

**key ideas:** { Current Stellar Composition  
reaction rates/cross-sections for key reaction pathways  
effortless simulation techniques

## 1.2 Nomenclature.

→ { Z: atomic/proton number  
N: neutron number  
A: nucleon/mass number

$$A = Z + N$$

→ Same  $Z$ , same  $N$ , same nuclear property.

→ symbol:  $A_z X_N$  (where  $X$  is the element symbol)

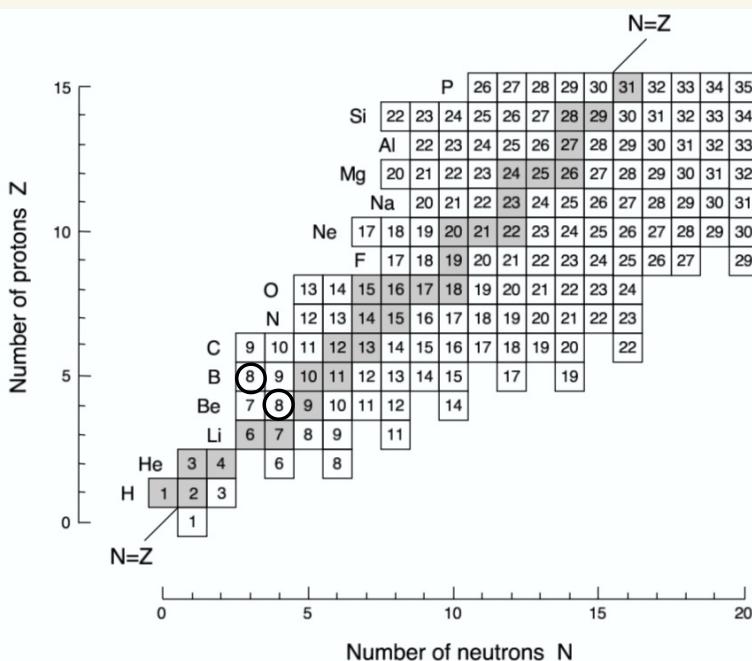
→ A distinct kind of atom, defined by specific atomic and neutron number, is called a nuclide.

nuclide with  $\rightarrow$  Same A, different Z and N  $\rightarrow$  isobars

different A, same Z  $\longrightarrow$  isotopes

different A, same N  $\longrightarrow$  isotones

## → chart of the nuclides



**Fig. 1.1** Section of the chart of the nuclides, showing the lightest species with  $Z \leq 15$  and  $N \leq 20$ . The shaded squares represent stable nuclides, while the open squares correspond to unstable nuclides with half-lives in excess of 1 ms. The only exceptions are the nuclides  $^8\text{Be}$  and  $^9\text{B}$  which have much shorter half-lives. Note that no stable nuclides exist with a mass number of  $A = 5$  or  $8$ .

→ Shaded box are considered  
has +ve half lives.

### 1.3 Solar system abundances

- Solar system formed from the collapse of a gaseous nebula that has uniform chemical and isotopic distribution.
- Method 1 : Solar photosphere → outer layer of a star, 5800K, most light can escape.
- Method 2 : primitive meteorites (CI carbonaceous chondrites)  
 (These two methods give result with uncertainties less than 10% for the most elements)

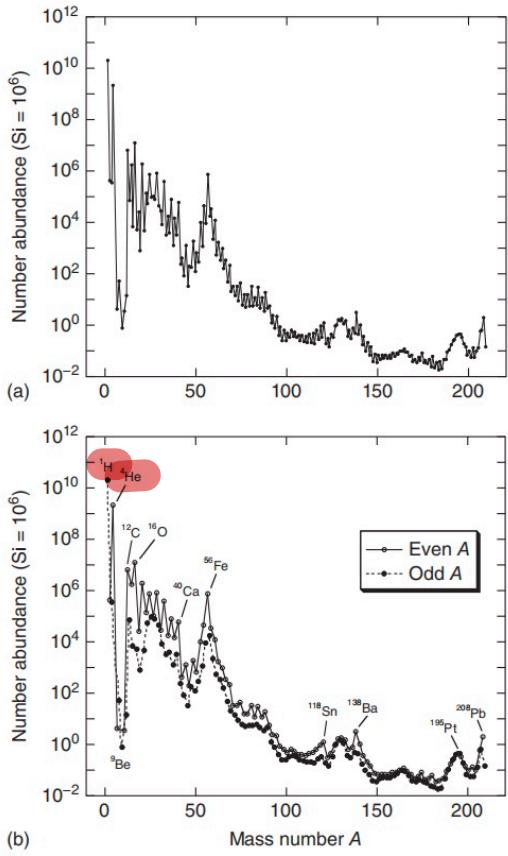


Figure 1.2 Abundances of the nuclides in the solar system at its birth. Number abundances are normalized to the number of silicon atoms ( $\text{Si} = 10^6$ ). Data from Lodders (2003). (a) Sum of all nuclide abundances at a given value of  $A$  versus mass number. The

maximum in the  $A = 50-65$  region is referred to as the *iron peak*. (b) Separate abundance contributions from nuclides with an even or an odd value of  $A$  versus mass number. Even- $A$  nuclides are in general more abundant than odd- $A$  nuclides.

→ normalised ( $\text{Si} = 10^6$ )

→  ${}^1\text{H}$  (71.1%)  
 ${}^4\text{He}$  (27.4%)  
 ${}^{12}\text{C} + {}^{16}\text{O}$  (1.5%)

→ first two minimum:

$A = 5-11$  : Li, Be, B

$A = 41-49$  : Sc

→ iron peak:  $A = 50-65$ .

→ before iron, fusion.  $\text{H}_2/\text{D}_2$   
 → after iron, fission.  $\text{U}/\text{Pu}$

→ generally decreasing abundances as  $A$  increases.

→ significance: Since almost all nuclides are formed within the star, therefore analysing the Solar system abundance offers powerful clues for stellar history and evolution.

### 1.5 Nuclear mass and Binding energy

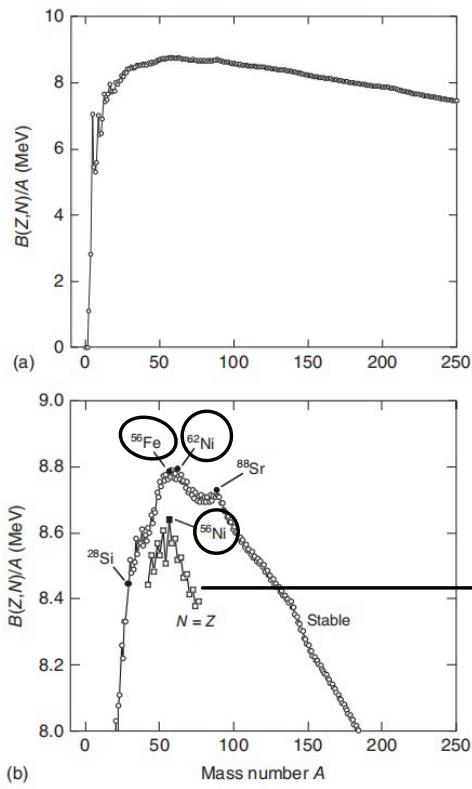
$$m_{\text{nuc}} = Z m_p + N m_n - \Delta m$$

↓  
 'nuclear mass'

mass defect, ( $\Delta E = \Delta m c^2$ )  
 Binding energy

(It's the energy released in assembling a given nucleus from its constituent nucleons)  
 or say it in the reverse way.

$$\Delta E = B(Z, N) = (Z m_p + N m_n - \Delta m) c^2$$



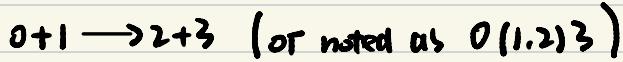
**Figure 1.9** Experimental binding energy per nucleon,  $B(Z,N)/A$ , versus mass number,  $A$ . (a) Value of  $B(Z,N)/A$  for the most tightly bound nuclide at a given mass number  $A$ . (b) Expanded section; the round symbols have the same meaning as in part (a), that is, they correspond to the most tightly bound nuclide of a given mass number;

the nuclide with the largest binding energy per nucleon is  $^{62}\text{Ni}$  ( $B/A = 8.795$  MeV); the square symbols show  $B(Z,N)/A$  values for  $N = Z$  nuclides above  $A = 40$ , which are all radioactive; the  $N = Z$  species with the largest binding energy per nucleon is  $^{56}\text{Ni}$  ( $B/A = 8.643$  MeV). Data from Wang et al. (2012).

→ squares are all radioactive.

## Energetics of Nuclear reactions

→ Commonly, we will encounter reactions:



If  $0, 1$  are identical to  $2, 3$ , then it is elastic/inelastic scattering.

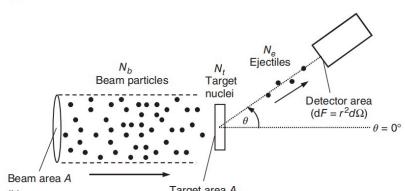
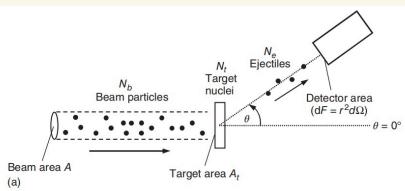
If light is involved in the reactants, then it is called photodisintegration reaction

light is involved in the reactants/products, then it is called reactive capture reaction

## 2.1 Cross Sections

→ Def: It's a quantitative measure of the probability that an interaction will occur.

(effective interaction area of a incident and target particles)



**Figure 2.1** Typical nuclear physics counting experiment, showing a beam of  $N_b$  particles per unit time,  $N_t$  nonoverlapping target nuclei within the beam area  $A$ ,  $N_e$  interaction products and a detector of area  $dF$ .

Assumptions here:  $N_e = N_t \underline{r}$

→ reaction number.

→ say if emitted product = incident product → elastic scattering  
if emitted product = incident product → reaction.

$$[G] = m^2$$

$G = \frac{\text{number of interaction per time}}{\text{number of incident particle per area per time}}$

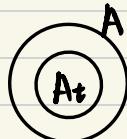
(number of incident particle per area per time) (number of target nuclei within the beam)

$$= \frac{N_p/t}{N_b/(tA) \cdot N_t}$$

① If  $A > A_t$

$$N_p/t = \frac{N_b N_t}{t A_t} G$$

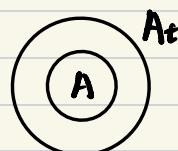
$$G = \frac{N_p A_t}{N_b N_t}$$



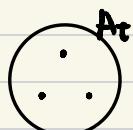
② If  $A < A_t$

$$N_p/t = \frac{N_b N_t}{t A} G$$

$$G = \frac{N_p A}{N_b N_t}$$



Plus, for a homogeneous target.



$\rightarrow$  "Distributed evenly"  $\frac{N_t}{A_t}$  is easy to obtain.

Then, for the emitted particles (assumed we used  $A_t A_t$  case)

$$\frac{N_p}{t} = \frac{N_e}{t}$$

$$= \frac{N_b N_t}{t A} G$$

$$\text{Then, } \frac{N_p(d\lambda)}{t} = \frac{N_b N_t}{t A} dG$$

$$= \frac{N_b N_t}{t A} \left( \frac{dG}{d\lambda} \right) d\lambda$$

$$\text{So, } \frac{dG}{d\lambda} = \frac{N_p(d\lambda) A}{N_b N_t d\lambda}$$

Also define  $N_{et} = \frac{N_e}{N_t}$ .

$$\text{then we can obtain: } \frac{dG}{d\lambda} = \frac{N_{et}(d\lambda) A}{N_b d\lambda}$$

$$= \frac{N_{et}(d\lambda)/t}{(N_b/t)(1/A)} \frac{1}{d\lambda}$$

Plus we know current density.

$$j_b = \frac{N_b}{t A}$$

$$j_{et} = \frac{N_{et}(d\lambda)}{t d\lambda} = \frac{N_{et}(d\lambda)}{dF} \quad (dF = d\lambda r^2, d\lambda = \frac{dF}{r^2})$$

One can finally gets:

$$G = \frac{N_{et}/t}{j_b}$$

$$\frac{dG}{d\lambda} = \frac{j_{et} r^2}{j_b}$$

$$G = \int \frac{dG}{d\lambda} d\lambda$$

### 3.1 Cross section and reaction rate.

→ reaction rate:

$$\begin{aligned} \text{reaction rate} &= \frac{N_p}{tV} \\ &= \frac{6NbN_t}{A+tV} \\ &= 6 \frac{N_t}{V} \frac{Nb}{A \cdot t} \end{aligned}$$

$$\text{reaction rate} = 6\gamma \left( \frac{N_t}{V} \right) \left( \frac{Nb}{V} \right)$$

where  $\gamma$  is the relative velocity of interacting particles.

→ particle-induced reactions.



where neither 4 particles is photon. (all particles have rest mass)

$$\Gamma_{01} = 6\gamma v N_0 N_1$$

where  $N_0 \equiv \frac{N_t}{V}$ ,  $N_1 \equiv \frac{Nb}{V}$ , named as number densities of interacting particles.

For thermodynamical equilibrium,  $v$  is not constant.

But we have  $P(v)dv$  in the small range.

with  $\int_0^\infty P(v)dv = 1$  to normalize.

→ non-identical

$$S_0, \langle \Gamma_{01} \rangle = N_0 N_1 \int_0^\infty v P(v) 6\gamma v dv = N_0 N_1 \langle 6\gamma v \rangle_{01}$$

(where  $N_0 N_1$  are non-identical number density, and  $N_0 N_1$  are total number of pairs)

→ identical

$$\langle \Gamma_{01} \rangle = \frac{N_0(N_0 - 1)}{2} \xrightarrow[N_0 \text{ is large}]{} \frac{N_0^2}{2}$$

In sum up: we obtain

$$\langle \Gamma_{01} \rangle = \frac{N_0 N_1 \langle 6\gamma v \rangle_{01}}{1 + S_0}$$

Now we made two assumptions:

{ ① the velocity distribution follows the Maxwell-Boltzmann distribution.

② we used the center-of-mass frame.

$$\text{So we have: } P(v)dv = \left( \frac{m_0}{2\pi kT} \right)^{\frac{3}{2}} e^{-m_0 v^2 / 2kT} 4\pi v^2 dv$$

(where  $m_{01} = \frac{m_0 m_1}{m_0 + m_1}$ )

Then we wanna express it in terms of energy.

we have:  $E = \frac{m_0 v^2}{2}$

$$\frac{dE}{dv} = m_0 v \quad (\text{we want to energy distribution})$$

$$P(v) dv = P(E) dE$$

$$= \left( \frac{m_0}{2\pi kT} \right)^{\frac{3}{2}} e^{-\frac{E}{kT}} 4\pi \frac{2E}{m_0} \frac{dE}{m_0} \sqrt{\frac{m_0}{2E}}$$

$$P(E) dE = \frac{2}{\sqrt{\pi}} \frac{1}{(kT)^{3/2}} E e^{-\frac{E}{kT}} dE$$

$$\text{Then } \langle Ev \rangle_{\text{av}} = \int_0^\infty v P(E) dE$$

$$= \left( \frac{8}{\pi m_0} \right)^{\frac{1}{2}} \frac{1}{(kT)^{3/2}} \int E g(E) e^{-\frac{E}{kT}} dE$$

Substitute this back to  $\Gamma_0$ , will result in the rate of reaction.

→ other types of reactions, see details back in textbook...