

¹ Paulie: A Python package to study Lie algebraic properties of quantum systems

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⁸ Summary

⁹ The control of quantum systems lies at the heart of quantum computation and is studied via ¹⁰ the dynamical Lie algebra (DLA), i.e. the Lie algebra that is generated by terms in the system's ¹¹ Hamiltonian. Tensor products of Pauli matrices (Pauli strings) form a basis of Hermitian ¹² operators and as such appear as generators in trotterized Hamiltonian simulation. Moreover, ¹³ they naturally arise as effective control Hamiltonians in trapped-ion systems and also occur in ¹⁴ measurement-based quantum computation. While a conceptually complete classification of ¹⁵ systems that are generated by Pauli strings is known, this classification was not packaged as an ¹⁶ explicit, scalable procedure, and therefore remained difficult to deploy in concrete design tasks. ¹⁷ Within the Python package `paulie` we present an efficient algorithm for this classification that ¹⁸ directly allows to study the controllability of Pauli string generated systems. We showcase the ¹⁹ utility of our framework by numerically providing the optimal universal generator sets.

²⁰ Statement of need

²¹ The dynamical Lie algebra (DLA) is generated via nested commutation of the interaction ²² terms in the system's Hamiltonian. It constitutes the tangent space of the associated Lie ²³ group which captures the possible trajectories of the evolving quantum system. Therefore, ²⁴ the DLA is a pivotal tool to study the symmetries and dynamics of quantum systems and ²⁵ occurs in various fields within quantum information theory such as quantum control ([Zeier & Schulte-Herbrüggen, 2011](#)), simulation ([Goh et al., 2025](#)) and quantum machine learning ²⁶ [Ragone et al. \(2024\)](#).

²⁷ Pauli strings form a basis for Hermitian operators and as such appear as generators in ²⁸ trotterized Hamiltonian simulation ([Trotter, 1959](#)). They also arise as native gate sets in ²⁹ trapped ion systems ([Nam et al., 2020](#)) and occur in measurement-based quantum computation ³⁰ ([Raussendorf et al., 2003](#)).

³¹ A recent line of work ([Kökçü et al., 2024](#); [Wiersema et al., 2024](#)) culminated in a classification ³² of DLAs generated by arbitrary sets of Pauli strings ([Aguilar et al., 2024](#)). While conceptually ³³ complete, this classification was not packaged as an explicit, scalable procedure, and therefore ³⁴ remained difficult to deploy in concrete design and compilation tasks. This work provides a ³⁵ bridge as we provide an efficient algorithm that takes any set of Pauli strings and outputs its ³⁶ DLA. This allows the direct study of controllability of quantum systems generated by Pauli ³⁷ strings. In addition equivalences of generating sets can be easily checked and transitions ³⁸ between algebras by changing the generator set studied. We demonstrate the utility of the ³⁹ framework by constructing universal generator sets with optimal generation rate.

41 Preliminaries

42 A state of an isolated physical system is represented at a fixed time t by a vector $|\psi\rangle$ in a
 43 Hilbert space \mathcal{H} , that is, a complex vector space endowed with an inner product. The time
 44 evolution of the state is determined by the Schrödinger equation $i\frac{d}{dt}|\psi\rangle = H|\psi\rangle$, where we
 45 set $\hbar = 1$. H is the Hamiltonian that describes the observable corresponding to the total
 46 energy of the system. An observable is a Hermitian linear map acting on \mathcal{H} . The solution
 47 to this differential equation with initial state $|\psi_0\rangle$ and time-independent Hamiltonian H is
 48 $|\psi(t)\rangle = e^{-iHt}|\psi_0\rangle$. We specify a set of operators $\mathcal{G} = \{h_j\}_{j \in J}$ that describe the ways in
 49 which we can control the quantum system and that corresponds to the individual interaction
 50 terms of the Hamiltonian. The reachable states are then of the form

$$|\psi(t)\rangle = e^{-ih_{l_m}t_{l_m}} \dots e^{-ih_{l_1}t_{l_1}} |\psi_0\rangle \quad (1)$$

51 for some index set $\{l_i\}_{i=1}^m \subset J$. The associated matrix Lie algebra \mathfrak{g} , that is, a real subspace of
 52 complex matrices that is closed under the matrix commutator $i[h, g] = i(hg - gh)$, is dubbed
 53 dynamical Lie algebra (DLA). Hence, the elements of \mathcal{G} span all elements in \mathfrak{g} through nested
 54 commutators, $\mathfrak{g} = \text{span}_{\mathbb{R}}(\langle \mathcal{G} \rangle_{Lie})$. In case the generator set \mathcal{G} consists of tensor products of
 55 Pauli matrices (Paulistrings), \mathfrak{g} is termed Pauli Lie algebra.

56 The classical compact simple Lie algebras are

$$\mathfrak{su}(d) = \{x \in \mathbb{C}^{d \times d} | x = -x^\dagger, \text{tr } x = 0\} \quad (2)$$

$$\mathfrak{so}(d) = \{x \in \mathbb{R}^{d \times d} | x = -x^T\} \quad (3)$$

$$\mathfrak{sp}(2d) = \{x \in \mathfrak{su}(2d) | x = -\Omega x^T \Omega^T\} \quad (4)$$

57 where the symplectic form is denoted as

$$\Omega = \begin{pmatrix} 0 & \mathbb{I}_d \\ -\mathbb{I}_d & 0 \end{pmatrix}. \quad (5)$$

58 Considering a system out of n qubits, the Hilbert space is $\mathcal{H} = \mathbb{C}^{2^n}$. A generator set that
 59 spans all of $\mathfrak{su}(2^n)$ is called universal, and the system is fully controllable. To characterise
 60 spin chains, we employ the following graph

61 **Definition 1** *The anticommutation graph has \mathcal{G} as a node set and the edges are between the
 62 generators that anticommute $E = \{(p, q) | [p, q] \neq 0\}$.*

63 Classification algorithm

64 Optimal universal generator set

65 With the feature `get_optimal_su_2_n_generators` we find among all universal generator sets
 66 those that have optimal generation rate, i.e. the sets that generate $\mathfrak{su}(2^n)$ the fastest. In
 67 (Smith et al., 2025) it is derived that these are exactly generating sets with a fraction of
 68 anticommuting pairs of generators out of the total number of pairs to be approximately 0.706.
 69 Analytically this fraction maximizes a particular q-Pochhammer symbol. We can reframe the
 70 problem in terms of the anticommutation graph. Optimal universal generating sets have an
 71 anticommutation graph with approximately $\lfloor 0.706 \cdot \binom{n}{2} \rfloor$ edges. To search for such graphs, we
 72 start with the canonical graph of $\mathfrak{su}(2^n)$. We can obtain the canonical graph by the algorithm
 73 described in the previous section. The canonical graph is the anticommutation graph with
 74 minimal connectivity, hence we are guaranteed to be below the target anticommutation fraction.
 75 We iterate over the edge set of the canonical graph and perform a contraction such that the
 76 anticommutation fraction increases and become closer to the target value until we reach the
 77 target value.

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