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# Computational Physics – Exercise 1

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## 1 What is $J$ ?

For the 1-D Ising Hamiltonian,

$$\mathcal{H}(s) = -J \sum_{\langle x,y \rangle} s_x s_y - h \sum_x s_x \quad (1)$$

Where  $\langle x, y \rangle$  denotes a sum over the nearest spins and  $h$  is the external coupling.  $J$  is a coupling constant that models the interaction between neighbouring spins. For  $J > 0$ , the similar nearest neighbour pairs are favoured as this minimizes the total energy. In this context of magnets, this corresponds to having ferromagnetic order (see figure 1) [6].

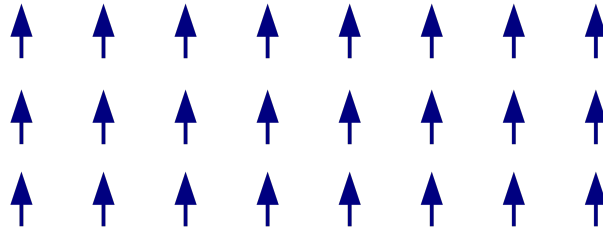


Figure 1: Ferromagnetic ordering in  $d = 2$ . Taken from [3]

For  $J < 0$ , opposite nearest neighbour pairs are favoured. In this context of magnets, this corresponds to having anti-ferromagnetic order (see figure 2).

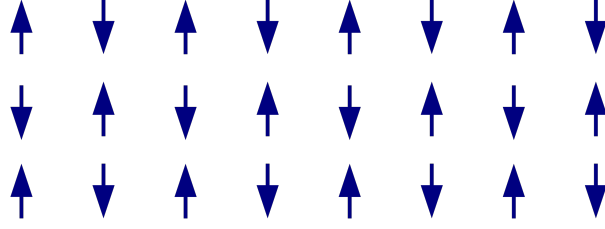


Figure 2: Antiferromagnetic ordering in  $d = 2$ . Taken from [2]

## 2 Periodic boundary conditions

Having periodic boundary conditions implies that all sites have the same number of nearest neighbours [5].

## 3 Simulating the Ising model in $d = 1$

The average magnetization is given by:

$$\langle m \rangle = \frac{T}{N} \frac{\partial}{\partial h} \ln Z \quad (2)$$

Where the partition function  $Z$  is given by,

$$Z = \lambda_+^N + \lambda_-^N \quad (3)$$

for

$$\lambda_{\pm} = e^{\frac{J}{T}} \left( \cosh \left( \frac{h}{T} \right) \pm \sqrt{\sinh \left( \frac{h}{T} \right)^2 + e^{-4\frac{J}{T}}} \right) \quad (4)$$

In this problem the relevant dimensionless ratios are as follows:

- $\frac{J}{T}$
- $\frac{h}{T}$

In order to calculate the average magnetization numerically, we use the formula

$$\langle m \rangle = \frac{1}{N} \sum_i S_i \quad (5)$$

We shall now simulate the Ising model in  $d = 1$ , the python code for that can be found at [\[1\]](#).

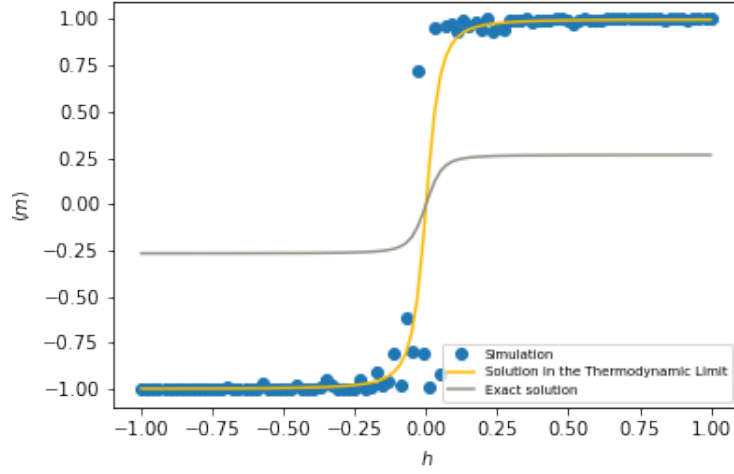


Figure 3: A plot of the  $\langle m \rangle$  versus external field  $h$  for a fixed  $J = 1$ ,  $T = 0.75$  and  $N = 20$ .

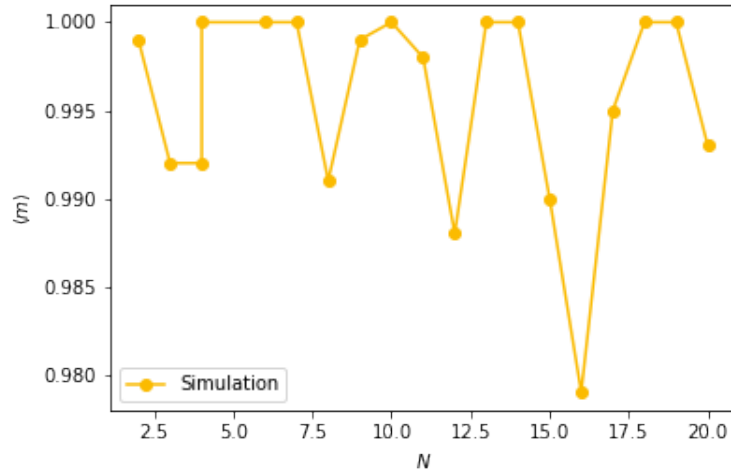


Figure 4: A plot of the  $\langle m \rangle$  versus external field  $N$  for a fixed  $J = 1$ ,  $T = 0.75$  and  $h = 0.5$ .

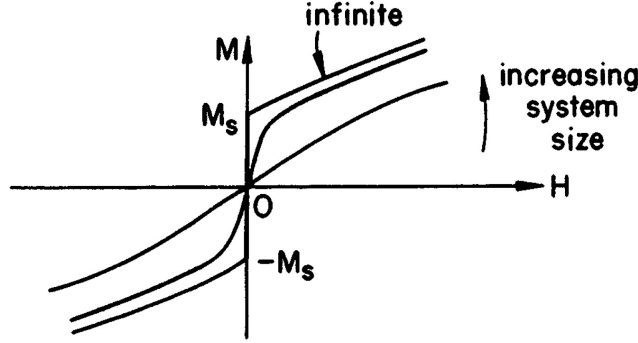


Figure 5: A plot of magnetisation  $M$  versus the external coupling parameter  $H$ . Taken from [4].

Numerically we were able to approximate  $\langle m \rangle$  close enough to the value we observe at the thermodynamic limit as shown in figure 3. We can compare this to a plot from the literature in figure 5, we find these to be in close agreement. We find that the value of  $\langle m \rangle$  is constrained heavily by the couplings  $h$  and  $J$ , however it does vary a little depending on the number of sites for  $d = 1$  i.e. if it is even or odd and so on (see figure 4).

We computed the error for simulation via the residual standard deviation, which came to about 0.0004.

## References

- [1] P. Anancia Devaneyan and R. K. Senthilkumar. Exercises for the Physics 760: Computational Physics course during the WS 22/23 Term.
- [2] W. Commons. File:antiferromagnetic ordering.svg — wikimedia commons, the free media repository, 2020. Online; accessed 23-October-2022.
- [3] W. Commons. File:ferromagnetic ordering.svg — wikimedia commons, the free media repository, 2020. [Online; accessed 24-October-2022].
- [4] N. Goldenfeld. *Lectures On Phase Transitions And The Renormalization Group*. CRC Press.
- [5] M. E. J. Newman, G. T. Barkema, M. E. J. Newman, and G. T. Barkema. *Monte Carlo Methods in Statistical Physics*. Oxford University Press.
- [6] J. Thijssen. *Computational Physics*. Cambridge University Press, 2 edition.