

# AA 242A Homework 8 Solution

December 5, 2021

## 1 Problem 1

a) It should be obvious that Lagrangian approach should be used here. In this solution I will present a concise approach which requires some intuitions and insights into the problem (as shall be seen later); a more complete approach can be found in a separate file in the homework section.

First we express the location of the center of mass of the disk:

$$\begin{aligned}x_{disk} &= x \cos \theta + r \sin \theta \\y_{disk} &= -x \sin \theta + r \cos \theta \\\Rightarrow \dot{x}_{disk} &= \dot{x} \cos \theta - x \dot{\theta} \sin \theta + r \dot{\theta} \cos \theta \\\dot{y}_{disk} &= -\dot{x} \sin \theta - x \dot{\theta} \cos \theta - r \dot{\theta} \sin \theta \\\Rightarrow T_{trans} &= \frac{1}{2}m(\dot{x}_{disk}^2 + \dot{y}_{disk}^2) \\&= \frac{1}{2}m\dot{x}^2 + \frac{1}{2}mx^2\dot{\theta}^2 + \frac{1}{2}mr^2\dot{\theta}^2 + mr\dot{x}\dot{\theta}\end{aligned}$$

Next is the rotational kinetic energy of the disk. Now here comes the little bit of insight: the rate of rotation of the disk in inertial frame is not only  $\frac{\dot{x}}{r}$ , but  $\frac{\dot{x}}{r} + \dot{\theta}$  since it also drops with the rod.

$$\begin{aligned}T_{rotation} &= \frac{1}{2}I_{center} \left( \frac{\dot{x}}{r} + \dot{\theta} \right)^2 \\&= \frac{1}{4}m\dot{x}^2 + \frac{1}{2}mr\dot{x}\dot{\theta} + \frac{1}{4}mr^2\dot{\theta}^2\end{aligned}$$

The rest is just routine work:

$$\begin{aligned}
T_{rod} &= \frac{1}{2} I_{end} \dot{\theta}^2 \\
&= \frac{8}{3} m r^2 \dot{\theta}^2 \\
V_{rod} &= -2 r m g \sin \theta \\
V_{disk} &= -m g x \sin \theta + m g r \cos \theta \\
L &= \frac{3 m \dot{x}^2}{4} + \frac{m x^2 \dot{\theta}^2}{2} + \frac{41 m r^2 \dot{\theta}^2}{12} + \frac{3 m r \dot{x} \dot{\theta}}{2} - (m g r \cos \theta - m g x \sin \theta - 2 r m g \sin \theta)
\end{aligned}$$

Because no non-conservative forces are acting on the system (frictionless pivot), we shall use the form of Lagrangian equations with no generalized forces:

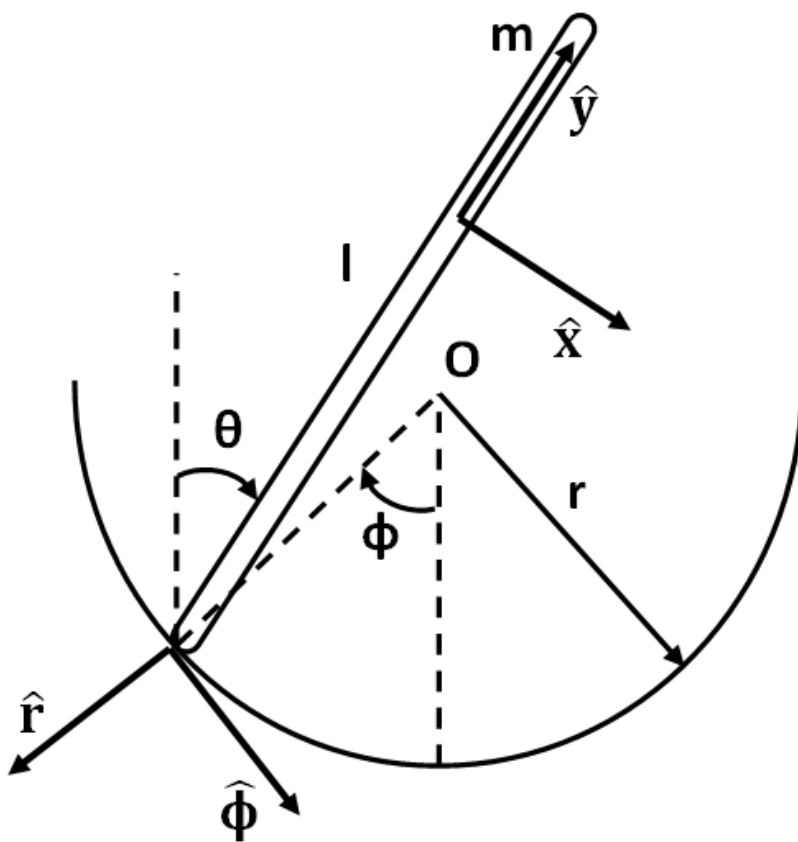
$$\begin{aligned}
\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}} \right) - \frac{\partial L}{\partial x} &= 0 \\
\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} &= 0 \\
\Rightarrow \text{x-EOM: } 0 &= \frac{3}{2} m \ddot{x} + \frac{3}{2} m r \ddot{\theta} - m x \dot{\theta}^2 - m g \sin \theta \\
\theta\text{-EOM: } 0 &= \frac{41}{6} m r^2 \ddot{\theta} + \frac{3}{2} m r \ddot{x} + m x^2 \ddot{\theta} + 2 m x \dot{x} \dot{\theta} - m g (x + 2r) \cos \theta - m g r \sin \theta
\end{aligned}$$

b) at  $t = 0^+$ , we know  $\theta(0) = 0$ ,  $x(0) = 2r$ ,  $\dot{\theta}(0) = 0$ ,  $\dot{x}(0) = 0$ :

$$\begin{aligned}
\text{x-EOM} &\rightarrow \ddot{x} = -\ddot{\theta} r \\
\theta\text{-EOM} &\rightarrow \left( \frac{41}{6} r^2 + 4r^2 - \frac{3}{2} r^2 \right) \ddot{\theta} = 4gr \\
&\Rightarrow \ddot{\theta} = \frac{3}{7} \frac{g}{r} \\
&\Rightarrow \frac{M_{total}}{I_{end}} = \ddot{\theta} = \frac{3}{7} \frac{g}{r} \\
\frac{2mgr + 2F_{disk}r}{\frac{16}{3}mr^2} &= \frac{3g}{7r} \\
&\Rightarrow F_{disk} = \frac{1}{7}mg
\end{aligned}$$

Notice that we have implicitly defined the downward direction as the positive direction of force. In inertial frame we would have  $\vec{F}_{disk} = \left(-\frac{1}{7}mg\right) \hat{y}$ .

## 2 Problem 2



Find position and velocity vectors for COM of rod:

$$\vec{r}_{CM} = r\hat{r} + \frac{l}{2}\hat{y}$$

$$\vec{v}_{CM} = -r\dot{\phi}\hat{\phi} + \frac{1}{2}\dot{\theta}\hat{x}$$

Form Lagrangian:

$$\begin{aligned}
T &= \frac{1}{2}m\vec{v}_{CM} \cdot \vec{v}_{CM} + \frac{1}{2}I\omega^2 \\
&= \frac{1}{2}m \left( r^2\dot{\phi}^2 + \frac{1}{4}l^2\dot{\theta}^2 - r l \dot{\theta} \dot{\phi} \cos(\phi - \theta) \right) + \frac{1}{2} \frac{1}{12} m l^2 \dot{\theta}^2 \\
&= \frac{1}{6} m l^2 \dot{\theta}^2 + \frac{1}{2} m r^2 \dot{\phi}^2 - \frac{1}{2} m r l \dot{\theta} \dot{\phi} \cos(\phi - \theta)
\end{aligned}$$

$$\begin{aligned}
V &= m g \left( -r \cos \phi + \frac{l}{2} \cos \theta \right) \\
L &= \frac{1}{6} m l^2 \dot{\theta}^2 + \frac{1}{2} m r^2 \dot{\phi}^2 - \frac{1}{2} m r l \dot{\theta} \dot{\phi} \cos(\phi - \theta) + m g r \cos \phi - m g \frac{l}{2} \cos \theta
\end{aligned}$$

Lagrangian Math:

$$\begin{aligned}
\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) &= \frac{d}{dt} \left( \frac{1}{3} m l^2 \dot{\theta} - \frac{1}{2} m r l \dot{\phi} \cos(\phi - \theta) \right) \\
&= \frac{1}{3} m l^2 \ddot{\theta} - \frac{1}{2} m r l \left( \ddot{\phi} \cos(\phi - \theta) - \dot{\phi}(\dot{\phi} - \dot{\theta}) \sin(\phi - \theta) \right) \\
\frac{\partial L}{\partial \theta} &= -\frac{1}{2} m r l \dot{\theta} \dot{\phi} \sin(\phi - \theta) + m g \frac{l}{2} \sin \theta
\end{aligned}$$

$$\begin{aligned}
\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\phi}} \right) &= \frac{d}{dt} \left( m r^2 \dot{\phi} - \frac{1}{2} m r l \dot{\theta} \cos(\phi - \theta) \right) \\
&= m r^2 \ddot{\phi} - \frac{1}{2} m r l \ddot{\theta} \cos(\phi - \theta) + \frac{1}{2} m r l \dot{\theta}(\dot{\phi} - \dot{\theta}) \sin(\phi - \theta) \\
\frac{\partial L}{\partial \phi} &= \frac{1}{2} m r l \dot{\theta} \dot{\phi} \sin(\phi - \theta) - m g r \sin \phi
\end{aligned}$$

Hence the final equations are:

$$\begin{aligned}
\frac{1}{3} m l^2 \ddot{\theta} - \frac{1}{2} m r l \ddot{\phi} \cos(\phi - \theta) + \frac{1}{2} m r l \dot{\phi}^2 \sin(\phi - \theta) - m g \frac{l}{2} \sin \theta &= 0 \\
m r^2 \ddot{\phi} - \frac{1}{2} m r l \ddot{\theta} \cos(\phi - \theta) - \frac{1}{2} m r l \dot{\theta}^2 \sin(\phi - \theta) + m g r \sin \phi &= 0
\end{aligned}$$

### 3 Problem 3

On the y-direction:

$$\begin{aligned} L_{disk} &= Iw \\ L_{balls} &= 2m(r + l \sin \theta)^2 w \end{aligned}$$

with no external torque acting on the system, the y angular momentum is conserved:

$$\begin{aligned} Iw_0 + 2mr^2w_0 &= Iw + 2m(r + l \sin \theta)^2 w \\ w &= \frac{(I + 2mr^2)w_0}{I + 2m(r + l \sin \theta)^2} \end{aligned}$$

When solving for  $\dot{\theta}$  in terms of  $\theta$  it is important to realize that as the balls swing up and down, the disk itself, due to conservation of y linear momentum, is moving as well! If we describe the location of the disk as  $[0; y]^T$  we have the following:

$$\begin{aligned} |\dot{x}_{ball}| &= \dot{\theta} l \cos \theta \\ |\dot{y}_{ball}| &= \dot{y} + \dot{\theta} l \sin \theta \\ \text{conservation of } \vec{p}_y &\Rightarrow 0 = 2m\dot{y} + 2m(\dot{y} + l\dot{\theta} \sin \theta) \\ \dot{y} &= -\frac{1}{2}l\dot{\theta} \sin \theta \end{aligned}$$

Finally we apply conservation of energy to the entire system:

$$\frac{1}{2}Iw_0^2 + mr^2w_0^2 = m\dot{y}^2 + \frac{1}{2}Iw^2 + m[(\dot{y} + l\dot{\theta} \sin \theta)^2 + (l\dot{\theta} \cos \theta)^2 + (r + l \sin \theta)^2 w^2]$$

substituting in the previous expression for  $w$  and isolate  $\dot{\theta}$ :

$$\dot{\theta} = \omega_0 \sqrt{\frac{I + 2mr^2}{ml^2(1 + \cos^2 \theta)} \left[ 1 - \frac{I + 2mr^2}{I + 2m(r + l \sin \theta)^2} \right]}$$

## 4 Problem 4

a) We just need to find the impulsive moment imparted to the gimbal to make it rotate with the same angular velocity as the rocket:

$$\begin{aligned}
 \vec{Y}_g &= \begin{bmatrix} 0.05 & 0 & 0 \\ 0 & 0.05 & 0 \\ 0 & 0 & 0.1 \end{bmatrix} \\
 \vec{H}_g &= \vec{Y}_g (\vec{\omega}_f - \vec{\omega}_i) \\
 &= \begin{bmatrix} 0.05 & 0 & 0 \\ 0 & 0.05 & 0 \\ 0 & 0 & 0.1 \end{bmatrix} \left( \begin{bmatrix} 10 \\ 0 \\ 0 \end{bmatrix} - \begin{bmatrix} 0 \\ 0 \\ 10^4 \end{bmatrix} \right) \\
 &= \begin{bmatrix} 0.5 \\ 0 \\ -1000 \end{bmatrix}
 \end{aligned}$$

The rocket experiences the negative of this, so:

$$\Delta \vec{H} = -0.5\hat{x} + 1000\hat{z}$$

b) Formulate this the same way as part a:

$$\begin{aligned}
 \vec{Y}_r &= \begin{bmatrix} 10^3 & 0 & 0 \\ 0 & 10^4 & 0 \\ 0 & 0 & 10^4 \end{bmatrix} \\
 \vec{H}_{rf} &= \vec{H}_{ri} + \Delta \vec{H} \\
 &= \begin{bmatrix} 10^3 & 0 & 0 \\ 0 & 10^4 & 0 \\ 0 & 0 & 10^4 \end{bmatrix} \begin{bmatrix} 10 \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} -0.5 \\ 0 \\ 1000 \end{bmatrix} \\
 &= \begin{bmatrix} 9999.5 \\ 0 \\ 1000 \end{bmatrix}
 \end{aligned}$$

We can then find the angle between the original angular momentum and the new angular momentum:

$$\begin{aligned}
 \beta &= \tan^{-1} \frac{H_z}{H_x} = \tan^{-1} \frac{1000}{9999.5} \\
 &= 5.71^\circ \\
 2\beta &= 11.42^\circ
 \end{aligned}$$

c) We want to use the equation  $\dot{\psi} = -\frac{H}{I_t}$  to solve for the precession rate. With this in mind we need:

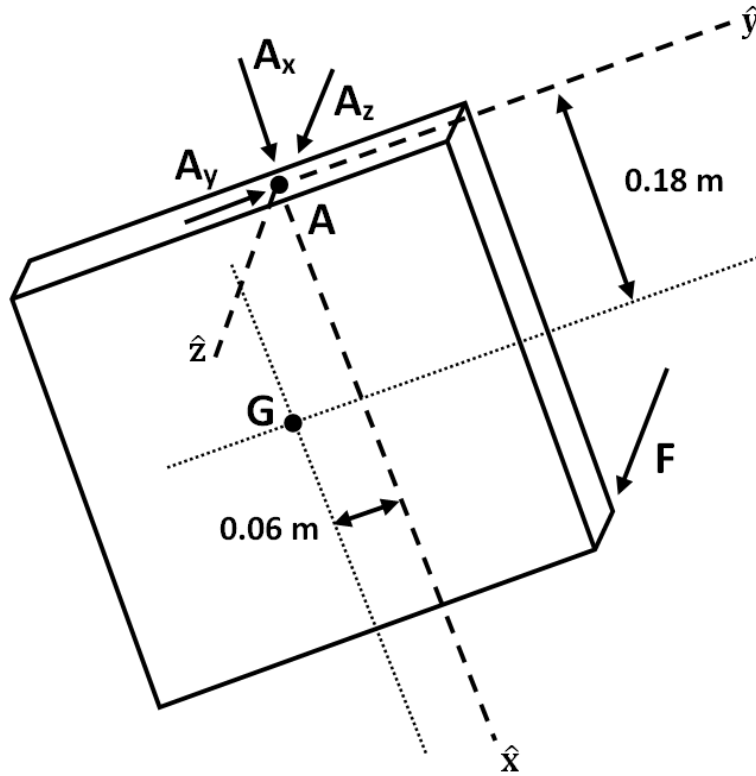
$$H = \sqrt{(9999.5)^2 + (1000)^2} = 10050$$

So then we have:

$$\dot{\psi} = -\frac{10050}{10000} = -1.005 \frac{rad}{s}$$
$$T = \frac{2\pi}{|\dot{\psi}|} = 6.25s$$

## 5 Problem 5

Original problem from Ginsberg: Advanced Engineering Dynamics 2nd edition



a) The force  $\vec{F}$  is much larger than the weight of the plate, so the latter is omitted from the free-body diagram. In contrast, the reaction exerted by the ball-and-socket joint is impulsive, because it must be as large as necessary to prevent movement of point A. We place the origin of  $xyz$  at point A in order to eliminate the angular impulse of this reaction. The coordinates of point A relative to parallel centroidal axes are  $(-0.18, 0.06, 0)$  meters, so the inertia properties are

$$I_{xx} = \frac{1}{12}(10)(0.36^2) + 10(0.06^2) = 0.144$$

$$I_{yy} = \frac{1}{12}(10)(0.36^2) + 10(0.18^2) = 0.432$$

$$I_{zz} = \frac{1}{12}(10)(0.36^2 + 0.36^2) + 10(0.18^2 + 0.06^2) = 0.576$$

$$I_{xy} = 0 + 10(-0.18)(0.06) = -0.108$$

$$I_{xz} = I_{yz} = 0$$



The angular velocity is initially zero. Let  $\vec{\omega}_2 = \omega_x \hat{x} + \omega_y \hat{y} + \omega_z \hat{z}$  denote the angular velocity at the termination of the impulsive action. Then the corresponding velocity of the center of mass is

$$\begin{aligned} (\vec{v}_G)_2 &= \vec{\omega}_2 \times \vec{r}_{G/A} \\ &= 0.06\omega_z \hat{x} + 0.18\omega_z \hat{y} - (0.06\omega_x + 0.18\omega_y) \hat{z} \end{aligned}$$

The final angular momentum about pivot A is

$$(\vec{H}_A)_2 = (0.144\omega_x + 0.108\omega_y) \hat{x} + (0.432\omega_y + 0.108\omega_x) \hat{y} + 0.576\omega_z \hat{z}$$

Applying the angular impulse-momentum principle to the 4-ms interval of the force leads to

$$\begin{aligned} (\vec{H}_A)_2 &= (\vec{r}_{F/A} \times \vec{F}) \Delta t = [(0.36\hat{x} + 0.12\hat{y}) \times 5000\hat{z}](0.004) \\ &= 2.4\hat{x} - 7.2\hat{z} \end{aligned}$$

The result of matching like components of  $(\vec{H}_A)_2$  is

$$\begin{aligned} (\vec{H}_A)_2 \cdot \hat{x} &= 0.144\omega_x + 0.108\omega_y = 2.4 \\ (\vec{H}_A)_2 \cdot \hat{y} &= 0.432\omega_y + 0.108\omega_x = -7.2 \\ (\vec{H}_A)_2 \cdot \hat{z} &= 0.576\omega_z = 0 \end{aligned}$$

from which we obtain

$$\vec{\omega}_2 = 35.90\hat{x} - 25.64\hat{y} \text{ rad/s}$$

b) Form the linear impulse-momentum principle in order to determine the reaction. Using the earlier expression for  $(\vec{v}_G)_2$  leads to

$$\begin{aligned} m(\vec{v}_G)_2 &= (\vec{A} + \vec{F}) \Delta t \\ 10[0.06\omega_z \hat{x} + 0.18\omega_z \hat{y} - (0.06\omega_x + 0.18\omega_y) \hat{z}] &= [A_x \hat{x} + A_y \hat{y} + (A_z + F) \hat{z}] \Delta t \end{aligned}$$

After substitution of the result for the angular velocity  $\vec{\omega}_2$ , the components of this equation yield

$$\begin{aligned} A_x &= A_y = 0 \\ A_z &= -F - \frac{10(0.06\omega_x + 0.18\omega_y)}{\Delta t} = 1153 \text{ N} \end{aligned}$$