
PyRsw Documentation

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Contents:

PLAN

List or problems to solve

1. solve 1d 1-layer SW (periodic/boundaries)
2. solve 1d n-layer SW
3. solve 2d 1-layer SW
4. solve 2d n-layer SW
5. Linear Stability Calculations

Equations and latex to include:

1. derivation of SW
2. governing equations for different models (see above)
3. different numerical methods:
 - (a) FD Sardouney
 - (b) spectral
 - (c) WENO
 - (d) f2py
 - (e) openmp
 - (f) openmp + mpi

ONE-LAYER ROTATING SHALLOW WATER MODEL

2.1 Classical Form

The one-layer, two-dimensional rotating shallow water model on a rotating f -plane can be written as,

$$\begin{aligned}\frac{\partial u}{\partial t} + (\mathbf{u} \cdot \nabla) u - f v &= -g \frac{\partial h}{\partial x}, \\ \frac{\partial v}{\partial t} + (\mathbf{u} \cdot \nabla) v + f u &= -g \frac{\partial h}{\partial y}, \\ \frac{\partial h}{\partial t} + \nabla \cdot (h \mathbf{u}) &= 0.\end{aligned}$$

This describes the motion of a pancake like fluid in that it is thin in the vertical and much longer in the horizontal. It contains pressure forces due to the free surface and a Coriolis pseudo-force because of the rotating frame of reference.

The fluid moves as columns that can be translated in the horizontal and stretched/contracted in the vertical. If the height of a column changes then the vorticity must change, as can be reflected in the fact that, in the absence of forcing and dissipation, Potential Vorticity is conserved following the motion,

$$\frac{D}{Dt} \left(\frac{\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} + f}{h} \right) = 0.$$

2.2 Conservation Form

There are a variety of forms, these equations can be written. One of them is conservation where the three fields that have time derivatives are $U = hu$, $V = hv$ and h .

$$\begin{aligned}\frac{\partial U}{\partial t} + \frac{\partial}{\partial x} \left(\frac{U^2}{h} + \frac{gh^2}{2} \right) + \frac{\partial}{\partial y} \left(\frac{UV}{h} \right) - fV &= 0, \\ \frac{\partial V}{\partial t} + \frac{\partial}{\partial x} \left(\frac{UV}{h} \right) + \frac{\partial}{\partial y} \left(\frac{V^2}{h} + \frac{gh^2}{2} \right) + fU &= 0, \\ \frac{\partial h}{\partial t} + \nabla \cdot (h \mathbf{u}) &= 0.\end{aligned}$$

Conservation form is attractive because it can help to ensure that, using a clever numerical scheme, some quantities are conserved. Note that if you have topography then there are source terms that appear on the right-hand side and this causes some problems.

2.3 Vorticity-Bernoulli Form

A third form arises from rewriting the nonlinear acceleration terms as a gradient and a cross product term. Using a vector identity, it can be shown that the above system is mathematically equivalent to the following,

$$\begin{aligned}\frac{\partial u}{\partial t} - qhv &= -\frac{\partial B}{\partial x}, \\ \frac{\partial v}{\partial t} + qhu &= -\frac{\partial B}{\partial y}, \\ \frac{\partial h}{\partial t} + \nabla(h\mathbf{u}) &= 0.\end{aligned}$$

Note that above we have defined the vorticity and the Bernoulli function,

$$\begin{aligned}q &= \frac{\zeta + f}{h} = \frac{\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} + f}{h}, \\ B &= gh + \frac{1}{2}(u^2 + v^2).\end{aligned}$$

Before we look at solving this complicated set of equations we consider the one and a half dimensional limit.

2.4 $1\frac{1}{2}$ Dimensional Limit

We assume that none of the variables depend on one horizontal direction, say y . But, it is very important to realize that the velocity in that direction is not necessarily zero. Indeed, if you have flow in the x -direction, the Coriolis force will deflect it to the right, which will then generate a flow that is perpendicular. This will continue and often give rise to inertial oscillations in the horizontal. So we can have motion in either direction but the motion only changes with respect to x .

If we simplify the governing equations we get

$$\begin{aligned}\frac{\partial u}{\partial t} - qhv &= -\frac{\partial B}{\partial x}, \\ \frac{\partial v}{\partial t} + qhu &= 0, \\ \frac{\partial h}{\partial t} + \frac{\partial}{\partial x}(hu) &= 0.\end{aligned}$$

where the vorticity and the Bernoulli function simplify to,

$$\begin{aligned}q &= \frac{\zeta + f}{h} = \frac{\frac{\partial v}{\partial x} + f}{h}, \\ B &= gh^2 + \frac{1}{2}(u^2 + v^2).\end{aligned}$$

2.5 Conserved Quantities

In the purely conservative or nondissipative limit, there are three quantities that are exactly conserved.

1. Mass:

$$M = \int_D h dA$$

2. Total Energy: sum of potential and kinetic energies

$$E = \frac{1}{2} \int_D (gh^2 + h(u^2 + v^2)) \, dA$$

3. Potential Enstrophy:

$$Q = \frac{1}{2} \int_D h q^2 \, dA$$

In a numerical model we cannot expect these to be conserved but we would like them to be close to be conserved. If they are very badly conserved then this could reflect that the numerical scheme is behaving badly. However, just because these are conserved that does not guarantee that the solution is correct. But they are usually good indicators as to how our method is doing in the conservative limit.

Of course when nonconservative forces are introduced things will change. One might argue that since the world is non-dissipative then we don't need to worry about conserving these. However, it is desirable to know that basis of your model is well behaved and therefore why we should worry about conserved quantities.

INDICES AND TABLES

- `genindex`
- `modindex`
- `search`