

RB-SFA: High Harmonic Generation in the Strong Field Approximation via *Mathematica*

© Emilio Pisanty 2014-2016. Licensed under GPL and CC-BY-SA.

■ Introduction

Readme

RB-SFA is a compact and flexible *Mathematica* package for calculating High Harmonic Generation emission within the Strong Field Approximation. It combines *Mathematica*'s analytical integration capabilities with its numerical calculation capacities to offer a fast and user-friendly plug-and-play solver for calculating harmonic spectra and other properties. In addition, it can calculate first-order nondipole corrections to the SFA results to evaluate the effect of the driving laser's magnetic field on harmonic spectra.

The name RB-SFA comes from its first application (as Rotating Bicircular High Harmonic Generation in the Strong field Approximation) but the code is general so RB-SFA just stands for itself now. This first application was used to calculate the polarization properties of the harmonics produced by multi-colour circularly polarized fields, as reported in the paper

Spin conservation in high-order-harmonic generation using bicircular fields. E. Pisanty, S. Sukiasyan and M. Ivanov. *Phys. Rev. A* **90**, 043829 (2014), arXiv:1404.6242.

This code is dual-licensed under the GPL and CC-BY-SA licenses. If you use this code or its results in an academic publication, please cite the paper above or the GitHub repository where the latest version will always be available. An example citation is

E. Pisanty. RB-SFA: High Harmonic Generation in the Strong Field Approximation via *Mathematica*. <https://github.com/episanty/RB-SFA> (2014).

This software consists of this notebook, which contains the code and its documentation, a corresponding auto-generated package file. This notebook also contains a Usage and Examples section which explains how to use the code and documents the calculations used in the original publication.

Specifications

This code calculates the harmonic dipole as per the original Lewenstein *et al.* paper,

M. Lewenstein, Ph. Balcou, M.Yu. Ivanov, A. L'Huillier and P.B. Corkum. Theory of high-harmonic generation by low-frequency laser fields. *Phys. Rev. A* **49** no. 3, 2117 (1994).

The time-dependent dipole is given by

$$\mathbf{D}(t) = i \int_0^t dt' \int d^3 \mathbf{p} \, \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}, t)) \times \mathbf{F}(t') \cdot \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}, t')) \times \exp[-i S(\mathbf{p}, t, t')] + c.c.$$

Here $\mathbf{F}(t)$ is the time-dependent applied electric field, with vector potential $\mathbf{A}(t)$; the electron charge is taken to be -1 . The integration time t' can be interpreted as the ionization time and the integration limit t represents the recollision time. (Upon applying the saddle-point approximation to the temporal integrals these notions apply exactly to the

corresponding quantum orbits.) The ionization and recollision events are governed by the corresponding dipole matrix elements, which are taken for a hydrogenic 1s-type orbital with characteristic momentum κ and ionization potential $I_p = \frac{1}{2} \kappa^2$:

$$\mathbf{d}(\mathbf{p}) = \langle \mathbf{p} | \hat{\mathbf{d}} | g \rangle = \frac{8i}{\pi} \frac{\sqrt{2\kappa^5} \mathbf{p}}{(\mathbf{p}^2 + \kappa^2)^3}.$$

(Taken from B. Podolsky & L. Pauling. *Phys. Rev.* **34** no. 1, 109 (1929).) The displaced momentum $\boldsymbol{\pi}(\mathbf{p}, t) = \mathbf{p} + \mathbf{A}(t)$ is used to aid in the generalization to nondipole cases.

The action

$$S(\mathbf{p}, t, t') = \int_{t'}^t \left(\frac{1}{2} \kappa^2 + \frac{1}{2} \boldsymbol{\pi}(\mathbf{p}, \tau)^2 \right) d\tau$$

is the governing factor of the integral. It is highly oscillatory and must be dealt with carefully. The momentum integral is approximated using the saddle point method; this is valid and straightforward since there is always one unique saddle point,

$$\mathbf{p}_s(t, t') = - \frac{\int_{t'}^t \mathbf{A}(\tau) d\tau}{t - t'},$$

as long as the action's dependence on the momentum is quadratic. After the momentum has been integrated via the saddle point method, the time-dependent dipole is given by

$$\mathbf{D}(t) = i \int_0^t dt' \left(\frac{2\pi}{\epsilon + i(t-t')} \right)^{3/2} \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}_s(t, t'), t)) \times \mathbf{F}(t') \cdot \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}_s(t, t'), t')) \times \exp[-i S(\mathbf{p}_s(t, t'), t, t')] + c.c.$$

A regularization factor ϵ has been introduced to prevent divergence of the prefactor as t' approaches t , representing the fact that at $t = t'$ the \mathbf{p} integral no longer has a quadratic exponent.

The integral is performed with respect to the excursion time $\tau = t - t'$ for simplicity. To reduce integration time, a gating function $\text{gate}[\omega \tau]$ has been introduced, and the integration time has been cut short to a controllable length. This is physically reasonable as the wavepacket spreading (which scales as $\tau^{-3/2}$) makes the contributions after the gate negligible.

More physically, suppressing the contributions from large excursion times represents the fact that the extra-long trajectories they represent produce harmonics which are much harder to phase match (as their harmonic phase is very sensitively dependent on many factors) which means that they are rarely present in experimental spectra unless a concerted effort is performed to observe them.

Quantum-orbit functions suite

The package also includes an experimental suite of functions for calculating spectra via quantum-orbit methods. That is, the goal is to calculate the Fourier transform directly,

$$\mathbf{D}(\Omega) = i \int_{-\infty}^{\infty} dt \int_{-\infty}^t dt' \left(\frac{2\pi}{\epsilon + i\tau} \right)^{3/2} \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}_s(t, t - t'), t)) \times \mathbf{F}(t') \cdot \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}_s(t, t'), t')) \times \exp[-i S(\mathbf{p}_s(t, t'), t, t')] \times e^{+i\Omega t},$$

via a saddle-point approximation on the ionization and recombination times t' and t , which gives a harmonic dipole of the form

$$D(\Omega) = i \sum_s \sqrt{\frac{(2\pi)^2}{i^2 \det(\partial_{(t,t')}^2 S(t, t'_s))}} \left(\frac{2\pi}{\epsilon + i(t_s - t'_s)} \right)^{3/2} \times \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}_s(t_s, t_s - t'_s), t_s)) \times \mathbf{F}(t'_s) \cdot \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}_s(t_s, t'_s), t'_s)) \times \exp[-i S(\mathbf{p}_s(t_s, t'_s), t_s, t'_s)] \times e^{+i\Omega t_s}$$

where the times t_s and t'_s (and equivalently $\tau_s = t_s - t'_s$) are obtained via the saddle-point equations

$$\frac{\partial S}{\partial t'}(t_s, t'_s) = 0, \quad \frac{\partial S}{\partial t}(t_s, t'_s) = \Omega,$$

as described in detail in the original paper by Lewenstein et al. (*Phys. Rev. A* **49** no. 3, 2117 (1994)), or in e.g. M. Ivanov and O. Smirnova. *Multielectron high harmonic generation: simple man on a complex plane*. In T. Schultz and M. Vrakking (eds.), *Attosecond and XUV Physics: Ultrafast Dynamics and Spectroscopy*, pp. 201–256 (Wiley-VCH, Weinheim, 2014), arXiv:1304.2413.

First-order nondipole calculations

This code can also be used to calculate the first-order effect of nondipole terms, which become relevant at high intensities or long wavelengths, when the electron's speed from the laser-driven oscillations (which scales as F/ω) becomes comparable to the speed of light (equal to $c = 1/\alpha \approx 137$ in atomic units). (More specifically, the nondipole effects on the harmonic generation become relevant when the magnetic pushing per cycle, which offsets the returning wavepacket by $F^2/c\omega^3$, becomes comparable with the wavepacket spread upon recollision.)

This code implements a generalization of the theory presented in various forms in the papers N.J. Kylstra et al. *J. Phys B: At. Mol. Opt. Phys.* **34** no. 3, L55 (2001); N.J. Kylstra et al. *Laser Phys.* **12** no. 2, 409 (2002); M.W. Walser et al. *Phys. Rev. Lett.* **85** no. 24, 5082 (2000); V. Averbukh et al. *Phys Rev. A* **65**, 063402 (2002); and C.C. Chirilă et al. *Phys. Rev. A* **66**, 063411 (2002).

The nondipole-HHG code calculates the harmonic dipole of exactly the same form as the standard SFA,

$$D(t) = i \int_0^t dt' \int d^3\mathbf{p} \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}, t, t')) \times \mathbf{F}(t') \cdot \mathbf{d}(\boldsymbol{\pi}(\mathbf{p}, t', t')) \times \exp[-i S(\mathbf{p}, t, t')] + c.c.$$

with $S(\mathbf{p}, t, t') = \int_t^{t'} dt'' \left(\frac{1}{2} \kappa^2 + \frac{1}{2} \boldsymbol{\pi}(\mathbf{p}, t'', t')^2 \right)$ as before, but with a nondipole term in the displaced momentum, which now reads

$$\boldsymbol{\pi}(\mathbf{p}, t, t') = \mathbf{p} + \mathbf{A}(t) - \int_t^{t'} \nabla \mathbf{A}(\tau) \cdot (\mathbf{p} + \mathbf{A}(\tau)) d\tau,$$

or in component notation

$$\pi_j(\mathbf{p}, t, t') = p_j + A_j(t) - \int_t^{t'} \partial_j A_k(\tau) (p_k + A_k(\tau)) d\tau.$$

■ Implementation

Supporting functions

Initialization

```
BeginPackage["RBSFA`"];
```

Version number

The command `RBSFAversion` prints the version of the RB-SFA package currently loaded and its timestamp

```
$RBSFAversion::usage = "$RBSFAversion prints the
  current version of the RB-SFA package in use and its timestamp.";
$RBSFAtimestamp::usage = "$RBSFAtimestamp prints the timestamp of
  the current version of the RB-SFA package.";
Begin["`Private`"];
$RBSFAversion := "RB-SFA v2.1.0, " <> $RBSFAtimestamp;
End[];
```

Old syntax (in functional form `RBSFAversion[]`), deprecated

```
RBSFAversion::usage = "RBSFAversion[] has been deprecated in favour of $RBSFAversion.";
RBSFAversion::dprc = "RBSFAversion[] has been deprecated in favour of $RBSFAversion.";
Begin["`Private`"];
RBSFAversion[] := (Message[RBSFAversion::dprc]; $RBSFAversion);
End[];
```

The timestamp is updated every time the notebook is saved via an appropriate notebook option, which is set by the code below.

```
SetOptions[
  EvaluationNotebook[],
  NotebookEventActions → {{"MenuCommand", "Save"} → (
    NotebookWrite[
      Cells[CellTags → "version-timestamp"][[1]],
      Cell[
        BoxData[RowBox[
          {"Begin[\"`Private`\"]; \n$RBSFAtimestamp= \"\" <> DateString[] <> \"\"; \nEnd[];"}]]],
        "Input", InitializationCell → True, CellTags → "version-timestamp"
      ], None, AutoScroll → False];
    NotebookSave[]
  ), PassEventsDown → True}
];
```

To reset this behaviour to normal, evaluate the cell below

```
SetOptions[EvaluationNotebook[],
  NotebookEventActions → {{"MenuCommand", "Save"} → (NotebookSave[], PassEventsDown → True)}]
```

Timestamp

```
Begin["`Private`"];
$RBSFAtimestamp = "Wed 11 May 2016 15:26:26";
End[];
```

Git commit hash and message

```

$RBSFAdirectory::usage = "$RBSFAdirectory is the
    directory where the current RB-SFA package instance is located.";
$RBSFACommit::usage = "$RBSFACommit returns the git commit log at
    the location of the RB-SFA package if there is one.";
$RBSFACommit::OS = "$RBSFACommit has only been tested on Linux.";

Begin["`Private`"];
With[{softLinkTestString = StringSplit[StringJoin[ReadList[
    "! ls -la " <> StringReplace[$InputFileName, {" " -> "\\ "}], String]], " -> "]},
    If[Length[softLinkTestString] > 1, (*Testing in case $InputFileName
        is a soft link to the actual directory.*)
        $RBSFAdirectory = StringReplace[DirectoryName[softLinkTestString[[2]], {" " -> "\\ "}],
        $RBSFAdirectory = StringReplace[DirectoryName[$InputFileName], {" " -> "\\ "};
    ]]
$RBSFACommit := (If[$OperatingSystem != "Unix", Message[$RBSFACommit::OS];
    StringJoin[
        Riffle[ReadList["!cd " <> $RBSFAdirectory <> " && git log -1", String], {"\n"}]]];
End[];

```

Dipole transition matrix elements

Default DTME is for a hydrogenic 1s state.

```

hydrogenicDTME::usage =
    "hydrogenicDTME[p,κ] returns the dipole transition matrix element for
    a 1s hydrogenic state of ionization potential  $I_p = \frac{1}{2}\kappa^2$ .";
hydrogenicDTMERegularized::usage = "hydrogenicDTMERegularized[p,κ] returns
    the dipole transition matrix element for a 1s hydrogenic state of
    ionization potential  $I_p = \frac{1}{2}\kappa^2$ , regularized to remove the denominator
    of  $1/(p^2 + \kappa^2)^3$ , where the saddle-point solutions are singular."
Begin["`Private`"];

```

$$\text{hydrogenicDTME}[p_List, \kappa_] := \frac{8 \, i}{\pi} \frac{\sqrt{2} \, \kappa^5 \, p}{(\text{Total}[p^2] + \kappa^2)^3}$$

$$\text{hydrogenicDTME}[p_?NumberQ, \kappa_] := \frac{8 \, i}{\pi} \frac{\sqrt{2} \, \kappa^5 \, p}{(p^2 + \kappa^2)^3}$$

$$\text{hydrogenicDTMERegularized}[p_List, \kappa_] := \frac{8 \, i}{\pi} \frac{\sqrt{2} \, \kappa^5 \, p}{1}$$

$$\text{hydrogenicDTMERegularized}[p_?NumberQ, \kappa_] := \frac{8 \, i}{\pi} \frac{\sqrt{2} \, \kappa^5 \, p}{1}$$
End[];

hydrogenicDTMERegularized[p,κ] returns the dipole transition matrix element for
 a 1s hydrogenic state of ionization potential $I_p = \frac{1}{2}\kappa^2$, regularized to remove
 the denominator of $1/(p^2 + \kappa^2)^3$, where the saddle-point solutions are singular.

Another possibility is for a gaussian orbital

```

gaussianDTME::usage = "gaussianDTME[p,κ] returns the dipole transition
    matrix element for a gaussian state of characteristic size 1/κ.";
Begin["`Private`"];

gaussianDTME[p_List, κ_] := -i (4 π)3/4 κ-7/2 p Exp[- $\frac{\text{Total}[p^2]}{2 \kappa^2}$ ]

gaussianDTME[p_?NumberQ, κ_] := -i (4 π)3/4 κ-7/2 p Exp[- $\frac{p^2}{2 \kappa^2}$ ]

End[];

```

Various field envelopes

flatTopEnvelope

```

flatTopEnvelope::usage =
    "flatTopEnvelope[ω,num,nRamp] returns a Function object representing
    a flat-top envelope at carrier frequency ω lasting a total
    of num cycles and with linear ramps nRamp cycles long.";
Begin["`Private`"];
flatTopEnvelope[ω_, num_, nRamp_] := Function[t,
    Piecewise[{
        {0, t < 0}, {Sin[ $\frac{\omega t}{4 \text{nRamp}}$ ]2, 0 ≤ t <  $\frac{2 \pi}{\omega} \text{nRamp}$ }, {1,  $\frac{2 \pi}{\omega} \text{nRamp} \leq t < \frac{2 \pi}{\omega} (\text{num} - \text{nRamp})$ },
        {Sin[ $\frac{\omega (\frac{2 \pi}{\omega} \text{num} - t)}{4 \text{nRamp}}$ ]2,  $\frac{2 \pi}{\omega} (\text{num} - \text{nRamp}) \leq t < \frac{2 \pi}{\omega} \text{num}$ }, {0,  $\frac{2 \pi}{\omega} \text{num} \leq t$ }
    }]]
End[];

```

cosPowerFlatTop

```

cosPowerFlatTop::usage =
    "cosPowerFlatTop[ω,num,power] returns a Function object representing
    a smooth flat-top envelope of the form 1-Cos(ω t/2 num)power";
Begin["`Private`"];
cosPowerFlatTop[ω_, num_, power_] := Function[t, 1 - Cos[ $\frac{\omega t}{2 \text{num}}$ ]power]
End[];

```

Field duration standard options

The standard options for the duration of the pulse and the resolution are

```

PointsPerCycle::usage =
    "PointsPerCycle is a sampling option which specifies the number of sampling
    points per cycle to be used in integrations.";
TotalCycles::usage = "TotalCycles is a sampling option which specifies
    the total number of periods to be integrated over.";
CarrierFrequency::usage = "CarrierFrequency is a sampling option
    which specifies the carrier frequency to be used.";
Protect[PointsPerCycle, TotalCycles, CarrierFrequency];

standardOptions = {PointsPerCycle → 90, TotalCycles → 1, CarrierFrequency → 0.057};

PointsPerCycle dictates how many sampling points are used per laser cycle (at frequency CarrierFrequency, of the

```

infrared laser), and it should be at least twice the highest harmonic of interest. The total duration is TotalCycles cycles. CarrierFrequency is the frequency of the fundamental laser, in atomic units.

harmonicOrderAxis

harmonicOrderAxis produces a list that can be used as a harmonic order axis for the given pulse parameters.

The length can be fine-tuned (to match exactly a spectrum, for instance, and get a matrix of the correct shape) using the correction option, or a TargetLength can be directly specified.

```
harmonicOrderAxis::usage =
  "harmonicOrderAxis[opt→value] returns a list of frequencies which can be
    used as a frequency axis for Fourier transforms, scaled in units of
    harmonic order, for the provided field duration and sampling options.";
TargetLength::usage = "TargetLength is an option for harmonicOrderAxis
  which specifies the total length required of the resulting list.";
LengthCorrection::usage = "LengthCorrection is an option for harmonicOrderAxis
  which allows for manual correction of the length of the resulting list.";
Protect[LengthCorrection, TargetLength];
Begin["`Private`"];
Options[harmonicOrderAxis] =
  standardOptions~Join~{TargetLength → Automatic, LengthCorrection → 1};
harmonicOrderAxis::target =
  "Invalid TargetLength option `1`. This must be a positive integer or Automatic.";
harmonicOrderAxis[OptionsPattern[]] :=
  Module[{num = OptionValue[TotalCycles], npp = OptionValue[PointsPerCycle]},
    Piecewise[{
      {

$$\frac{1}{\text{num}} \text{Range}\left[0., \text{Round}\left[\frac{\text{npp num} + 1}{2.}\right] - 1 + \text{OptionValue[LengthCorrection]}\right],$$

        OptionValue[TargetLength] === Automatic},
      {

$$\frac{\text{Round}\left[\frac{\text{npp num} + 1}{2.}\right]}{\text{num}} \frac{\text{Range}\left[0, \text{OptionValue[TargetLength]} - 1\right]}{\text{OptionValue[TargetLength]}},$$

        IntegerQ[OptionValue[TargetLength]] && OptionValue[TargetLength] ≥ 0}
    },
    Message[harmonicOrderAxis::target, OptionValue["TargetLength"]]; Abort[]
  ]
End[];
```

frequencyAxis

frequencyAxis produces a list that can be used as a harmonic order axis for the given pulse parameters. Identical to harmonicOrderAxis but produces a frequency axis (in atomic units) instead.

```

frequencyAxis::usage =
  "frequencyAxis[opt→value] returns a list of frequencies which can be used
  as a frequency axis for Fourier transforms, in atomic units of
  frequency, for the provided field duration and sampling options.";
Begin["`Private`"];
Options[frequencyAxis] = Options[harmonicOrderAxis];
frequencyAxis[options:OptionsPattern[]] :=
  OptionValue[CarrierFrequency] harmonicOrderAxis[options]
End[];

```

timeAxis

timeAxis produces a list that can be used as a time axis for the given pulse parameters.

Quit

```

timeAxis::usage =
  "timeAxis[opt→value] returns a list of times which can be used as a time axis ";
TimeScale::usage = "TimeScale is an option for timeAxis which specifies the units
  the list should use: AtomicUnits by default, or LaserPeriods if required.";
AtomicUnits::usage = "AtomicUnits is a value for the option TimeScale of timeAxis
  which specifies that the times should be in atomic units of time.";
LaserPeriods::usage = "LaserPeriods is a value for the option TimeScale of timeAxis
  which specifies that the times should be in multiples of the carrier laser period.";
Protect[TimeScale, AtomicUnits, LaserPeriods];
Begin["`Private`"];
Options[timeAxis] =
  standardOptions~Join~{TimeScale→AtomicUnits, PointNumberCorrection→0};
timeAxis::scale =
  "Invalid TimeScale option `1`. Available values are AtomicUnits and LaserPeriods";
timeAxis[OptionsPattern[]] := Block[{T = 2  $\pi$  /  $\omega$ ,  $\omega$  = OptionValue[CarrierFrequency],
  num = OptionValue[TotalCycles], npp = OptionValue[PointsPerCycle]},
  Piecewise[{
    {1, OptionValue[TimeScale] === AtomicUnits},
    { $\frac{1}{T}$ , OptionValue[TimeScale] === LaserPeriods}
  },
  ],
  Message[timeAxis::scale, OptionValue[TimeScale]]; Abort[]
] × Table[t
  , {t, 0, num  $\frac{2 \pi}{\omega}$ ,  $\frac{\text{num}}{\text{num} \times \text{npp} + \text{OptionValue[PointNumberCorrection]}} \frac{2 \pi}{\omega}$ }
]
]
End[];

```



```

tInit = 0;
tFinal =  $\frac{2 \pi}{\omega}$  num;

 $\delta t = \frac{tFinal - tInit}{num \times npp + OptionValue[PointNumberCorrection]}$ ; (*integration and looping timestep*)

```

getSpectrum

getSpectrum takes a time-dependent dipole list and returns its Fourier transform in absolute-value-squared. It takes as options

- pulse parameters ω , TotalCycles and PointsPerCycle,
- a polarization parameter ϵ , which gives an unpolarized spectrum when given False, or polarizes along an ellipticity vector ϵ (this is meant primarily to select right- and left-circularly polarized spectra using $\epsilon = \{1, i\}$ and $\epsilon = \{1, -i\}$ respectively),
- a DifferentiationOrder, which can return the dipole value (default, = 0), velocity (= 1), or acceleration (= 2),
- a power of ω , ω Power, with which to multiply the spectrum before returning it (which should be equivalent to DifferentiationOrder except for pathological cases), and
- a ComplexPart function to apply immediately after differentiation (default is the identity function, but Re, Im, or Abs[#]^2 & are reasonable choices).

If no option is passed to ω Power and DifferentiationOrder, the pulse parameters do not really affect the output, except by a global factor of TotalCycles.

```

getSpectrum::usage = "getSpectrum[DipoleList] returns the power spectrum of DipoleList.";
Polarization::usage =
  "Polarization is an option for getSpectrum which specifies a polarization
  vector along which to polarize the dipole list. The default,
  Polarization->False, specifies an unpolarized spectrum.";
ComplexPart::usage = "ComplexPart is an option for getSpectrum which
  specifies a function (like Re, Im, or by default #&) which should
  be applied to the dipole list before the spectrum is taken.";
 $\omega$ Power::usage = " $\omega$ Power is an option for getSpectrum which specifies
  a power of frequency which should multiply the spectrum.";
DifferentiationOrder::usage = "DifferentiationOrder is an option for
  getSpectrum which specifies the order to which the dipole
  list should be differentiated before the spectrum is taken.";
Protect[Polarization, ComplexPart,  $\omega$ Power, DifferentiationOrder];

Begin["`Private`"];
Options[getSpectrum] = {Polarization -> False, ComplexPart -> (# &),
   $\omega$ Power -> 0, DifferentiationOrder -> 0} ~Join~ standardOptions;

getSpectrum::diffOrd = "Invalid differentiation order `1`.";
getSpectrum:: $\omega$ Pow = "Invalid  $\omega$  power `1`.";

getSpectrum[dipoleList_, OptionsPattern[]] := Block[
  {polarizationVector, differentiatedList, depth, dimensions,
    num = OptionValue[TotalCycles],
    npp = OptionValue[PointsPerCycle],  $\omega$  = OptionValue[CarrierFrequency],  $\delta t = \frac{2 \pi / \omega}{npp}$ 
  },

```

$$\text{polarizationVector} = \frac{\text{OptionValue}[\text{Polarization}]}{\text{Norm}[\text{OptionValue}[\text{Polarization}]}];$$

```

differentiatedList = OptionValue[ComplexPart][Piecewise[{
  {dipoleList, OptionValue[DifferentiationOrder] == 0},
  {1/2 dt (Most[Most[dipoleList]] - Rest[Rest[dipoleList]]),
   OptionValue[DifferentiationOrder] == 1},
  {1/dt^2 (Most[Most[dipoleList]] - 2 Most[Rest[dipoleList]] + Rest[Rest[dipoleList]]),
   OptionValue[DifferentiationOrder] == 2}],
  Message[getSpectrum::diffOrd, OptionValue[DifferentiationOrder]];
  Abort[]
]];

```

```

If[NumberQ[OptionValue[ωPower]], Null, Message[getSpectrum::ωPow, OptionValue[ωPower]];
  Abort[]];

```

```

num Table[
  (ω/num k)^2 OptionValue[ωPower], {k, 1, Round[Length[differentiatedList]/2]}
] × If[
  OptionValue[Polarization] == False, (*unpolarized spectrum*)
  (*funky depth thing so this can take lists of numbers and lists of vectors,
  of arbitrary length. Makes for easier benchmarking.*)
  depth = Length[Dimensions[dipoleList]];
  dimensions = If[Length[#] > 1, #[[2]], 1 (*#[[1]]*)] &[Dimensions[dipoleList]];
  Sum[Abs[
    Fourier[
      If[depth > 1, Re[differentiatedList[[All, i]]], Re[differentiatedList[[All]]]]
      , FourierParameters → {-1, 1}
    ][[1 ;; Round[Length[differentiatedList]/2]]]]
  ]^2, {i, 1, dimensions}]
, (*polarized spectrum*)
Abs[
  Transpose[Table[
    Fourier[
      Re[differentiatedList[[All, i]]]
      , FourierParameters → {-1, 1}
    ]
    , {i, 1, 2}]]][[1 ;; Round[Length[differentiatedList]/2]]].polarizationVector
  ]^2
]
End[];

```

spectrumPlotter

spectrumPlotter takes a spectrum and a list of options and returns a plot of the spectrum. The available options are

- a FrequencyAxis option, which will give the harmonic order as a horizontal axis by default, and an arbitrary scale with any other option,
- all the options of harmonicOrderAxis, which will be passed to the call that makes the horizontal axis, and
- all the options of ListLinePlot, which will be used to format the plot.

```
spectrumPlotter::usage = "spectrumPlotter[spectrum] plots
  the given spectrum with an appropriate axis in a log10 scale.";
FrequencyAxis::usage = "FrequencyAxis is an option for spectrumPlotter
  which specifies the axis to use.";
Protect[FrequencyAxis];
Begin["`Private`"];
Options[spectrumPlotter] = Join[{FrequencyAxis → "HarmonicOrder"},
  Options[harmonicOrderAxis], Options[ListLinePlot]];
spectrumPlotter[spectrum_, options : OptionsPattern[]] := ListPlot[
  {Which[
    OptionValue[FrequencyAxis] === "HarmonicOrder",
    harmonicOrderAxis["TargetLength" → Length[spectrum], Sequence@@
      FilterRules[{options}~Join~Options[spectrumPlotter], Options[harmonicOrderAxis]]],
    OptionValue[FrequencyAxis] === "Frequency",
    frequencyAxis["TargetLength" → Length[spectrum], Sequence@@
      FilterRules[{options}~Join~Options[spectrumPlotter], Options[harmonicOrderAxis]]],
    True, Range[Length[spectrum]]
  ]},
  Log[10, spectrum]
]
, Sequence@@FilterRules[{options}, Options[ListLinePlot]]
, Joined → True
, PlotRange → Full
, PlotStyle → Thick
, Frame → True
, Axes → False
, ImageSize → 800
]
End[];
```

biColorSpectrum

biColorSpectrum takes a time-dependent dipole list and produces overlaid plots of the right- and left-circular components of the spectrum, in red and blue respectively. It takes all the options of getSpectrum and spectrumPlotter, which are passed directly to the corresponding calls, as well as the options of Show, which can be used to modify the plot appearance.

Quit

```

biColorSpectrum::usage =
  "biColorSpectrum[DipoleList] produces a two-colour spectrum of DipoleList,
    separating the two circular polarizations.";
Begin["`Private`"];
Options[biColorSpectrum] = Join[{PlotRange → All}, Options[Show],
  Options[spectrumPlotter], DeleteCases[Options[getSpectrum], Polarization → False]];
biColorSpectrum[dipoleList_, options : OptionsPattern[]] := Show[{
  spectrumPlotter[
    getSpectrum[dipoleList, Polarization → {1, +i},
      Sequence@@FilterRules[{options}, Options[getSpectrum]]],
    PlotStyle → Red, Sequence@@FilterRules[{options}, Options[spectrumPlotter]]],
  spectrumPlotter[
    getSpectrum[dipoleList, Polarization → {1, -i},
      Sequence@@FilterRules[{options}, Options[getSpectrum]]],
    PlotStyle → Blue, Sequence@@FilterRules[{options}, Options[spectrumPlotter]]]
  }, PlotRange → OptionValue[PlotRange],
  Sequence@@FilterRules[{options}, Options[Show]]
]
End[];

```

Various gate functions

Gate functions are used to suppress the contributions of extra-long trajectories with long excursion times, partly to reflect the effect of phase matching but mostly to keep integration times reasonable. They are provided to the main numerical integrator `makeDipoleList` via its `Gate` option.

```

SineSquaredGate::usage =
  "SineSquaredGate[nGateRamp] specifies an integration gate with a sine-squared
    ramp, such that SineSquaredGate[nGateRamp][ωt,nGate]
    has nGate flat periods and nGateRamp ramp periods.";
LinearRampGate::usage = "LinearRampGate[nGateRamp] specifies an integration
  gate with a linear ramp, such that SineSquaredGate[nGateRamp][ωt,nGate]
  has nGate flat periods and nGateRamp ramp periods.";
Begin["`Private`"];
SineSquaredGate[nGateRamp_][ωτ_, nGate_] := Piecewise[{{1, ωτ ≤ 2 π (nGate - nGateRamp)}},
  {Sin[ $\frac{2 \pi nGate - \omega \tau}{4 nGateRamp}$ ]2, 2 π (nGate - nGateRamp) < ωτ ≤ 2 π nGate}, {0, nGate < ωτ}}]
LinearRampGate[nGateRamp_][ωτ_, nGate_] := Piecewise[{{1, ωτ ≤ 2 π (nGate - nGateRamp)}},
  {- $\frac{\omega \tau - 2 \pi (nGate + nGateRamp)}{2 \pi nGateRamp}$ , 2 π (nGate - nGateRamp) < ωτ ≤ 2 π nGate}, {0, nGate < ωτ}}]
End[];

```

getIonizationPotential

```
getIonizationPotential::usage =
  "getIonizationPotential[Target] returns the ionization potential
    of an atomic target, e.g. \"Hydrogen\", in atomic units.

getIonizationPotential[Target,q] returns the ionization
  potential of the q-th ion of the specified Target, in atomic units.

getIonizationPotential[{Target,q}] returns the ionization
  potential of the q-th ion of the specified Target, in atomic units.";
Begin["`Private`"];
getIonizationPotential[Target_, Charge_: 0] :=
  UnitConvert[ElementData[Target, "IonizationEnergies"][[Charge + 1]] /
    (Quantity[1, "AvogadroConstant"] Quantity[1, "Hartrees"])]
getIonizationPotential[{Target_, Charge_: 0}] := getIonizationPotential[Target, Charge]
End[];
```

makeDipoleList: main numerical integrator

The main integration function is `makeDipoleList`, and its basic syntax is of the form `makeDipoleList[VectorPotential→A]`. Here the vector potential `A` must be a function object, such that for numeric `t` the construct `A[t]` returns a list of numbers after the appropriate field parameters have been introduced: thus the criterion is that, for a call of the form `makeDipoleList[VectorPotential→A, FieldParameters→pars]`, a call of the form `A[t]//.pars` returns a list of numbers for numeric `t`. To see the available options use `Options[makeDipoleList]`, and to get information on each option use the `?VectorPotential` construct.

```
makeDipoleList::usage = "makeDipoleList[VectorPotential→A]
  calculates the dipole response to the vector potential A.";

VectorPotential::usage =
  "VectorPotential is an option for makeDipole list which specifies the field's vector
    potential. Usage should be VectorPotential→A, where A[t]//.pars must yield a
    list of numbers for numeric t and parameters indicated by FieldParameters→pars.";
VectorPotentialGradient::usage = "VectorPotentialGradient is an option for makeDipole
  list which specifies the gradient of the field's vector potential. Usage should be
  VectorPotentialGradient→GA, where GA[t]//.pars must yield a square matrix of the
  same dimension as the vector potential for numeric t and parameters indicated by
  FieldParameters→pars. The indices must be such that GA[t][[i,j]] returns  $\partial_i A_j[t]$ .";
FieldParameters::usage = "FieldParameters is an option for makeDipole list which ";
Preintegrals::usage =
  "Preintegrals is an option for makeDipole list which specifies whether the
    preintegrals of the vector potential should be \"Analytic\" or \"Numeric\".";
ReportingFunction::usage = "ReportingFunction is an option for makeDipole
  list which specifies a function used to report the results, either
  internally (by the default, Identity) or to an external file.";
Gate::usage = "Gate is an option for makeDipole list which specifies the
  integration gate to use. Usage as Gate→g, nGate→n will gate the integral
  at time  $\omega t/\omega$  by  $g[\omega t, n]$ . The default is Gate→SineSquaredGate[1/2].";
nGate::usage = "nGate is an option for makeDipole list which specifies
  the total number of cycles in the integration gate.";
```

```

IonizationPotential::usage = "IonizationPotential is an option for makeDipoleList
    which specifies the ionization potential  $I_p$  of the target.";
Target::usage = "Target is an option for makeDipoleList which specifies
    chemical species producing the HHG emission, pulling the ionization
    potential from the Wolfram ElementData curated data set.";
DipoleTransitionMatrixElement::usage = "DipoleTransitionMatrixElement is an option
    for makeDipoleList which specifies a function f to use as the dipole transition
    matrix element, or a pair of functions {fion,frec} to be used separately for the
    ionization and recombination dipoles, to be used in the form f[p,κ]=f[p,√2 Ip ].";
eCorrection::usage = "eCorrection is an option for makeDipoleList which specifies
    the regularization correction ε, i.e. as used in the factor  $\frac{1}{(t - t_t + i \epsilon)^{3/2}}$ .";
PointNumberCorrection::usage = "PointNumberCorrection is an option for
    makeDipoleList and timeAxis which specifies an extra number of points
    to be integrated over, which is useful to prevent Indeterminate errors
    when a Piecewise envelope is being differentiated at the boundaries.";
IntegrationPointsPerCycle::usage = "IntegrationPointsPerCycle is an option for
    makeDipoleList which controls the number of points per cycle to use for the
    integration. Set to Automatic, to follow PointsPerCycle, or to an integer.";
RunInParallel::usage = "RunInParallel is an option for makeDipoleList which
    controls whether each RB-SFA instance is parallelized. It accepts False
    as the (Automatic) option, True, to parallelize each instance, or a pair
    of functions {TableCommand, SumCommand} to use for the iteration and
    summing, which could be e.g. {Inactive[ParallelTable], Inactive[Sum]}.";
Simplifier::usage = "Simplifier is an option for makeDipoleList which specifies a
    function to use to simplify the intermediate and final analytical results.";

Protect[VectorPotential, VectorPotentialGradient, FieldParameters,
    Preintegrals, ReportingFunction, Gate, nGate, IonizationPotential, Target,
    eCorrection, PointNumberCorrection, DipoleTransitionMatrixElement,
    IntegrationPointsPerCycle, RunInParallel, Simplifier];

Begin["`Private`"];
Options[makeDipoleList] = standardOptions~Join~{
    VectorPotential → Automatic, FieldParameters → {}, VectorPotentialGradient → None,
    Preintegrals → "Analytic", ReportingFunction → Identity,
    Gate → SineSquaredGate[1/2], nGate → 3/2, eCorrection → 0.1,
    IonizationPotential → 0.5,
    Target → Automatic, DipoleTransitionMatrixElement → hydrogenicDTME,
    PointNumberCorrection → 0, Verbose → 0,
    RunInParallel → Automatic, IntegrationPointsPerCycle → Automatic,
    Simplifier → Identity
};
makeDipoleList::gate =
    "The integration gate g provided as Gate→`1` is incorrect. Its usage as
    g[`2`,`3`] returns `4` and should return a number.";
makeDipoleList::pot = "The vector potential A provided as VectorPotential→`1`
    is incorrect or is missing FieldParameters. Its usage as

```

```

A[2] returns 3 and should return a list of numbers.";
makeDipoleList::gradpot = "The vector potential GA provided as
  VectorPotentialGradient→1 is incorrect or is missing FieldParameters.
  Its usage as GA[2] returns 3 and should return a square matrix
  of numbers. Alternatively, use VectorPotentialGradient→None.";
makeDipoleList::preint = "Wrong Preintegrals option 1. Valid
  options are \"Analytic\" and \"Numeric\".";
makeDipoleList::runpar = "Wrong RunInParallel option 1.";

makeDipoleList[OptionsPattern[]] := Block[
{
  num = OptionValue[TotalCycles],
  npp = OptionValue[PointsPerCycle],  $\omega$  = OptionValue[CarrierFrequency],
  dipoleRec, dipoleIon,  $\kappa$ ,
  A, F, GA, pi, ps, S,
  gate, tGate, setPreintegral,
  tInit, tFinal,  $\delta t$ ,  $\delta tint$ ,  $\epsilon$  = OptionValue[ $\epsilon$ Correction],
  AInt, A2Int, GAInt, GAdotAInt, AdotGAInt, GAIntInt, bigPSCorrectionInt, AdotGAdotAInt,
  simplifier, prefactor, integrand, dipoleList,
  TableCommand, SumCommand
},

A[t_] = OptionValue[VectorPotential][t] //. OptionValue[FieldParameters];
F[t_] = -D[A[t], t];
GA[t_] = If[
  TrueQ[OptionValue[VectorPotentialGradient] == None],
  Table[0, {Length[A[tInit]]}, {Length[A[tInit]]}],
  OptionValue[VectorPotentialGradient][t] //. OptionValue[FieldParameters]
];

tInit = 0;
tFinal =  $\frac{2\pi}{\omega}$  num;
(*looping timestep*)
 $\delta t = \frac{tFinal - tInit}{num \times npp + OptionValue[PointNumberCorrection]}$ ;
(*integration timestep*)
 $\delta tint = If[OptionValue[IntegrationPointsPerCycle] === Automatic, \delta t, (tFinal - tInit) /$ 
   $(num \times OptionValue[IntegrationPointsPerCycle] + OptionValue[PointNumberCorrection])]$ ;

tGate = OptionValue[nGate]  $\frac{2\pi}{\omega}$ ;
(*Check potential and potential gradient for correctness.*)
With[{ $\omega tRandom$  = RandomReal[{ $\omega tInit$ ,  $\omega tFinal$ }]},
  If[OptionValue[Verbose]  $\neq$  3
    (*Turn off numerics check if quantum-orbits verbose option is in use.*),
    If[! And@@ (NumberQ/@A[ $\omega tRandom/\omega$ ]), Message[makeDipoleList::pot,
      OptionValue[VectorPotential],  $\omega tRandom$ , A[ $\omega tRandom$ ]];

```

```

Abort[]];
If[! And @@ (NumberQ /@ Flatten[GA[wtRandom/ω]]), Message[makeDipoleList::gradpot,
  OptionValue[VectorPotentialGradient], wtRandom, GA[wtRandom]];
Abort[]];
]];

gate[ωτ_] := OptionValue[Gate][ωτ, OptionValue[nGate]];
With[{wtRandom = RandomReal[{ω tInit, ω tFinal}]},
  If[! TrueQ[NumberQ[gate[wtRandom]]],
    Message[makeDipoleList::gate,
      OptionValue[Gate], wtRandom, OptionValue[nGate], gate[wtRandom]];
    Abort[]];
];

(*Target setup*)
Which[
  OptionValue[Target] === Automatic, κ =  $\sqrt{2 \text{OptionValue[IonizationPotential]}}$ ,
  True, κ =  $\sqrt{2 \text{getIonizationPotential[OptionValue[Target]}}$ 
];
With[{dim = Length[A[RandomReal[{ω tInit, ω tFinal}]]]},
  (*Explicit conjugation of the
  recombination matrix element to keep the integrand analytic.*)
  Which[
    Head[OptionValue[DipoleTransitionMatrixElement]] === List,
    dipoleIon[{p1_, p2_, p3_}][1 ;; dim], κκ_] =
      First[OptionValue[DipoleTransitionMatrixElement]][{p1, p2, p3}][1 ;; dim], κκ];
    dipoleRec[{p1_, p2_, p3_}][1 ;; dim], κκ_] = Assuming[{p1, p2, p3, κκ} ∈ Reals], Simplify[
      Conjugate[
        Last[OptionValue[DipoleTransitionMatrixElement]][{p1, p2, p3}][1 ;; dim], κκ]]
    ];
  , True,
  dipoleIon[{p1_, p2_, p3_}][1 ;; dim], κκ_] =
    OptionValue[DipoleTransitionMatrixElement][{p1, p2, p3}][1 ;; dim], κκ];
  dipoleRec[{p1_, p2_, p3_}][1 ;; dim], κκ_] = Assuming[{p1, p2, p3, κκ} ∈ Reals], Simplify[
    Conjugate[OptionValue[DipoleTransitionMatrixElement][{p1, p2, p3}][1 ;; dim], κκ]]
  ];
];
];
simplifier = OptionValue[Simplifier];

setPreintegral[integralVariable_, preintegrand_,
  dimensions_, integrateWithoutGradient_, parametric_] := Which[
  OptionValue[VectorPotentialGradient] != None || TrueQ[integrateWithoutGradient],
  (*Vector potential gradient specified,
  or integral variable does not depend on it, so integrate*)
  Which[
    OptionValue[Preintegrals] == "Analytic",
    integralVariable[t_, tt_] =

```



```

simplifier[(# /. {τ → t}) - (# /. {τ → tt})] & [Integrate[preintegrand[τ, tt], τ]]];

, OptionValue[Preintegrals] == "Numeric",
Which[
  TrueQ[Not[parametric]],
  Block[{innerVariable},
    integralVariable[t_, tt_] = (innerVariable[t] - innerVariable[tt] /. First[
      NDSolve[{innerVariable'[τ] == preintegrand[τ],
        innerVariable[tInit] == ConstantArray[0, dimensions]},
        innerVariable, {τ, tInit, tFinal}, MaxStepSize → 0.25 / ω]
    ])
  ];
, True,
Block[{matrixpreintegrand, innerVariable, τpre},
  matrixpreintegrand[indices_, t_?NumericQ, tt_?NumericQ] :=
    preintegrand[t, tt][[## & @@ indices]];
  integralVariable[t_, tt_] = Array[(
    innerVariable[##][t - tt, tt] /. First@NDSolve[{
      D[innerVariable[##][τpre, tt], τpre] == Piecewise[
        {{matrixpreintegrand[##, tt + τpre, tt], tt + τpre ≤ tFinal}}, 0],
      innerVariable[##][0, tt] == 0
    }, innerVariable[##]
    ], {τpre, 0, tFinal - tInit}, {tt, tInit, tFinal}
    ], MaxStepSize → 0.25 / ω
  ]
  ) &, dimensions];
];
];
];
, OptionValue[VectorPotentialGradient] === None,
(*Vector potential gradient has not been specified,
and integral variable depends on it, so return appropriate zero matrix*)
integralVariable[t_] = ConstantArray[0, dimensions];
integralVariable[t_, tt_] = ConstantArray[0, dimensions];
];
Apply[setPreintegral,
{
  AInt          A[#1] & {Length[A[tInit]]}
  A2Int         A[#1].A[#1] & {}
  GAInt         GA[#1] & {Length[A[tInit]], Length[A[tI
  GAdotAInt     GA[#1].A[#1] & {Length[A[tInit]]}
  AdotGAInt     A[#1].GA[#1] & {Length[A[tInit]]}
  GAIntInt      GAInt[#1, #2] & {Length[A[tInit]], Length[A[tI
  AdotGAdotAInt A[#1].GAdotAInt[#1, #2] & {}
  bigPSCorrectionInt GAdotAInt[#1, #2] + A[#1].GAInt[#1, #2] & {Length[A[tInit]]}
}

```

```

)
, {1}];

(*{

$$\int_{t_0}^t A(\tau) d\tau, \int_{t_0}^t A(\tau)^2 d\tau, \int_{t_0}^t \nabla A(\tau) d\tau, \int_{t_0}^t \nabla A(\tau) \cdot A(\tau) d\tau, \int_{t_0}^t A(\tau) \cdot \nabla A(\tau) d\tau, \int_{t_0}^t \int_{t_0}^t \nabla A(\tau') d\tau' d\tau,$$


$$\int_{t_0}^t A_k(\tau) \cdot \int_{t_0}^t \partial_k A_j(\tau') A_j(\tau') d\tau' d\tau, \int_{t_0}^t \int_{t_0}^t (A_k(\tau') \partial_j A_k(\tau') + A_k(\tau) \partial_k A_j(\tau')) d\tau' d\tau \}; *)$$


(*Displaced momentum*)
pi[p_, t_, tt_] := p + A[t] - GAInt[t, tt].p - GAdotAInt[t, tt];

(*Stationary momentum and action*)
ps[t_, tt_] := ps[t, tt] = -  $\frac{1}{t - tt - i\epsilon}$  Inverse[IdentityMatrix[Length[A[tInit]]]] -  $\frac{1}{t - tt - i\epsilon}$ 
  (GAIntInt[t, tt] + GAIntInt[t, tt]T).(AInt[t, tt] - bigPSCorrectionInt[t, tt]);
S[t_, tt_] := simplifier[
   $\frac{1}{2}$  (Total[ps[t, tt]2] +  $\kappa^2$ ) (t - tt) + ps[t, tt].AInt[t, tt] +  $\frac{1}{2}$  A2Int[t, tt] - (
    ps[t, tt].GAIntInt[t, tt].ps[t, tt] +
    ps[t, tt].bigPSCorrectionInt[t, tt] + AdotGAdotAInt[t, tt]
  )
];

prefactor[t_,  $\tau$ ] :=  $i \left( \frac{2\pi}{\epsilon + i\tau} \right)^{3/2}$  dipoleRec[pi[ps[t, t -  $\tau$ ], t, t -  $\tau$ ],  $\kappa$ ]  $\times$ 
  dipoleIon[pi[ps[t, t -  $\tau$ ], t -  $\tau$ , t -  $\tau$ ],  $\kappa$ ].F[t -  $\tau$ ];
integrand[t_,  $\tau$ ] := prefactor[t,  $\tau$ ] Exp[-i S[t, t -  $\tau$ ]] gate[ $\omega \tau$ ];

(*Debugging constructs. Verbose→
1 prints information about the internal functions. Verbose→2 returns all the relevant
internal functions and stops. Verbose→3 for quantum-orbit constructs.*)
Which[
  OptionValue[Verbose] == 1, Information/@{A, GA, ps, pi, S, AInt, A2Int,
    GAInt, GAdotAInt, AdotGAInt, GAIntInt, bigPSCorrectionInt, AdotGAdotAInt},
  OptionValue[Verbose] == 2, Return[With[{t = Global`t,
    tt = Global`tt, p = Global`t,  $\tau$  = Global` $\tau$ },
    {A[t], GA[t], ps[t, tt], pi[p, t, tt], S[t, tt], AInt[t], AInt[t, tt], A2Int[t],
    A2Int[t, tt], GAInt[t], GAInt[t, tt], GAdotAInt[t], GAdotAInt[t, tt], AdotGAInt[t],
    AdotGAInt[t, tt], GAIntInt[t], GAIntInt[t, tt], bigPSCorrectionInt[t],
    bigPSCorrectionInt[t, tt], AdotGAdotAInt[t], AdotGAdotAInt[t, tt], integrand[t,  $\tau$ ]}]],
  OptionValue[Verbose] == 3,
  Return[{
    Function[{Global`t, Global`tt}, Evaluate[prefactor[Global`t, Global`t - Global`tt]]],

```

```

Function[{Global`t, Global`tt}, Evaluate[S[Global`t, Global`tt]]]
}]
];

(*Single-run parallelization*)
Which[
  OptionValue[RunInParallel] == Automatic ||
  OptionValue[RunInParallel] == False, TableCommand = Table;
  SumCommand = Sum;;
  OptionValue[RunInParallel] == True, TableCommand = ParallelTable;
  SumCommand = Sum;;
  True, TableCommand = OptionValue[RunInParallel][[1]];
  SumCommand = OptionValue[RunInParallel][[2]];
];

(*Numerical integration loop*)
dipoleList = TableCommand[
  OptionValue[ReportingFunction][
     $\delta t$  SumCommand[(
      integrand[t,  $\tau$ ]
    ),
    { $\tau$ , 0, If[OptionValue[Preintegrals] == "Analytic", tGate, Min[t - tInit, tGate]},  $\delta t$ ]}]
  ], {t, tInit, tFinal,  $\delta t$ }
];
dipoleList

]
End[];

```

Quantum orbit functions suite

GetSaddlePoints

GetSaddlePoints::usage =

"GetSaddlePoints[Ω , *S*, {*tmin*, *tmax*}, { τ min, τ max}] finds a list of solutions {*t*, τ } of the HHG temporal saddle-point equations at harmonic energy Ω for action *S*, in the range {*tmin*, *tmax*} of recombination time and { τ min, τ max} of excursion time, where both ranges should be the lower-left and upper-right corners of rectangles in the complex plane.

GetSaddlePoints[Ω Range, *S*, {*tmin*, *tmax*}, { τ min, τ max}] finds solutions of the HHG temporal saddle-point equations for a range of harmonic energies Ω Range, and returns an Association with each harmonic energy Ω indexing a list of saddle-point solution pairs {*t*, τ }.

GetSaddlePoints[Ω spec, *S*, {{*tmin*₁, *tmax*₁}, { τ min₁, τ max₁}}, {{*tmin*₂, *tmax*₂}, { τ min₂, τ max₂}}, ...]] uses multiple time domains and combines the solutions.

GetSaddlePoints[Ω spec, *S*, {{*urange*, *vrange*}, ...}, IndependentVariables→{*u*, *v*}] uses the explicit

```

independent variables u and v to solve the equations and over the given
ranges, where u and v can be any of "RecombinationTime", "IonizationTime"
and "ExcursionTime", or their shorthands "t", "tt" and "τ" resp.";
SortingFunction::usage = "SortingFunction is an option of GetSaddlePoints
which sets a function f, to be used as f[t,τ,S,Ω], to be
used to sort the solutions, or a list of such functions.";
SelectionFunction::usage = "SelectionFunction is an option of
GetSaddlePoints that sets a function f, to be used as
f[t,τ,S,Ω], such that roots are only kept if f returns True.";
IndependentVariables::usage = "IndependentVariables is an option for
GetSaddlePoints that specifies the two independent variables, out of
"RecombinationTime", "IonizationTime" and "ExcursionTime" (or their
shorthands "t", "tt" and "τ", respectively), to be used in solving
the saddle-point equations, and which range over the given regions.";

FiniteDifference::usage =
  "FiniteDifference is a value for the option Jacobian of FindRoot, FindComplexRoots,
  GetSaddlePoints, and related functions, which specifies that the Jacobian at
  each step should be evaluated using numerical finite difference procedures.";

GetSaddlePoints::error = "Errors encountered for frequency Ω=`1`";

Begin["`Private`"];
Options[GetSaddlePoints] =
  Join[{SortingFunction → (#2 &), SelectionFunction → (True &), IndependentVariables →
    {"RecombinationTime", "ExcursionTime"}}, Options[FindComplexRoots]];
Protect[SortingFunction, SelectionFunction, IndependentVariables, FiniteDifference];

GetSaddlePoints[Ωspec_, S_, {tmin_, tmax_}, {τmin_, τmax_}, options:OptionsPattern[]] :=
  GetSaddlePoints[Ωspec, S, {{tmin, tmax}, {τmin, τmax}}, options]

GetSaddlePoints[Ω_, S_, timeRanges_, options:OptionsPattern[]] :=
  Block[{equations, roots, t, tt, τ, indVars, depVar, depVarRule, tolerances},
    indVars = OptionValue[IndependentVariables] /.
      {"RecombinationTime" → "t", "ExcursionTime" → "τ", "IonizationTime" → "tt"};
    depVar = First[DeleteCases[{"t", "τ", "tt"}, Alternatives @@ indVars]];
    depVarRule = depVar /. {"tt" → {tt → t - τ}, "t" → {t → tt + τ}, "τ" → {τ → t - tt}};
    equations = {D[S[t, tt], t] == Ω, D[S[t, tt], tt] == 0} /. depVarRule;
    tolerances = Which[
      ListQ[OptionValue[Tolerance]], OptionValue[Tolerance],
      True, ConstantArray[
        Which[
          NumberQ[OptionValue[Tolerance]], OptionValue[Tolerance],
          True, 10^If[NumberQ[OptionValue[WorkingPrecision]],
            2 - OptionValue[WorkingPrecision], 2 - $MachinePrecision]
        ],
        2];

SortBy[
  DeleteDuplicates[
    Flatten[Table[

```

```

Select[
  Check[
    roots = ({t,  $\tau$ } /. depVarRule) /. (FindComplexRoots[
      equations
      , Evaluate[Sequence[{Symbol[indVars[[1]]], range[[1, 1]],
        range[[1, 2]]}, {Symbol[indVars[[2]]], range[[2, 1]], range[[2, 2]]}]]
      , Evaluate[Sequence@@FilterRules[{options}, Options[FindComplexRoots]]]
      , SeedGenerator  $\rightarrow$  RandomSobolComplexes
      , Seeds  $\rightarrow$  50
    ] /. {{}}  $\rightarrow$  ({t,  $\tau$ } /. depVarRule)  $\rightarrow$  {}}) (*to deal with empty results*)
    , Message[GetSaddlePoints::error,  $\Omega$ ]; roots
  ]
  ,
  Function[timesPair, OptionValue[SelectionFunction][timesPair[[1]], timesPair[[2]], S,  $\Omega$ ]]
]
, {range, timeRanges}], 1]
, Function[{timesPair1, timesPair2},
  And@@Thread[Abs[timesPair1 - timesPair2] < tolerances]
]
]
, If[
  ListQ[OptionValue[SortingFunction]],
  Table[Function[timesPair, f[timesPair[[1]], timesPair[[2]], S,  $\Omega$ ]],
    {f, OptionValue[SortingFunction]}],
  Function[timesPair, OptionValue[SortingFunction][timesPair[[1]], timesPair[[2]], S,  $\Omega$ ]]
]
]
]
GetSaddlePoints[ $\Omega$ Range_List, S_, timeRanges_, options : OptionsPattern[]] :=
  Association[ParallelTable[
     $\Omega$   $\rightarrow$  GetSaddlePoints[ $\Omega$ , S, timeRanges, options]
    , { $\Omega$ , Sort[ $\Omega$ Range]}]]
End[];

```

ClassifyQuantumOrbits

```

ClassifyQuantumOrbits::usage =
  "ClassifyQuantumOrbits[saddlePoints,f] sorts an indexed set of saddle
    points of the form  $\langle \Omega_1 \rightarrow \{t_{11}, \tau_{11}\}, \{t_{12}, \tau_{12}\}, \dots \rangle$  using a function f,
    which should turn  $f[t, \tau, \Omega]$  into an appropriate label, and returns an
    association of the form  $\langle \text{label}_1 \rightarrow \langle \Omega_1 \rightarrow \langle 1 \rightarrow \{t, \tau\}, 2 \rightarrow \{t, \tau\}, \dots \rangle, \dots \rangle, \dots \rangle$ .

ClassifyQuantumOrbits[saddlePoints,f,sortFunction]
  uses the function sortFunction to sort the sets of saddle points
   $\{\{t_{11}, \tau_{11}\}, \{t_{12}, \tau_{12}\}, \dots\}$  for each label and harmonic energy.

ClassifyQuantumOrbits[saddlePoints,f,sortFunction,DiscardedLabels $\rightarrow$ {label1,label2,...}]
  specifies a list of labels to discard from the final output.";
DiscardedLabels::usage = "DiscardedLabels is an option for ClassifyQuantumOrbits
  which specifies a list of labels to discard from the final output.";

Begin["`Private`"];

Options[ClassifyQuantumOrbits] = {DiscardedLabels  $\rightarrow$  {}};
Protect[DiscardedLabels];

ClassifyQuantumOrbits[saddlePointList_,
  classifierFunction_, sortingFunction_: Sort, OptionsPattern[]] := Map[
  Association,
  MapIndexed[
    (#2[[1]]  $\rightarrow$  #1 &), (*index after sorting each pair*)
    Map[
      sortingFunction,
      Delete[DeleteMissing[
        Query[Transpose][
          MapIndexed[
            GroupBy[classifierFunction@@# &][
              Flatten /@ Transpose[{#1, ConstantArray[#2[[1]], 1], Length[#1]]}] &
            , saddlePointList][All, All, All, {1, 2}]
          ]
        , 2], List /@ OptionValue[DiscardedLabels]]
      , {2}], {3}], {2}]
End[];

```

ReperiodSaddles

```
ClearAll[ReperiodSaddles]
ReperiodSaddles::usage =
  "ReperiodSaddles[{{t1,τ1},{t2,τ2},...},f] readjusts the assigned cycle of
  the saddle points {ti,τi}, returning the list {{t1+f[t1,τ1],τ1},...}."

ReperiodSaddles[<|Ω1→{{t11,τ11},...},Ω2→...|>,f] reperiods
  saddle-point pairs in a harmonic-energy-indexed association.

ReperiodSaddles[<|label1→<|Ω1→{{t11,τ11},...},...|>,...|>,f]
  reperiods saddle-point pairs of a classified set of saddle points.";

Begin["`Private`"];

ReperiodSaddles[pair_ /; Depth[pair] == 2, f_] := {pair[[1]] + f[pair[[1]], pair[[2]]], pair[[2]]}
ReperiodSaddles[association_, f_] := Apply[f, association, {Depth[association] - 2}]

End[];
```

HessianRoot

```
HessianRoot::usage = "HessianRoot[S,t,τ] calculates the Hessian root  $\sqrt{\frac{(2\pi)^2}{\hbar^2 \text{Det}[\partial_{\{t,\tau\}}^2 S]}}$ .";

Begin["`Private`"];

HessianRoot[S_, t_, τ_] := 
$$\sqrt{\frac{2\pi}{\hbar \text{Derivative}[0, 2][S][t, t - \tau]}}$$


$$\sqrt{\left( (2\pi \text{Derivative}[0, 2][S][t, t - \tau]) / (\hbar (\text{Derivative}[2, 0][S][t, t - \tau] \text{Derivative}[0, 2][S][t, t - \tau] - \text{Derivative}[1, 1][S][t, t - \tau]^2)) \right)}$$


End[];
```

FindStokesTransitions

```
FindStokesTransitions::usage =
  "FindStokesTransitions[S,<|Ω1→<|1→{t11,τ11},2→{t12,τ12}|>,Ω2→<|1→{t21,τ21},2→{t22,τ22}|>,...|>]
  finds the set {{ΩS},{ΩAS},n} of the Stokes and anti-Stokes transition
  energies for the given set of saddle points, where Re(S) changes sign after the
  ΩS and Im(S) changes sign after the ΩAS, and n is the index of the member of
  the pair that should be chosen after the transition (taken as the member with
  a positive imaginary part of the action at the largest Ωi in the given keys).

FindStokesTransitions[S,<|label1→<|Ω1→...|>|>] finds the Stokes transitions for the given set
  of saddle-point curve pairs, and returns them labeled with the label1.";
FindStokesTransitions::saddleno = "FindStokesTransitions called with `1`
  of `2` saddle-point sets of length different from 2, with set
  length structure `3`. Excluding those sets from the calculation.";
FindStokesTransitions::multipleS = "FindStokesTransitions found multiple
  Stokes transitions; using `1` to return a single transition.";
```

```

FindStokesTransitions::multipleAS = "FindStokesTransitions found multiple
  anti-Stokes transitions; using `1` to return a single transition.";
ChooserFunction::usage = "ChooserFunction is an option for FindStokesTransitions
  that specifies which transition to take if there are multiple transitions
  in the given dataset. The default is Last and gives the one with higher
  energy; to get the full set of transitions found use Full or Identity.";
ReperiodingFunction::usage = "ReperiodingFunction is an option for FindStokesTransitions,
  SPAdipole and UAdipole which specifies a function f[t,τ] of recombination
  time t and excursion time τ that will be used to re-period the pairs
  {t,τ} into the form {t+f[t,τ],τ}. The default is Function[0], but
  if pairs are split it can be useful to set ReperiodingFunction to
  Function[{t,τ},Floor[-Re[t-τ], $\frac{2\pi}{\omega}$ ]] for ω the carrier frequency. In general,
  however, it is preferable to do this in a single go using ReperiodSaddles.";

Begin["`Private`"];

Protect[ReperiodingFunction, ChooserFunction];
Options[FindStokesTransitions] =
  {ReperiodingFunction → Function[{t, τ}, 0], ChooserFunction → Automatic};

FindStokesTransitions[S_,
  deeperAssociation_ /; Depth[deeperAssociation] == 5, options : OptionsPattern[]] := Map[
  FindStokesTransitions[S, #, options] &,
  deeperAssociation
]

FindStokesTransitions[S_, saddlesAssociation_, options : OptionsPattern[]] :=
Block[{reducedSaddlesAssociation, actionList, signsList, s, processor},
  reducedSaddlesAssociation = KeySort[Select[saddlesAssociation, Length[#] == 2 &]];
  If[Length[saddlesAssociation] - Length[reducedSaddlesAssociation] > 0,
    Message[FindStokesTransitions::saddleno,
      Length[saddlesAssociation] - Length[reducedSaddlesAssociation],
      Length[saddlesAssociation], First /@ Tally /@ Split[Values[Length /@ saddlesAssociation]]
    ]
  ];
  actionList = ReIm[
    Map[(#reduces each  $\Omega \rightarrow |1 \rightarrow S_1, 2 \rightarrow S_2\rangle$  to  $\Omega \rightarrow (S_1 - S_2) *$ )
      Apply[Subtract],
    MapIndexed[(#reduces each  $\Omega \rightarrow |1 \rightarrow \{t_1, \tau_1\}, 2 \rightarrow \{t_2, \tau_2\}\rangle$  to  $\Omega \rightarrow |1 \rightarrow S_1, 2 \rightarrow S_2\rangle *$ )
      With[
        {t = #1[[1]] + OptionValue[ReperiodingFunction][#1[[1]], #2[[2]]], τ = #1[[2]], Ω = #2[[1, 1]]},
        S[t, t - τ] - Ω t
      ] &
    , reducedSaddlesAssociation, {2}
  ]
];
  signsList = Sign[Times[
    Rest[actionList],
    AssociationThread[Rest[Keys[actionList]], Most[Values[actionList]]]
  ]];

```



```

processor = OptionValue[ChooserFunction] /. {Automatic → Last, Full → Identity};
If[Length[Keys[Select[signsList, #[[1]] < 0 &]]] > 1,
  Message[FindStokesTransitions::multipleS, processor]];
If[Length[Keys[Select[signsList, #[[2]] < 0 &]]] > 1,
  Message[FindStokesTransitions::multipleAS, processor]];
{
  processor[Keys[Select[signsList, #[[1]] < 0 &]] /. {{} → {Missing["No transition"]}}],
  processor[Keys[Select[signsList, #[[2]] < 0 &]] /. {{} → {Missing["No transition"]}}],
  Sign[Last[actionList][[2]]] /. {1 → 2, -1 → 1}
}
]

End[];

```

SPAdipole

```

SPAdipole::usage =
  "SPAdipole[S,prefactor, $\Omega$ , {t, $\tau$ }] returns the saddle-point approximation amplitude
  corresponding to action  $S[t,t-\tau]-\Omega t$  and the given prefactor.

SPAdipole[S,prefactor, $\Omega$ , <|1→{t1, $\tau$ 1},2→{t2, $\tau$ 2},...|>] returns the total harmonic-dipole
  contribution in the saddle-point approximation from the specified saddle points.

SPAdipole[S,prefactor, $\Omega$ , <|1→{t1, $\tau$ 1},2→{t2, $\tau$ 2}|>,transition] uses the given Stokes
  transition set to drop the relevant saddle after the anti-Stokes transition.";
SPAdipole::wrongno = "SPAdipole called with a Stokes transition but
  with an input association of length `1` at harmonic
  energy  $\Omega=\tau$ . Reverting to unstructured evaluation.";
SPAdipole::invldtrns = "SPAdipole called with invalid Stokes transition
  set `1`. Reverting to unstructured evaluation.";

Begin["`Private`"];

Options[SPAdipole] = {ReperiodingFunction → Function[{t,  $\tau$ ], 0}};

SPAdipole[S_, prefactor_,  $\Omega$ _, {t_,  $\tau$ _}, options:OptionsPattern[]] :=
  Block[{tr = t + OptionValue[ReperiodingFunction][t,  $\tau$ ]},
    HessianRoot[S, tr,  $\tau$ ] prefactor[tr, tr -  $\tau$ ] Exp[- $i$  S[tr, tr -  $\tau$ ] +  $i$   $\Omega$  tr]
  ]
SPAdipole[S_, prefactor_,  $\Omega$ _, times_Association, options:OptionsPattern[]] := Block[{},
  Total[SPAdipole[S, prefactor,  $\Omega$ , #, options] & /@ times]
]

SPAdipole[S_, prefactor_,  $\Omega$ _, times_Association,
  transition_, options:OptionsPattern[]] := Block[{},
  If[!NumberQ[transition[[2]]], Message[SPAdipole::invldtrns, transition];
    Return[SPAdipole[S, prefactor,  $\Omega$ , times] ]];
  If[Length[times] ≠ 2, Message[SPAdipole::wrongno, Length[times],  $\Omega$ ];
    Return[SPAdipole[S, prefactor,  $\Omega$ , times, options] ]];
  If[ $\Omega$  < transition[[2]],
    SPAdipole[S, prefactor,  $\Omega$ , times, options],
    SPAdipole[S, prefactor,  $\Omega$ , KeySelect[times, # == transition[[3]] &], options]
  ]
]

End[];

```

UAdipole

```

UAdipole::usage =
  "UAdipole[S,prefactor,Ω,<|1→{t1,τ1},2→{t2,τ2},...|>,transition] returns the total
    harmonic-dipole contribution in the uniform approximation from the specified
    saddle points, taking the given Stokes transition set as a reference.";
UAdipole::saddleno = "UAdipole called with `1` time pairs at Ω=`2`.
  Reverting to the saddle-point approximation for this set.";
UAdipole::invldtrns = "UAdipole called with invalid Stokes transition set
  `1`. Reverting to the saddle-point approximation for this set.";

Begin["`Private`"];

Options[UAdipole] = {ReperiodingFunction → Function[{t, τ}, 0]};

UAdipole[S_, prefactor_, Ω_, times_, transition_, options : OptionsPattern[]] := (
  If[Length[times] ≠ 2, Message[UAdipole::saddleno, Length[times], Ω];
    Return[SPAdipole[S, prefactor, Ω, times] ]];
  If[! NumberQ[transition[[2]], Message[UAdipole::invldtrns, transition];
    Return[SPAdipole[S, prefactor, Ω, times] ]];
  Block[
    {A1, A2, S1, S2, Ss, Sm, z,
     t1 = times[1][[1]] + OptionValue[ReperiodingFunction][times[1][[1]], times[1][[2]]],
     τ1 = times[1][[2]],
     t2 = times[2][[1]] + OptionValue[ReperiodingFunction][times[2][[1]], times[2][[2]]],
     τ2 = times[2][[2]]},
    A1 = HessianRoot[S, t1, τ1] prefactor[t1, t1 - τ1];
    S1 = S[t1, t1 - τ1] - Ω t1;
    A2 = HessianRoot[S, t2, τ2] prefactor[t2, t2 - τ2];
    S2 = S[t2, t2 - τ2] - Ω t2;
    Ss =  $\frac{S1 + S2}{2}$ ; Sm =  $\frac{S1 - S2}{2}$ ;
    If[Ω < transition[[2]], z =  $\left(-\frac{3}{2} Sm\right)^{2/3}$ ,
      z =  $\left(-\frac{3}{2} Sm\right)^{2/3} \text{Exp}\left[i \left(\text{transition}[[3]] /. \{2 \rightarrow -1, 1 \rightarrow 1\} \frac{2\pi}{3}\right)\right]$ ;
       $\sqrt{6\pi Sm} \text{Exp}\left[-i Ss + i \frac{\pi}{4}\right] \left(\frac{A1 - i A2}{2} \frac{\text{AiryAi}[-z]}{\sqrt{z}} + i \frac{A1 + i A2}{2} \frac{\text{AiryAi}'[-z]}{z}\right)$ 
    ]
  )

End[];

```

Package closure

End of package

```
EndPackage[];
```

Add to distributed contexts.

```
DistributeDefinitions["RBSFA`"];
```