \right| E_{0I}, & &= \left(\frac{2n_1}{n_1 + n_2} \right) E_{0I}. \end{aligned} \quad \begin{array}{l} \text{(in terms of} \\ \text{refraction indices)} \end{array}$$

Let us see, what fraction of the incident energy is reflected / transmitted?

Intensity (average power per unit area), $I = \frac{1}{2} \epsilon v E_0^2$.

Reflection coefficient (R)

$$R \equiv \frac{I_R}{I_I} = \left(\frac{E_{0R}}{E_{0I}} \right)^2 = \left(\frac{n_1 - n_2}{n_1 + n_2} \right)^2,$$

Transmission coefficient (T)

$$T \equiv \frac{I_T}{I_I} = \frac{\epsilon_2 v_2}{\epsilon_1 v_1} \left(\frac{E_{0T}}{E_{0I}} \right)^2 = \frac{4n_1 n_2}{(n_1 + n_2)^2}.$$

Example. light passes from air ($n_1 = 1$) into glass ($n_2 = 1.5$),

R = 0.04 and T = 0.96.

No Surprise! R + T = 1