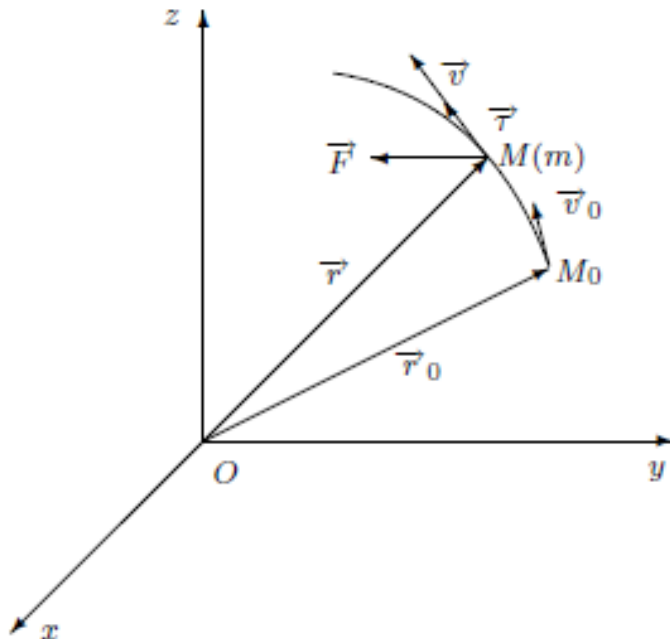


Theoretical Mechanics

9. Dynamics of the particle. General theorems.

The differential equations of the motion

Consider the material point $M(m)$ moving in the Cartesian frame $Oxyz$ under the action of the force $\vec{F} = \vec{F}(\vec{r}, \vec{v}, t)$ where $\vec{F}(X, Y, Z)$.



The differential equation of the motion is

$$m \frac{d^2 \vec{r}}{dt^2} = \vec{F}(\vec{r}, \vec{v}, t), \quad t \in (t_0, T] \quad (9.1)$$

along with the initial conditions

$$\vec{r}(t_0) = \vec{r}_0, \quad \vec{v}(t_0) = \vec{v}_0, \quad (9.2)$$

Theoretical Mechanics

In the Cartesian frame $Ox\ yz$, Eqs. (9.1) and (9.2) are

$$m \frac{d^2 x}{dt^2} = X, \quad m \frac{d^2 y}{dt^2} = Y, \quad m \frac{d^2 z}{dt^2} = Z, \quad (9.3)$$

$$\left| \begin{array}{l} x(t_0) = x_0, \quad y(t_0) = y_0, \quad z(t_0) = z_0 \\ \dot{x}(t_0) = \dot{x}_0, \quad \dot{y}(t_0) = \dot{y}_0, \quad \dot{z}(t_0) = \dot{z}_0, \end{array} \right. \quad (9.4)$$

where

$$\vec{r}_0 = (x_0, y_0, z_0), \quad \vec{v}_0 = (\dot{x}_0, \dot{y}_0, \dot{z}_0)$$

A first integral of the differential system of motion (9-3) – (9.4) is a (non-constant) continuously-differentiable function (class C^1)

$$\boxed{\mathcal{F}(x, y, z, \dot{x}, \dot{y}, \dot{z}, t) = c \in \mathbb{R}, \quad \forall t \geq t_0} \quad (9.5)$$

which reduces to a constant when the functions

$$x = x(t), \quad y = y(t), \quad z = z(t) \quad (9.6)$$

satisfy (9.3).

Theoretical Mechanics

If we determine $k \leq 6$ first integrals, for which

$$f_j(x_1, x_2, x_3, v_1, v_2, v_3; t) = C_j, \quad C_j = \text{const}, \quad j = 1, 2, \dots, k,$$

the matrix

$$\mathbf{M} \equiv \left[\frac{\partial (f_1, f_2, \dots, f_k)}{\partial (x_1, x_2, x_3, v_1, v_2, v_3)} \right]$$

being of rank k , then all the first integrals are functionally independent (for the sake of simplicity, further we say *independent first integrals*) :

Thus, the problem is reduced to the integration of a system of equations (9.3) with only $6 - k$ unknowns (hence, a smaller number of unknowns). If $k = 6$, then all the first integrals are independent, so that the system (9.5) of first integrals determines all the unknown functions. We notice that for $k > 6$ the first integrals (9.5) are no more independent; we may thus set up at the most six independent first integrals.

P.P. Teodorescu, *Mechanical Systems, Classical Models. Vol. I: Particle Mechanics*, Springer, 2007

Theoretical Mechanics

Momentum of a particle

Newton has introduced the notion of *momentum* (which he called *quantity of motion*) representing the product of the mass by the velocity of a point.

$$\vec{H} := m \vec{v} \quad (9.7)$$

- momentum is a vector parallel to the velocity vector

From (9.1) we have

$$\frac{d}{dt}(m \vec{v}) = \vec{F} \Rightarrow \frac{d\vec{H}}{dt} = \vec{F}, \quad \frac{d}{dt}(m\dot{x}) = X, \dots \quad \vec{H}(t_0) = \vec{H}_0 \quad (9.8)$$

Equation (9.8) we have is called the momentum equation. From (9.8) we can formulate the theorem of momentum:

Theorem (*theorem of momentum*). The derivative with respect to time of the momentum of a free particle is equal to the resultant of the given forces which act upon it.

Another usual notation for momentum is: $\vec{P} = m\vec{v}$

Average Force, Momentum, and Impulse

Suppose you are pushing a cart with a force that is non-uniform, but has an average value $\vec{\mathbf{F}}_{\text{ave}}$ during the time interval Δt . We can find the average acceleration according to Newton's Second Law,

$$\vec{\mathbf{F}}_{\text{ave}} = m \vec{\mathbf{a}}_{\text{ave}} \quad \vec{\mathbf{a}}_{\text{ave}} = \frac{\Delta \vec{\mathbf{v}}}{\Delta t}.$$

Therefore Newton's Second Law can be recast as

$$\vec{\mathbf{F}}_{\text{ave}} = m \vec{\mathbf{a}}_{\text{ave}} = \frac{m \Delta \vec{\mathbf{v}}}{\Delta t}.$$

The change in momentum is the product of the mass and the change in velocity,

$$\Delta \vec{\mathbf{p}} = m \Delta \vec{\mathbf{v}}.$$

Theoretical Mechanics

Newton's Second Law can be restated as follows: the product of the average force acting on an object and the time interval over which the force acts will produce a change in momentum of the object,

$$\vec{\mathbf{F}}_{\text{ave}} \Delta t = \Delta \vec{\mathbf{p}}.$$

This change in momentum is called the *impulse*,

$$\vec{\mathbf{I}} = \vec{\mathbf{F}}_{\text{ave}} \Delta t = \Delta \vec{\mathbf{p}}.$$

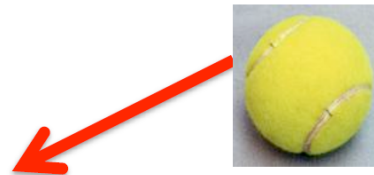
Force is a vector quantity; impulse is obtained by multiplying a vector by a scalar, and so impulse is also a vector quantity. The SI units for impulse are $[\text{N} \cdot \text{s}] = [\text{kg} \cdot \text{m} \cdot \text{s}^{-1}]$, which are the same units as momentum.

<https://sites.ualberta.ca/~ygu/courses/phys144/notes/momentum1.pdf>

Example

Impulse and momentum:

A tennis ball (mass=0.1 kg) comes in from the right at 40 m/s at an angle 10 deg below horizontal. Then on a bad shot Rafa hits it at 30 deg above the horizontal at a speed of 60 m/s. What is the impulse of the net force and the average net force, assuming the contact (collision) duration is 0.01 sec.



Solution: Take positive X to the Right and positive Y as Up. The velocity components before (subscript 1) and after (2) the tennis ball is hit are:

$$v_{1x} = -40 \cos(10) = -39.39 \text{ m/s} \quad v_{1y} = -40 \sin(10) = -6.95 \text{ m/s}$$

$$v_{2x} = 60 \cos(30) = 51.96 \text{ m/s} \quad v_{2y} = 60 \sin(30) = 30 \text{ m/s}$$

Theoretical Mechanics

X-component of impulse is equal to x-component of momentum change

$$\begin{aligned} J_x &= p_{2x} - p_{1x} = m(v_{2x} - v_{1x}) \\ &= 0.1 \times (51.96 - (-39.39)) = 9.14 \text{ kg} \cdot \text{m/s} \end{aligned}$$

$$\begin{aligned} J_y &= p_{2y} - p_{1y} = m(v_{2y} - v_{1y}) \\ &= 0.1 \times (30 - (-6.95)) = 3.70 \text{ kg} \cdot \text{m/s} \end{aligned}$$

By vector addition

$$J = \sqrt{J_x^2 + J_y^2} = \sqrt{9.14^2 + 3.70^2} = 9.86 \text{ kg} \cdot \text{m/s}$$

Average net force: $\bar{F} = \frac{J}{\Delta t} = \frac{9.86}{0.01} = 986 \text{ N}$

Alternatively, we could use components of J to find components of average net force. The components of average net force are:

$$\bar{F}_x = \frac{J_x}{\Delta t} = \frac{9.14}{0.01} = 914 \text{ N} \quad \bar{F}_y = \frac{J_y}{\Delta t} = \frac{3.70}{0.01} = 370 \text{ N}$$

Theoretical Mechanics

The average net force is:

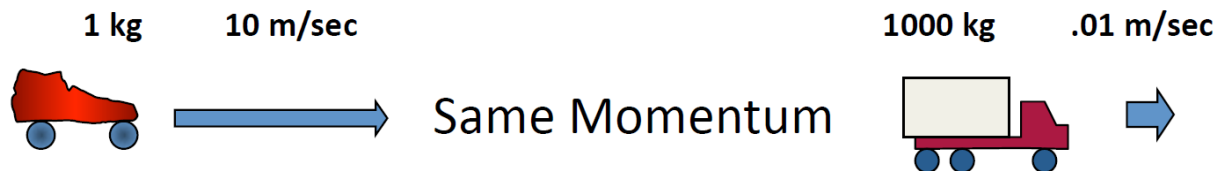
$$\bar{F} = \sqrt{\bar{F}_x^2 + \bar{F}_y^2} = \sqrt{914^2 + 370^2} = \sqrt{835,396 + 136,900} = 986.05 \text{ N}$$

Average Force Direction:

$$\theta = \tan^{-1}\left(\frac{\bar{F}_y}{\bar{F}_x}\right) = 22.04^\circ$$

Remarks: (1) Keep signs straight due to vector use (2) Because the ball was Not at rest initially, average force direction is NOT the same as ball direction (analogy: projectile motion)

Example



First integrals

- If $\vec{F} = 0$ then $\vec{H} = \overrightarrow{\text{const}}$, $\forall t \geq t_0$. This first integral is the **law of momentum conservation**.

$$\mathbf{H} = m\mathbf{v} = \mathbf{C}, \quad \mathbf{C} = \overrightarrow{\text{const}}, \quad H_i = C_i, \quad i = 1, 2, 3;$$

Theorem (*conservation theorem of momentum*). The momentum (and the velocity) of a free particle is conserved in time if and only if the resultant of the given forces which act upon it vanishes.

We notice that the relation $m\mathbf{v} = m\dot{\mathbf{r}} = \mathbf{C}$ leads to

$$m\mathbf{r} = \mathbf{C}t + \mathbf{C}', \quad \mathbf{C}, \mathbf{C}' = \overrightarrow{\text{const}}, \quad mx_i = C_i t + C'_i, \quad i = 1, 2, 3;$$

Taking into account the initial conditions we obtain:

$$\begin{aligned} \vec{v} &= \vec{v}_0, \\ \vec{r}(t) &= \vec{v}_0(t - t_0) + \vec{r}_0. \end{aligned}$$

Theoretical Mechanics

The motion of the particle P is thus rectilinear and uniform. Besides, **this result corresponds to the principle of inertia**, which appears thus as a particular case of the principle of action of forces (lex secunda).

- If $\vec{F} \neq 0$, $\forall t \geq t_0$ and exist a fixed direction $\vec{u} = (\alpha, \beta, \gamma)$ such that $\vec{F} \cdot \vec{u} = 0$, then using this in (9.1) we obtain

$$\frac{d}{dt}(\vec{H} \cdot \vec{u}) = \vec{F} \cdot \vec{u} = 0 \Rightarrow \vec{H} \cdot \vec{u} = c \in \mathbb{R}.$$

Thus,

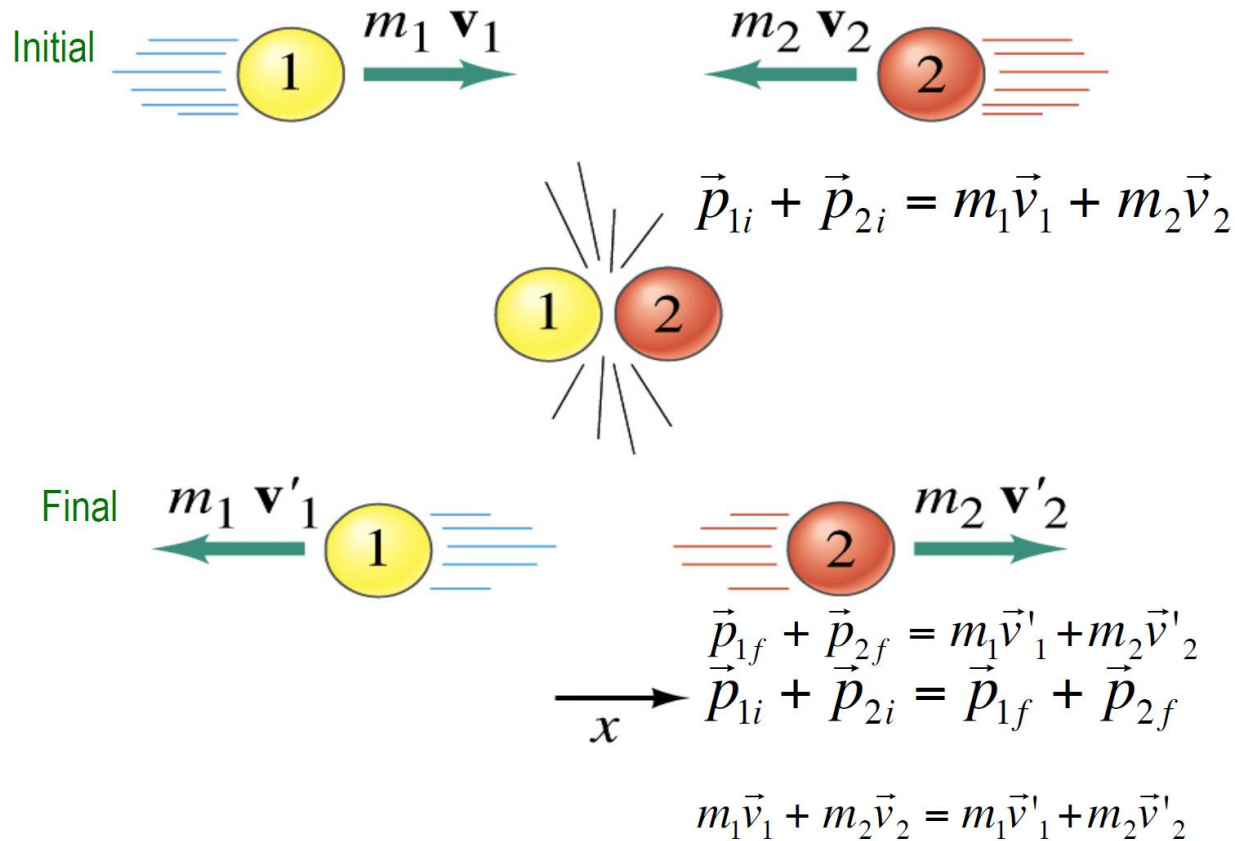
$$\alpha \dot{x} + \beta \dot{y} + \gamma \dot{z} = c, \quad \forall t \geq t_0, \quad (9.9)$$

is a first integral of the differential equations of motion.

Hence, *if the force \mathbf{F} is parallel to a fixed plane, then the projection of the velocity of the free particle P on the normal to this plane is conserved (is constant) in time*

Example

Linear Momentum Conservation (1D)



Moment of momentum (angular momentum)

The moment of the momentum with respect to the pole O (origin of the coordinate system) is called *moment of momentum (angular momentum) of the particle*, with respect to this pole, and is given by

$$\vec{K}_0 := \vec{r} \times m \vec{v} = \vec{r} \times \vec{H} \quad (9.10)$$

$$K_x = m(y\dot{z} - z\dot{y}), \quad K_y = m(z\dot{x} - x\dot{z}), \quad K_z = m(x\dot{y} - y\dot{x}) \quad (9.11)$$

From the equation of motion (9.1) we have

$$m \frac{d\vec{v}}{dt} = \vec{F} \quad | \times \vec{r}$$

$$\vec{r} \times \frac{d}{dt}(m \vec{v}) = \vec{r} \times \vec{F} \Rightarrow \frac{d}{dt}(\vec{r} \times m \vec{v}) = \vec{r} \times \vec{F}$$

or

$$(9.12) \quad \frac{d\vec{K}_0}{dt} = \vec{M}_0(\vec{F})$$

Moment (torque) of force \vec{F} with respect to the pole O:
 $\vec{M}_O(\vec{F}) = \vec{r} \times \vec{F}$

Theorem (*theorem of angular momentum*) *The derivative with respect to time of the moment of momentum of a particle, with respect to a fixed pole, is equal to the moment of the force which act upon it, with respect to the same pole.*

On projection on the $Oxyz$ axes:

$$m \frac{d}{dt}(y\dot{z} - z\dot{y}) = yZ - zY, \dots \quad (9.13)$$

First integrals

- If $\vec{M}_O(\vec{F}) = 0$ we deduce the **law of the angular momentum conservation**.

$$\vec{K}_0 = \vec{C} \in \mathbb{R}^3, \quad \forall t \geq t_0 \quad (9.14)$$

Thus,

$$\begin{aligned} \vec{K}_0 &= \vec{r} \times m \vec{v} = \vec{c}, \quad \forall t \geq t_0 \\ \vec{r} \times \vec{v} &= \vec{c} = \vec{r}_0 \times \vec{v}_0 \end{aligned} \quad (9.15)$$

Equation (9.15) is equivalent with

$$\frac{d\vec{A}}{dt} = \overrightarrow{const} \quad (9.16)$$

where $\frac{d\vec{A}}{dt} = \frac{1}{2}(\vec{r} \times \vec{v})$ is the areal velocity of the particle.

From (9.15) the following first integrals are obtained

$$y\dot{z} - z\dot{y} = c_1, \dots \quad (9.17)$$

In the particular case when $\vec{F} \parallel \vec{r}$ (in this case \vec{F} is called central force) we have

$$\vec{r} \times \vec{F} = 0 \Rightarrow \vec{r} \times \vec{v} = \vec{c} = \vec{r}_0 \times \vec{v}_0.$$

Notice that:

- i) the motion is rectilinear if $\vec{r}_0 \parallel \vec{v}_0 \Rightarrow \vec{r}_0 \times \vec{v}_0 = 0 = \vec{r} \times \vec{v} \Rightarrow \vec{r} \parallel \vec{v}$,
- ii) the motion takes place in a plane if $\vec{r}_0 \nparallel \vec{v}_0 \Rightarrow \vec{r} \cdot (\vec{r}_0 \times \vec{v}_0) = 0$
 $\Rightarrow \vec{r} \in (\vec{v}_0, \vec{r}_0)$

Theoretical Mechanics

- If $\vec{M}_O(\vec{F}) \neq 0$, $\forall t \geq t_0$ and exist a fixed direction $\vec{u} = (\alpha, \beta, \gamma)$ such that $\vec{M}_O(\vec{F}) \cdot \vec{u} = 0$, then

$$\frac{d}{dt}(\vec{K}_0 \cdot \vec{u}) = 0 \Rightarrow \vec{K}_0 \cdot \vec{u} = c \in \mathbb{R}.$$

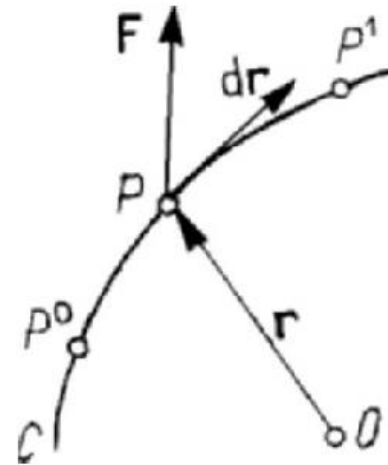
Thus, we obtain the first integral

$$(\vec{r} \times \vec{v}) \cdot \vec{u} = (\vec{r}_0 \times \vec{v}_0) \cdot \vec{u}. \quad (9.18)$$

Work

Consider the force \vec{F} acting on the point P ($\vec{r} = \overrightarrow{OP}$) which effects a real displacement $d\vec{r}$. The elementary work of the force \vec{F} is the scalar

$$\delta L = \vec{F} \cdot d\vec{r} = xdx + Ydy + Zdz \quad (9.19)$$



Theoretical Mechanics

Next, we suppose \vec{F} depends only on \vec{r} ($\vec{F} = \vec{F}(X, Y, Z)$) and exist a function $U = U(x, y, z)$ with the property:

$$X = \frac{\partial U}{\partial x}, \quad Y = \frac{\partial U}{\partial y}, \quad Z = \frac{\partial U}{\partial z} \quad (9.20)$$

or

$$\mathbf{F} = \text{grad } U = \nabla U = U_{,j} \mathbf{i}_j, \quad F_j = U_{,j}, \quad (9.21)$$

where $U = U(\mathbf{r}) = U(x, y, z)$ is the force function (potential function or potential). Function $V := -U$ is the potential energy and we say that \vec{F} is conservative.

Theorem. The force \vec{F} is potential (conservative) if and only if

$$\frac{\partial X}{\partial y} = \frac{\partial Y}{\partial x}, \quad \frac{\partial Y}{\partial z} = \frac{\partial Z}{\partial y}, \quad \frac{\partial Z}{\partial x} = \frac{\partial X}{\partial z} \quad (9.22)$$

This is equivalent with the fact that

$$\vec{F} \cdot d\vec{r} = Xdx + Ydy + Zdz \quad (9.23)$$

is an exact differential.

Theoretical Mechanics

A force field $\vec{F}: D \subset R^3 \rightarrow R^3$ is a conservative force field if $\exists U: D \rightarrow R$ with the property:

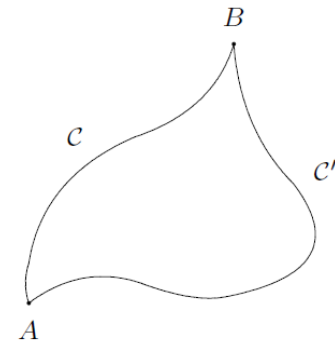
$$\vec{F} = \text{grad}U = -\text{grad}V$$

Theorem. Consider $D \subset R^3$ a simply connected domain.

A force field $\vec{F}: D \rightarrow R^3$ is potential (conservative) if and only if one of the following equivalent conditions takes place:

- $\text{rot } \vec{F} = \nabla \times \vec{F} = 0$
- $\int_{\Gamma} \underbrace{\vec{F} \cdot d\vec{r}}_{=Xdx+Ydy+Zdz} = 0, \quad \forall \gamma \text{ closed curve}$
- $\int_C \vec{F} \cdot d\vec{r} = \int_{C'} \vec{F} \cdot d\vec{r},$

$\forall C, C' \in D$ arbitrary curved from A to B .



Theoretical Mechanics

- the differential form

$$\vec{F} \cdot d\vec{r} = Xdx + Ydy + Zdz$$

is an exact differential, i.e.

$$\frac{\partial X}{\partial y} = \frac{\partial Y}{\partial x}, \quad \frac{\partial Y}{\partial z} = \frac{\partial Z}{\partial y}, \quad \frac{\partial Z}{\partial x} = \frac{\partial X}{\partial z}$$

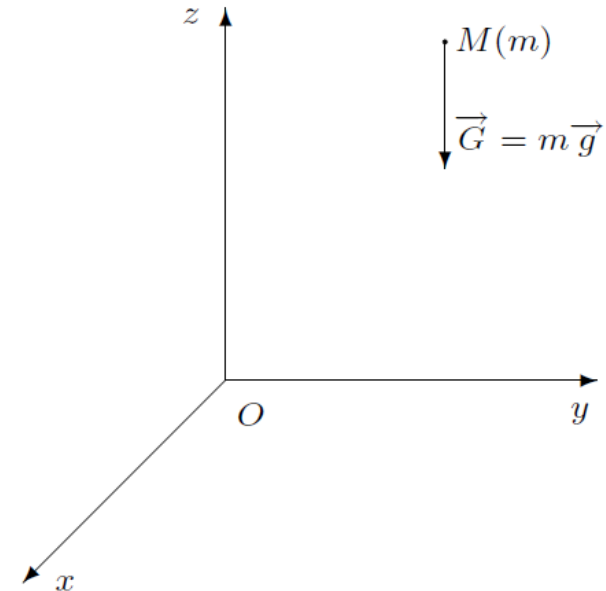
Example

$$\vec{F} = m\vec{g} = -\text{grad}V \Leftrightarrow \frac{\partial V}{\partial x} = 0, \quad \frac{\partial V}{\partial y} = 0, \quad \frac{\partial V}{\partial z} = mg$$

$$\Rightarrow dV = d(mgz)$$

$$\Rightarrow V = mgz + c, \quad U = -mgz$$

$\Rightarrow \vec{G} = m\vec{g}$ is a potential (conservative)
force



Force function in the plane

Consider $\vec{F} : D \rightarrow \mathbb{R}^2$, $\vec{F} = (X, Y)$, $X = X(x, y)$, $Y = Y(x, y)$

such that

$$\frac{\partial X}{\partial y} = \frac{\partial Y}{\partial x} \quad (9.24)$$

$$\text{Thus, } \exists U : D \rightarrow \mathbb{R} \text{ such that } \vec{F} = \text{grad}U \Leftrightarrow X = \frac{\partial U}{\partial x}, \quad Y = \frac{\partial U}{\partial y} \quad (9.25)$$

We integrate (9.25)₁ with respect to x

$$U(x, y) = \int_{x_0}^x X(s, y) ds + \varphi(y), \quad (9.26)$$

where $\varphi(y)$ is determined from (9.25)₂:

$$\int_{x_0}^x \frac{\partial U(s, y)}{\partial y} ds + \varphi'(y) = Y(x, y)$$

Taking into account (9.24) we obtain

$$\int_{x_0}^x \frac{\partial Y(s, y)}{\partial s} ds + \varphi'(y) = Y(x, y)$$

$$Y(x, y) - Y(x_0, y) + \varphi'(y) = Y(x, y)$$

$$\varphi'(y) = Y(x_0, y),$$

or

$$\varphi(y) = \int_{y_0}^y Y(x_0, u) du + \text{const.} \quad (9.27)$$

Using (9.26) and (9.27) we have

$$U(x, y) = \int_{x_0}^x X(s, y) ds + \int_{y_0}^y Y(x_0, u) du. \quad (9.28)$$

Theoretical Mechanics

Next, we suppose that the force field $\vec{F} = \vec{F}(\vec{r}) = \vec{F}(x, y, z)$ is conservative (potential), i.e. $\exists U: D \rightarrow \mathbb{R}$ with the property:

$$\vec{F} = \text{grad}U = -\text{grad}V$$

or

$$X = \frac{\partial U}{\partial x}, \quad Y = \frac{\partial U}{\partial y}, \quad Z = \frac{\partial U}{\partial z}$$

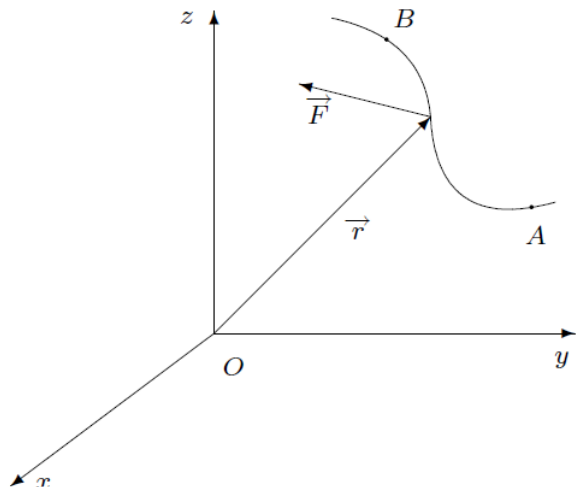
Thus, the elementary work is an exact differential

$$\delta L = Xdx + \dots = dU = -dV$$

Total work along an arc AB of a curve

Consider AB an arc of a curve defined by the equations:

$$x = x(q), \quad y = y(q), \quad z = z(q), \quad q \in [q_0, q_1]. \quad (9.29)$$



The curvilinear integral

$$\begin{aligned} L_{AB} = L &:= \int_{AB} \vec{F} \cdot d\vec{r} = \int_{AB} X dx + Y dy + Z dz = \int_{AB} (X(x(z), \dots) dx(q) \\ &= \int_{q_0}^{q_1} (X(x(q), y(q), z(q))x'(q) + \dots) dq \end{aligned} \quad (9.30)$$

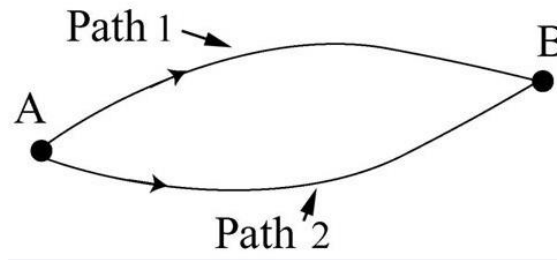
is the total work of the force \vec{F} along the arc AB.

Theoretical Mechanics

If the force is conservative, $\vec{F} = \text{grad}U$, then

$$L = \int_{AB} \vec{F} \cdot d\vec{r} = \int_{AB} dU = U(B) - U(A) \quad (9.31)$$

and the work is independent of the path from A to B.



If the force is not conservative we use:

$$L = \int_{AB} \vec{F} \cdot d\vec{r} = \int_{q_0}^{q_1} [X(x(q), y(q), z(q))x'(q) + \dots] dq \quad (9.32)$$

and if the force is a function of velocity and time we can use:

$$L = \int_{AB} X dx + Y dy + Z dz = \int_{t_0}^{t_1} X(x(t), \dots) \dot{x}(t) + \dots \quad (9.33)$$

Theoretical Mechanics

Kinetic energy

The scalar quantity

$$T = \frac{1}{2}mv^2 \quad (9.34)$$

is the *kinetic energy of the particle P* ; this quantity depends on the mass and the velocity of the particle.

From $m \frac{d\vec{v}}{dt} = \vec{F} \mid \cdot d\vec{r}$ we deduce

$$m \underbrace{\frac{d\vec{r}}{dt}}_{=\vec{v}} \cdot \underbrace{\frac{d\vec{v}}{dt}}_{=d\vec{v}} dt = \vec{F} \cdot d\vec{r} \Rightarrow m \vec{v} \cdot d\vec{v} = \vec{F} \cdot d\vec{r} \Rightarrow dT = \delta L \quad (9.35)$$

Equation (9.35) expresses the following theorem

Theorem (theorem of kinetic energy). *The differential of the kinetic energy of a free particle is equal to the elementary work of the resultant of the given forces which act upon it.*

Theoretical Mechanics

If the motion of the particle takes place in a conservative field of forces, $\vec{F} = \vec{F}(\vec{r})$, then exist $V = V(\vec{r})$ such that $\vec{F} = -\text{grad}V$.

Thus, $\delta L = -dV$ and Eq. (9.35) becomes

$$d(T + V) = 0 \Rightarrow T + V = h, \forall t \geq t_0. \quad (9.36)$$

The first integral (9.36) is the *energy integral*, V is the *potential energy*, while h is the *energy constant* and has to be calculated from the initial conditions.

$$E = T + V \quad \text{total (mechanical) energy}$$

Theorem (*mechanical energy conservation theorem*). *The mechanical energy of a free particle is conserved in time if and only if the resultant of the given forces which act upon it is conservative.*

<https://sites.ualberta.ca/~ygu/courses/phys144/notes/momentum2.pdf>

Ex. A Ballistic Pendulum

The mass of the block of wood is 2.50-kg and the mass of the bullet is 0.0100-kg. The block swings to a maximum height of 0.650 m above the initial position. Find the initial speed of the bullet.

What kind of collision? **Perfectly inelastic**

No net external force \rightarrow momentum conserved

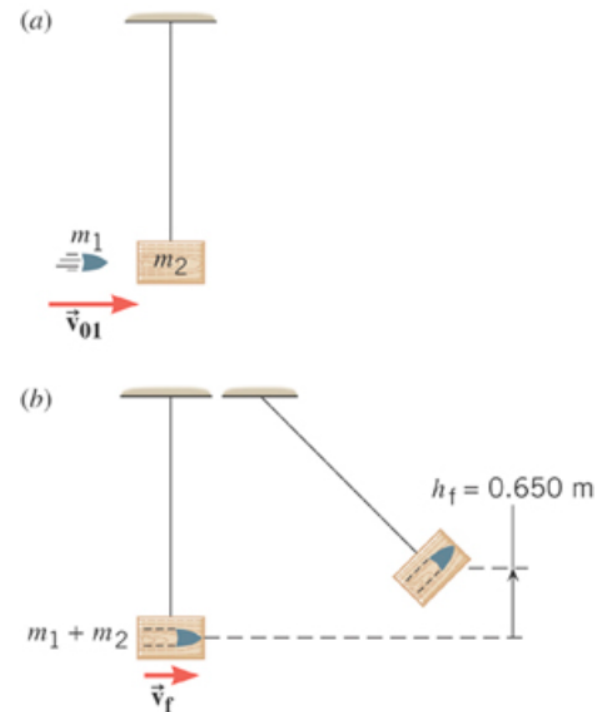
$$m_1 v_{f1} + m_2 v_{f2} = m_1 v_{o1} + m_2 v_{o2}$$
$$(m_1 + m_2) v_f = m_1 v_{o1}$$

Solve for v_{o1}

$$v_{o1} = \frac{(m_1 + m_2) v_f}{m_1}$$

What do we not know? The final speed!!

How can we get it? Using the mechanical energy conservation!



Theoretical Mechanics

Now using the mechanical energy conservation

$$\frac{1}{2}mv^2 = mgh$$

$$\cancel{(m_1 + m_2)}gh_f = \frac{1}{2}\cancel{(m_1 + m_2)}v_f^2$$

$$gh_f = \frac{1}{2}v_f^2$$

Solve for V_f

$$v_f = \sqrt{2gh_f} = \sqrt{2(9.80 \text{ m/s}^2)(0.650 \text{ m})}$$

Using the solution obtained previously, we obtain

$$\begin{aligned} v_{o1} &= \frac{(m_1 + m_2)v_f}{m_1} = \frac{(m_1 + m_2)\sqrt{2gh_f}}{m_1} \\ &= \left(\frac{0.0100 \text{ kg} + 2.50 \text{ kg}}{0.0100 \text{ kg}} \right) \sqrt{2(9.80 \text{ m/s}^2)(0.650 \text{ m})} \\ &= +896 \text{ m/s} \end{aligned}$$

