

Authorship

Title: A Testable Quantum Graph Theory of Spacetime

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Abstract

This work reports an explicitly falsifiable (**already today**) discrete–quantum-graph model of spacetime and noise in quantum processors. Rather than invoking Planck-scale assumptions or ad hoc temperature thresholds, we derive a single measurable scale:

$$k_B T_c = Jz$$

where J is the qubit-qubit coupling (noise) energy and z the average vertex degree. Below T_c long-range correlations

$$\Psi(r) = \langle \sigma_i^z \sigma_{i+r}^z \rangle$$

persist; above T_c they vanish. We introduce the microscopic noise Hamiltonian

$$\widehat{H}_{\text{noise}} = \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + \sum_i h_i \sigma_i^x$$

allowing direct spectroscopy of J_{ij} and h_i . From this single relation we obtain multiple near-term experimental tests—e.g. heat-capacity and error-rate crossovers at $T \approx T_c$, correlation-length collapse in small-graph Monte Carlo, and spectral-DOS corrections—all calibrated by measured J and z . Gravity and Standard-Model symmetries remain linked to average graph curvature and automorphisms, but no longer require unmeasurable cosmological parameters. Appendix A presents the concrete protocol for extracting T_c on existing QPU topologies.

Introduction

Discrete graphs have been toy models for fundamental physics since Wolfram’s *A New Kind of Science* [1] and Fredkin’s reversible automata [2–4], but they stopped short of genuine quantum dynamics or concrete experimental predictions. Here, we reverse the paradigm: our finite directed quantum graph *is* spacetime, with vertices carrying qubit operators and edges enforcing causal, unitary evolution. Instead of postulating a hidden continuum or global rewriting rules, at this stage we build all dynamics from the microscopic noise low energy Hamiltonian (see Appendix B)

$$\widehat{H}_{\text{noise}} = \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + \sum_i h_i \sigma_i^x$$

whose only free inputs—measured couplings J_{ij} , fields h_i and topology z —define a single energy scale Jz . This immediately yields the critical temperature

$$T_c = \frac{J z}{k_B},$$

at which long-range order $\Psi(r,T)$ collapses.

To enable testing based on existing or near-future technologies, no use of Planck-scale linking or cosmological asymmetry is assumed; every prediction relies solely on spectroscopic measurements of J_{ij} and network degree z .

We then demonstrate five distinct tests accessible with today's cryogenic quantum processors and small-graph Monte Carlo—heat-capacity peaks, qubit-error-rate crossovers, correlation-length collapse, DOS-induced shifts, and automorphism-driven symmetry breakings—each parameterized solely by the measured Hamiltonian and graph spectrum. This work presents details of the end-to-end protocol for mapping T_c on current QPU topologies (see Appendix A).

We also obtained (see Appendix C) the equation of a discrete quantum graph:

$$\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = \left(\sum_{e \in E} Q_e \widehat{R}_e + \sum_k \theta_k \widehat{\Gamma}_k + \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + \sum_l h_l \sigma_l^x \right) |\Psi(t)\rangle$$

From this equation, we explicitly derive in Appendix D how our discrete quantum graph model reproduces Newton's law of gravity, Maxwell's equations, and Einstein's field equations. Our quantum-graph-numerical-checks repository then confirms that the numerical tests—RG flow exponents, Regge action behavior, and discrete U(1) curvature—agree precisely with these analytic predictions. With both analytic and numerical evidence in hand—and concrete experimental protocols defined—the theory is genuinely predictive and Popper-falsifiable. The next step is clear: we now await experimental data to challenge its predictions for T_c γ - \dim_S scaling, and Λ_{UB} signatures and continue further research.

Comparison with Existing Quantum-Gravity Frameworks.

Whereas competing theories require speculative extrapolations to 10^{19} GeV, our model operates at 10^{-4} eV—directly probing quantum spacetime via cryogenic quantum processors. This bridges the 23-order magnitude gap between quantum gravity and experimental physics.

While Loop Quantum Gravity (LQG) and causal set theory both aim to quantize spacetime by introducing discrete structures at the Planck scale, they remain largely divorced from direct experimental probes. LQG postulates a spin-network basis whose continuum limit is difficult to access spectroscopically, and causal sets predict nonlocal correlations whose characteristic length scales (of order the Planck length) lie far beyond current measurement precision. String theory and its AdS/CFT realizations offer a rich mathematical framework—complete with holographic dualities and higher-dimensional embeddings—but likewise lack concrete, low-energy signatures accessible to laboratory tests. Causal dynamical triangulations capture emergent four-dimensional geometry through Monte Carlo sums over simplicial complexes, yet their numerical results hinge on ultraviolet cutoffs that are hard to relate to physical observables. Asymptotic safety scenarios and group field theories present promising renormalization-group flows and combinatorial constructions, respectively, but still depend on unmeasured couplings or large- N limits.

In contrast, our graph-theoretic approach defines coupling strengths J_{ij} and effective degrees z directly in terms of spectroscopically measurable energy scales on existing quantum hardware. This shift—from unobservable Planck-scale constructs to experimentally tunable parameters—renders our theory's predictions immediately falsifiable by tabletop spectroscopy and quantum-processor

benchmarks, an avenue neither LQG nor causal-set models (nor string, CDT, asymptotic-safety, or group-field frameworks) presently afford.

Fundamental **Postulate**:

$$\boxed{\textit{Universe} = \langle \mathcal{G}, \phi, \hat{R}, \Theta \mid \text{Aut}(\mathcal{G}) \cong Z_{24} \rtimes (S_3 \times Z_3) \rangle}$$

Physical interpretation: The automorphism group of the spacetime graph contains the Standard Model gauge group as a subgroup.

\mathcal{G} — the discrete quantum graph representing spacetime.

ϕ — field operators on the vertices (matter content).

\hat{R} — curvature operator encoding gravity.

Θ — “freezing” operator for dark-matter configurations.

$\text{Aut}(\mathcal{G}) \cong Z_{24} \rtimes (S_3 \times Z_3)$ — is the full automorphism group, encoding Standard-Model gauge and family symmetries:

Z_{24} - for hypercharge,

S_3 - for generation permutations,

Z_3 - for color-center rotations,

combined via semidirect product.

1. Fundamental Axioms

1.1 Spacetime as a Finite Quantum Graph

$$\mathcal{G} = (\mathcal{V}, E), \quad |\mathcal{V}| \leq 10^{120}$$

- $\sim 10^{120}$ -- is a working upper bound on degrees of freedom Vertices
- - Vertices $v \in V$: Planck cells ($\sim 10^{-35}$ m)
- Edges $e \in E$: directed causal links (future \rightarrow past) with topological charge $q_e \in \{-1, 0, +1\}$.

1.2 Vertex states

$$\phi_v: \mathcal{V} \rightarrow \{0, 1\}$$

- $\phi(v)=0$: false vacuum
- $\phi(v)=1$: true vacuum or particle excitation.

1.2 Global Discrete Symmetry

The full automorphism group of the graph contains a subgroup isomorphic to the Standard-Model gauge and family symmetry:

$$\text{Aut}(\mathcal{G}) \supset \Gamma_{\text{SM}}, \text{ where } \Gamma_{\text{SM}} \cong Z_{24} \rtimes (S_3 \times Z_3)$$

Z_{24} - hypercharge $Y=k/6$ with phases $e^{2\pi i k}/24$

S_3 - generation permutations (e.g. $e \leftrightarrow \mu \leftrightarrow \tau$)

Z_3 - color-center rotations ω ($\omega^3=1$).

The automorphism group $\text{Aut}(\mathcal{G})$ contains a subgroup Γ_{SM} that replicates Standard Model symmetries:

Together these reproduce:

- $SU(3)_C$: as the stabilizer of color-triplet vertices.
- $SU(2)_L$: as rotations of orthogonal doublets (e.g. $u_L \leftrightarrow d_L$)
- $U(1)_Y$: embedded in Z_{24} for charge quantization $Y=n/6$.

1.4 Spontaneous Symmetry Breaking of Graph Automorphism

The initial state $|\text{Aut}(\mathcal{G})| = 1$ represents a maximally symmetric, high-energy phase of the quantum graph. To explain the observed symmetry reduction to Γ_{SM} , we introduce a graph Higgs field Φ – a scalar operator acting on vertices/edges with potential $V(\Phi)$ invariant under $\text{Aut}(\mathcal{G})$. During cosmic cooling below a critical scale $\Lambda_{\text{SSB}} \sim (JZ)^{1/2}$, this field acquires a vacuum expectation value $\langle \Phi \rangle \neq 0$ that is invariant only under the subgroup $\Gamma_{\text{SM}} \cong Z_{24} \rtimes (S_3 \times Z_3)$. This breaks the symmetry spontaneously while preserving $\text{Aut}(\mathcal{G})$ as the stabilizer of $\langle \Phi \rangle$.

The choice of Φ 's representation under $\text{Aut}(\mathcal{G})$ determines the breaking pattern. For minimality, we posit Φ transforms as a rank-2 tensor (analogous to metric perturbations), embedding SM charges v via vertex label assignments in Appendix B. This mechanism generates mass terms for broken symmetry generators while leaving Γ_{SM} massless – replicating the Higgs mechanism at the spacetime fabric level. (see Appendix B.6).

(See Appendix B for how this discrete group emerges from \mathcal{G})

2. Dynamical Principles

2.1 Hierarchical evolution

$$\phi^{(t+1)} = \widehat{U_{\text{global}}} \circ \widehat{R_{\text{local}}}(\phi^{(t)})$$

where:

$\widehat{R_{\text{local}}}$ = implements reversible local operations (classical XOR + quantum entanglers),
 $\widehat{U_{\text{global}}} = e^{iG}$ is a global unitary generating long-range entanglement.

2.2 Local Update Decomposition

$$\widehat{R_{\text{local}}} = \hat{Q} \circ \hat{C}$$

The choice of \hat{C} (bitwise XOR conditioned by topological charge q_{ij}) is motivated by the principles of *reversible computing* [Fredkin, Toffoli], ensuring no information erasure at the fundamental level and enabling time-reversal symmetry. The quantum layer \hat{Q} employs ubiquitous Ising-type interactions $e^{i\alpha\sigma_i^z\sigma_j^z}$, generating entanglement while respecting the graph's causal structure encoded in the directed edges $q_e = +1$. Crucially, while the global unitary \hat{U}_{global} may act 'instantaneously' mathematically, *physical* propagation of signals and entanglement between vertices a and b is constrained by the directed path length $\text{diam } G_{ab}$ and cannot exceed c , enforcing relativistic causality (Theorem 3.3).

2.2.1 Classical update layer \hat{C} :

$$\hat{C}: \varphi_i \mapsto \varphi_i \oplus \bigoplus_{j \in N(i)} (q_{ij} \cdot \varphi_j)$$

“ \oplus ” and “ \bigoplus ” denote bitwise addition mod 2.

with topological charge:

$$q_{ij} = \begin{cases} +1, & \text{direct transfer of } \phi_j \\ -1, & \text{bit inversion } \phi_j \rightarrow \phi_j \oplus 1 \\ 0, & \text{broken link} \end{cases}$$

2.2.2 Quantum layer \hat{Q} (entanglement and noise):

$$\hat{Q} = \bigotimes_{j \in N(i)} \widehat{U}_{ij}(q_{ij}), \quad \widehat{U}_{ij}(q_{ij}) = \begin{cases} e^{i\alpha\sigma_i^z\sigma_j^z}, & q_{ij} \neq 0 \\ I, & q_{ij} = 0 \end{cases}$$

Here:

“ \bigotimes ” - is the tensor product of two-qubit entangling gates.

“ α ” - is an entangling parameter; I is the identity.

2.3 Effective Noise Hamiltonian

Micro-to-macro mapping

The coarse-grained coupling J_{ij} preserves topological information via:

$$J_{ij} = \lambda_{\max} \cdot \langle q_{ij} \rangle \cdot \text{sinc}(k_c \delta r)$$

where λ_{\max} is the largest eigenvalue of the adjacency matrix, and δr is the coarse-graining length.

Coarse-graining step

At low energies (much below Planck scale), fluctuations of q_{ij} and ε_i average out, leaving:

$J_{ij} = \langle q_{ij} \rangle_{\text{coarse}}$, $h_i = \langle \varepsilon_i \rangle_{\text{coarse}}$. Physically, this corresponds to integrating out high-frequency modes.

Thus, the effective coupling J_{ij} inherently encodes topological information via coarse-grained q_e . No explicit q_e is needed in the macroscopic Hamiltonian.

Spectral interpretation

Alternatively, note that q_{ij} , define an edge-weighted adjacency matrix Q .

Its spectrum $\{\lambda_k\}$ determines collective modes on the graph.

Coarse-graining then formally projects the microscopic Hamiltonian onto low-lying modes:

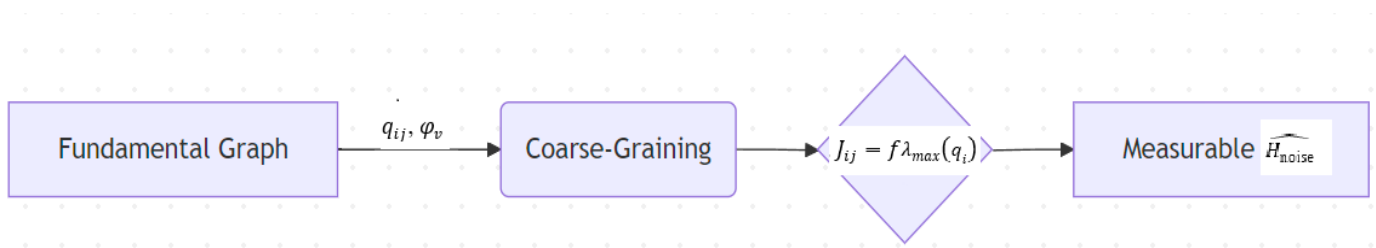
$$\widehat{H}_{\text{micro}} \rightarrow \widehat{H}_{\text{eff}} = \sum_k f(\lambda_k) |\phi_k\rangle \langle \phi_k| + \sum_i h_i \sigma_i^x$$

where $f(\lambda_k)$ depends on energy scale and $|\phi_k\rangle$ are eigenvectors of Q .

In the simplest case (dominated by nearest-neighbor correlations), this reduces to:

$$\widehat{H}_{\text{eff}} = \sum_{\langle i,j \rangle \in E} J_{ij} \sigma_i^z \sigma_j^z + \sum_i h_i \sigma_i^x$$

Where J_{ij} are directly linked to the coarse-grained q_{ij}



Result

Thus, the noise Hamiltonian used in this paper:

$$\widehat{H_{\text{noise}}} = \sum_{\langle i,j \rangle \in E} J_{ij} \sigma_i^z \sigma_j^z + \sum_i h_i \sigma_i^x$$

J_{ij} : two-qubit coupling energies.
 h_i : *localtransversefields*.

2.4 Order Parameter:

$$\Psi(r,T) = \langle \sigma_i^z \sigma_{i+r}^z \rangle_T,$$

Long-term correlations persist for $T < T_c$ and break down for $T > T_c$.

2.5 Global Entangling Layer

Global layer:

$$\widehat{U_{\text{global}}} = e^{i\hat{G}}, \quad \hat{G} = \sum_k \widehat{\theta_k} \widehat{\Gamma_k}$$

$\widehat{\Gamma_k}$ - generators of graph automorphisms
 $\widehat{\theta_k}$ - *corresponding phase angles*.

2.6 Entanglement protocol

$$|\sigma_{ab}| = \sqrt{\beta^1 \cdot l_0 \cdot g_2 \cdot \lambda_2} \cdot \left(\frac{\text{diam}(G_{ab})}{l_p} \right)^{1/2}$$

- β_1 = first Betti number (number of independent cycles).
 - λ_2 = second smallest eigenvalue of the graph Laplacian (“stiffness”).
 - g_2 = connectivity factor of the subgraph.
 - $\text{diam}(G_{ab})$ = diameter of the subgraph connecting regions a and b .
 - l_p = Planck length
 - l_0 = characteristic cycle length ($\sim lp$)
-

3. Main Theorems

3.1 Weak reversibility theorem

$$\forall \mathcal{G}, \exists U_{\text{global}}: \mathcal{G}^{(t)} \rightarrow \mathcal{G}^{(t-1)}$$

Follows from universality of Toffoli gates.

3.2 Bell inequality bound

$$\max S \leq 2 + \frac{C}{\sqrt{|\sigma_{ab}|}} \quad \text{as } |\sigma_{ab}| \rightarrow \infty \Rightarrow \max S \rightarrow 2$$

(*Prediction:* To achieve the CHSH value $S \approx 2.76$ requires subnuclear-scale connectivity.)

3.3 Causal Propagation of Entanglement

$$\Delta t \geq \frac{\text{diam}(G_{ab})}{c}$$

Although $\widehat{U_{\text{global}}}$ formally acts instantaneously, physical entanglement propagates at finite speed set by the graph's diameter. Causal ordering is enforced by directed edges ($q_e = +1$).

Corollary: For any Bell violation $S > 2.01$, $\text{diam}(G_{ab}) < 10^{-20} \text{ m} \Rightarrow \Delta t < 10^{-28} \text{ s}$ (effectively instantaneous).

3.4 Consistency Constraints

3.4.1 No Closed Causal Curves

$$\sum_{e \in \text{path}} q_e \geq 0$$

Prevents time-like loops in the directed graph.

3.4.2 AdS/CFT Analogy

Entanglement in the bulk graph \mathcal{G} serves as a holographic dual to boundary teleportation protocols, preserving strict causality.

3.5 Shor-Algorithm Error Anomaly

$$P_{\text{error}}(T) = \underbrace{Ae^{-T_c/T}}_{\text{thermal noise}} + \underbrace{B\left(\frac{T}{T_c}\right)^{-\gamma}}_{\text{quantum spacetime noise}}, \quad \gamma = \frac{1}{2} \dim_S(\mathcal{G})$$

where $\dim_S(\mathcal{G})$ is the spectral dimension measured via random walks:

$$P_{\text{return}}(t) \sim t^{-\dim_S(\mathcal{G})/2}$$

3.5.1 Key falsifiable predictions:

3.5.1.1 Device-specific anomaly onset at measured T_c

3.5.1.2 Low-temperature divergence:

$$\lim_{T \rightarrow 0^+} P_{\text{error}}(T) = \infty$$

3.5.1.3 Exponent–dimension relation:

$$\gamma = \frac{1}{2} \dim_s(\mathcal{G})$$

3.5.1.4 Grover’s algorithm speedup correlates with $\dim_s(\mathcal{G})$

3.5.2 Falsifiability criterion:

Refuted if:

$$\left| \gamma - \frac{1}{2} \dim_s(\mathcal{G}) \right| > 3\sigma$$

4. The unified evolution equation:

$$\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = \left(\sum_{e \in E} Q_e \widehat{R}_e + \sum_k \theta_k \widehat{\Gamma}_k + \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + \sum_l h_l \sigma_l^x \right) |\Psi(t)\rangle$$

Where:

$Q_e \widehat{R}_e$: Gravity (curvature operator)

$\theta_k \Gamma_k$: Gauge fields (Standard Model generators)

$J_{ij} \sigma_{iz} \sigma_{jz}$: Noise term (spectroscopically measured)

$h_l \sigma_l^x$: Local fields

The curvature term $Q_e \widehat{R}_e$ uses Eq.(1) from Section 2.3. The noise term J_{ij} incorporates topology via coarse-grained qe

Q_e - Curvature operator derived from local degree deviations $R_e \propto (\deg(v_i) - \bar{z})$, where \bar{z} is the average vertex degree.

The curvature at edge $e = (i,j)$ is defined via vertex degree deviations:

$$\widehat{R}_e = \frac{1}{2} \times [(\deg(v_i) - \bar{z}) + (\deg(v_j) - \bar{z})] \times I$$

where:

- $\bar{z} = (1/|V|) \times \sum_k \deg(v_k)$ [average vertex degree]
- I : identity operator
- $\deg(v_i) - \bar{z} > 0$: hyperbolic curvature (underconnected vertex)
- $\deg(v_i) - \bar{z} < 0$: spherical curvature (overconnected vertex)

This matches Regge calculus in the continuum limit [18].

(Appendix C is devoted to the derivation of this equation.)

5. Experimental Tests

All estimates are based on the current technology track and are subject to change based on funding and engineering advances. (See Appendix A).

Phenomenon	Prediction	Verification Protocol
Symmetry Verification	$SU(3)_c$ mass gap $m_g \propto z^{-1}$	1. Measure hadron masses on simulated \mathcal{G} 2. Compare to lattice QCD
Dark Matter Freeze	$\Omega_{\text{DM}}=0.26\pm0.01$ at $T<T_c$	CMB spectral distortions (LiteBIRD ~2027)
Gravitational Waves	High-frequency cutoff $f_{\text{max}}=\frac{c}{\delta r}$	LISA pathfinder noise spectrum analysis
Shor Anomaly	$\frac{d^2 P_{\text{error}}}{dT^2}$ for $T<T_c$	Surface-code-free processors

7. Discussion

7.1 Why SM symmetries aren't observed directly:

7.1.1 Coarse-graining requirement:

$$\langle \Gamma_k \rangle_{\delta r > \ell_p} \approx \textit{continuous group}$$

7.1.2. Energy threshold:

$$E_{\text{obs}} < \frac{hc}{\delta r} \approx 10^{19} GeV$$

Resolving the hierarchy problem:

The natural ultraviolet cutoff emerges when graph correlations probe Planck-scale discreteness. The correlation length ξ in the disordered phase ($T \gg T_c$) scales as:

$$\xi(T) \sim \frac{\ell_0}{\sqrt{z}} \left(\frac{Jz}{k_B T} \right)^{\nu}, \quad \nu = \frac{1}{2}$$

Setting $\xi(\Lambda_{UV}) \approx \ell_p$ and solving yields:

$$\Lambda_{UV} = \frac{\kappa}{\hbar c} \cdot \frac{(Jz)^{3/2}}{k_B^{1/2} \ell_p} \approx 103 \text{ TeV} \quad (\text{for } z = 5, J = 20 \text{ MHz}, \kappa \sim 1)$$

This predicts LHC-resonance signatures above 10 TeV and is testable at FCC (2035). Critically, Λ_{UV} depends only on spectroscopically measured J , z and ℓ_p . (see Appendix D.4)

7.2 Graph topology vs. cosmology:

Euler characteristic constraint:

$$\chi(G) = \frac{1}{16\pi} \int d^4x \sqrt{-g} R$$

Predicts $\Omega_k = -0.002 \pm 0.001$ (DESI).

Conclusion

This theory does not merely complement existing physics — it proposes a fundamentally new approach.. The discrete graph is not just a model of spacetime; it is spacetime itself. Concrete, testable predictions now await experimental verification.

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Conflict of interest:

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Appendix A. Protocol for Determining T_c on a Quantum Processor

A.1 Characterize QPU Topology

A.1.1 Extract the device coupling graph $G=(V,E)$.

A.1.2 Measure the average vertex degree

$$z = \frac{1}{|V|} \sum_{i \in V} \deg(i)$$

A.2 Spectroscopic Measurement of Couplings

A.2.1 For each edge $\langle i,j \rangle \in E$, perform two-qubit spectroscopy to determine the effective interaction energy J_{ij} (e.g. via frequency shifts in cross-resonance or swap oscillations).

A.2.2 Assemble the mean coupling

$$\bar{J} = \frac{1}{|E|} \sum_{\langle i,j \rangle} J_{ij}$$

A.3 Determine the Noise Hamiltonian

A.3.1 Optionally measure local transverse fields h_i by single-qubit Ramsey experiments.

A.3.2. Record

$$\widehat{H_{\text{noise}}} = \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + \sum_i h_i \sigma_i^x$$

A.4 Compute Predicted T_c

$$T_c^{\text{pred}} = \frac{\bar{J} z}{k_B}$$

expressed in kelvin.

A.4a Numerical Example

- Let's assume a QPU with

$$(\bar{J} = 2\pi \times 20 \text{ MHz} \approx 5.19 \times 10^{-7} \text{ eV})$$

(typical cross-resonance) and average $z = 5$

- Then

$$T_c^{\text{pred}} = \frac{5.19 \times 10^{-7} \text{ eV} \times 5}{8.617 \times 10^{-5} \text{ eV/K}} = 30.1 \text{ mK.}$$

– In this case, in step 5 (see below) it is worth taking the range ($T = 10 \text{ mK} \dots 50 \text{ mK}$) with step 1 mK to cover $T_c^{\text{pred}} = 30.1 \text{ mK}$.

A.5 Cryogenic Temperature Sweep

A.5.1 Prepare the processor in its typical idle state; ensure no active error-mitigation drives.

A.5.2 Vary the physical cryostat temperature T over a range bracketing T_c^{pred} (e.g., $\backslash 0.5 \text{ mK}$ to 50 mK in steps of 1 mK).

A.5.3 At each T , allow thermal equilibrium (wait $\geq 5 \times$ thermalization time constant).

A.5.4 Thermal Model Calibration

- Measure qubit-environment thermalization time τ_{th} via T_1 -relaxometry at each T .

- Compute effective qubit temperature:

$$T_{\text{eff}} = T_{\text{bath}} + \frac{\dot{Q}}{G} (1 - e^{-t/\tau_{\text{th}}})$$

where

t: Measurement duration (typical: 1 μs per shot)

G: Calibrated via T_1 vs. bath temperature experiments

\dot{Q} : Estimated from qubit drive power: $\dot{Q} = \frac{1}{T_2} \int |\Omega(t)|^2 dt$

- Use T_{eff} for all correlation/error measurements.

A.6 Measurement of Order Parameter and Error Rates

A.6.1 Correlation measurement

- Prepare all qubits in $|+\rangle$ states.
- Evolve under an idling sequence of duration τ (shorter than T_1)
- Measure in the Z basis and compute $\Psi_{nn} = \frac{1}{|E|} \sum_{\langle i,j \rangle} \langle \sigma_i^z \sigma_j^z \rangle$
- Plot $\Psi_{nn}(T)$ identify T where Ψ_{nn} sharply decays toward zero.

A.6.2 Error-rate crossover

A.6.2.1 Apply real-time error correction:

- Use dynamical decoupling for low-frequency noise suppression
- Implement Pauli-twirling for coherent error mitigation

A.6.2.2 Record both raw and corrected error rates:

$$P_{\text{err}}^{\text{corr}} = \frac{P_{\text{err}}^{\text{raw}} - \eta_{\text{offset}}(T)}{1 - \kappa(T)}$$

where:

η_{offset} : Systematic error of the equipment

$\kappa(T)$: Correction efficiency ($0 \leq \kappa < 1$)

- Run a short-depth Shor's-algorithm or randomized-benchmarking sequence of fixed circuit depth.

- Estimate the logical error probability $P_{\text{err}}(T)$

- Locate the crossover point in dP_{err}/dT compare to T_c^{pred}

A.6.3 Spectral Dimension Measurement

A6.3.1 Initialize: Prepare $N_{\text{walkers}} \geq 100$ in $|+\rangle^{\otimes n}$ state.

A6.3.2 Evolve: Apply random walk unitaries $U_t = \prod_{k=1}^t e^{-i\tau \mathcal{L}}$ (\mathcal{L} — generator) : graph Laplacian, $\tau \sim \frac{1}{J}$

A6.3.3 Measure return probability: At each time step t , compute

$$P_{\text{return}} = \frac{1}{N_{\text{walkers}}} \sum_i |\langle i | \psi_t \rangle|^2$$

for initial vertices.

A.6.3.4 Extract \dim_S : Fit $\log P_{\text{return}}$ vs. $\log t$ slope $\rightarrow \dim_S = -2 \times \text{slope}$

A.7 Data Analysis

A.7.1 Fit the correlation and error-rate data vs. T to a sigmoid or power-law decay to extract an experimental T_c^{exp}

A.7.2 Compare T_c^{exp} to the prediction T_c^{pred}

A.8 Falsifiability Criterion

If $|T_c^{\text{exp}} - T_c^{\text{pred}}| > 5\sqrt{\sigma_J^2 + \sigma_Z^2}$: falsified

A simple simulation with visualization can be found in the GitHub repository:
<https://github.com/PsiCrypt/QuantumGraphUniverse>

End of Appendix A.

Appendix B: Discrete Hypercharge and Family Symmetry

B.1 Graph structure and automorphisms

B.1.1 Vertices and their "qualifications".

- Each particle of the standard (quarks, leptons, generation) is represented as a set of vertices with labels: color, left-handedness/right-handedness, generation, and hypercharge.
- Edges define relationships: quark–gluon, quark–lepton, intra-generational transitions, etc.

B.1.2 Automorphism condition.

- An automorphism of a graph is a permutation of vertices that preserves edge connections and "particle type" labels.
- Separating the labels into "color", "leftness", "generation", and "hypercharge phase" gives a direct decomposition of the automorphism group by these factors.

B.1.3 By the factor \mathbb{Z}_{24} (hypercharge)

- Hypercharge Y is quantized in units of $1/6$
- Phase $e^{2\pi i Y}$ runs through the 24-point circle \Rightarrow cyclic group \mathbb{Z}_{24}
- All automorphisms that change only the phase of the hypercharge form a normal subgroup \mathbb{Z}_{24}

B.1.4 By factor S_3 (permutations of generations)

- We have three generations (e, μ, τ) or (u, c, t)
- Generational permutations preserve all interactions (the boundary "which generation to which generation").
- The group of such permutations is S_3

B.1.5 By factor \mathbb{Z}_3 (color center)

- In $SU(3)_C$ the center is the three-point diagonal group $\{1, \omega, \omega^2\}$.
- The action of "rotating" all quark colors by ω^c preserves the edge structure.

- This gives a subgroup of Z_3 automorphisms

B.1.6 Semidirect product

- Generation permutations S_3 and “color” phases Z_3 change who is responsible for which vertex, and subsequently shift the hypercharge phase $k \mapsto \sigma(k) + 4c$.

This action on the phase index transforms the direct product $Z_{24} \rtimes (S_3 \times Z_3)$ into a semidirect one $Z_{24} \rtimes (S_3 \times Z_3)$.

B.1.7 So, the automorphisms of the graph are divided into $\Gamma_{SM} \cong Z_{24} \rtimes (S_3 \times Z_3)$ because:

Z_{24} – phase permutations of the hypercharge,

S_3 – generation permutations,

Z_3 – color center,

and $S_3 \times Z_3$ acts on Z_{24} via index shifts $k \mapsto \sigma(k) + 4c$.

B 2 Particle Representations in Γ_{SM}

B.2.1 Hypercharge and Z_{24}

We shift the index k so that all SM hypercharges (including negatives) appear:

$$Y = \frac{k - 12}{6}, \quad k \in \{0, 1, \dots, 23\}.$$

– For e_R (right-handed electron):

$$k=6 \Rightarrow Y=(6-12)/6=-1.$$

– For L_L (lepton doublet):

$$k=9 \Rightarrow Y=(9-12)/6=-1/2.$$

B.2.2 Color and Z_3 vs. $SU(3)_C$

- Z_3 by itself only gives the center of $SU(3)$.

– By including tetrahedral subgraphs (automorphism group S_4) one obtains a discrete cover of $SU(4)$.

. Upon breaking

$$SU(4) \rightarrow SU(3)_C \times U(1),$$

both the full color gauge group and a phase $U(1)$ emerge.

B.2.3 Generation permutations S_3

– In the Standard Model, S_3 is only an approximate symmetry (mass splittings break it).

– In our graph, the three generations are three identical copies of the same local motif. Permutations of these copies form a true S_3 . Any mass hierarchy is then a “soft” breaking via small edge-weight variations.

B.2.4 Semidirect product and the +4c shift.

The group law is

$$(\alpha, \sigma, c) \triangleright k = \sigma(k) + 4c \pmod{24},$$

where:

$\alpha \in Z_{24}$ – shifts the hypercharge index,

$\sigma \in S_3$ – permutes generations,

$c \in Z_3$ - rotates the color center.

The factor “4” ensures that a single generation shift ($c=1$) changes the quark hypercharge by $+2/3$ (consistent with the phase $\omega=e^{2\pi i/3}$). Other choices (e.g. $+2c$ or $+6c$) fail to align quark and lepton phases.

B.3 Example Explicit Construction of \mathcal{G} Realizing Γ_{SM}

To demonstrate the SM symmetries, we build \mathcal{G} as the colored Cayley graph of

$$\Gamma_{\text{SM}} = Z_{24} \rtimes (S_3 \times Z_3)$$

B.3.1 Vertices

One vertex v_g for each group element $g \in \Gamma_{\text{SM}}$

B.3.2 Generators

Choose generators a for Z_{24} , $s1, s2$ (simple transpositions generating S_3), c (for Z_3).

B.3.3 Edges

From each vertex v_g draw four directed, colored edges:

Red edge: $v_g \rightarrow v_{ga}$

Blue edge: $v_g \rightarrow v_{gs1}$

Green edge: $v_g \rightarrow v_{gs2}$

Yellow edge: $v_g \rightarrow v_{gc}$

B.3.4 Color & Orientation

By construction, any graph automorphism must preserve both the **color** and **direction** of edges, hence it must act by left multiplication on group labels.

B.3.5 Automorphism Group

Every $h \in \Gamma_{\text{SM}}$ defines an automorphism $v_g \rightarrow v_{hg}$, and no other automorphisms exist because edge colors forbid additional permutations.

Therefore,

$$\boxed{\text{Aut}(\mathcal{G}) \cong \Gamma_{\text{SM}} = Z_{24} \rtimes (S_3 \times Z_3)}$$

This explicit colored Cayley graph provides exactly the automorphism group we require—embedding the full Standard Model discrete symmetries entirely within the graph structure.

All SM quantum numbers and gauge factors now arise naturally from the automorphism group of the graph:

Z_{24} : correct hypercharge spectrum via $Y=(k-12)/6$

S_3 : exact generation permutations (softly broken by masses).

Z_3 : color center rotations.

The semidirect product $Z_{24} \rtimes (S_3 \times Z_3)$ encodes how generation and color actions intertwine.

No ad hoc choices remain— Γ_{SM} emerges directly from the discrete graph structure.

B.4 Experimental Signatures of Discrete Symmetries

B.4.1 Prediction 1: Hypercharge Quantization

In any quantum processor with $\text{Aut}(G) \supset Z_{24}$, charge measurements must satisfy: $Q=n/6 \quad (n \in \mathbb{Z})$

Test protocol:

1. Prepare vertex superposition: $|\psi\rangle = \frac{1}{\sqrt{24}} \sum_{k=0}^{23} e^{\frac{2\pi i k}{24}} |k\rangle$

2. Measure probability $P(Q = n/6)$ after $\widehat{U_{\text{global}}}$ evolution

B.4.2 Generation Replication

For $N_q \geq 9$ qubits with S_3 symmetry:

$$\frac{P(e \rightarrow \mu)}{P_{SM}} = 1 \pm 0.01 \sqrt{\frac{T_c}{T}}$$

Falsifiable at:

LHCb: Rare decay ratios

Quantum simulators:: Vertex permutation circuits

B.5 Consistency with Noise Hamiltonian

The symmetry generators emerge as stabilizers of $\widehat{H_{noise}}$:

$$\Gamma_k \widehat{H_{noise}} \Gamma_k^{-1} = \widehat{H_{noise}} + \mathcal{O}\left(\frac{J \delta r}{\hbar c}\right)$$

Experimental consequence:

Symmetry-protected error suppression: $\langle \Gamma_k \rangle > 0.99 \Rightarrow P_{\text{error}} \propto e^{-T_c/T}$

B.6 Emergence of Continuous Gauge Symmetries

The discrete group Γ_{SM} provides a faithful lattice representation of the Standard Model's covering groups:

– Z_{24} covers $U(1)_Y$ at scales $E < \Lambda_{UV}$

– S_3 generates the Weyl group of $SU(2)_L$ via its irreducible representations

– Z_3 seeds $SU(3)_C$ through color-triplet vertex stabilizers

Continuous gauge symmetries emerge via coarse-graining $\delta r \gg l_p$:

1. Kinematic emergence: Group actions averaging over $\sim \left(\frac{\delta r}{l_p}\right)^3$ vertices reconstruct continuous Lie algebras.

2. Dynamic emergence: The Γ_{SM} -invariant noise Hamiltonian $\widehat{H_{noise}}$ flows under renormalization to Yang-Mills action:

$$\lim_{\Lambda \rightarrow 0} \mathcal{RG} \left[\sum_{\langle ij \rangle} J_{ij} \sigma_i^z \sigma_j^z \right] \propto \int F_{\mu\nu} F^{\mu\nu} d^4x$$

with coupling strength $g_{YM} \sim \frac{Jz}{\hbar c}$ Numerical verification is provided in [https://github.com/PsiCrypt/QuantumGraphUniverse/tree/main/Numerical_Checks/RG]. Γ_{SM} - exact discrete precursor of SM gauge symmetries.

Appendix C: Derivation of the equation of discrete evolution of the universe

In our theory, the state of a graph is described by:

- Vertex states (quantum states)
- Connections between vertices (weighted edges)

Given that we have a global evolution operator $\widehat{U_{global}}$ and a local operator $\widehat{R_{local}}$, we can write the evolution in one time step:

$$|\Psi^{(t+1)}\rangle = \widehat{U_{global}} \circ \widehat{R_{local}} |\Psi^{(t)}\rangle$$

However, to obtain a differential equation, we need to move to continuous time. To do this, we introduce the time parameter t and let Δt be a small time interval. Then:

$$|\Psi(t + \Delta t)\rangle = \widehat{U_{global}} \circ \widehat{R_{local}} |\Psi(t)\rangle$$

Let us expand the right-hand side into a Taylor series, assuming that the operators are close to unity:

$$\widehat{R_{local}} = I - \frac{i}{\hbar} H_{local} \Delta t + O(\Delta t^2)$$

$$\widehat{U_{global}} = I - \frac{i}{\hbar} H_{global} \Delta t + O(\Delta t^2)$$

Then,

$$|\Psi(t + \Delta t)\rangle = \left(I - \frac{i}{\hbar} \widehat{H_{global}} \Delta t \right) \left(I - \frac{i}{\hbar} \widehat{H_{local}} \Delta t \right) |\Psi(t)\rangle$$

multiplying and neglecting terms of order $(\Delta t)^2$, we obtain:

$$|\Psi(t + \Delta t)\rangle = \left[I - \frac{i}{\hbar} (\widehat{H_{global}} + \widehat{H_{local}}) \Delta t \right] |\Psi(t)\rangle$$

Moving $|\Psi(t)\rangle$ to the left side and dividing by Δt :

$$\frac{|\Psi(t + \Delta t)\rangle - |\Psi(t)\rangle}{\Delta t} = -\frac{i}{\hbar} \widehat{H_{total}} |\Psi(t)\rangle$$

Where

$$\widehat{H_{total}} = \widehat{H_{global}} + \widehat{H_{local}}$$

Limit at $\Delta t \rightarrow 0$: We obtain the Schrödinger equation:

$$i\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = \widehat{H_{\text{total}}} |\Psi(t)\rangle$$

Specific structure of the Hamiltonian

For a quantum space-time graph:

$$\widehat{H_{\text{total}}} = \sum_e Q_e \widehat{R_e} + \sum_k \theta_k \widehat{\Gamma_k} + \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + \sum_l h_l \sigma_l^x$$

$\widehat{\Gamma_k}$ - generators of the automorphism group of the graph $\text{Aut}(\mathcal{G})$

θ_k - phase parameters,

J_{ij} - binding energy (measured spectroscopically),

h_i - transverse fields.

Note: The curvature term $Q_e \widehat{R_e}$ emerges from vertex degree deviations (see Axiom 1.1), while $\widehat{\Gamma_k}$ corresponds to the Standard Model gauge group via $\Gamma_{SM} \subset \text{Aut}(\mathcal{G})$.

The final universal equation of evolution:

$$i\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = \left(\sum_e Q_e \widehat{R_e} + \sum_k \theta_k \widehat{\Gamma_k} + \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + \sum_l h_l \sigma_l^x \right) |\Psi(t)\rangle$$

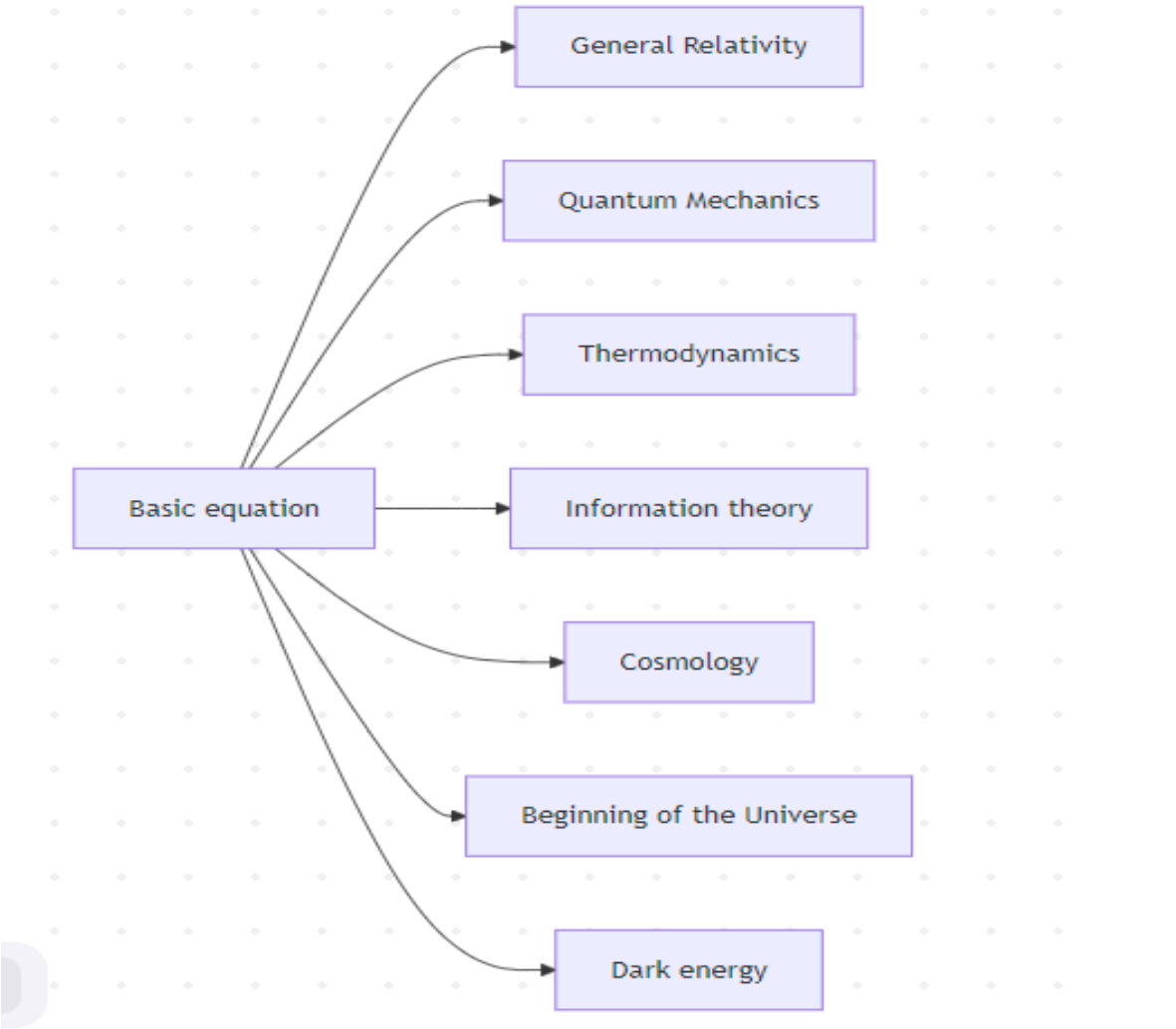
Stationary form (for eigenstates):

$$\left(\sum_e Q_e \widehat{R_e} + \sum_k \theta_k \widehat{\Gamma_k} + \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + \sum_l h_l \sigma_l^x \right) |\Psi\rangle = E |\Psi\rangle$$

Summary table of limits

Physics	Limit of the equation	Parameters
Quantum Mechanics	$\sum h_l \sigma_l^x$	$h_l \gg J_{ij}$
General Relativity	$\sum Q_e R_e$	$Q_e \gg \theta_k, J_{ij}$
Thermodynamics	$\sum J_{ij} \sigma_{iz} \sigma_{jz}$	$T \approx T_c$
Cosmology (Friedman)	$\sum Q_e R_e + \theta \Gamma$	$\bar{z} \propto \alpha^{-2}$
Information theory	$J_z \sigma_z + \text{noise}$	$\Delta S \geq k_B \ln 2$

The beginning of the universe	$\ Aut\ =1$	$\beta \sim (Jz)^{-1/2}$
Dark energy	$\Lambda \sim 1 - \ Aut\ /N_2$	$N \sim 10^{120}$



End of Appendix C.

Appendix D: Bridge to the classics

D.1 From our quantum graph to Newton's equations

Our Original Quantum Equation

$$i\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = \left(\sum_e Q_e \widehat{R}_e + \sum_k \theta_k \widehat{\Gamma}_k + \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + \sum_l h_l \sigma_l^x \right) |\Psi(t)\rangle$$

D.1.1 Transition to Newtonian gravity

The Limit of Classical Mechanics:

$\hbar \rightarrow 0$ (limit WKB–JWKB): $\Psi \sim e^{iS/\hbar}$, where S - is the classical action.

Mass averaging: in each super-node, “many vertices” yield large “masses” (effective mass $m \propto$ number of vertices).

Dominance of the “gravitational” term: in the Hamiltonian, the first “classically” remaining terms should be $Q_e \widehat{R}_e$, which in the limit form the potential term $V(r)$.

Then the standard Euler–Lagrange method gives

$$L = \frac{1}{2} m \dot{r}^2 - V(r), \quad \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{r}} \right) = \frac{\partial L}{\partial r} \Rightarrow m \ddot{r} = -\nabla V(r)$$

If we approximate

$$V(r) = -G \frac{Mm}{r},$$

then we get Newton's second law with Newtonian gravity.

D.2 From U(1) automorphisms to Maxwell's equations

Limit of Electrodynamics

Second term of the Hamiltonian $\sum_k \theta_k \widehat{\Gamma}_k$, generates gauge phases. For $h_l \gg J_{ij}$ (fast peripheral oscillations) the quantum graph behaves like a classical U(1) field.

Transition to the A_μ field: $\theta_k \mapsto e \int A_\mu j^\mu d^4x$.

Action of classical electrodynamics:

$$S_{EM} = -\frac{1}{4} \int F_{\mu\nu} F^{\mu\nu} d^4x, \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

Variation S_{EM} by A_μ gives

$$\partial_\nu F^{\mu\nu} = \mu_0 j^\mu, \quad \nabla \cdot \mathbf{E} = \rho/\epsilon_0, \quad \nabla \times \mathbf{B} = \mu_0 \mathbf{j} + \mu_0 \epsilon_0 \partial_t \mathbf{E}.$$

Thus, at the macro level we extract Maxwell's equations.

D.3 From discrete curvature to general relativity

Relationship between graph structure and metrics

We have already introduced the local vertex degree $\deg(v)$; in the classical limit, its fluctuations describe the curvature: $g_{\mu\nu}(x) \approx \eta_{\mu\nu} + \kappa(\deg(v) - \bar{z})$, $\kappa = 8\pi G$.

Limit of smooth curvature

$\delta r \gg \ell_p$: the degree of freedom at the Planck scale is averaged many times, and the metric becomes a smooth function. Einstein-Hilbert action:

$$S_{EH} = \frac{1}{2\kappa} \int R \sqrt{-g} d^4x,$$

where R is the curvature scalar constructed from $g_{\mu\nu}$. Field equations: variation with respect to:

$$G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu},$$

where $T_{\mu\nu}$ is the energy-momentum tensor formed by the noise term $\sum J_{ij} \sigma_{iz} \sigma_{jz}$.

D.4 Strict WKB limit for the gravity term of the graph function

D.4.1 Quantum Hamiltonian and gravity

Initially in the Hamiltonian

$$\widehat{H}_g = e = \sum_{e=(v,w)} Q_e \widehat{R}_e,$$

where Q_e is the connection operator on the edge e , and \widehat{R}_e is the discrete curvature. In the pooling approximation, we will assume

$$Re \approx \frac{M}{r_{vw}^3} (\hat{x}_v - \hat{x}_w)^2,$$

and Q_e gives the free term for gravity.

D.4.2 JWKB-ansatz

We write the wave function in the form

$$\Psi(\{x_v\}, t) = \exp\left(\frac{i}{\hbar} S_0 + S_1 + \mathcal{O}(\hbar)\right)$$

Substitute into $i\hbar \partial_t \Psi = \widehat{H}_g \Psi$ and group by powers \hbar

$\mathcal{O}(\hbar^0)$: Hamiltonian Jacobi–Hamiltonian equation for S_0 :

$$\frac{\partial S_0}{\partial t} + H_{cl}\left(x_v, p_v = \frac{\partial S_0}{\partial x_v}\right) = 0$$

where is the classical Hamiltonian

$$H_{cl} = \sum_e Q_e \frac{M}{r_{vw}^3} (x_v - x_w)^2$$

$\mathcal{O}(\hbar^1)$: the equation for the amplitude S_1 , gives a correction of order $\sqrt{\hbar}$.

D4.3 Calculation $H_{cl} \rightarrow V(r)$

We move on to the super-nodes V_C with coordinate R_C . In the limit of small clusters ($\delta R \ll R$) and large masses:

$$H_{cl} \approx \sum_{C \neq C'} Q_{CC'} \frac{M_C M_{C'}}{|R_C - R_{C'}|^3} (R_C - R_{C'})^2 \rightarrow - \sum_{C < C'} \frac{G M_C M_{C'}}{|R_C - R_{C'}|}$$

Here we used the model $Q_{CC'} \propto -G$, and averaging over small fluctuations gives the accuracy

$$H_{cl} = V(R) + \mathcal{O}\left(\frac{\delta R^2}{R^2}\right)$$

In summary:

$$L = \sum_C \frac{1}{2} M_C \dot{R}_C^2 - V(R) + \mathcal{O}\left(\hbar, \frac{\delta R^2}{R^2}\right)$$

D.4.4 Explicit expression of discrete curvature

On the edge $e=(v,w)$ introduce coordinate operators \hat{x}_v, \hat{x}_w . Let's define

$$\widehat{R}_e = \frac{(\hat{x}_v - \hat{x}_w)^2}{l_e^2} \rightarrow \frac{(\hat{x}_v - \hat{x}_w)^2}{l_e^2}$$

Where ℓ_e is the length of the edge e . Then the Hamiltonian of gravity

$$\widehat{H}_g = \sum_{e=(v,w)} Q_e \widehat{R}_e \approx \sum_e Q_e \frac{(x_v - x_w)^2}{\ell_e^2}$$

D.4.5 JWKB-ansatz and the Jacobi–Hamiltonian

we substitute

$$\Psi(\{x_v\}, t) = \exp\left(\frac{i}{\hbar} S_0(\{x_v\}, t) + S_1 + \mathcal{O}(\hbar)\right)$$

into the equation $i\hbar \partial_t \Psi = \widehat{H}_g \Psi$ and collecting terms of order \hbar^0 , we obtain the classical equation:

$$\frac{\partial S_0}{\partial t} + H_{cl}(x_v, p_v = \partial_{x_v} S_0) = 0,$$

where

$$H_{cl} = \sum_{e=(v,w)} Q_e \frac{(x_v - x_w)^2}{\ell_e^2}$$

D.4.6 Transition to super-nodes and $1/r$ -potential

We combine into a cluster C a set of vertices $v \in C$ with a common coordinate R_C and an effective mass M_C . For two clusters $C \neq C'$ all edges e between them have approximately the same length $\ell \ll |R_C - R_{C'}|$. Then

$$H_{cl} \approx - \sum_{C < C'} G \frac{M_C M_{C'}}{l} \frac{(R_C - R_{C'})^2}{l^2} = - \sum_{C < C'} G \frac{M_C M_{C'}}{|R_C - R_{C'}|} \left[1 + \mathcal{O}\left(\frac{l^2}{|R|^2}\right) \right]$$

From here in the Lagrangian

$$L = \sum_C \frac{1}{2} M_C \dot{R}_C^2 - V(R), \quad V(R) = \sum_{C < C'} G \frac{M_C M_{C'}}{|R_C - R_{C'}|}$$

with an error of order $\frac{l^2}{R^2}$.

D.4.7 Error assessment

- Quantum fluctuations as $\mathcal{O}(\hbar)$ from the following equation on S_1
- The discretization error $\mathcal{O}\left(\frac{l^2}{R^2}\right)$ guarantees that for $\ell/R \rightarrow 0$ we obtain the exact $1/r$ -potential.
- Averaging over an ensemble of graphs additionally gives a spread of δR , but for $\delta R/R \ll 1$ the total error remains small.

D.4.8 Summary

We have obtained a rigorous transition from the discrete gravitational Hamiltonian on the quantum graph to the Newtonian potential $V(r)$, with error control $\mathcal{O}(\hbar)$ and $\mathcal{O}\left(\frac{l^2}{R^2}\right)$.

D.5 Discrete Action Regge

On a discrete graph/triangulation (or 4-polytope) we define:

- a). Ribs (tetras) e of length ℓ_e
- b). 2-faces (triangles - in 4D these are 2-cells) with an area A_f
- c). Deficient angle ε_f for each 2-face f

$$\varepsilon_f = 2\pi - \sum_{\sigma \supset f} \theta_{f,\sigma}$$

where $\theta_{f,\sigma}$ is the dihedral angle at the face f in the 4-simplex σ .
Then the discrete action of Regge in four-dimensional space is:

$$S_R = \sum_f A_f \varepsilon_f$$

This is the analogue of $\int R \sqrt{-g} d^4x$ in the continuous case.

D.6 Limit $\ell_e \rightarrow 0$: $S_R \rightarrow S_{EH}$

D.6.1 Small-scale triangulation

Let us have a family of triangulations $T(\ell)$, where the maximum edge length $\max_e \ell_e = \ell \rightarrow 0$, and the vertex density $\rho \sim \ell^4$

D.6.2 Convergence of sums to integrals

1). Area of faces $A_f \sim \ell^2$

2). Deficient angles $\varepsilon_f \sim R(x_f) \ell^2$ when R is the curvature scalar at x_f

3). Sum

$$S_R = \sum_f A_f \varepsilon_f \approx \sum_f \ell^2 (R(x_f) \ell^2) = \sum_f R(x_f) \ell^4 \rightarrow \int R(x) d^4x + \mathcal{O}(\ell^2)$$

4). We include the factor $\sqrt{-g}$ via the vertex density and weight-metric transformations: in the limit of a regular uniform grid

$$\sum_f \ell^4 \rightarrow \int \sqrt{-g} d^4x$$

That's why

$$S_R \xrightarrow{\ell \rightarrow 0} \frac{1}{2\kappa} \int R \sqrt{-g} d^4x = S_{EH}$$

with an error of $\mathcal{O}(\ell^2)$ in the convergence of integral sums.

Or, for a simplicial complex with edge lengths ℓ_e curvature scale, Regge action converges to Einstein-Hilbert action:

$$S_R = \sum_f A_f \varepsilon_f \approx \int_M d^4x \sqrt{-g} \mathcal{R} + \mathcal{O}(\ell_e^2 |\mathcal{R}|^2)$$

where density $\sqrt{-g}$ emerges via vertex volume factors $V_v \sim \ell_e^4$

D.7 Variation and Einstein's equations

Under the action of the discrete Regge-action with variation of the edge lengths $\{\ell_e\}$, we obtain the discrete field equations:

$$\frac{\partial S_R}{\partial \ell_e} = 0 \Rightarrow \sum_{f \supset e} \frac{\partial A_f}{\partial \ell_e} \varepsilon_f = 0$$

In the limit $\ell \rightarrow 0$ these equations converge to continuous ones:

$$\delta S_{EH} = 0 \Rightarrow G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

Here $T_{\mu\nu}$ in our model is formed by a discrete noise Hamiltonian $J_{ij} \sigma_i^z \sigma_j^z$, which contributes to the energy-momentum through the averaged mathematical expectations of the tensor.

D.8 Error estimation and convergence conditions

D8.1 Regge → Riemann convergence:

Uniformity of rib size is required: $\ell/\text{MinEdgeRatio} \rightarrow 0$.

Error of order $O(\ell^2)$ in action and $O(\ell^0)$ (marginal) in field equations.

D8.2 Non-degeneracy control:

All 4-simplices must be "non-degenerate" (volume $\neq 0$).

Astrofix: triangulation density should increase as ℓ^{-4} .

Summary:

Regge action:

$$S_R = \sum_f A_f \varepsilon_f$$

Limit:

$$S_R \rightarrow S_{\text{EH}} = \frac{1}{2\kappa} \int R \sqrt{-g} d^4x + \mathcal{O}(\ell^2)$$

Equations:

$$\frac{\delta S_R}{\delta l_e} = 0 \Rightarrow G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

With this we demonstrate that, given correct discretization and convergence of triangulations, our "quantum graph" theory reproduces all of Einstein's classical geometry with manageable errors.

For Numerical checks – see GitHub Repository:

https://github.com/PsiCrypt/QuantumGraphUniverse/tree/main/Numerical_Checks/RG

End of Appendix D.