Solutions to problem Set - 2

Introduction to quantum computing using QSim

Problem 1: Multi-qubit transformations

The following questions are based on the properties of the Kronecker (Tensor) product and transformations on multi-qubit systems.

a) If A and B are two single qubit gates, prove the following relations: [I is the 2×2 identity matrix and \circ denotes matrix multiplication. All other symbols carry their usual meanings]

(This problem uses concepts discussed in slides 9-11, Week 2 Session 1-2.)

i.
$$(A \otimes I) \circ (I \otimes B) = (A \otimes B)$$

Solution: The aim of this problem is to prove the distributive property of the tensor product. Let the action of the A and B on the single-qubit standard basis given by the following relations:

$$\begin{array}{ll} A \left| 0 \right\rangle = \left| a_0 \right\rangle & B \left| 0 \right\rangle = \left| b_0 \right\rangle \\ A \left| 1 \right\rangle = \left| a_1 \right\rangle & B \left| 1 \right\rangle = \left| b_1 \right\rangle \end{array}$$

The proof can be proven by showing that the action of the operators $(A \otimes I) \circ (I \otimes B)$ and $(A \otimes B)$ on the two-qubit basis is the same:

Let us first consider the action of the two operators on the basis element $|00\rangle$

$$(A \otimes I) \circ (I \otimes B) |00\rangle = (A \otimes I) \circ (I |0\rangle \otimes B |0\rangle)$$

$$= (A \otimes I)(|0\rangle \otimes |b_0\rangle)$$

$$= A |0\rangle \otimes I |b_0\rangle$$

$$\therefore (A \otimes I) \circ (I \otimes B) |00\rangle = |a_0\rangle \otimes |b_0\rangle$$

and

$$(A \otimes B) |00\rangle = (A |0\rangle \otimes B |0\rangle)$$
$$= |a_0\rangle \otimes |b_0\rangle$$

Similarly, the equality can be proven for all other basis elements and this completes the proof.

ii.
$$(A \circ B)^{-1} = B^{\dagger} \circ A^{\dagger}$$

Solution: The proof of this question is based on the fact that A and B are single qubit gates and are therefore unitary transformations.

Therefore, from the definition of matrix inverse:

$$(A \circ B)^{-1}(A \circ B) = I$$

Now, considering the product of the operators:

$$(B^{\dagger} \circ A^{\dagger})(A \circ B) = (B^{\dagger} \circ (A^{\dagger}A) \circ B)$$
$$= (B^{\dagger} \circ B) :: A \text{ is unitary}$$
$$= I :: B \text{ is unitary}$$

Therefore, $B^{\dagger} \circ A^{\dagger}$ is the inverse of $A \circ B$ and the proof is complete.

b) Let U be a single qubit gate with actions on the standard basis given as follows:

$$U |0\rangle = |a\rangle$$
$$U |1\rangle = |b\rangle$$

Given V be a single qubit gate related to U through the following relation: $V^{\dagger}XV = U$ Find the action of the following circuit (i.e. find $|\psi\rangle$), on the two-qubit computational basis $\{|00\rangle, |01\rangle, |10\rangle, |11\rangle\}$.

(There was an error in the circuit diagram for this question, the correct figure is given below:.)

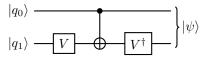


Figure 1: Circuit diagram for Problem 1 b).

Solution: Since the circuit uses a CX gate, it may be worthwhile examining the effect of the circuit in the cases, $|q_0\rangle = |0\rangle$ and $|q_0\rangle = |1\rangle$.

The matrix corresponding to the given circuit is given by:

$$M = (I \otimes V^{\dagger}) \circ CX \circ (I \otimes V)$$

(The method for this has been discussed in slides 11-13, Week 2 Session 2-3.)

When the qubit q_0 is in $|0\rangle$, the action of the above circuit is given by:

$$\begin{split} \left(I\otimes V^{\dagger}\right)\circ CX\circ \left(I\otimes V\right)\,|0\rangle\otimes |q_{1}\rangle &=\left(I\otimes V^{\dagger}\right)\circ CX\,\,|0\rangle\otimes V\,|q_{1}\rangle\\ &=\left(I\otimes V^{\dagger}\right)\,|0\rangle\otimes V\,|q_{1}\rangle\,\,, \text{ the control qubit is set to }|0\rangle\\ &=|0\rangle\otimes V^{\dagger}V\,|q_{1}\rangle\\ &=|0\rangle\otimes |q_{1}\rangle\,\,,\,V \text{ is unitary.} \end{split}$$

When the qubit q_0 is in $|1\rangle$, the X gate is applied to second qubit. Therefore, the action of the above circuit is given by:

$$\begin{split} \left(I\otimes V^{\dagger}\right)\circ CX\circ \left(I\otimes V\right)\,|1\rangle\otimes |q_{1}\rangle &=\left(I\otimes V^{\dagger}\right)\circ CX\,\,|1\rangle\otimes V\,|q_{1}\rangle\\ &=\left(I\otimes V^{\dagger}\right)\,|1\rangle\otimes XV\,|q_{1}\rangle\\ &=|1\rangle\otimes V^{\dagger}XV\,|q_{1}\rangle\\ &=|1\rangle\otimes U\,|q_{1}\rangle\,\,, \text{ from the definition of }V. \end{split}$$

Therefore, the action of the circuit on the two-qubit standard basis is given by;

$$\left\{\left|00\right\rangle,\left|01\right\rangle,\left|10\right\rangle,\left|11\right\rangle\right\} \xrightarrow{M} \left\{\left|00\right\rangle,\left|01\right\rangle,\left|1\right\rangle \otimes \left|a\right\rangle,\left|1\right\rangle \otimes \left|b\right\rangle\right\}$$

The above circuit is therefore a representation of the controlled U gate.

c) Show that there exists no pair of single qubit gates, A, B that satisfy the following relation:

$$A \otimes B = \text{CNOT}_1^0$$

Solution: Assuming that it is possible to write $CNOT_1^0$ in the form $A \otimes B$. Where A and B are single qubit gates given by the matrices:

$$A = \begin{bmatrix} A_{00} & A_{01} \\ A_{10} & A_{11} \end{bmatrix} \qquad B = \begin{bmatrix} B_{00} & B_{01} \\ B_{10} & B_{11} \end{bmatrix}$$

Therefore the expansion of $CNOT_1^0$ may be expressed as:

$$A \otimes B = \begin{bmatrix} A_{00} \begin{bmatrix} B_{00} & B_{01} \\ B_{10} & B_{11} \\ A_{10} \begin{bmatrix} B_{00} & B_{01} \\ B_{10} & B_{11} \end{bmatrix} & A_{01} \begin{bmatrix} B_{00} & B_{01} \\ B_{10} & B_{11} \\ B_{00} & B_{01} \\ B_{10} & B_{11} \end{bmatrix} \end{bmatrix}$$
$$= \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

This further implies that the following relations are true:

$$A_{00} \begin{bmatrix} B_{00} & B_{01} \\ B_{10} & B_{11} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \qquad A_{11} \begin{bmatrix} B_{00} & B_{01} \\ B_{10} & B_{11} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\Rightarrow A_{00}B_{01} = A_{00}B_{10} = 0 \qquad \Rightarrow A_{11}B_{01} = A_{00}B_{10} = 1$$

$$\Rightarrow A_{00}B_{00} = A_{00}B_{11} = 1 \qquad \Rightarrow A_{11}B_{00} = A_{00}B_{11} = 0$$

$$\therefore B_{01} = B_{10} = 0 \qquad \therefore B_{01} = B_{10} \neq 0$$
and $B_{00} = B_{11} \neq 0$
and $B_{00} = B_{11} = 0$

The above results are a contradiction, therefore it is impossible to express CNOT_1^0 as $A \otimes B$.

d) If a two-qubit gate has the following action on the two-qubit standard basis:

$$U |00\rangle = |v_0\rangle$$

$$U |01\rangle = |v_1\rangle$$

$$U |10\rangle = |v_2\rangle$$

$$U |11\rangle = |v_3\rangle$$

i. Prove the following relation:

$$|v_0\rangle\langle 00| + |v_1\rangle\langle 01| + |v_2\rangle\langle 10| + |v_3\rangle\langle 11| = U$$

Solution: This relation can be verified by multiplying the left hand side of the above equation with elements of the two-qubit standard basis. As an example, the product with $|00\rangle$ is shown.

$$(|v_0\rangle\langle 00| + |v_1\rangle\langle 01| + |v_2\rangle\langle 10| + |v_3\rangle\langle 11|) |00\rangle = |v_0\rangle \langle 00|00\rangle + |v_1\rangle \langle 01|00\rangle + |v_2\rangle \langle 10|00\rangle + |v_3\rangle \langle 11|00\rangle = |v_0\rangle$$

The other terms vanish because of the orthonormality of the two-qubit standard basis. working out the expression for the other elements of the basis will complete the proof.

ii. Using the above relation, find the matrix form of $CNOT_0^1$.

Solution: The action of $CNOT_0^1$ on the two-qubit standard basis is given as:

$$\begin{aligned} & \text{CNOT}_0^1 |00\rangle = |00\rangle \\ & \text{CNOT}_0^1 |01\rangle = |11\rangle \\ & \text{CNOT}_0^1 |10\rangle = |10\rangle \\ & \text{CNOT}_0^1 |11\rangle = |01\rangle \end{aligned}$$

The operator form of $CNOT_0^1$ is given by:

$$CNOT_0^1 = |00\rangle\langle 00| + |11\rangle\langle 01| + |10\rangle\langle 10| + |01\rangle\langle 11|$$

if one were to try to expand this to the matrix form, we get the following result.

$$CNOT_0^1 = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix}$$

Problem 2: Boolean functions and quantum gates

The following questions are based on quantum implementations of Boolean functions.

- a) Find the quantum circuits (bit oracles) for the following Boolean functions: $[\oplus$ XOR, \vee OR, $^-$ NOT]
 - i. $F(x_0, x_1) = x_0 \oplus x_1$

Solution: This is a two variable function. Therefore the bit oracle requires a 3-qubit operation and is represented by the following circuit:

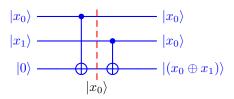


Figure 2: Circuit diagram for Problem 2 a) i.

(This is also explained in the python notebook titled, "2. Boolean function oracles.ipynb".)

ii.
$$F(x_0, x_1) = x_0 \vee x_1$$

Solution: This is a two variable function. Therefore the bit oracle requires a 3-qubit operation. It will also be of benefit to represent the OR gate in terms of the AND gate.

$$x_0 \vee x_1 = \overline{\overline{x}_0 \wedge \overline{x}_1}$$

This is due to the De Morgan's law.

The bit oracle of the AND gate can be represented by a single Toffoli (CCX) gate. Additionally, the negation of the bit value stored in any qubit can be performed using the Pauli X gate.

Therefore, the circuit for this function is as shown below:

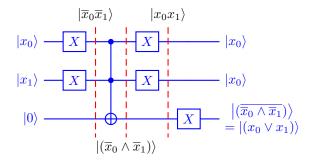


Figure 3: Circuit diagram for Problem 2 a) ii.

iii. $F(x) = \overline{x}$

Solution: It should be remembered that what is required for this question is a bit oracle, which will be a two-qubit circuit. This circuit will be given by:

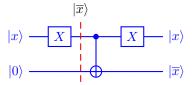


Figure 4: Circuit diagram for Problem 2 a) iii.

b) Prove the following statements for a *n*-bit Boolean function bit oracle: ['+' denotes the XOR operation]

$$U_F |\mathbf{x}\rangle \qquad |y\rangle \rightarrow |\mathbf{x}\rangle |F(\mathbf{x}) + y\rangle$$
_{n-qubit} _{1-qubit}

i. The bit oracle causes entanglement between the input qubits, $|\mathbf{x}\rangle$ and the target qubit, $|y\rangle$.

Solution: To prove this statement, we can consider two distinct input values \mathbf{x}_1 and \mathbf{x}_2 for the Boolean function such that:

$$F(\mathbf{x}_1) \neq F(\mathbf{x}_2)$$

The action of the the bit-oracle U_F on the state, $|\mathbf{x}_1\rangle + |\mathbf{x}_2\rangle$ is given by:

$$U_F(|\mathbf{x}_1\rangle + |\mathbf{x}_2\rangle)|0\rangle = |\mathbf{x}_1\rangle|F(\mathbf{x}_1)\rangle + |\mathbf{x}_2\rangle|F(\mathbf{x}_2)\rangle$$

In the above expression, normalization has been temporarily disregarded. Now since it is known that $\mathbf{x}_1 \neq \mathbf{x}_2$ and $F(\mathbf{x}_1) \neq F(\mathbf{x}_2)$. It is impossible to separate the target qubit from the *n*-qubit input states. This implies that there is entanglement between the input qubits and the target qubit. The following statement is a property of all Boolean function oracles.

"As long as the Boolean function $F(\mathbf{x})$ is **not constant** the quantum bit oracle U_F causes entanglement between the input and the output qubits."

(This was also discussed in slides 11 - 12, Week 3 Session 2-3.)

ii.
$$U_F^{\dagger} = U_F$$

Solution: The proof of this statement involves applying the bit oracle of the Boolean function twice. This is done as shown below:

$$|\mathbf{x}\rangle |y\rangle \xrightarrow{U_F} |\mathbf{x}\rangle |y + F(\mathbf{x})\rangle \xrightarrow{U_F} |\mathbf{x}\rangle |y + F(\mathbf{x}) + F(\mathbf{x})\rangle$$

Using the property of the XOR gate, y + y = 0, the resultant state of the above equation can be written as:

$$|\mathbf{x}\rangle |y + F(\mathbf{x}) + F(\mathbf{x})\rangle = |\mathbf{x}\rangle |y + 0\rangle = |\mathbf{x}\rangle |y\rangle$$

This implies that applying the Boolean function bit oracle, U_F twice is the same as applying the identity gate (or no gate). This implies that:

$$U_F^{-1} = U_F$$

and since, U_F is also unitary (as it is a quantum gate), $U_F^{-1} = U_F^{\dagger}$, this when used with the previous result proves the statement.

$$U_F^{\dagger} = U_F$$