

# Linear Stability Analysis

## One-Layer Shallow Water

In this subsection we consider the one-layer reduced gravity RSW model with topography below. We define the following:

- $H$ : mean depth of the layer
- $z = \eta$ : height of the free surface
- $z = -H + \eta_B$ : height of the topography.
- $h = H + \eta - \eta_B$ : total depth of layer
- $(u, v)$ : horizontal velocity
- $g'$ : reduced gravity
- $\rho_0$ : reference density

The governing nonlinear equations are,

$$\begin{aligned}\frac{\partial u}{\partial t} + \vec{u} \cdot \vec{\nabla} u - fv &= -g \frac{\partial}{\partial x} (h + \eta_B), \\ \frac{\partial v}{\partial t} + \vec{u} \cdot \vec{\nabla} v + fu &= -g \frac{\partial}{\partial y} (h + \eta_B), \\ \frac{\partial h}{\partial t} + \vec{\nabla} \cdot (h \vec{u}_1) &= 0.\end{aligned}$$

## Basic State

To study shear flows in a meridional channel we consider solutions of the form,

$$\begin{aligned}u &= U_B(y), \\ v &= 0, \\ h &= H_B(y).\end{aligned}$$

For this to be an exact solution we require that the flow is in geostrophic balance,

$$fU_B = -g \frac{d}{dy} (H_B + \eta_B).$$

## Perturbation

We perturb the basic state with infinitesimal quantities,

$$\begin{aligned}u &= U_B(y) + u', \\ v &= 0 + v', \\ h &= H_B(y) + h'.\end{aligned}$$

We substitute our perturbation into the governing equations and drop the primes (for brevity) and cancelling out the geostrophic terms

$$\begin{aligned}\frac{\partial u}{\partial t} + (u + U_B) \frac{\partial u}{\partial x} + v \frac{\partial}{\partial y} (u + U_B) - fv &= -g \frac{\partial h}{\partial x}, \\ \frac{\partial v}{\partial t} + (u + U_B) \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + fu &= -g \frac{\partial h}{\partial y}, \\ \frac{\partial h}{\partial t} + (u + U_B) \frac{\partial h}{\partial x} + v \frac{\partial}{\partial y} (H_B + h) + (H_B + h) \left( \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) &= 0.\end{aligned}$$

Now we neglect the quadratic terms to obtain the linearized equations,

$$\begin{aligned}\frac{\partial u}{\partial t} &= -U_B \frac{\partial u}{\partial x} + \left(f - \frac{dU_B}{dy}\right) v - g \frac{\partial h_1}{\partial x}, \\ \frac{\partial v}{\partial t} &= -f u - U_B \frac{\partial v}{\partial x} - g \frac{\partial h_1}{\partial y}, \\ \frac{\partial h}{\partial t} &= -H_B \frac{\partial u}{\partial x} - v \frac{dH_B}{dy} - H_B \frac{\partial v}{\partial y} - U_B \frac{\partial h}{\partial x}.\end{aligned}$$

Finally, we assume a normal mode decomposition in the zonal direction and time,

$$[u, v, h] = \text{Re} \left\{ e^{ik(x-ct)} [\hat{u}, ik\hat{v}, \hat{h}] \right\}, \quad (1)$$

which we can substitute into the above equations to yield in the inviscid limit

$$\begin{aligned}c\hat{u} &= U_B \hat{u} - \left(f - \frac{dU_B}{dy}\right) \hat{v} + g\hat{h}, \\ c\hat{v} &= -\frac{f}{k^2} \hat{u} + U_B \hat{v} - \frac{g}{k^2} \frac{d\hat{h}}{dy}, \\ c\hat{h} &= H_B \hat{u} + \frac{d}{dy} (H_B \hat{v}) + U_B \hat{h}.\end{aligned}$$