

Lecture 8

Conservation laws, D'Alembert's principle of virtual work,
Derivation of Lagrange's equation from Newton's law using
D'Alembert's principle

Symmetry and conservation laws

Cyclic coordinates, symmetries and conservation laws

Example: For planetary motion

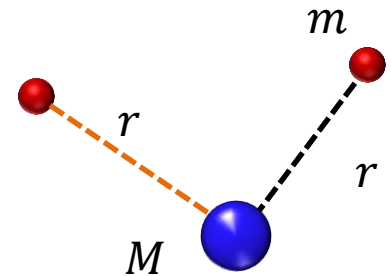
$$L = \frac{1}{2}m(\dot{r}^2 + r^2\dot{\theta}^2) + \frac{GMm}{r}$$

Here, θ is cyclic coordinate ($i, e. \frac{\partial L}{\partial \theta} = 0$) corresponding generalized (canonical) momentum $\mathbf{p}_\theta = \frac{\partial L}{\partial \dot{\theta}} = \mathbf{mr}^2\dot{\theta} = \mathbf{Constant}$; [Using Lagrange's eqn. $\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{\theta}}\right) - \frac{\partial L}{\partial \theta} = 0$]

This generalized momentum is nothing but angular momentum.

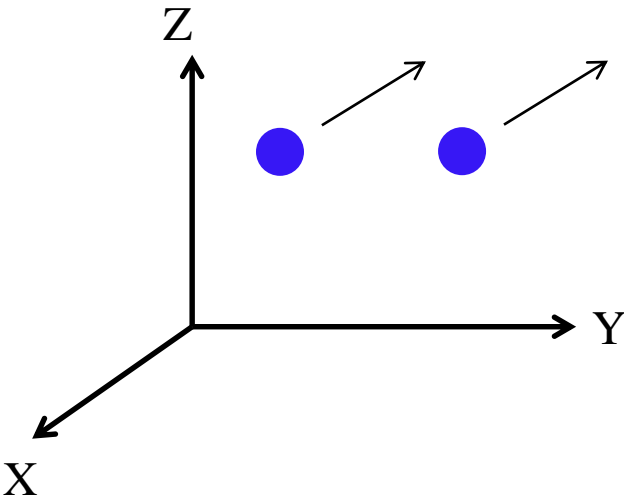
Another way of saying the same statement ' θ is cyclic':

L is independent of rotation angle θ , the system has rotational symmetry,
the system remains the same after any change in θ .



Conclusion: Conservation of angular momentum is related to rotational symmetry of the system.

Cyclic coordinates, symmetries and conservation laws



Lagrangian of a point mass moving in general direction under gravitational field.

$$L = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - mgz$$

x and y are cyclic coordinates, thus $p_x = \frac{\partial L}{\partial \dot{x}} = m\dot{x}$ and $p_y = \frac{\partial L}{\partial \dot{y}} = m\dot{y}$ both are constant of motion .

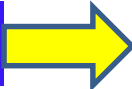
In this case, generalized momentum is the linear momentum.

In other words the system has translation symmetry in x and y coordinates. Any change in either x or y will not change the Lagrangian.

Conclusion: Conservation of linear momentum is associated with translational symmetry of the system.

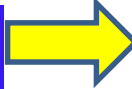
Symmetry and conservation laws

Translational symmetry



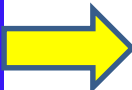
Conservation of linear momentum

Rotational symmetry



Conservation of angular momentum

?



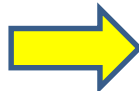
Conservation of energy

If L does not explicitly depend on time, then energy of the system is conserved, provided potential energy is velocity independent.

If L does not have explicit time dependence, change in time does not cause any change in Lagrangian i.e., $\frac{\partial L}{\partial t} = 0$,

$E = T + U = \text{constant}$ (looking for proof !)

Translational symmetry in time



Conservation of energy

Time-translation symmetry leads to energy conservation: Proof

Lagrangian $L = L(q_1, \dots, q_n, \dot{q}_1, \dots, \dot{q}_n, t)$

Using the chain rule of partial differentiation

$$dL = \sum_j \frac{\partial L}{\partial \dot{q}_j} d\dot{q}_j + \sum_j \frac{\partial L}{\partial q_j} dq_j + \frac{\partial L}{\partial t} dt$$

$$\frac{dL}{dt} = \sum_j \frac{\partial L}{\partial \dot{q}_j} \frac{d\dot{q}_j}{dt} + \sum_j \frac{\partial L}{\partial q_j} \frac{dq_j}{dt} + \frac{\partial L}{\partial t} = \sum_j \frac{\partial L}{\partial \dot{q}_j} \ddot{q}_j + \sum_j \frac{\partial L}{\partial q_j} \dot{q}_j + \frac{\partial L}{\partial t}$$

If L does not contain explicit time dependence, i.e. $L = L(q_1, \dots, q_n, \dot{q}_1, \dots, \dot{q}_n)$, Hence $\frac{\partial L}{\partial t} = 0$

Thus

$$\frac{dL}{dt} = \sum_j \frac{\partial L}{\partial \dot{q}_j} \ddot{q}_j + \sum_j \frac{\partial L}{\partial q_j} \dot{q}_j = \sum_j \frac{\partial L}{\partial \dot{q}_j} \ddot{q}_j + \sum_j \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_j} \right) \dot{q}_j$$

$$\frac{dL}{dt} = \sum_j \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_j} \dot{q}_j \right)$$

$$\frac{d}{dt} \left(L - \sum_j \frac{\partial L}{\partial \dot{q}_j} \dot{q}_j \right) = 0$$

Using Lagrange's
eqn
 $\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_j} \right) = \frac{\partial L}{\partial q_j}$

Equation 1

Proof continue...

Now, $L = T - U$

If U does not depend on generalized velocity ($\frac{\partial U}{\partial \dot{q}_j} = 0$) and only function of generalized coordinate, then $\frac{\partial L}{\partial \dot{q}_j} = \frac{\partial T}{\partial \dot{q}_j}$

$$\text{Thus from eqn 1. } \frac{d}{dt} \left(L - \sum_j \frac{\partial L}{\partial \dot{q}_j} \dot{q}_j \right) = 0; \frac{d}{dt} \left(L - \sum_j \frac{\partial T}{\partial \dot{q}_j} \dot{q}_j \right) = 0$$

Now, check that for a single free particle $T = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$

$$\sum_j \frac{\partial T}{\partial \dot{q}_j} \dot{q}_j = m(\dot{x}\dot{x} + \dot{y}\dot{y} + \dot{z}\dot{z}) = 2T$$

The relationship is true for general case as well, which can be obtained immediately from Euler's theorem. If $f(x_i)$ is a homogeneous function of the n_{th} degree of set of variables x_i , then $\sum_j \frac{\partial f}{\partial x_j} x_j = nf$.

Here, T is a function of 2nd degree of generalized velocities \dot{q}_j , thus $\sum_j \frac{\partial T}{\partial \dot{q}_j} \dot{q}_j = 2T$

$$\frac{d}{dt}(L - 2T) = 0; \frac{d}{dt}(T - U - 2T) = 0 \text{ hence } \frac{d}{dt}(T + U) = 0; \quad T + U = \text{constant}$$

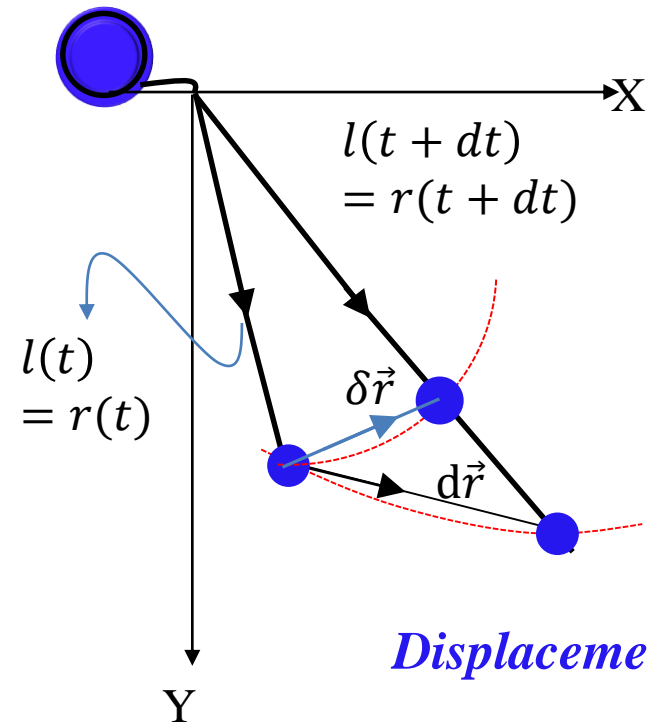
Conclusion, $T + U = E = \text{constant}$,

if L does not explicitly depend on time and potential is velocity independent

D'Alembert's principle of virtual work and proof of newton's law

Real vs Virtual displacement

Simple pendulum with a variable string length $l(t)$ [Time dependent constraint]



Real displacement of the bob in time dt is given by
$$d\vec{r} = \vec{r}(t + dt) - \vec{r}(t)$$

Let's imagine any instantaneous displacement at time t (that it, *without allowing the real time to change* $dt = 0$) which is consistent with the constrain relation at time t ?

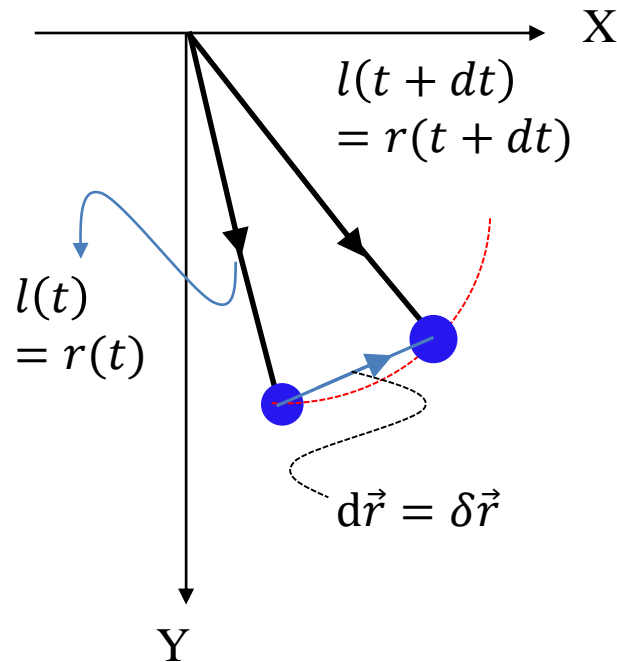
Displacement without allowing time to change -Seems absurd?

As you are making an imagination of displacement at time t ($dt = 0$), this displacement must not change the constrain condition at time t [i.e, $l(t)$ remains same] as shown.

Imaginary, instantaneous displacement which is consistent with the constrain relation at a given instant (i.e. without allowing real time to change) is called *Virtual displacement* and denoted by $\delta\vec{r}$ (for infinitesimal case)

Real vs Virtual displacement

- If the constrain is not time dependent, the real and virtual displacements matches each other, only the difference you need to remember that **virtual displacement is imagined without allowing time to change.**



Virtual displacement in generalized coordinates

- Consider a system of N particles with k constraints, DOF, $n = 3N - k$
- Cartesian coordinates, $\vec{r}_i = \vec{r}_i(x_i, y_i, z_i) \quad (i = 1, \dots, N)$
- Generalized coordinates $q_j \quad (j = 1, \dots, n)$
- Virtual displacements of the particles $\delta\vec{r}_1, \delta\vec{r}_2, \dots, \delta\vec{r}_N$

- Virtual displacements of the particles in the generalized coordinates $\delta q_1, \delta q_2, \dots, \delta q_n$ can be found from given transformation relations

$$\begin{aligned}\vec{r}_1 &= \vec{r}_1(q_1, q_2, \dots, q_n, t) \\ \vec{r}_2 &= \vec{r}_2(q_1, q_2, \dots, q_n, t) \\ &\dots\dots\dots \\ \vec{r}_N &= \vec{r}_N(q_1, q_2, \dots, q_n, t)\end{aligned}$$

$$\delta\vec{r}_i = \sum_{j=1}^n \frac{\partial \vec{r}_i}{\partial q_j} \delta q_j$$

$3N$ coordinates,
not independent

$n = 3N - k$ generalized
coordinates, independent

Note: There is no $\frac{\partial \vec{r}_i}{\partial t} \delta t$, as virtual displacement is instantaneous without allowing time to change, $\delta t=0$

Virtual work done

- ❑ **Real work done:** Work done due to real displacement ($d\vec{r}$) of a particle acted on by total force \vec{F} is given by

$$dW = \vec{F} \cdot d\vec{r}$$

- ❑ As you can always **imagine** an instantaneous displacement (without allowing time to change), known as virtual displacement ($\delta\vec{r}$), and hence you can always define a scalar function

$$\delta W = \vec{F} \cdot \delta\vec{r}$$

This scalar function is called as **Virtual work**.

Note: ‘**Virtual work**’ is different from ‘**Real work**’, as virtual displacement is imagined without allowing time to change.

If time is not allowed to change, no real displacement ($d\vec{r}$) is possible and hence no real work (dW) is possible without allowing time to change

Virtual work done for a system of particles

□ Consider a system of particles and $\vec{F}_1, \vec{F}_2, \dots, \vec{F}_N$ are the forces on 1, 2, ..., N_{th} particles, then

Total virtual work done

$$\delta W = \sum_{i=1}^N \vec{F}_i \cdot \delta \vec{r}_i$$

□ Here, force on each particle, \vec{F}_i is the sum of external force and also forces of constraints.

$$\vec{F}_i = \vec{F}_{ie} + \vec{f}_{ic}$$

Where,

- ❖ \vec{F}_{ie} is the external applied force on i_{th} particle.
- ❖ \vec{f}_{ic} is the constrain force

Virtual work for a dynamical system

□ Newton's second law for a single particle reads as

$$m\ddot{\vec{r}} = \vec{F}$$

$$m\ddot{\vec{r}} = \vec{F}_e + \vec{f}_c$$

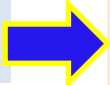
□ Taking dot product with an infinitesimal virtual displacement $\delta\vec{r}$

$$m\ddot{\vec{r}} \cdot \delta\vec{r} = (\vec{F}_e + \vec{f}_c) \cdot \delta\vec{r}$$

$$(\vec{F}_e - m\ddot{\vec{r}}) \cdot \delta\vec{r} + \vec{f}_c \cdot \delta\vec{r} = 0 \quad \boxed{1}$$

$$\begin{aligned} \text{Total force}(\vec{F}) &= \\ &\text{Applied force}(\vec{F}_e) \\ &+ \text{constrain force}(\vec{f}_c) \end{aligned}$$

Now, consider a general case of system of N particles having virtual displacements, $\delta\vec{r}_1, \delta\vec{r}_2, \dots, \delta\vec{r}_N$,

Eqn.1 for i_{th} particle  $(\vec{F}_{ie} - m_i\ddot{\vec{r}}_i) \cdot \delta\vec{r}_i + \vec{f}_{ic} \cdot \delta\vec{r}_i = 0$

Summing for all the particles of the system

$$\sum_{i=1}^N (\vec{F}_{ie} - m_i\ddot{\vec{r}}_i) \cdot \delta\vec{r}_i + \sum_{i=1}^N \vec{f}_{ic} \cdot \delta\vec{r}_i = 0$$

$\vec{F}_{ie} \rightarrow$ Applied force on i_{th} particle

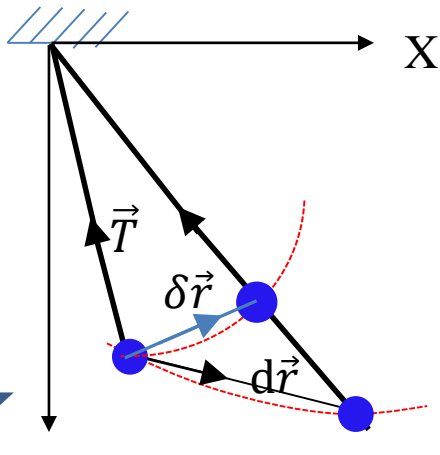
$\vec{f}_{ic} \rightarrow$ Constrain force on i_{th} particle

$$\sum_{i=1}^N \vec{f}_{ic} \cdot \delta \vec{r}_i$$

Total virtual work done by all the constrain forces

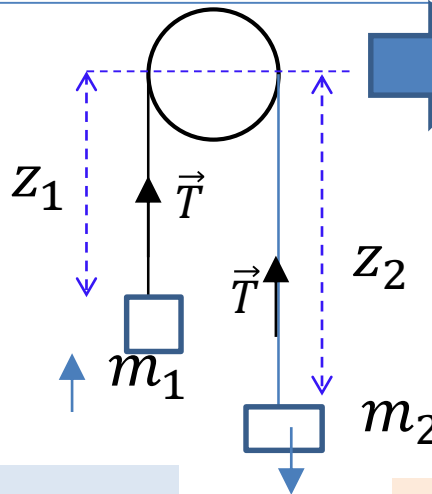
❑ Now, ‘virtual displacement’ is instantaneous, imaginary displacement **consistence with constrain relation**.

To keep the imaginary displacement, consistence with constrains at a given moment, either **each individual terms within this summation is zero** or **summation is zero as a whole** (even if individual terms are non-zero)



To keep the imaginary displacement, **consistence with constrains at a given moment**, the virtual displacement must be chosen perpendicular to constrain force.

$$\vec{f}_{ic} \cdot \delta \vec{r}_i = \vec{T} \cdot \delta \vec{r} = 0$$



To keep the imaginary displacement, **consistence with constrains at a given moment**, the virtual displacement must be chosen along *z - direction* and $\delta \vec{z}_2 = -\delta \vec{z}_1$

$$\begin{aligned} \sum_{i=1}^N \vec{f}_{ic} \cdot \delta \vec{r}_i &= \vec{T} \cdot \delta \vec{z}_1 + \vec{T} \cdot \delta \vec{z}_2 \\ &= T \delta z_1 - T \delta z_1 = 0 \end{aligned}$$


Note: Individual terms are not zero

D'Alembert's principle of virtual work

Since total virtual work done by the all the constraint forces is zero, I,e

$$\sum_{i=1}^N \vec{f}_{ic} \cdot \delta \vec{r}_i = 0$$

Then from equation 1

$$\sum_{i=1}^N (\vec{F}_{ie} - m_i \ddot{\vec{r}}_i) \cdot \delta \vec{r}_i = 0$$


D'Alembert's principle of Virtual work

$\vec{F}_{ie} \rightarrow$ Applied force on i_{th} particle

Does not necessarily means that individual terms of the summation are zero as \vec{r}_i **are not independent**, they are connected by constrain relation

Want to express this relation in such a way where all the terms in the summation becomes individually zero.

How to do?

Converting this relationship in terms of generalized coordinates

Quick recap of basic mathematics

If $u_1 \delta x_1 + u_2 \delta x_2 = 0$; does this always mean $u_1 = 0$ and $u_2 = 0$?

If x_1 and x_2 are independent then $u_1 = 0$ and $u_2 = 0$ for all possible variation of x_1 and x_2 ,

If x_1 and x_2 are independent then you can vary one without changing other. If you fix x_1 and vary x_2 , and still the relation is always giving zero, then only possibility is u_1 and u_2 must be zero.

If x_1 and x_2 are not independent, changing one will change the other.

Generalization:

$$\text{If, } \sum u_i \delta x_i = 0,$$

*then all u_i will be individually zero for all possible variation of the x_i only when x_i are **independent to each other**.*