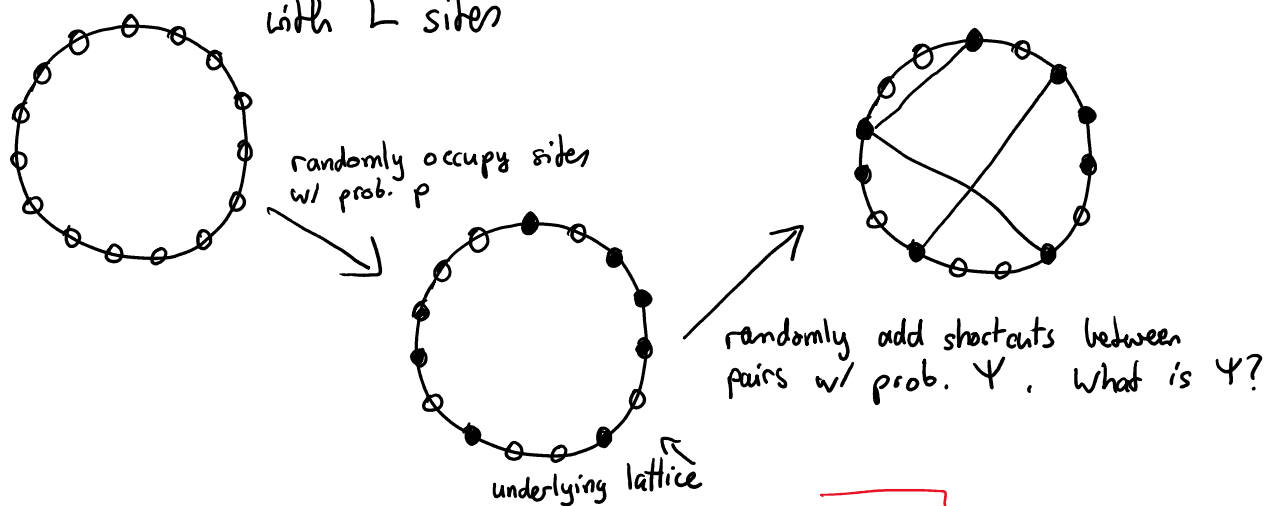


Site percolation

take a periodic 1D lattice

with L sites



$$\phi = \frac{\text{average number of shortcuts}}{\text{number of bonds in underlying lattice}} = \frac{\langle n_s \rangle}{n_b} = \frac{\langle n_s \rangle}{kL}$$

Ψ = prob. of a randomly chosen pair having a shortcut

$$\Psi = 1 - \bar{\Psi} = 1 - \left(1 - \frac{2!}{L^2}\right)^{\langle n_s \rangle}$$

$\bar{\Psi}$ = prob. of not picking this pair when adding the shortcuts = $\left(1 - \frac{2!}{L^2}\right)^{\langle n_s \rangle}$

prob. of picking a given pair

We assume that $L \gg 1$, then

$$\Psi = 1 - (1 - x)^{\langle n_s \rangle} \stackrel{x \rightarrow 0}{\approx} 1 - 1 + \langle n_s \rangle x = \langle n_s \rangle \frac{2}{L^2} = \frac{2\phi k}{L}$$

Taylor Expansion

What is the average number of local clusters of size i :

$$\langle N_i \rangle = \text{prob. of local cluster having size } i \cdot L$$

$$= (1-p)^{2k} \cdot p \cdot (1 - (1-p)^k)^{i-1} \cdot L$$

prob. of having breaks on the two edges

prob. of having one occupied

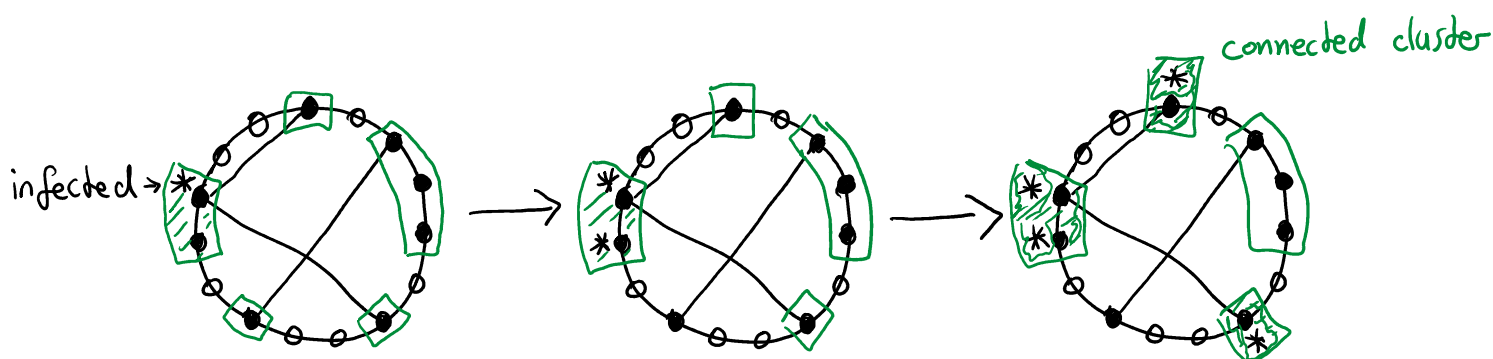
prob. of not having a break $i-1$ times

break local cluster break local cluster

$k=3$

Example

Spreading of the disease



Adding a local cluster of size i to the connected cluster at step n .

$$\vec{v}_n = \begin{pmatrix} v_{n1} \\ v_{n2} \\ \vdots \\ v_{nL} \end{pmatrix} = \begin{pmatrix} \text{prob. of adding a local cluster of size 1 to the connected cluster} \\ \text{" " " " " " " " 2 " " " " " " " " } \\ \vdots \end{pmatrix}$$

for $p \leq p_c$ the prob. increases linearly with a transition Matrix M

$$\vec{v}_{n+1} = M \vec{v}_n \quad \text{where } M_{ij} = \text{number of local clusters of size } i \text{ connected to local clusters of size } j$$

$$\Rightarrow M_{ij} = N_i \cdot (\text{prob. of a size } i \text{ cluster having a shortcut to a } j \text{ size cluster})$$

$$= N_i \cdot (1 - \underbrace{(1 - \Psi)^{ij}}_{\text{prob. of having no shortcut between a pair } i-j \text{ times}})$$

$\vec{v}_{n+1} = M \vec{v}_n$ is an iterative eq.

$$= M M \vec{v}_{n-1} = M^2 \vec{v}_{n-1}$$

$$\vdots \\ = M^n \vec{v}_0$$

express \vec{v}_0 in terms of Eigenvectors \vec{w}_j of M

$$\vec{v}_{n+1} = M^n \sum_{j=1}^L \alpha_j \vec{w}_j = \sum_{j=1}^L \alpha_j \lambda_j^n \vec{w}_j$$

(w/ positive Eigenvector)

Perron-Frobenius theorem $\rightarrow M$ has a unique largest Eigenvalue λ_{\max} and for large n , we have

$$\vec{v}_{n+1} \approx \lambda_{\max}^n \cdot \alpha_{\max} \vec{w}_{\max}$$

if $\lambda_{\max} < 1 \Rightarrow \vec{v}_{n+1} \rightarrow 0$ no spreading of disease

if $\lambda_{\max} > 1 \Rightarrow \vec{v}_{n+1}$ grows exponentially in $\lambda_{\max} \rightarrow$ Epidemic

therefore

$\lambda_{\max} = 1$ corresponds to the critical value at the separation between an epidemic and no epidemic.

Remember $\Psi = \frac{2\phi k}{L}$ for $L \gg 1$

If we fix ϕ and let $L \rightarrow \infty$, then $\Psi \rightarrow 0$ and we obtain via Taylor expansion

$$M_{ij} = N_i (1 - (1 - \Psi)^{ij}) \approx N_i (1 - 1 + ij\Psi) = N_i ij\Psi$$

$$\Rightarrow M_{ij} \text{ can be written as } M = \vec{a} \overset{\text{outer product}}{\circ} \vec{b}^T \text{ with } \vec{a} = \begin{pmatrix} N_1 1 \\ N_2 2 \\ \vdots \end{pmatrix} \Psi \text{ and } \vec{b} = \begin{pmatrix} 1 \\ 2 \\ \vdots \end{pmatrix}$$

$\Rightarrow M$ has rank 1 \Rightarrow one non-zero Eigenvalue

$\Rightarrow M$ has rank 1 \Rightarrow one non-zero Eigenvalue

since $\vec{a} \otimes \vec{b}^T = \vec{a} \vec{b}^T = \begin{pmatrix} b_1 \vec{a} & b_2 \vec{a} & \dots & b_L \vec{a} \end{pmatrix} \Rightarrow$ all columns are linearly dependent

The eigenvector corresponding to the non-zero eigenvalue is \vec{a} since

$$\begin{pmatrix} b_1 \vec{a} & b_2 \vec{a} & \dots & b_L \vec{a} \end{pmatrix} \cdot \vec{a} = b_1 a_1 \vec{a} + b_2 a_2 \vec{a} + \dots = \underbrace{(\vec{a} \cdot \vec{b})}_{=\lambda} \vec{a}$$

So we have the Eigenvector \vec{a} with $a_i = \Psi N_i$ to the

only non-zero Eigenvalue $\lambda = \vec{a} \cdot \vec{b} = \Psi \sum_{j=1}^L N_j j^2$

we insert $N_j = (1-p)^{2k} p (1-(1-p)^k)^{j-1} L$

and $\Psi = \frac{2\phi k}{L}$ for large L

$$\Rightarrow \lambda = \frac{2\phi k}{L} (1-p)^{2k} p \sum_{j=1}^L j^2 q^{j-1} \quad \text{where we set } q = 1 - (1-p)^k$$

$$= 2\phi k (1-q)^2 p \sum_{j=1}^L j^2 q^{j-1} \quad \text{using geometric series (and "Mathematica")}$$

$$\stackrel{L \rightarrow \infty}{\approx} 2\phi k \cancel{(1-q)^2} p \frac{1+q}{(1-q)^3}$$

$$= 2\phi k p \frac{1 + 1 - (1-p)^k}{1 - 1 + (1-p)^k} = 2\phi k p \frac{2 - (1-p)^k}{(1-p)^k}$$

$$\Rightarrow \lambda = 2\phi k p \frac{2 - (1-p)^k}{(1-p)^k}$$

$\lambda = 1$ corresponds to the critical probability p_c

$$\Rightarrow 1 = 2\phi k p_c \frac{2 - (1-p_c)^k}{(1-p_c)^k}$$