

(6) optimal filtering for cross-correlation statistic

$$\hat{S}_h \approx \int_{-\infty}^{\infty} df \int_{-\infty}^{\infty} df' \int_T (f-f') \tilde{d}_1(f) \tilde{d}_2^*(f') \tilde{Q}^*(f')$$

Expected value

$$\mu \equiv \langle \hat{S}_h \rangle$$

$$= \int_{-\infty}^{\infty} df \int_{-\infty}^{\infty} df' \delta_T(f-f') \langle \tilde{d}_1(f) \tilde{d}_2^*(f') \rangle \tilde{Q}^*(f')$$

For uncorrelated noise

$$\langle \tilde{d}_1(f) \tilde{d}_2^*(f') \rangle = \langle \tilde{h}_1(f) \tilde{h}_2^*(f') \rangle$$

$$= \frac{1}{2} \delta(f-f') \Gamma_{12}(f) \tilde{S}_h(f)$$

$$\rightarrow \mu = \frac{1}{2} \int_{-\infty}^{\infty} df \underbrace{\delta_T(0)}_T \Gamma_{12}(f) \tilde{S}_h(f) \tilde{Q}^*(f)$$

$$= \frac{T}{2} \int_{-\infty}^{\infty} df \Gamma_{12}(f) \tilde{S}_h(f) \tilde{Q}^*(f)$$

Variance:

$$\sigma^2 \equiv \langle \hat{s}_h - \langle \hat{s}_h \rangle^2 \rangle$$

$$= \langle \hat{s}_h^2 \rangle - \mu^2$$

$$= \int dt \int dt' \int dp' \int dp \delta_T(t-t') \delta_T(p-p') \tilde{Q}^*(t') \tilde{Q}(p') \tilde{Q}(p)$$

$$\langle \tilde{d}_1(t) \tilde{d}_2^*(t') \tilde{d}_1^*(p) \tilde{d}_2(p') \rangle$$

$$= \langle \tilde{d}_1(t) \tilde{d}_2^*(t') \tilde{d}_1^*(p) \tilde{d}_2(p') \rangle$$

Use $\langle abcd \rangle = \langle ab \rangle \langle cd \rangle + \langle ac \rangle \langle bd \rangle + \langle ad \rangle \langle bc \rangle$
for zero-mean Gaussian random variables

so

$$\langle \tilde{d}_1 \tilde{d}_2^* \tilde{d}_1^* \tilde{d}_2 \rangle = \langle \tilde{d}_1 \tilde{d}_2^* \rangle \langle \tilde{d}_1^* \tilde{d}_2 \rangle$$

$$= \langle \tilde{d}_1(t) \tilde{d}_1^*(p) \rangle \langle \tilde{d}_2^*(t') \tilde{d}_2(p') \rangle + \langle \tilde{d}_1(t) \tilde{d}_2^*(p') \rangle \langle \tilde{d}_1^*(p) \tilde{d}_2(t') \rangle$$

$$\approx \frac{1}{2} \delta(t-p) P_1(t) \pm \delta(t'-p') P_2(t')$$

for uncorrelated noise, this is proportional to $S_h^2(t)$

total auto-correlated power

$$P_1(t) = P_{in}(t) + P_{gw}(t) \approx P_{in}(t) \text{ for weak signal}$$

Thus,

$$\sigma^2 \approx \frac{1}{4} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} 1 f' \int_{-\infty}^{\infty} \underbrace{\delta_T(t-t') \delta_T(t-t'') \tilde{Q}^*(t) \tilde{Q}(t')}_{\approx \delta(t-t')} P_1(t) P_2(t) dt$$

$$= \frac{1}{4} \int_{-\infty}^{\infty} \underbrace{\delta_T(t)}_{\approx \delta(t-t')} |\tilde{Q}(t)|^2 P_1(t) P_2(t) dt$$

$$= \frac{1}{4} \int_{-\infty}^{\infty} 1 f |\tilde{Q}(t)|^2 P_1(t) P_2(t) dt$$

$$\rightarrow \text{SNR} = \frac{\mu}{\sigma} = \frac{\frac{1}{4} \int_{-\infty}^{\infty} 1 f \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} 1 f \int_{-\infty}^{\infty} \tilde{Q}^*(t) \tilde{Q}(t) dt dt}{\sqrt{\frac{1}{4} \int_{-\infty}^{\infty} 1 f |\tilde{Q}(t)|^2 P_1(t) P_2(t) dt}}$$

$$= \sqrt{T} \frac{\int_{-\infty}^{\infty} 1 f \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \tilde{Q}^*(t) \tilde{Q}(t) dt dt}{\int_{-\infty}^{\infty} 1 f P_1(t) P_2(t) dt}$$

$$\sqrt{\int_{-\infty}^{\infty} 1 f \int_{-\infty}^{\infty} \tilde{Q}^*(t) \tilde{Q}(t) dt dt} / \int_{-\infty}^{\infty} 1 f P_1(t) P_2(t) dt$$

$$= \sqrt{T} \left(\frac{\int_{-\infty}^{\infty} \tilde{Q}^*(t) \tilde{Q}(t) dt}{\int_{-\infty}^{\infty} P_1(t) P_2(t) dt} \right)$$

where $(A, B) = \int_{-\infty}^{\infty} A(t) B^*(t) P_1(t) P_2(t) dt$

$\rightarrow \text{SNR} = \sqrt{T} \frac{(A, \tilde{Q})}{\sqrt{(Q, Q)}}$

where $A(t) = \frac{\int_{12}(t) \int_1(t)}{P_1(t) P_2(t)}$

Recall: $\vec{A} \cdot \vec{B} = |\vec{A}| |\vec{B}| \cos \theta$
 $\vec{B} \cdot \vec{B} = |\vec{B}|^2$

$\frac{\vec{A} \cdot \vec{B}}{\sqrt{B \cdot B}} = |\vec{A}| \cos \theta$

is maximized by have $\theta = 0$
 $\rightarrow \vec{A}$ should point in same direction as \vec{B} .

Thus, SNR is maximized when $\tilde{Q}(t) \propto \frac{\int_{12}(t) \int_1(t)}{P_1(t) P_2(t)}$ just different by overall amplitude
 $\propto \frac{\int_{12}(t) H(t)}{P_1(t) P_2(t)}$

$$(7) (4) p(d|a, \sigma) \propto \exp \left[-\frac{1}{2} \sum_{i=1}^N \frac{(d_i - a)^2}{\sigma_i^2} \right]$$

$$N_{\text{maximize}} \text{ likelihood} \iff N_{\text{maximize}} \ln(\text{likelihood})$$

$$\text{Derive } L(a) = \ln[p(d|a, \sigma)] = -\frac{1}{2} \sum_{i=1}^N \frac{(d_i - a)^2}{\sigma_i^2}$$

$$0 = \frac{dL}{da} = - \sum_{i=1}^N \frac{2(d_i - a)(-1)}{\sigma_i^2} \Big|_{a=a}$$

$$= 2 \sum_{i=1}^N \frac{(d_i - a)}{\sigma_i^2}$$

$$= 2 \left(\sum_{i=1}^N \frac{d_i}{\sigma_i^2} - a \sum_{i=1}^N \frac{1}{\sigma_i^2} \right)$$

$$\rightarrow a = \frac{\sum_{i=1}^N \frac{d_i}{\sigma_i^2}}{\sum_{i=1}^N \frac{1}{\sigma_i^2}}$$

$$(b) \quad p(d|A, c) \propto \exp \left[-\frac{1}{2} (d - mA)^T C^{-1} (d - mA) \right]$$

Dagger \equiv complex-conjugate transpose
allow for complex data d , parameter A

Define: $L(A, A^T) = \ln p(d|A, c)$

$$= -\frac{1}{2} (d - mA)^T C^{-1} (d - mA)$$

Treat A, A^T as independent variables for the variations

$$\delta L = \frac{1}{2} \delta A^T m^T C^{-1} (d - mA) + \frac{1}{2} (d - mA)^T C^{-1} m \delta A$$

$$\delta L = 0 \text{ for } \delta A^T \rightarrow 0 = m^T C^{-1} (d - mA) \\ = m^T C^{-1} d - m^T C^{-1} m \hat{A}$$

$$\rightarrow \hat{A} = (m^T C^{-1} m)^{-1} m^T C^{-1} d \\ = F^{-1} X$$

where $F \equiv m^T C^{-1} m$ (Fisher matrix)

$$X \equiv m^T C^{-1} d \quad ("dirty map")$$

$$\{L=0 \quad f_0 \quad SA$$

$$\frac{1}{2} (I - MA) \hat{C}^{-1} M = 0$$

$$d C^{-1} M - A^T M^T + C^{-1} M = 0$$

Take + of above equation using $(C^{-1})^T = C^{-1}$

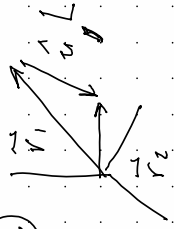
$$\rightarrow M^T C^{-1} d - M^T C^{-1} M \hat{A} = 0$$

$$\rightarrow \hat{A} = (M^T C^{-1} M)^{-1} M^T C^{-1} d$$

$$= F^{-1} X$$

(as before)

8.



$$r(t) = \Delta T(t) = \frac{1}{z_c} u u^5 \int_0^L d_s h_{45}(t/s, \vec{x}/s)$$

where $t(s) = (t - \frac{L}{c}) + \frac{z}{c}$

$$\vec{x}(s) = \vec{r}_1 + s \vec{u}$$

$$s \in [0, L]$$

plane-wave expansion

$$h_{45}(t, \vec{x}) = \int_{-\infty}^{\infty} d\omega \int d^3\Omega \sum_A h_4(f, \vec{k}) e^{i\vec{k} \cdot \vec{x}} e^{i\omega t} e^{i2\pi f(t - \frac{L}{c})}$$

Replace t, \vec{x} by $t(s), \vec{x}(s)$

Exponential becomes

$$e^{i2\pi f \left((t - \frac{L}{c}) + \frac{z}{c} - \frac{1}{k} (\vec{r}_1 + s \vec{u}) \right)} e^{i2\pi f \left((t - \frac{L}{c}) - \frac{1}{k} \cdot \vec{r}_1 \right)} e^{i2\pi f \frac{z}{c} (1 - \frac{\vec{k} \cdot \vec{u}}{k})}$$

$$= e^{i2\pi f \left(t - \frac{L}{c} - \frac{1}{k} \cdot \vec{r}_1 \right)} e^{i2\pi f \frac{z}{c} (1 - \frac{\vec{k} \cdot \vec{u}}{k})}$$

$$= e^{i2\pi f \left(t - \frac{L}{c} - \frac{1}{k} \cdot \vec{r}_1 \right)} e^{i2\pi f \frac{z}{c} (1 - \frac{\vec{k} \cdot \vec{u}}{k})}$$

$$D_0 \text{ integral} \int_0^L e^{i 2 \pi f \frac{L}{c} (1 - \hat{k} \cdot \hat{u})} = \frac{e^{i 2 \pi f \frac{L}{c} (1 - \hat{k} \cdot \hat{u})}}{i 2 \pi f (1 - \hat{k} \cdot \hat{u})}$$

$$= \frac{c}{(i 2 \pi f) (1 - \hat{k} \cdot \hat{u})} \left(e^{i 2 \pi f \frac{L}{c} (1 - \hat{k} \cdot \hat{u})} - 1 \right)$$

Therefore,

$$\Delta T(t) = \frac{1}{2c} \int d\Omega_A \sum_A h_A(t, \hat{k}) e^{i \frac{A}{c} (\hat{k})} \frac{1}{(i 2 \pi f) (1 - \hat{k} \cdot \hat{u})}$$

$$\left[e^{i 2 \pi f (t_1 - \hat{k} \cdot \frac{\vec{r}_1}{c})} e^{i 2 \pi f \frac{L}{c} (1 - \hat{k} \cdot \hat{u})} - e^{i 2 \pi f (t_1 - \hat{k} \cdot \frac{\vec{r}_1}{c})} \right]$$

$$e^{i 2 \pi f \left((t_1 + \frac{L}{c}) - \underbrace{\hat{k} \cdot (\vec{r}_1 + \hat{u} L)}_{\frac{1}{r_2}} \right)}$$

$$= \frac{1}{2} \int d\Omega_A \sum_A h_A(t, \hat{k}) e^{i \frac{A}{c} (\hat{k})} \left(\frac{1}{i 2 \pi f} \right) \left(\frac{1}{1 - \hat{k} \cdot \hat{u}} \right) e^{i 2 \pi f (t_1 - \hat{k} \cdot \frac{\vec{r}_1}{c})}$$

NOTE:

$$e^{i2\pi f(t_2 - \hat{r}_2/c)}$$

e

$$e^{i2\pi f(t_1 - \hat{r}_1/c)}$$

e

corresponds to reception of pulse at \vec{r}_2 at time t_2

corresponds to emission of pulse at \vec{r}_1 at time t_1

Rewrite:

$$\Delta T(t) = \int d\Omega_A \sum_A h_A(t, \hat{u}) \left(\frac{1}{i2\pi f} \right) \left(\frac{1}{1 - \hat{u} \cdot \hat{u}} \right) e^{A(\hat{u})} u_{ab}$$

$$e^{i2\pi f t - i2\pi f \hat{r}_1 \cdot \vec{r}_2/c} \left[1 - e^{i2\pi f t_1 - i2\pi f \hat{r}_1 \cdot \vec{r}_2/c} \right]$$

$$e^{i2\pi f t - i2\pi f \hat{r}_1 \cdot \vec{r}_2/c}$$

$$\left. \begin{aligned} & e^{i2\pi f t_1 - i2\pi f \hat{r}_1 \cdot \vec{r}_2/c} \\ & e^{i2\pi f \frac{L}{c} - i2\pi f \frac{L}{c} (1 - \hat{u} \cdot (\vec{r}_2 - \vec{r}_1))} \\ & e^{-i2\pi f \frac{L}{c} (1 - \hat{u} \cdot \hat{u})} \end{aligned} \right\} t - \frac{L}{c}$$

Thru,

$$\Delta T(t) = \int_{-\infty}^{\infty} dt e^{i2\pi ft} \int d^2 \Omega_H \sum_A h_A(t, \hat{k}) \left(\frac{1}{i2\pi f} \right) u^a u^b e_{ab}^A(\hat{k}) \frac{1}{1 - \hat{k} \cdot \hat{u}}$$

$$e^{-i2\pi f \hat{k} \cdot \hat{r}_2 \frac{1}{c}} \left[1 - e^{-i2\pi f \frac{L}{c} (1 - \hat{k} \cdot \hat{u})} \right]$$

So

$$R_A(t, \hat{k}) = \frac{1}{i2\pi f} \frac{u^a u^b e_{ab}^A(\hat{k})}{2(1 - \hat{k} \cdot \hat{u})} \left[1 - e^{-i2\pi f \frac{L}{c} (1 - \hat{k} \cdot \hat{u})} \right] e^{-i2\pi f \hat{k} \cdot \hat{r}_2 \frac{1}{c}}$$

NOTE: for polar timing $\hat{p} = -\hat{u}$, take $\hat{r}_2 = \vec{0}$ (SSB)

$\frac{1}{i2\pi f} \rightarrow 1$ for redshift $\left(\frac{\Delta \nu}{\nu}\right)$ measurement

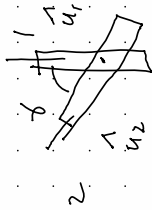
$$R_A(t, \hat{k}) = \underbrace{\frac{1}{2} \left(\frac{p^a p^b e_{ab}^A(\hat{k})}{1 + \hat{k} \cdot \hat{p}} \right)}_{\text{Earth term}} \left[1 - e^{-i2\pi f \frac{L}{c} (1 + \hat{k} \cdot \hat{p})} \right] \underbrace{\quad}_{\text{polar term}}$$

$$\equiv F^A(\hat{k})$$

(9) ORF for colorless electric dipole antennae

$$r_I(t) = u_I, \vec{E}(t, \vec{x}=0)$$

$$\vec{E}(t, \vec{x}) = \int d\kappa \int d\Omega_{\kappa} \sum_{\alpha=1,2} \vec{E}_{\alpha}(t, \kappa) \hat{e}_{\alpha}(\kappa) e^{i\kappa(t - \vec{x} \cdot \vec{x}/c)}$$



$$\vec{u}_1 = \hat{z}$$

$$\vec{u}_2 = \sin \gamma \hat{x} + \cos \gamma \hat{z}$$

$$\hat{e}_1(\kappa) = -\hat{\phi}$$

$$\hat{e}_2(\kappa) = -\hat{\theta}$$

$$\vec{E}(t, \vec{x}=0)$$

$$\begin{aligned} \text{Now: } r_I(t) &= u_I, \vec{E}(t, \vec{x}=0) \\ &= \int_{-\infty}^{\infty} d\kappa \int d\Omega_{\kappa} \sum_{\alpha=1,2} \vec{E}_{\alpha}(\kappa, \kappa) u_I \cdot \hat{e}_{\alpha}(\kappa) e^{i2\pi f t} \\ &= \int_{-\infty}^{\infty} d\kappa e^{i2\pi f t} \sum_{\alpha} \vec{E}_{\alpha}(\kappa, \kappa) u_I \cdot \hat{e}_{\alpha}(\kappa) \end{aligned}$$

$$s.o. \quad R_I^{\alpha}(\kappa, \kappa) = u_I \cdot \hat{e}_{\alpha}(\kappa)$$

$$\chi_{501} \theta_{415} + \phi_{102} \theta_{101} \chi_{415} =$$

$$\theta \cdot \left(\chi_{102} + \chi_{415} \right) = \begin{pmatrix} 1 \\ 1 \end{pmatrix}^2 \otimes \begin{pmatrix} 1 \\ 1 \end{pmatrix}^2$$

$$\theta_{415} = \theta \cdot \chi = \begin{pmatrix} 1 \\ 1 \end{pmatrix}^2 \otimes \begin{pmatrix} 1 \\ 1 \end{pmatrix}^2$$

$$\chi_{415} \phi_{415} =$$

$$\phi \cdot \left(\chi_{102} + \chi_{415} \right) = \begin{pmatrix} 1 \\ 1 \end{pmatrix}^2 \otimes \begin{pmatrix} 1 \\ 1 \end{pmatrix}^2$$

$$0 = \phi \cdot \chi = \begin{pmatrix} 1 \\ 1 \end{pmatrix}^2 \otimes \begin{pmatrix} 1 \\ 1 \end{pmatrix}^2$$

$$\chi_{102} - \chi_{415} \theta_{415} - \chi_{102} \theta_{101} = \chi_{415}$$

$$\chi_{102} + \chi_{415} \theta_{415} - \chi_{102} \theta_{101} = \chi_{415}$$

$$\chi_{102} - \chi_{415} = \chi_{415}$$

$$\chi_{102} + \chi_{415} = \chi_{415}$$

