# Physics II (PH 102) Electromagnetism (Lecture 8 & 9)

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Feb 2020

### Electrostatic Potential Energy of a Point Test Charge in Electric Field

Consider a stationary configuration of source charge distribution.

- ▶ Let E(r) be a pre-existing electric field with potential V(r) at a point P.
- ightharpoonup The electrostatic force  $\mathbf{F}_{field}$  on a positive test charge Q is

$$\mathbf{F}_{field} = Q\mathbf{E} = -\mathbf{F}_{ext}$$
.

#### Definition

Electrostatic Potential Energy of a test charge Q at point  $P(\mathbf{r})$  is equal to the amount of work done by an external agent against the electrostatic field to bring the charge Q from  $\infty$  (or ref) to the point  $P(\mathbf{r})$ :

$$U_{E}(\mathbf{r}) = \int_{-\infty}^{r} \mathbf{F}_{ext} \cdot d\mathbf{r}' = -\int_{-\infty}^{r} \mathbf{F}_{field} \cdot d\mathbf{r}' = -Q \int_{-\infty}^{r} \mathbf{E} \cdot d\mathbf{r}' = QV(\mathbf{r}).$$

**Note**: The ambiguity in the absolute value of  $U_E$  at a point like V!

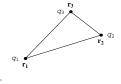
► Electrostatic Energy difference between two points **a** and **b** can be unambiguously expressed in terms of the Potential difference:

$$U_{E}(\mathbf{b}) - U_{E}(\mathbf{a}) = Q[V(\mathbf{b}) - V(\mathbf{a})] = -Q \int_{a}^{\mathbf{b}} \mathbf{E}(\mathbf{r}') \cdot d\mathbf{r}'.$$

# Potential Energy due to a system of Point Charges in free Field

Consider bringing in source charges  $q_1, q_2, q_3, \dots, q_n$ , one by one from  $\infty$ .

- No work done in placing first charge  $q_1$  at  $\mathbf{r_1}$ , i.e.,  $W_1=0$
- ► Total Work done is placing up to the second charge q<sub>2</sub> at r<sub>2</sub>:



$$W_2 = W_1 + \delta w_2 = 0 + \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{|\mathbf{r}_1 - \mathbf{r}_2|} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r_{12}}$$

► Total Work done in placing up to the third charge q<sub>3</sub> at r<sub>3</sub>:

$$W_3 = W_2 + \delta w_3 = W_2 + \frac{q_3}{4\pi\epsilon_0} \left( \frac{q_1}{r_{13}} + \frac{q_2}{r_{23}} \right)$$

• q<sub>5</sub> • q<sub>2</sub> • q<sub>3</sub> • q<sub>1</sub> • q<sub>4</sub>

Generalize formula up to n charges:  $(j \neq i \text{ Self interactions excluded!})$ 

$$W_n = W_{n-1} + \delta w_n = rac{1}{4\pi\epsilon_0} \sum_{i=1}^{n-1} \sum_{j=2,j>i}^n \left(rac{q_i q_j}{r_{ij}}
ight) 
ightarrow ext{No Double Counting}$$
  $\equiv rac{1}{8\pi\epsilon_0} \sum_{i=1}^n \sum_{j=2,j>i}^n \left(rac{q_i q_j}{r_{ij}}
ight) 
ightarrow ext{With Double Counting}$ 



### Potential Energy of a system of Point Charges in free Field (contd.)

▶ The Electrostatic Potential Energy is equal to the total work done  $W_n$  to assemble the configuration of n point charges at  $\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_n$ :

$$U_{E}(\mathbf{r}_{1}, \mathbf{r}_{2}, \dots, \mathbf{r}_{n}) \equiv W_{n} = \frac{1}{8\pi\epsilon_{0}} \sum_{i=1}^{n} \sum_{j=1, j \neq i}^{n} \left(\frac{q_{i}q_{j}}{r_{ij}}\right)$$

$$= \frac{1}{2} \sum_{i=1}^{n} q_{i} \left(\sum_{j=1, j \neq i}^{n} \frac{1}{4\pi\epsilon_{0}} \frac{q_{j}}{r_{ij}}\right)$$

$$= \frac{1}{2} \sum_{i=1}^{n} q_{i} V(\mathbf{r}_{i}) \equiv \frac{1}{2} \sum_{i=1}^{n} q_{i} V_{i} = \frac{1}{2} \sum_{i=1}^{n} u_{Ei}$$

- $V_i \equiv V(\mathbf{r}_i)$ : Potential at the  $i^{th}$  charge site  $\mathbf{r}_i$  due to other n-1 charges.
- ▶ The result is independent of the order/sequence in which the charges are assembled at the respective EXACT locations,  $\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_n$ .

$$\textit{U}_{\textit{E}}(r_1, r_2, \cdots, r_n) = \textit{U}_{\textit{E}}(r_2, r_n, \cdots, r_1) = \cdots = \textit{U}_{\textit{E}}\left[ \mathrm{Permute}(r_1, r_2, \cdots, r_n) \right]$$

► The result OBVIOUSLY, depends on the respective locations of the charges. Thus, *U<sub>E</sub>* is called the CONFIGURATION ENERGY.

# Configuration Energy of Point Charges in free space

► Superposition Principle is <u>invalid</u> (1/2 factor avoids **double counting!**):

$$U_E = \frac{1}{2} \sum_{i=1}^n u_{Ei}$$

#### Recall:

- $\blacktriangleright$  While defining  $U_E$ , we said "No work done in placing first charge  $q_1$  at  $r_1$ ."
- $\blacktriangleright$  All terms i = j were absent in  $U_E$ .
- ► SELF-ENERGIES of the individual point charges were excluded!

#### Definition

SELF-ENERGY: This is the amount of energy needed to fabricate or build-up the individual point charges by bringing their respective differential amounts of constituent charges from  $\infty$  to the specific locations.

# Electrostatic Potential Energy of General Distribution of Source Charges



(a) Discrete charges



(b) Line charge, λ



(c) Surface charge, σ



(d) Volume charge, ρ

#### Localized charge distributions:

- ▶ Volume  $\mathcal{V}$  with volume charge density  $\rho(\mathbf{r})$
- Surface S with surface charge density  $\sigma(\mathbf{r})$
- ► Curve  $\Gamma$  with linear charge density  $\lambda(\mathbf{r})$
- ▶ Discrete point charges  $q_i$  at  $\mathbf{r}'_i$
- ▶ Potential V(r)
- ► Configuration Energy of system of charges:

$$U_{E} \approx \frac{1}{2} \iiint_{\mathcal{V}} \rho(\mathbf{r}') V(\mathbf{r}') d\tau' + \frac{1}{2} \iint_{S} \sigma(\mathbf{r}') V(\mathbf{r}') da'$$
$$+ \frac{1}{2} \int_{\Gamma} \lambda(\mathbf{r}') V(\mathbf{r}') dl' + \frac{1}{2} \sum_{i=1}^{n} q_{i} V(\mathbf{r}'_{i})$$

- Note: Self-energies of continuous distributions are included but not for discrete point charges!
- What else is missing here? INTERACTION ENERGIES
- ► Superposition Principle is <u>invalid</u> in general!

## True Configuration Energy of General Localized Charge Distribution

Consider the most general localized charge distribution  $\rho_{\rm tot}$  in a region V, bounded by a closed surface S:

$$U_E = rac{1}{2} \iiint\limits_{\mathcal{V}} 
ho_{\mathrm{tot}}(\mathbf{r}') V(\mathbf{r}') \, d au'.$$

▶ Using Gauss's differential law <u>at source</u>:  $\rho_{\text{tot}}(\mathbf{r}') = \epsilon_0 \nabla' \cdot \mathbf{E}(\mathbf{r}')$ ,

$$\begin{split} U_{\mathsf{E}} &= \frac{1}{2} \iiint\limits_{\mathcal{V}} \rho(\mathbf{r}') V(\mathbf{r}') \, d\tau' = \frac{\epsilon_0}{2} \iiint\limits_{\mathcal{V}} \left[ \nabla' \cdot \mathbf{E}(\mathbf{r}') \right] V(\mathbf{r}') \, d\tau' \\ &= \frac{\epsilon_0}{2} \iiint\limits_{\mathcal{V}} \left[ \left| \mathbf{E}(\mathbf{r}') \right|^2 + \nabla' \cdot \left[ V(\mathbf{r}') \, \mathbf{E}(\mathbf{r}') \right] \right]_{\mathrm{source}} d\tau' \\ &= \frac{\epsilon_0}{2} \iiint\limits_{\mathcal{V}} \left| \mathbf{E}(\mathbf{r}') \right|_{\mathrm{source}}^2 d\tau' + \frac{\epsilon_0}{2} \iint\limits_{\mathcal{S}} \left[ V(\mathbf{r}') \, \mathbf{E}(\mathbf{r}') \right]_{\mathrm{boundary}} \cdot d\mathbf{a}'. \end{split}$$

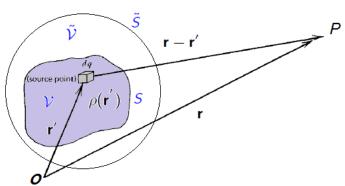
▶ In the last step, we applied the Gauss' Divergence Theorem to obtain the surface integral over the bounding surface S.

Recall Identity:  $\nabla \cdot (V \mathbf{E}) = V (\nabla \cdot \mathbf{E}) + (\nabla V) \cdot \mathbf{E} = V (\nabla \cdot \mathbf{E}) - |\mathbf{E}|^2$ 



Let us EXTEND the integration over a very large SPHERICAL volume  $\tilde{\mathcal{V}}$ , with bounding surface  $\tilde{\mathcal{S}}$ , enclosing the LOCALIZED distribution  $\mathcal{V}$ :

$$U_{E} = rac{1}{2} \iiint\limits_{\mathcal{V}} 
ho_{ ext{tot}}(\mathbf{r}') V(\mathbf{r}') \, d au' = rac{1}{2} \iiint\limits_{\mathcal{V}} 
ho_{ ext{tot}}(\mathbf{r}') V(\mathbf{r}') \, d au'$$



- lacktriangle For  ${
  m localized}$  distribution at a distant point P,  $V({
  m r}) \propto 1/r, \ |{
  m E}({
  m r})| \propto 1/r^2$
- For points  $\mathbf{r}'$  on this very large spherical surface  $\tilde{S}$ ,

$$\iint [V(\mathbf{r}') \mathbf{E}(\mathbf{r}')]_{\text{boundary}} \cdot d\mathbf{a}' \propto \left(\frac{1}{r'} \cdot \frac{1}{r'^2} \cdot r'^2\right)_{\tilde{S}} \sim \#\left(\frac{1}{r'}\right)_{\tilde{S}} \to 0, \ r' \to \infty.$$

## True Configuration Energy of General Charge Distribution (contd.)

Extending to include ALL SPACE:  $\tilde{\mathcal{V}} \to \mathcal{V}_{\infty} \equiv \mathbb{R}^3$  and  $\tilde{S} \to S_{\infty}$ , then the surface integral vanishes!

$$\begin{array}{ll} \textit{U}_{E} & = & \frac{1}{2} \iiint\limits_{\mathcal{V}} \rho_{\mathrm{tot}}(\mathbf{r}') \textit{V}(\mathbf{r}') \, d\tau' = \frac{1}{2} \iiint\limits_{\tilde{\mathcal{V}}} \rho_{\mathrm{tot}}(\mathbf{r}') \textit{V}(\mathbf{r}') \, d\tau' \\ \\ & = & \frac{\epsilon_{0}}{2} \iiint\limits_{\tilde{\mathcal{V}} \rightarrow \mathcal{V}_{\infty}} |\mathbf{E}|_{\mathrm{source}}^{2} \, d\tau' + \frac{\epsilon_{0}}{2} \iint\limits_{\tilde{\mathcal{S}} \rightarrow \tilde{\mathcal{S}}_{\infty}} (\textit{V} \, \mathbf{E})_{\mathrm{boundary}} \cdot d\mathbf{a}' \\ \\ & = & \frac{\epsilon_{0}}{2} \iiint\limits_{\mathrm{All \ Space}} |\mathbf{E}(\mathbf{r}')|_{\mathrm{source}}^{2} \, d\tau' \end{array}$$

All inclusive formula: SELF-ENERGIES and INTERACTION ENERGIES of ALL localized charge distributions.

**In Summary:** 2 Prescriptions to determine the Configuration Energy of a Localized distribution (depending on convenience) in a given problem

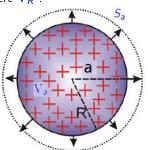
$$U_E = rac{1}{2} \iiint \int 
ho_{
m tot}({f r}) V({f r}) \, d au = rac{\epsilon_0}{2} \iiint \int 
ho_{
m pace} |{f E}({f r})|^2 \, d au$$

## Configuration Energy of a Charged Sphere

#### Example

Determine the total configuration/self energy of a uniformly charged solid sphere  $V_R$  of radius R and charge q.

▶ AUGMENTED VOLUME: Consider a very large <u>concentric</u> spherical volume  $V_a$  of radius  $a\gg R$ , with its bounding surface  $S_a$ , enclosing the original charged sphere  $V_R$ .



$$U_{E} = \frac{1}{2} \iiint\limits_{V_{R} \rightarrow V_{g}} \rho(\mathbf{r}) V(\mathbf{r}) d\tau = \frac{\epsilon_{0}}{2} \iiint\limits_{V_{g}} \left| \mathbf{E}(\mathbf{r}) \right|^{2} d\tau + \frac{\epsilon_{0}}{2} \oiint\limits_{S_{g}} V(\mathbf{r}) \mathbf{E}(\mathbf{r}) \cdot d\mathbf{a}$$

ightharpoonup Electric fields due to **original** charged sphere  $V_R$  (by Using Gauss's Law):

$$\mathsf{E}(\mathsf{r}) = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{qr}{R^3} \hat{\mathsf{r}} & r < R \\ \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathsf{r}} & r > R \end{cases}$$

Potentials due to **original** charged sphere  $V_R$  (using  $V(r) = -\int_{-\infty}^{r} \mathbf{E} \cdot d\mathbf{I}$ ):

$$V(\mathbf{r}) = \begin{cases} \frac{q}{8\pi\epsilon_0 R} \left( 3 - \frac{r^2}{R^2} \right) & r < R \\ \frac{1}{4\pi\epsilon_0} \frac{q}{r} & r > R \end{cases}$$

▶ Volume integral over the **augmented** volume  $V_a$  of radius  $a \gg R$ :

$$\frac{\epsilon_0}{2} \iiint_{\mathbf{V}_a} |\mathbf{E}(\mathbf{r})|^2 d\tau = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0}\right)^2 \left\{ \int_0^R \left(\frac{r^2}{R^6}\right) r^2 dr + \int_R^a \left(\frac{1}{r^4}\right) r^2 dr \right\} \int_0^{\pi_A} d\Omega$$

$$= \frac{q^2}{8\pi\epsilon_0} \left\{ \left(\frac{1}{5R}\right)_{r \le R} + \left(\frac{1}{R} - \frac{1}{a}\right)_{R < r \le a} \right\}$$

▶ Surface integral over the **augmented** sphere  $S_a$  of radius  $a \gg R$ :

$$\underbrace{\frac{\epsilon_0}{2} \oiint \left[ V(a\hat{\mathbf{r}}) \, \mathbf{E}(a\hat{\mathbf{r}}) \right] \cdot d\mathbf{a}}_{\epsilon} = \underbrace{\frac{\epsilon_0}{2} \left( \frac{q}{4\pi\epsilon_0 a} \right) \left( \frac{q}{4\pi\epsilon_0 a^2} \right) \left( 4\pi a^2 \right) = \left( \frac{q^2}{8\pi\epsilon_0 r} \right)_{r=a}$$

Extend  $V_a$  to include ALL SPACE:

$$V_a o \mathcal{V}_\infty \equiv \mathbb{R}^3$$
 &  $S_a o S_\infty \Longrightarrow a o \infty$ 

 $\triangleright$  Configuration energy (self-energy) of the charged sphere  $V_R$ :

$$\begin{split} U_E &= \frac{\epsilon_0}{2} \iiint_{V_a \to V_\infty} |\mathbf{E}(\mathbf{r})|^2 \ d\tau + \frac{\epsilon_0}{2} \iiint_{S_a \to S_\infty} V(\mathbf{r}) \, \mathbf{E}(\mathbf{r}) \cdot d\mathbf{a} \\ &= \lim_{a \to \infty} \left[ \frac{q^2}{8\pi\epsilon_0} \left\{ \frac{1}{5R} + \left( \frac{1}{R} - \frac{1}{A} \right) \right\} + \frac{q^2}{8\pi\epsilon_0 a} \right]^0 \\ &= \frac{1}{4\pi\epsilon_0} \left( \frac{3q^2}{5R} \right) \\ U_E &\equiv \frac{\epsilon_0}{2} \iiint_{\text{All Space}} |\mathbf{E}(\mathbf{r})|^2 \ d\tau = \frac{1}{4\pi\epsilon_0} \left( \frac{3q^2}{5R} \right). \end{split}$$

 $\triangleright$  Check by <u>direct</u> integration over the **original** charged sphere  $V_R$ :

$$U_{E} = \frac{1}{2} \iiint \rho(\mathbf{r}) V(\mathbf{r}) d\tau = \frac{\epsilon_{0}}{2} \left(\frac{q}{4\pi\epsilon_{0}}\right)^{2} \int \limits_{s}^{R} \left(\frac{r^{2}}{R^{6}}\right) r^{2} dr \int \limits_{s}^{4\pi} d\Omega = \frac{1}{4\pi\epsilon_{0}} \left(\frac{3q^{2}}{5R}\right).$$

## Self-energy of a Point Charge

► Electric field of a point charge placed at origin:

$$\mathsf{E}(\mathsf{r}) = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathsf{r}}$$

► Self-Energy:

$$\begin{split} W &= \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |\mathbf{E}(\mathbf{r})|^2 d\tau &= \frac{\epsilon_0}{2} \left( \frac{q}{4\pi\epsilon_0} \right)^2 \left\{ \int_0^\infty \left( \frac{1}{r^4} \right) r^2 dr \int\limits_0^\pi \sin\theta d\theta \int\limits_0^{2\pi} d\phi \right\} \\ &= \frac{\epsilon_0}{2} \left( \frac{q}{4\pi\epsilon_0} \right)^2 \left\{ \int_0^\infty \left( \frac{1}{r^4} \right) 4\pi r^2 dr \right\} \\ &= \frac{q^2}{8\pi\epsilon_0} \int_0^\infty \frac{1}{r^2} dr \\ &= \frac{q^2}{8\pi\epsilon_0} \lim_{\delta R \to 0} \int_{\epsilon_R}^\infty \frac{1}{r^2} dr = \frac{q^2}{8\pi\epsilon_0} \lim_{\delta R \to 0} \frac{1}{\delta R} \Rightarrow \infty \end{split}$$

- ightharpoonup Since the radius  $\delta R$  of the point charge vanishes, the self-energy blow up!
- Point charges are <u>idealized</u> concepts. In reality  $\delta R \neq 0$  (say, for electrons), so **Self-Energies** (classically) of spherical objects of finite radius  $\delta R$  is:

$$W_{\rm sphere} = \frac{1}{4\pi\epsilon_0} \left( \frac{3q^2}{5\delta R} \right) \rightarrow \text{finite}$$

### Interaction Energy of two Point Charges

#### Example

Find the Interaction energy of two charges,  $q_1$  and  $q_2$  located at  $r_1$  and  $r_2$ , respectively.

Net Electric field at any point r (Superposition Principle)

$$\mathsf{E}(\mathsf{r}) = \mathsf{E}_1(\mathsf{r}) + \mathsf{E}_2(\mathsf{r})\,,$$

#### Total Configuration Energy:

$$\begin{split} W &= \frac{\epsilon_0}{2} \iiint\limits_{\mathrm{All~Space}} |\mathsf{E}(\mathsf{r})|^2 \, d\tau = \frac{\epsilon_0}{2} \iiint\limits_{\mathrm{All~Space}} |\mathsf{E}_1(\mathsf{r}) + \mathsf{E}_2(\mathsf{r})|^2 \, d\tau \\ &= \frac{\epsilon_0}{2} \iiint\limits_{\mathrm{All~Space}} |\mathsf{E}_1(\mathsf{r})|^2 \, d\tau + \frac{\epsilon_0}{2} \iiint\limits_{\mathrm{Space}} |\mathsf{E}_2(\mathsf{r})|^2 \, d\tau + \frac{\epsilon_0}{2} \iiint\limits_{\mathrm{All~Space}} 2\mathsf{E}_1(\mathsf{r}) \cdot \mathsf{E}_2(\mathsf{r}) \, d\tau \end{split}$$

#### Interaction Energy:

$$W^{\rm int} = \epsilon_0 \iiint\limits_{\rm All~Space} {\bf E_1(r) \cdot E_2(r)} \ d\tau = \frac{q_1 q_2}{16 \pi^2 \epsilon_0} \iiint\limits_{\rm All~Space} \frac{({\bf r} - {\bf r_1}) \cdot ({\bf r} - {\bf r_2})}{|{\bf r} - {\bf r_1}|^3 |{\bf r} - {\bf r_2}|^3} d\tau = \frac{q_1 q_2}{4 \pi \epsilon_0 r_{12}}$$

 $\Rightarrow$  This is exactly the work done by an external agent in bringing  $q_2$  from  $\infty$  to  $r_2$  with  $q_1$  already present at  $r_1$ .



### Total Configuration Energy & Density

- ► Interesting Question: Where is the total energy stored?
- ...in the charges?...in the fields? No unique answer to that question!
- ► The equation involves integration over all charge distributions:

$$U_E = \frac{1}{2} \iiint\limits_{\mathcal{V}} \rho(\mathbf{r}) V(\mathbf{r}) \ d^3 r$$

⇒ suggests that the energy may be stored in the charges.

► The equation involves integration over all field configurations:

$$U_E = \frac{\epsilon_0}{2} \iiint_{\text{All Space}} E^2(\mathbf{r}) \ d^3 r$$

 $\Longrightarrow$  suggests that the energy may be stored in the fields.

▶ It is conventional to define an **ENERGY DENSITY**:

$$u(\mathbf{r}) = \frac{\epsilon_0}{2} \left| \mathbf{E} \left( \mathbf{r} \right) \right|^2$$

 $\Longrightarrow$  a volume dv will contain Electrostatic Potential Energy equal to  $u(\mathbf{r})dv$ .

