

QM3

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1 Tensors: Introduction

Tensors are used to describe how quantities change on coordinate transforms. Consider a vector(?) space mapped by the coordinate system $\{x^1, x^2 \dots\}$ that smoothly maps to a new set of coordinates $\{\bar{x}^1, \bar{x}^2 \dots\}$.

If both spaces are N dimensional, we can write

$$dx^i = \sum_{\alpha=1}^N \frac{\partial x^i}{\partial \bar{x}^\alpha} d\bar{x}^\alpha$$
$$d\bar{x}^i = \sum_{\alpha=1}^N \frac{\partial \bar{x}^i}{\partial x^\alpha} dx^\alpha$$

In the Einstein summation notation we can skip the \sum sign and sum over all possible values of repeated (dummy) indices. All the free indices give us a tensor quantity. For example

$$\frac{dx^i}{dx^j} = \delta_j^i$$
$$\frac{d\bar{x}^\alpha}{d\bar{x}^\beta} = \delta_\beta^\alpha$$

Where δ is the usual kronecker delta tensor.

The way mathematical objects 'transform' is used to classify them:

Rank 1 contravariant tensor

A group of N functions that depend on coordinates A^i is a contravariant tensor if it transforms on changing coordinates as

$$\bar{A}^\alpha = \frac{\partial \bar{x}^\alpha}{\partial x^j} A^j$$

Observe that on multiplying by $\frac{\partial x^i}{\partial \bar{x}^\alpha}$ and summing over all α we get

$$\frac{\partial x^i}{\partial \bar{x}^\alpha} \bar{A}^\alpha = \frac{\partial x^i}{\partial \bar{x}^\alpha} \frac{\partial \bar{x}^\alpha}{\partial x^j} A^j$$
$$\frac{\partial x^i}{\partial \bar{x}^\alpha} \bar{A}^\alpha = \delta_j^i A^j$$
$$\frac{\partial x^i}{\partial \bar{x}^\alpha} \bar{A}^\alpha = A^i$$

Using the chain rule expansion to show $\frac{\partial x^i}{\partial x^j} = \frac{\partial x^i}{\partial \bar{x}^\alpha} \frac{\partial \bar{x}^\alpha}{\partial x^j}$. This shows us the transformation is consistent for both coordinate changes $\{x^1, x^2 \dots\} \rightarrow \{\bar{x}^1, \bar{x}^2 \dots\}$ and vice versa.

Rank 1 covariant tensor

A group of N functions that depend on coordinates A_i is a covariant tensor if it transforms on changing coordinates as

$$\begin{aligned}\bar{A}_\alpha &= \frac{\partial x^j}{\partial \bar{x}^\alpha} A_j \\ A_i &= \frac{\partial \bar{x}^\alpha}{\partial x^i} \bar{A}_\alpha\end{aligned}$$

The gradient of a scalar field for example is a covariant vector. Let $\phi(\{x^i\}) = \phi(\{\bar{x}^i\})$ be a scalar field that is left unchanged on transforming to a new coordinate system. Its gradient is given by a covariant tensor.

$$A_i = \frac{\partial \phi}{\partial x^i}$$

It is easy to observe on transforming coordinates, using the chain rule that

$$\begin{aligned}\bar{A}_\alpha &= \frac{\partial \phi}{\partial \bar{x}^\alpha} \\ \bar{A}_\alpha &= \frac{\partial x^i}{\partial \bar{x}^\alpha} \frac{\partial \phi}{\partial x^i} \\ \bar{A}_\alpha &= \frac{\partial x^i}{\partial \bar{x}^\alpha} A_i\end{aligned}$$

Thus verifying that it is a covariant tensor.

General tensor

A is a general mixed tensor if it transforms as follows:

$$\bar{A}_{\beta_1 \beta_2 \dots \beta_q}^{\alpha_1 \alpha_2 \dots \alpha_p} = \frac{\partial \bar{x}^{\alpha_1}}{\partial x^{i_1}} \frac{\partial \bar{x}^{\alpha_2}}{\partial x^{i_2}} \dots \frac{\partial x^{j_1}}{\partial \bar{x}^{\beta_1}} \dots A_{j_1 j_2 \dots j_q}^{i_1 i_2 \dots i_p}$$

2 tensors can only be added or subtracted if they are of the same type ie they both have the same no of co and contravariant indices