

Density-functional perturbation theory

forces, response functions, phonons, and all that

Stefano Baroni

Scuola Internazionale Superiore di Studi Avanzati
& DEMOCRITOS National Simulation Center
Trieste - Italy

Energy derivatives

$$H = H_0 + \sum_i \lambda_i v_i$$

Energy derivatives

$$H = H_0 + \sum_i \lambda_i v_i$$

$$E[\lambda] = E_0 - \sum_i f_i \lambda_i + \frac{1}{2} \sum_{ij} h_{ij} \lambda_i \lambda_j + \dots$$

Energy derivatives

$$H = H_0 + \sum_i \lambda_i v_i$$

$$E[\lambda] = E_0 - \sum_i f_{\textcolor{blue}{i}} \lambda_i + \frac{1}{2} \sum_{ij} h_{ij} \lambda_i \lambda_j + \dots$$

- structural optimization & molecular dynamics

Energy derivatives

$$H = H_0 + \sum_i \lambda_i v_i$$

$$E[\lambda] = E_0 - \sum_i f_{\textcolor{blue}{i}} \lambda_i + \frac{1}{2} \sum_{ij} h_{ij} \lambda_i \lambda_j + \dots$$

- structural optimization & molecular dynamics
- (static) response functions
 - elastic constants
 - dielectric tensor
 - piezoelectric tensor
 - Born effective charges
 - ...

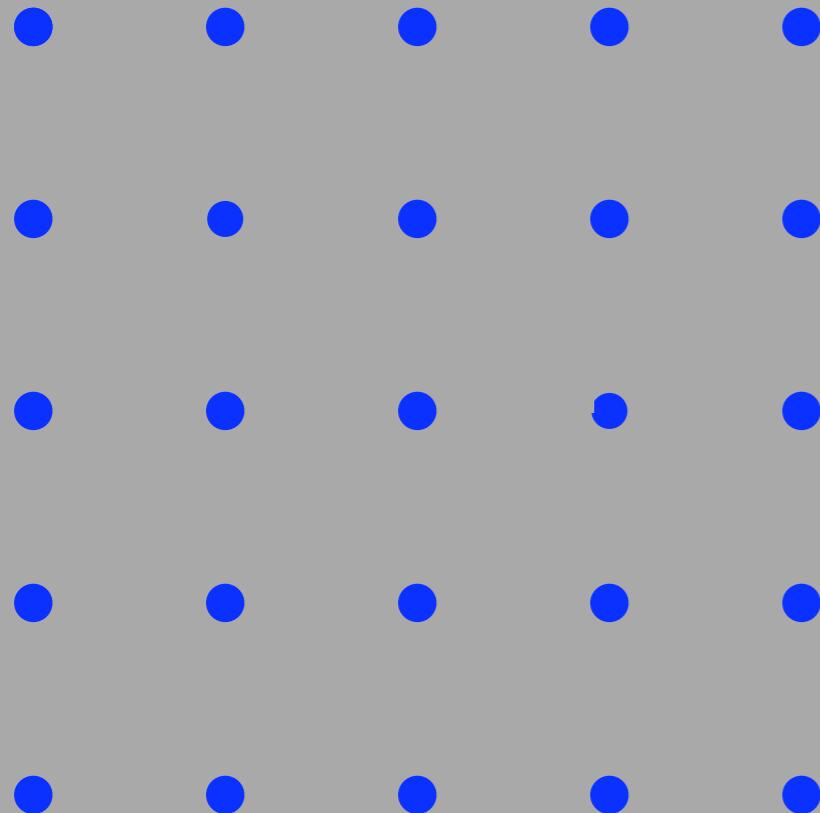
Energy derivatives

$$H = H_0 + \sum_i \lambda_i v_i$$

$$E[\lambda] = E_0 - \sum_i f_i \lambda_i + \frac{1}{2} \sum_{ij} h_{ij} \lambda_i \lambda_j + \dots$$

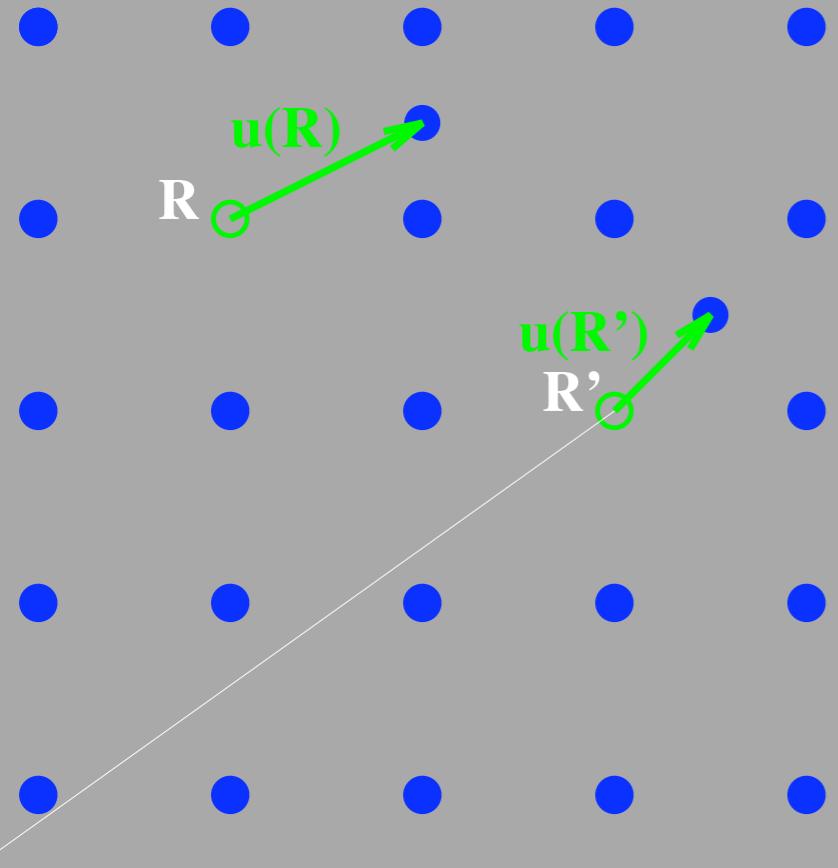
- structural optimization & molecular dynamics
- (static) response functions
 - elastic constants
 - dielectric tensor
 - piezoelectric tensor
 - Born effective charges
 - ...
- vibrational modes in the adiabatic approximaton

Lattice dynamics



$$\begin{aligned} V(\mathbf{r}) &= V_0(\mathbf{r}) \\ &= \sum_{\mathbf{R}} v(\mathbf{r} - \mathbf{R}) \\ E &= E_0 \end{aligned}$$

Lattice dynamics



$$\begin{aligned} V(\mathbf{r}) &= V_0(\mathbf{r}) \\ &\quad + \sum_{\mathbf{R}} \mathbf{u}(\mathbf{R}) \cdot \frac{\partial v(\mathbf{r} - \mathbf{R})}{\partial \mathbf{R}} \\ E &= E_0 \\ &\quad + \frac{1}{2} \sum_{\mathbf{R}, \mathbf{R}'} \mathbf{u}(\mathbf{R}) \cdot \frac{\partial^2 E}{\partial \mathbf{u}(\mathbf{R}) \partial \mathbf{u}(\mathbf{R}')} \cdot \mathbf{u}(\mathbf{R}') \\ &\quad + \dots \end{aligned}$$

Energy derivatives & perturbation theory

$$H = H_0 + \sum_i \lambda_i v_i \quad E[\lambda] = E_0 - \sum_i f_i \lambda_i + \frac{1}{2} \sum_{ij} h_{ij} \lambda_i \lambda_j + \dots$$

Energy derivatives & perturbation theory

$$H = H_0 + \sum_i \lambda_i v_i \quad E[\lambda] = E_0 - \sum_i f_i \lambda_i + \frac{1}{2} \sum_{ij} h_{ij} \lambda_i \lambda_j + \dots$$

$$f_i = - \left. \frac{\partial E}{\partial \lambda_i} \right|_{\lambda=0} = - \langle \Psi_0 | v_i | \Psi_0 \rangle$$

Energy derivatives & perturbation theory

$$H = H_0 + \sum_i \lambda_i v_i \quad E[\lambda] = E_0 - \sum_i f_i \lambda_i + \frac{1}{2} \sum_{ij} h_{ij} \lambda_i \lambda_j + \dots$$

$$f_i = - \left. \frac{\partial E}{\partial \lambda_i} \right|_{\lambda=0} = - \langle \Psi_0 | v_i | \Psi_0 \rangle$$

$$h_{ij} = \left. \frac{\partial^2 E}{\partial \lambda_i \partial \lambda_j} \right|_{\lambda=0} = 2 \sum_n \frac{\langle \Psi_0 | v_i | \Psi_n \rangle \langle \Psi_n | v_j | \Psi_0 \rangle}{\epsilon_0 - \epsilon_n}$$

Energy derivatives & perturbation theory

$$H = H_0 + \sum_i \lambda_i v_i \quad E[\lambda] = E_0 - \sum_i f_i \lambda_i + \frac{1}{2} \sum_{ij} h_{ij} \lambda_i \lambda_j + \dots$$

$$f_i = - \left. \frac{\partial E}{\partial \lambda_i} \right|_{\lambda=0} = - \langle \Psi_0 | v_i | \Psi_0 \rangle$$

$$\begin{aligned} h_{ij} &= \left. \frac{\partial^2 E}{\partial \lambda_i \partial \lambda_j} \right|_{\lambda=0} = 2 \sum_n \frac{\langle \Psi_0 | v_i | \Psi_n \rangle \langle \Psi_n | v_j | \Psi_0 \rangle}{\epsilon_0 - \epsilon_n} \\ &= 2 \langle \Psi_0 | v_i | \Psi_0'^j \rangle \end{aligned}$$

Energy derivatives & perturbation theory

$$H = H_0 + \sum_i \lambda_i v_i \quad E[\lambda] = E_0 - \sum_i f_i \lambda_i + \frac{1}{2} \sum_{ij} h_{ij} \lambda_i \lambda_j + \dots$$

$$f_i = - \left. \frac{\partial E}{\partial \lambda_i} \right|_{\lambda=0} = - \langle \Psi_0 | v_i | \Psi_0 \rangle$$

$$\begin{aligned} h_{ij} &= \left. \frac{\partial^2 E}{\partial \lambda_i \partial \lambda_j} \right|_{\lambda=0} = 2 \sum_n \frac{\langle \Psi_0 | v_i | \Psi_n \rangle \langle \Psi_n | v_j | \Psi_0 \rangle}{\epsilon_0 - \epsilon_n} \\ &= 2 \langle \Psi_0 | v_i | \Psi_0'^j \rangle \\ &= 2 \langle \Psi_i' | v_j | \Psi_0 \rangle \end{aligned}$$

The Hellmann-Feynman theorem

$$E(\lambda) = \min \langle \Psi | H(\lambda) | \Psi \rangle$$
$$\langle \Psi | \Psi \rangle = 1$$

The Hellmann-Feynman theorem

$$E(\lambda) = \min \langle \Psi | H(\lambda) | \Psi \rangle$$
$$\langle \Psi | \Psi \rangle = 1$$

$$g(\lambda) = \min_x G[x, \lambda]$$

The Hellmann-Feynman theorem

$$E(\lambda) = \min \langle \Psi | H(\lambda) | \Psi \rangle$$

$$\langle \Psi | \Psi \rangle = 1$$

$$g(\lambda) = \min_x G[x, \lambda] \quad \longleftarrow \quad \frac{\partial G}{\partial x} \Big|_{x=x(\lambda)} = 0$$

The Hellmann-Feynman theorem

$$E(\lambda) = \min \langle \Psi | H(\lambda) | \Psi \rangle$$

$$\langle \Psi | \Psi \rangle = 1$$

$$g(\lambda) = \min_x G[x, \lambda] \quad \longleftrightarrow \quad \left. \frac{\partial G}{\partial x} \right|_{x=x(\lambda)} = 0$$

$$g(\lambda) = G[x(\lambda), \lambda]$$

The Hellmann-Feynman theorem

$$E(\lambda) = \min \langle \Psi | H(\lambda) | \Psi \rangle$$

$$\langle \Psi | \Psi \rangle = 1$$

$$g(\lambda) = \min_x G[x, \lambda] \quad \longleftarrow \quad \frac{\partial G}{\partial x} \Big|_{x=x(\lambda)} = 0$$

$$g(\lambda) = G[x(\lambda), \lambda] \longrightarrow g'(\lambda) = x'(\lambda) \frac{\partial G}{\partial x} \Big|_{x=x(\lambda)} + \frac{\partial G}{\partial \lambda}$$

The Hellmann-Feynman theorem

$$E(\lambda) = \min \langle \Psi | H(\lambda) | \Psi \rangle$$

$$\langle \Psi | \Psi \rangle = 1$$

$$g(\lambda) = \min_x G[x, \lambda] \quad \longleftrightarrow \quad \boxed{\frac{\partial G}{\partial x} \Big|_{x=x(\lambda)} = 0}$$
$$g(\lambda) = G[x(\lambda), \lambda] \longrightarrow g'(\lambda) = x'(\lambda) \cancel{\frac{\partial G}{\partial x} \Big|_{x=x(\lambda)}} + \frac{\partial G}{\partial \lambda}$$

The Hellmann-Feynman theorem

$$E(\lambda) = \min \langle \Psi | H(\lambda) | \Psi \rangle$$

$$\langle \Psi | \Psi \rangle = 1$$

$$g(\lambda) = \min_x G[x, \lambda] \quad \longleftrightarrow \quad \boxed{\frac{\partial G}{\partial x} \Big|_{x=x(\lambda)} = 0}$$
$$g(\lambda) = G[x(\lambda), \lambda] \longrightarrow g'(\lambda) = x'(\lambda) \cancel{\frac{\partial G}{\partial x} \Big|_{x=x(\lambda)}} + \frac{\partial G}{\partial \lambda}$$

$$E'(\lambda) = \langle \Psi_\lambda | H'(\lambda) | \Psi_\lambda \rangle$$

The “2n+1” theorem

$$\Phi = \Phi_0 + \mathcal{O}(\lambda) \Rightarrow E = E_0 + \mathcal{O}(\lambda^2)$$

The “2n+1” theorem

$$\Phi = \Phi_0 + \mathcal{O}(\lambda) \Rightarrow E = E_0 + \mathcal{O}(\lambda^2)$$

$$\Phi = \Phi_0 + \sum_{l=1}^n \lambda^l \Phi^{(l)} + \mathcal{O}(\lambda^{n+1})$$

The “2n+1” theorem

$$\Phi = \Phi_0 + \mathcal{O}(\lambda) \Rightarrow E = E_0 + \mathcal{O}(\lambda^2)$$

$$\Phi = \Phi_0 + \sum_{l=1}^n \lambda^l \Phi^{(l)} + \mathcal{O}(\lambda^{n+1}) \Rightarrow$$

$$E = E_0 + \sum_{l=1}^{2n+1} \lambda^l E^{(l)} + \mathcal{O}(\lambda^{2n+2})$$

The “2n+1” theorem

$$\Phi = \Phi_0 + \mathcal{O}(\lambda) \Rightarrow E = E_0 + \mathcal{O}(\lambda^2)$$

$$\Phi = \Phi_0 + \sum_{l=1}^n \lambda^l \Phi^{(l)} + \mathcal{O}(\lambda^{n+1}) \Rightarrow$$

$$E = E_0 + \sum_{l=1}^{2n+1} \lambda^l E^{(l)} + \mathcal{O}(\lambda^{2n+2})$$

$$E = \frac{\langle \Phi_0 + \Phi' | (H_0 + V') | \Phi_0 + \Phi' \rangle}{\langle \Phi_0 + \Phi' | \Phi_0 + \Phi' \rangle} + \mathcal{O}(V'^4)$$

The “2n+1” theorem

$$\Phi = \Phi_0 + \mathcal{O}(\lambda) \Rightarrow E = E_0 + \mathcal{O}(\lambda^2)$$

$$\Phi = \Phi_0 + \sum_{l=1}^n \lambda^l \Phi^{(l)} + \mathcal{O}(\lambda^{n+1}) \Rightarrow$$

$$E = E_0 + \sum_{l=1}^{2n+1} \lambda^l E^{(l)} + \mathcal{O}(\lambda^{2n+2})$$

$$E = \frac{\langle \Phi_0 + \Phi' | (H_0 + V') | \Phi_0 + \Phi' \rangle}{\langle \Phi_0 + \Phi' | \Phi_0 + \Phi' \rangle} + \mathcal{O}(V'^4)$$

$$E^{(3)} = \langle \Phi' | V' | \Phi' \rangle - \langle \Phi' | \Phi' \rangle \langle \Phi_0 | V' | \Phi_0 \rangle$$

Density-functional perturbation theory

$$V(\mathbf{r}) = V_0(\mathbf{r}) + \sum_i \lambda_i V'_i(\mathbf{r})$$

Density-functional perturbation theory

$$V(\mathbf{r}) = V_0(\mathbf{r}) + \sum_i \lambda_i V'_i(\mathbf{r})$$

$$E(\lambda) = \min_n \left(F[n] + \int V_\lambda(\mathbf{r}) n(\mathbf{r}) \right) \quad \int n(\mathbf{r}) d\mathbf{r} = N \quad \text{DFT}$$

Density-functional perturbation theory

$$V(\mathbf{r}) = V_0(\mathbf{r}) + \sum_i \lambda_i V'_i(\mathbf{r})$$

$$E(\lambda) = \min_n \left(F[n] + \int V_\lambda(\mathbf{r}) n(\mathbf{r}) \right) \quad \int n(\mathbf{r}) d\mathbf{r} = N \quad \text{DFT}$$

$$\frac{\partial E(\lambda)}{\partial \lambda_i} = \int n_\lambda(\mathbf{r}) V'_i(\mathbf{r}) d\mathbf{r} \quad \text{HF}$$

Density-functional perturbation theory

$$V(\mathbf{r}) = V_0(\mathbf{r}) + \sum_i \lambda_i V'_i(\mathbf{r})$$

$$E(\lambda) = \min_n \left(F[n] + \int V_\lambda(\mathbf{r}) n(\mathbf{r}) \right) \quad \int n(\mathbf{r}) d\mathbf{r} = N \quad \text{DFT}$$

$$\frac{\partial E(\lambda)}{\partial \lambda_i} = \int n_\lambda(\mathbf{r}) V'_i(\mathbf{r}) d\mathbf{r} \quad \text{HF}$$

$$\frac{\partial^2 E(\lambda)}{\partial \lambda_i \partial \lambda_j} = \int \frac{\partial n_\lambda(\mathbf{r})}{\partial \lambda_j} V'_i(\mathbf{r}) d\mathbf{r}$$

DFPT

DFPT II: the equations

DFT

$$V_0(\mathbf{r}) \leftrightarrows n_0(\mathbf{r})$$

DFPT II: the equations

DFT

$$V_0(\mathbf{r}) \leftrightharpoons n_0(\mathbf{r})$$

$$V_{SCF}(\mathbf{r}) = V_0(\mathbf{r}) + \int \frac{n_0(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu_{xc}(\mathbf{r})$$

DFPT II: the equations

DFT

$$V_0(\mathbf{r}) \rightleftharpoons n_0(\mathbf{r})$$

$$V_{SCF}(\mathbf{r}) = V_0(\mathbf{r}) + \int \frac{n_0(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu_{xc}(\mathbf{r})$$

↓

$$(-\Delta + V_{SCF}(\mathbf{r}))\phi_v(\mathbf{r}) = \epsilon_v \phi_v(\mathbf{r})$$

DFPT II: the equations

DFT

$$V_0(\mathbf{r}) \Leftarrow n_0(\mathbf{r})$$

$$V_{SCF}(\mathbf{r}) = V_0(\mathbf{r}) + \int \frac{n_0(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu_{xc}(\mathbf{r})$$

$$(-\Delta + V_{SCF}(\mathbf{r}))\phi_v(\mathbf{r}) = \epsilon_v \phi_v(\mathbf{r})$$

$$n_0(\mathbf{r}) = \sum_{\epsilon_v < E_F} |\phi_v(\mathbf{r})|^2$$

DFPT II: the equations

DFT

DFPT

$$V_0(\mathbf{r}) \rightleftharpoons n_0(\mathbf{r})$$

$$V'(\mathbf{r}) \rightleftharpoons n'(\mathbf{r})$$

$$V_{SCF}(\mathbf{r}) = V_0(\mathbf{r}) + \int \frac{n_0(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu_{xc}(\mathbf{r})$$

$$(-\Delta + V_{SCF}(\mathbf{r}))\phi_v(\mathbf{r}) = \epsilon_v \phi_v(\mathbf{r})$$

$$n_0(\mathbf{r}) = \sum_{\epsilon_v < E_F} |\phi_v(\mathbf{r})|^2$$

DFPT II: the equations

DFT

DFPT

$$V_0(\mathbf{r}) \rightleftharpoons n_0(\mathbf{r})$$

$$V'(\mathbf{r}) \rightleftharpoons n'(\mathbf{r})$$

$$V_{SCF}(\mathbf{r}) = V_0(\mathbf{r}) + \int \frac{n_0(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu_{xc}(\mathbf{r})$$
$$\downarrow$$
$$(-\Delta + V_{SCF}(\mathbf{r}))\phi_v(\mathbf{r}) = \epsilon_v \phi_v(\mathbf{r})$$
$$\downarrow$$
$$n_0(\mathbf{r}) = \sum_{\epsilon_v < E_F} |\phi_v(\mathbf{r})|^2$$
$$\longrightarrow$$
$$V'_{SCF}(\mathbf{r}) = V'(\mathbf{r}) + \int \frac{n'(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu'_{xc}(\mathbf{r})$$

DFPT II: the equations

DFT

DFPT

$$V_0(\mathbf{r}) \rightleftharpoons n_0(\mathbf{r})$$

$$V'(\mathbf{r}) \rightleftharpoons n'(\mathbf{r})$$

$$V_{SCF}(\mathbf{r}) = V_0(\mathbf{r}) + \int \frac{n_0(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu_{xc}(\mathbf{r})$$

$$(-\Delta + V_{SCF}(\mathbf{r}))\phi_v(\mathbf{r}) = \epsilon_v \phi_v(\mathbf{r})$$

$$n_0(\mathbf{r}) = \sum_{\epsilon_v < E_F} |\phi_v(\mathbf{r})|^2$$

$$V'_{SCF}(\mathbf{r}) = V'(\mathbf{r}) + \int \frac{n'(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu'_{xc}(\mathbf{r})$$

$$(-\Delta + V_{SCF}(\mathbf{r}) - \epsilon_v)\phi'_v(\mathbf{r}) = -P_c V'_{SCF} \phi_v(\mathbf{r})$$

DFPT II: the equations

DFT

DFPT

$$V_0(\mathbf{r}) \rightleftharpoons n_0(\mathbf{r})$$

$$V'(\mathbf{r}) \rightleftharpoons n'(\mathbf{r})$$

$$V_{SCF}(\mathbf{r}) = V_0(\mathbf{r}) + \int \frac{n_0(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu_{xc}(\mathbf{r})$$



$$(-\Delta + V_{SCF}(\mathbf{r}))\phi_v(\mathbf{r}) = \epsilon_v \phi_v(\mathbf{r})$$

$$n_0(\mathbf{r}) = \sum_{\epsilon_v < E_F} |\phi_v(\mathbf{r})|^2$$



$$V'_{SCF}(\mathbf{r}) = V'(\mathbf{r}) + \int \frac{n'(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + \mu'_{xc}(\mathbf{r})$$



$$(-\Delta + V_{SCF}(\mathbf{r}) - \epsilon_v)\phi'_v(\mathbf{r}) = -P_c V'_{SCF} \phi_v(\mathbf{r})$$



$$n'(\mathbf{r}) = 2 \operatorname{Re} \sum_{\epsilon_v < E_F} \phi_v^*(\mathbf{r}) \phi'_v(\mathbf{r})$$



More about perturbation theory

$$H = H_0 + V' \quad \left\{ \begin{array}{l} \phi_v^0 \rightarrow \phi_v^0 + \phi'_v \\ \epsilon_v^0 \rightarrow \epsilon_v^0 + \epsilon'_v \end{array} \right.$$

More about perturbation theory

$$H = H_0 + V' \quad \left\{ \begin{array}{l} \phi_v^0 \rightarrow \phi_v^0 + \phi'_v \\ \epsilon_v^0 \rightarrow \epsilon_v^0 + \epsilon'_v \end{array} \right.$$

$$\phi'_v = \sum_{u \neq v} \phi_u^0 \frac{\langle \phi_u^0 | V' | \phi_v^0 \rangle}{\epsilon_v^0 - \epsilon_u^0}$$

More about perturbation theory

$$H = H_0 + V' \quad \left\{ \begin{array}{l} \phi_v^0 \rightarrow \phi_v^0 + \phi'_v \\ \epsilon_v^0 \rightarrow \epsilon_v^0 + \epsilon'_v \end{array} \right.$$

$$\phi'_v = \sum_{u \neq v} \phi_u^0 \frac{\langle \phi_u^0 | V' | \phi_v^0 \rangle}{\epsilon_v^0 - \epsilon_u^0}$$

$$(H_0 - \epsilon_v^0) \phi'_v = - \sum_{u \neq v} \phi_u^0 \langle \phi_u^0 | V' | \phi_v^0 \rangle$$

More about perturbation theory

$$H = H_0 + V' \quad \left\{ \begin{array}{l} \phi_v^0 \rightarrow \phi_v^0 + \phi'_v \\ \epsilon_v^0 \rightarrow \epsilon_v^0 + \epsilon'_v \end{array} \right.$$

$$\phi'_v = \sum_{u \neq v} \phi_u^0 \frac{\langle \phi_u^0 | V' | \phi_v^0 \rangle}{\epsilon_v^0 - \epsilon_u^0}$$

$$\begin{aligned} (H_0 - \epsilon_v^0) \phi'_v &= - \sum_{u \neq v} \phi_u^0 \langle \phi_u^0 | V' | \phi_v^0 \rangle \\ &= -(1 - P_v) V' | \phi_v^0 \rangle \end{aligned}$$

More about perturbation theory

$$H = H_0 + V' \quad \left\{ \begin{array}{l} \phi_v^0 \rightarrow \phi_v^0 + \phi'_v \\ \epsilon_v^0 \rightarrow \epsilon_v^0 + \epsilon'_v \end{array} \right.$$

$$\phi'_v = \sum_{u \neq v} \phi_u^0 \frac{\langle \phi_u^0 | V' | \phi_v^0 \rangle}{\epsilon_v^0 - \epsilon_u^0}$$

$$(H_0 - \epsilon_v^0) \phi'_v = - \sum_{u \neq v} \phi_u^0 \langle \phi_u^0 | V' | \phi_v^0 \rangle = -(1 - P_v) V' | \phi_v^0 \rangle$$

$$(H_0 - \epsilon_v^0) P_c \phi'_v = -P_c V' | \phi_v^0 \rangle$$

More about perturbation theory

$$H = H_0 + V' \quad \left\{ \begin{array}{l} \phi_v^0 \rightarrow \phi_v^0 + \phi'_v \\ \epsilon_v^0 \rightarrow \epsilon_v^0 + \epsilon'_v \end{array} \right.$$

$$\phi'_v = \sum_{u \neq v} \phi_u^0 \frac{\langle \phi_u^0 | V' | \phi_v^0 \rangle}{\epsilon_v^0 - \epsilon_u^0}$$

$$(H_0 - \epsilon_v^0) \phi'_v = - \sum_{u \neq v} \phi_u^0 \langle \phi_u^0 | V' | \phi_v^0 \rangle$$

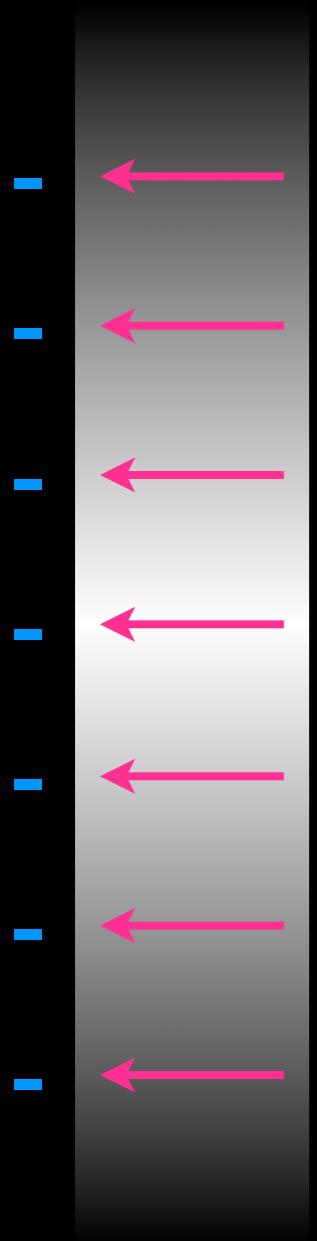
$$= -(1 - P_v) V' | \phi_v^0 \rangle$$

$$(H_0 - \epsilon_v^0) P_c \phi'_v = -P_c V' | \phi_v^0 \rangle$$

$$(H_0 - \epsilon_v^0 + \alpha P_v) \bar{\phi}'_v = -P_c V' | \phi_v^0 \rangle$$

Macroscopic electric fields

$$\vec{E} = \text{cnst}$$

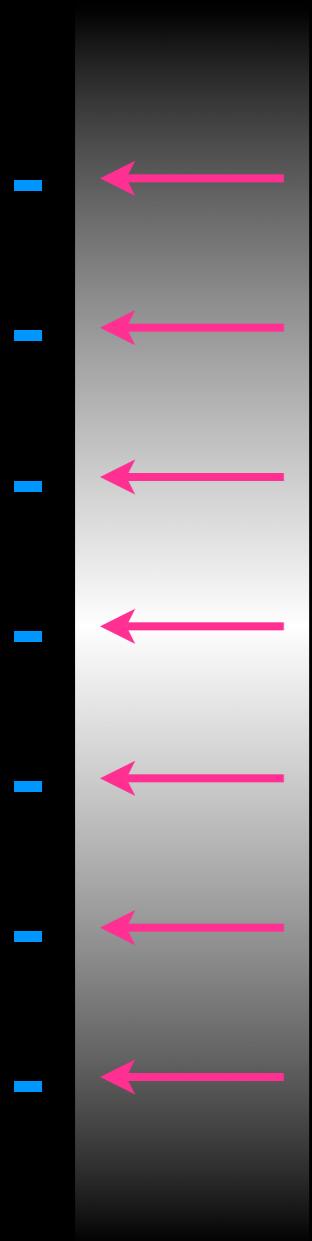


$$\nabla' = E \times$$

Macroscopic electric fields

$$\vec{E} = \text{cnst}$$

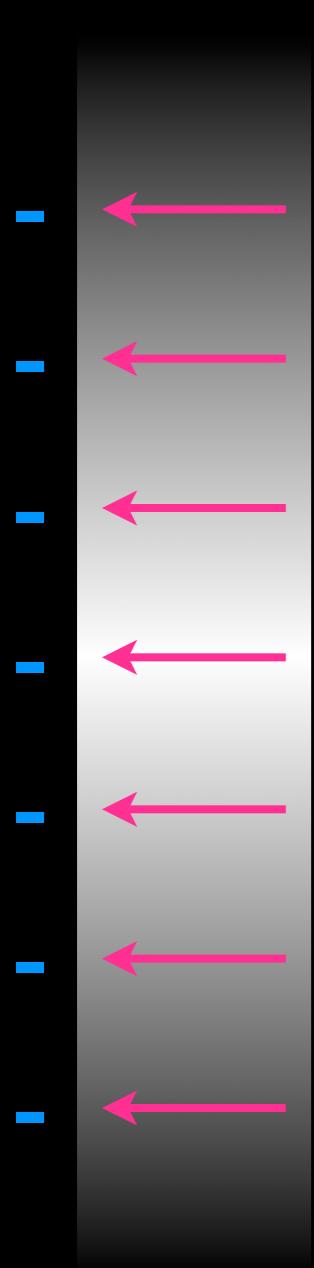
$$\phi_v^0(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}} u_{v,\mathbf{k}}(\mathbf{r})$$



$$\nabla' = \mathbf{E} \times$$

Macroscopic electric fields

$\vec{E} = \text{cnst}$



$$\begin{aligned}\phi_v^0(\mathbf{r}) &= e^{i\mathbf{k} \cdot \mathbf{r}} u_{v,\mathbf{k}}(\mathbf{r}) \\ V'(\mathbf{r}) \phi_v^0(\mathbf{r}) &= ??\end{aligned}$$

$\nabla' = \mathbf{E} \times$

Macroscopic electric fields

$\vec{E} = \text{cnst}$

$$\phi_v^0(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}} u_{v,\mathbf{k}}(\mathbf{r})$$
$$V'(\mathbf{r})\phi_v^0(\mathbf{r}) = ??$$
$$\langle \phi_v^0 | x | \phi_u^0 \rangle = \frac{\langle \phi_v^0 | [H, x] | \phi_u^0 \rangle}{\epsilon_v^0 - \epsilon_u^0} \quad [H, x] = -\frac{\hbar^2}{m} \frac{\partial}{\partial x}$$

$\nabla' = \mathbf{E} \times$

Macroscopic electric fields

$\vec{E} = \text{cnst}$

$$\phi_v^0(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}} u_{v,\mathbf{k}}(\mathbf{r})$$

$$V'(\mathbf{r})\phi_v^0(\mathbf{r}) = ??$$

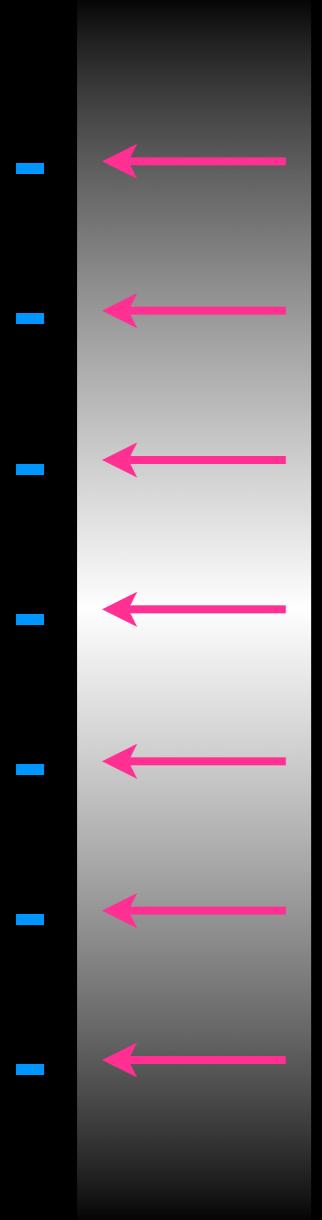
$$\langle \phi_v^0 | x | \phi_u^0 \rangle = \frac{\langle \phi_v^0 | [H, x] | \phi_u^0 \rangle}{\epsilon_v^0 - \epsilon_u^0} \quad [H, x] = -\frac{\hbar^2}{m} \frac{\partial}{\partial x}$$

$$-P_c V' \phi_v^0 = -E \sum_c \phi_c^0 \langle \phi_c^0 | x | \phi_v^0 \rangle$$

$\nabla' = E \times$

Macroscopic electric fields

$\vec{E} = \text{cnst}$



$$\phi_v^0(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}} u_{v,\mathbf{k}}(\mathbf{r})$$

$$V'(\mathbf{r})\phi_v^0(\mathbf{r}) = ??$$

$$\langle\phi_v^0|x|\phi_u^0\rangle = \frac{\langle\phi_v^0|[H,x]|\phi_u^0\rangle}{\epsilon_v^0 - \epsilon_u^0} \quad [H, x] = -\frac{\hbar^2}{m} \frac{\partial}{\partial x}$$

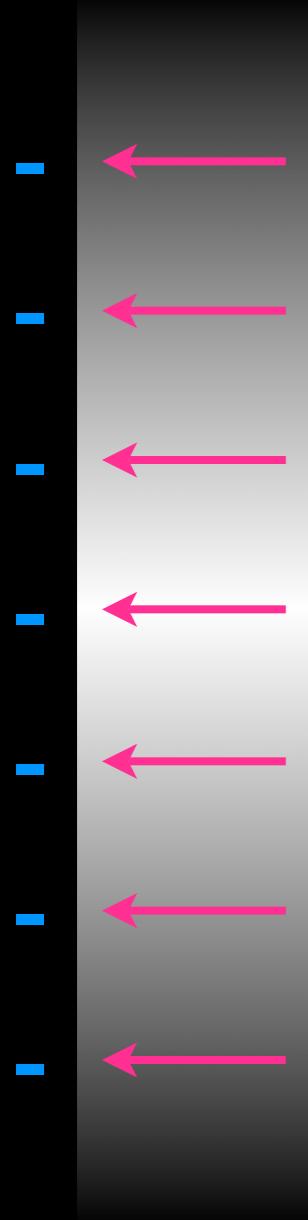
$$-P_c V' \phi_v^0 = -E \sum_c \phi_c^0 \langle\phi_c^0|x|\phi_v^0\rangle$$

$$= -E \sum_c \phi_c^0 \frac{\langle\phi_c^0|[H_0,x]|\phi_v^0\rangle}{\epsilon_c^0 - \epsilon_v^0} \equiv \psi'_v$$

$\nabla' = E \times$

Macroscopic electric fields

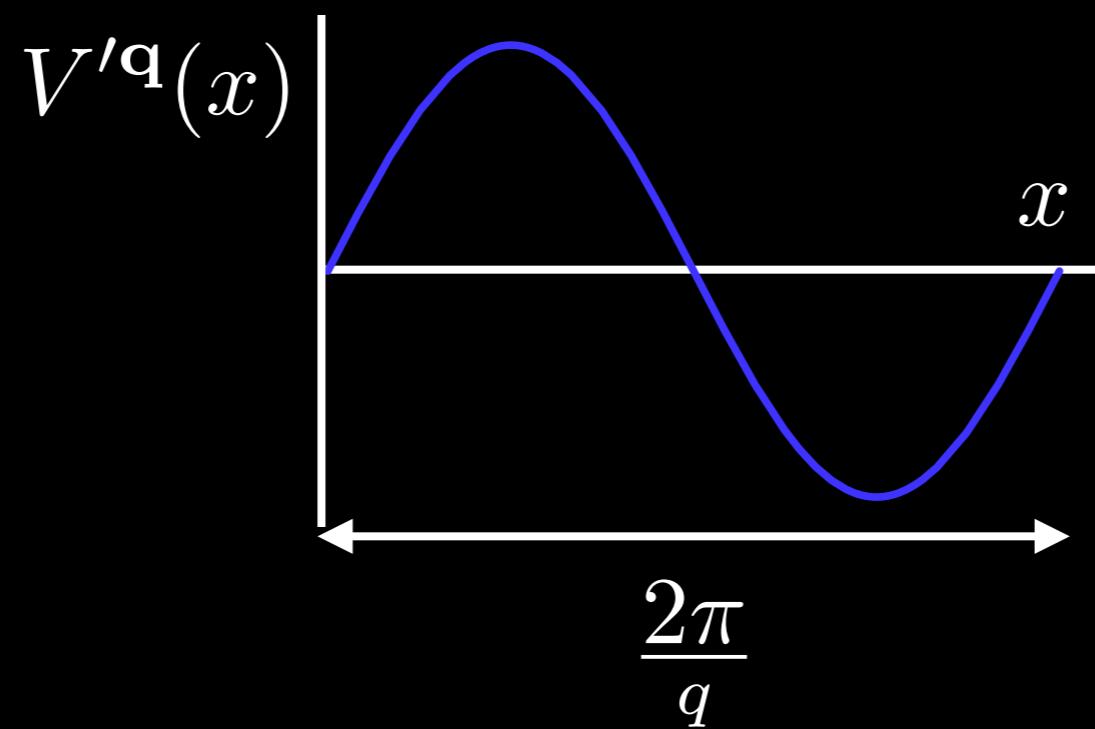
$\vec{E} = \text{cnst}$



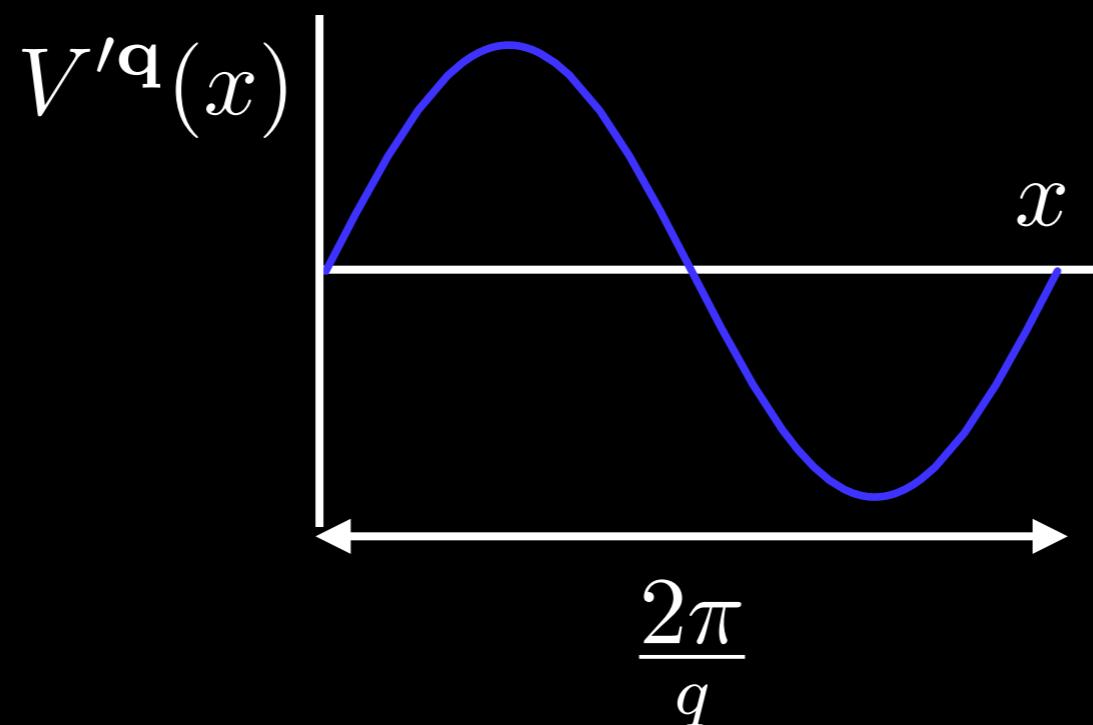
$$\begin{aligned} \phi_v^0(\mathbf{r}) &= e^{i\mathbf{k}\cdot\mathbf{r}} u_{v,\mathbf{k}}(\mathbf{r}) \\ V'(\mathbf{r})\phi_v^0(\mathbf{r}) &= ?? \\ \langle\phi_v^0|x|\phi_u^0\rangle &= \frac{\langle\phi_v^0|[H,x]|\phi_u^0\rangle}{\epsilon_v^0 - \epsilon_u^0} \quad [H, x] = -\frac{\hbar^2}{m} \frac{\partial}{\partial x} \\ -P_c V' \phi_v^0 &= -E \sum_c \phi_c^0 \langle\phi_c^0|x|\phi_v^0\rangle \\ &= -E \sum_c \phi_c^0 \frac{\langle\phi_c^0|[H_0,x]|\phi_v^0\rangle}{\epsilon_c^0 - \epsilon_v^0} \equiv \psi'_v \end{aligned}$$

$\nabla' = \mathbf{E} \times$ $(H_0 - \epsilon_c^0)\psi'_v = -EP_c[H_0,x]\phi_v^0$ DFPT RHS

Monochromatic perturbations

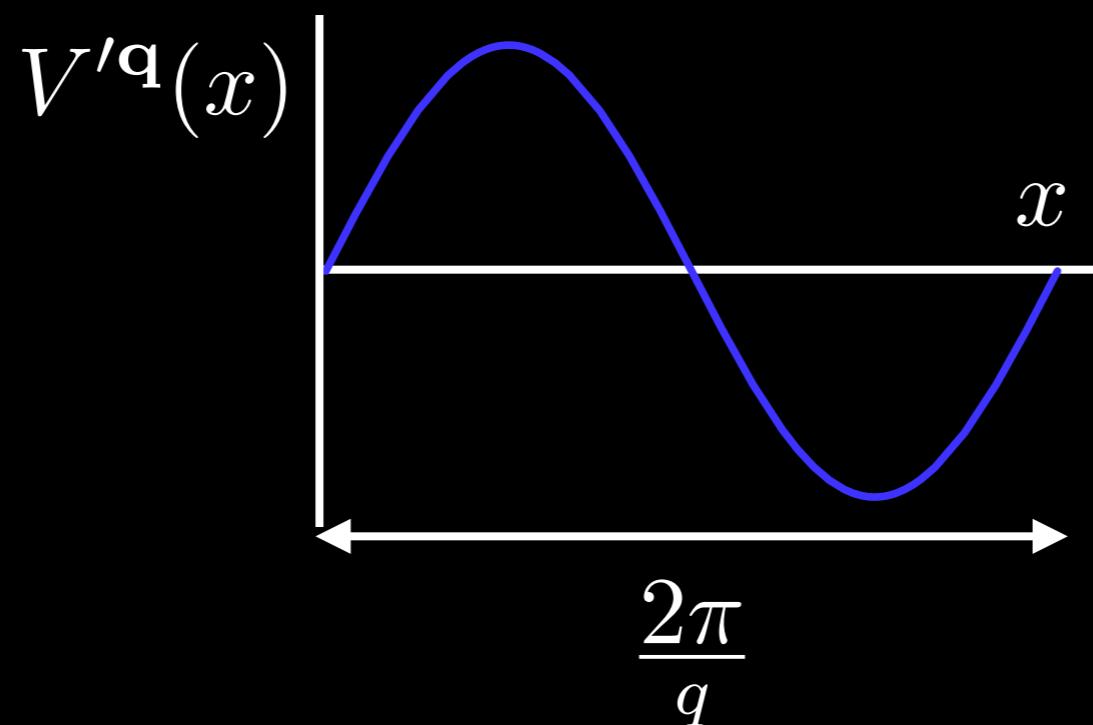


Monochromatic perturbations



$$(H_0 + \alpha P_v^{\mathbf{k}+\mathbf{q}} - \epsilon_v^{\mathbf{k}}) \phi_v'^{\mathbf{k}+\mathbf{q}}(\mathbf{r}) = -P_c V'^{\mathbf{q}} \phi_v^{\mathbf{k}}(\mathbf{r})$$

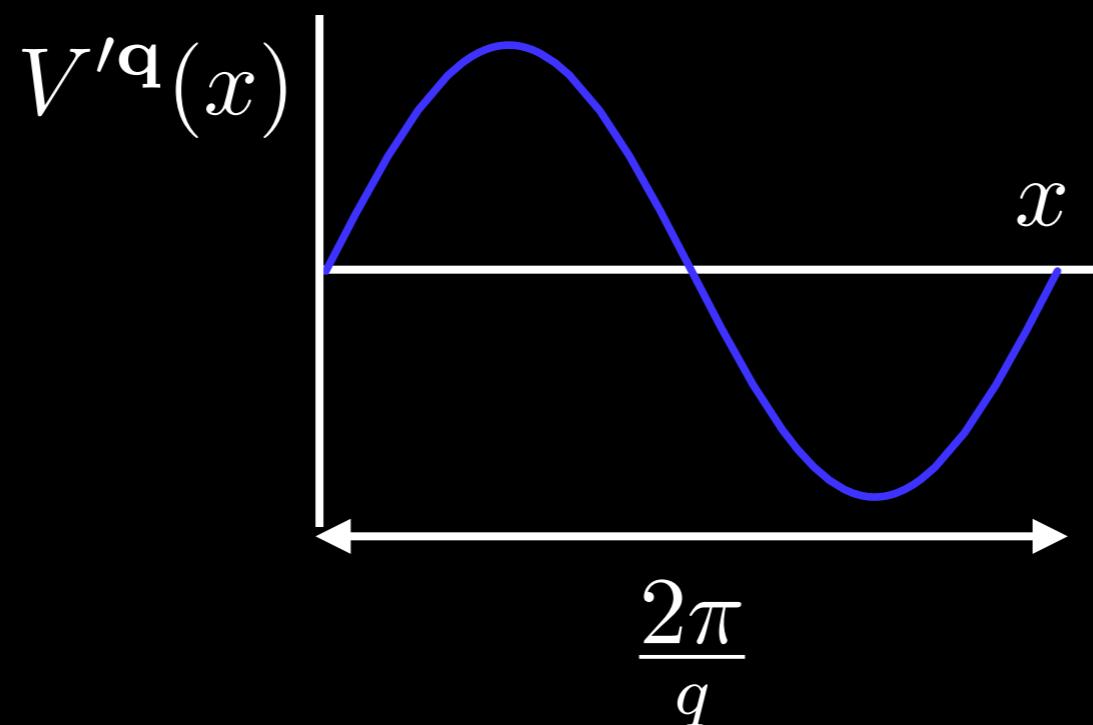
Monochromatic perturbations



$$(H_0 + \alpha P_v^{\mathbf{k}+\mathbf{q}} - \epsilon_v^{\mathbf{k}}) \phi_v'^{\mathbf{k}+\mathbf{q}}(\mathbf{r}) = -P_c V'^{\mathbf{q}} \phi_v^{\mathbf{k}}(\mathbf{r})$$

$$n'^{\mathbf{q}}(\mathbf{r}) = e^{i\mathbf{q}\cdot\mathbf{r}} \sum_{v,\mathbf{k}} u_v'^{\mathbf{k}*}(\mathbf{r}) u_v^{\mathbf{k}+\mathbf{q}}(\mathbf{r})$$

Monochromatic perturbations



$$(H_0 + \alpha P_v^{\mathbf{k}+\mathbf{q}} - \epsilon_v^{\mathbf{k}}) \phi_v'^{\mathbf{k}+\mathbf{q}}(\mathbf{r}) = -P_c V'^{\mathbf{q}} \phi_v^{\mathbf{k}}(\mathbf{r})$$

$$n'^{\mathbf{q}}(\mathbf{r}) = e^{i\mathbf{q}\cdot\mathbf{r}} \sum_{v,\mathbf{k}} u_v'^{\mathbf{k}*}(\mathbf{r}) u_v^{\mathbf{k}+\mathbf{q}}(\mathbf{r})$$

$$V'^{\mathbf{q}}(\mathbf{r}) = V'^{\mathbf{q}}_{ext}(\mathbf{r}) + \int \left(\frac{e^2}{|\mathbf{r} - \mathbf{r}'|} + \kappa_{xc}(\mathbf{r}, \mathbf{r}') \right) n'^{\mathbf{q}}(\mathbf{r}') d\mathbf{r}'$$

Phonons in polar materials

$$E(\mathbf{u}, \mathbf{E}) = \frac{1}{2} M \omega_0^2 u^2 - \frac{\Omega}{8\pi} \epsilon_\infty \mathbf{E}^2 - e Z^* \mathbf{u} \cdot \mathbf{E}$$

Phonons in polar materials

$$E(\mathbf{u}, \mathbf{E}) = \frac{1}{2} M \omega_0^2 u^2 - \frac{\Omega}{8\pi} \epsilon_\infty \mathbf{E}^2 - e Z^* \mathbf{u} \cdot \mathbf{E}$$

$$\mathbf{F} \equiv -\frac{\partial E}{\partial \mathbf{u}} = -M \omega_0^2 \mathbf{u} + Z^* \mathbf{E}$$

$$\mathbf{D} \equiv -\frac{4\pi}{\Omega} \frac{\partial E}{\partial \mathbf{E}} = \frac{4\pi}{\Omega} Z^* \mathbf{u} + \epsilon_\infty \mathbf{E}$$

Phonons in polar materials

$$E(\mathbf{u}, \mathbf{E}) = \frac{1}{2} M \omega_0^2 u^2 - \frac{\Omega}{8\pi} \epsilon_\infty \mathbf{E}^2 - e Z^* \mathbf{u} \cdot \mathbf{E}$$

$$\mathbf{F} \equiv -\frac{\partial E}{\partial \mathbf{u}} = -M \omega_0^2 \mathbf{u} + Z^* \mathbf{E}$$

$$\mathbf{D} \equiv -\frac{4\pi}{\Omega} \frac{\partial E}{\partial \mathbf{E}} = \frac{4\pi}{\Omega} Z^* \mathbf{u} + \epsilon_\infty \mathbf{E}$$

$$\text{rot } \mathbf{E} \sim i \mathbf{q} \times \mathbf{E} = 0$$

Phonons in polar materials

$$E(\mathbf{u}, \mathbf{E}) = \frac{1}{2} M \omega_0^2 u^2 - \frac{\Omega}{8\pi} \epsilon_\infty \mathbf{E}^2 - e Z^* \mathbf{u} \cdot \mathbf{E}$$

$$\mathbf{F} \equiv -\frac{\partial E}{\partial \mathbf{u}} = -M \omega_0^2 \mathbf{u} + Z^* \mathbf{E}$$

$$\mathbf{D} \equiv -\frac{4\pi}{\Omega} \frac{\partial E}{\partial \mathbf{E}} = \frac{4\pi}{\Omega} Z^* \mathbf{u} + \epsilon_\infty \mathbf{E}$$

$$\text{rot } \mathbf{E} \sim i \mathbf{q} \times \mathbf{E} = 0 \quad \mathbf{u} \perp \mathbf{q} \Rightarrow \mathbf{E} = 0 \quad (\text{T})$$

Phonons in polar materials

$$E(\mathbf{u}, \mathbf{E}) = \frac{1}{2} M \omega_0^2 u^2 - \frac{\Omega}{8\pi} \epsilon_\infty \mathbf{E}^2 - e Z^* \mathbf{u} \cdot \mathbf{E}$$

$$\mathbf{F} \equiv -\frac{\partial E}{\partial \mathbf{u}} = -M \omega_0^2 \mathbf{u} + Z^* \mathbf{E}$$

$$\mathbf{D} \equiv -\frac{4\pi}{\Omega} \frac{\partial E}{\partial \mathbf{E}} = \frac{4\pi}{\Omega} Z^* \mathbf{u} + \epsilon_\infty \mathbf{E}$$

$$\text{rot } \mathbf{E} \sim i \mathbf{q} \times \mathbf{E} = 0 \quad \mathbf{u} \perp \mathbf{q} \Rightarrow \mathbf{E} = 0 \quad (\text{T})$$

$$\text{div } \mathbf{D} \sim i \mathbf{q} \cdot \mathbf{D} = 0$$

Phonons in polar materials

$$E(\mathbf{u}, \mathbf{E}) = \frac{1}{2} M \omega_0^2 u^2 - \frac{\Omega}{8\pi} \epsilon_\infty \mathbf{E}^2 - e Z^* \mathbf{u} \cdot \mathbf{E}$$

$$\mathbf{F} \equiv -\frac{\partial E}{\partial \mathbf{u}} = -M \omega_0^2 \mathbf{u} + Z^* \mathbf{E}$$

$$\mathbf{D} \equiv -\frac{4\pi}{\Omega} \frac{\partial E}{\partial \mathbf{E}} = \frac{4\pi}{\Omega} Z^* \mathbf{u} + \epsilon_\infty \mathbf{E}$$

$$\text{rot } \mathbf{E} \sim i \mathbf{q} \times \mathbf{E} = 0 \quad \mathbf{u} \perp \mathbf{q} \Rightarrow \mathbf{E} = 0 \quad (\text{T})$$

$$\text{div } \mathbf{D} \sim i \mathbf{q} \cdot \mathbf{D} = 0 \quad \mathbf{u} \parallel \mathbf{q} \Rightarrow \mathbf{D} = 0 \quad (\text{L})$$

Phonons in polar materials

$$E(\mathbf{u}, \mathbf{E}) = \frac{1}{2} M \omega_0^2 u^2 - \frac{\Omega}{8\pi} \epsilon_\infty \mathbf{E}^2 - e Z^* \mathbf{u} \cdot \mathbf{E}$$

$$\mathbf{F} \equiv -\frac{\partial E}{\partial \mathbf{u}} = -M \omega_0^2 \mathbf{u} + Z^* \mathbf{E}$$

$$\mathbf{D} \equiv -\frac{4\pi}{\Omega} \frac{\partial E}{\partial \mathbf{E}} = \frac{4\pi}{\Omega} Z^* \mathbf{u} + \epsilon_\infty \mathbf{E}$$

$$\text{rot } \mathbf{E} \sim i \mathbf{q} \times \mathbf{E} = 0 \quad \mathbf{u} \perp \mathbf{q} \Rightarrow \mathbf{E} = 0 \quad (\text{T})$$

$$\text{div } \mathbf{D} \sim i \mathbf{q} \cdot \mathbf{D} = 0 \quad \mathbf{u} \parallel \mathbf{q} \Rightarrow \mathbf{D} = 0 \quad (\text{L})$$

$$\mathbf{F}_{\text{T}} = -M \omega_0^2 \mathbf{u}$$

Phonons in polar materials

$$E(\mathbf{u}, \mathbf{E}) = \frac{1}{2} M \omega_0^2 u^2 - \frac{\Omega}{8\pi} \epsilon_\infty \mathbf{E}^2 - e Z^* \mathbf{u} \cdot \mathbf{E}$$

$$\mathbf{F} \equiv -\frac{\partial E}{\partial \mathbf{u}} = -M \omega_0^2 \mathbf{u} + Z^* \mathbf{E}$$

$$\mathbf{D} \equiv -\frac{4\pi}{\Omega} \frac{\partial E}{\partial \mathbf{E}} = \frac{4\pi}{\Omega} Z^* \mathbf{u} + \epsilon_\infty \mathbf{E}$$

$$\text{rot } \mathbf{E} \sim i \mathbf{q} \times \mathbf{E} = 0 \quad \mathbf{u} \perp \mathbf{q} \Rightarrow \mathbf{E} = 0 \quad (\text{T})$$

$$\text{div } \mathbf{D} \sim i \mathbf{q} \cdot \mathbf{D} = 0 \quad \mathbf{u} \parallel \mathbf{q} \Rightarrow \mathbf{D} = 0 \quad (\text{L})$$

$$\mathbf{F}_{\text{T}} = -M \omega_0^2 \mathbf{u}$$

$$\mathbf{F}_{\text{L}} = -M \left(\omega_0^2 + \frac{4\pi Z^*}{M \Omega \epsilon_\infty} \right) \mathbf{u}$$

Interatomic force constants

$$\Phi(\mathbf{R} - \mathbf{R}') = -\frac{\partial^2 E}{\partial \mathbf{u}_R \partial \mathbf{u}_{R'}}$$

Interatomic force constants

$$\Phi(\mathbf{R} - \mathbf{R}') = -\frac{\partial^2 E}{\partial \mathbf{u}_R \partial \mathbf{u}_{R'}}$$

Short ranged +
dipole-dipole

Interatomic force constants

$$\begin{aligned}\Phi(\mathbf{R} - \mathbf{R}') &= -\frac{\partial^2 E}{\partial \mathbf{u}_R \partial \mathbf{u}_{R'}} && \text{Short ranged +} \\ &= \frac{\Omega}{(2\pi)^3} \int e^{i\mathbf{q} \cdot (\mathbf{R} - \mathbf{R}')} D(\mathbf{q}) d\mathbf{q} && \text{dipole-dipole}\end{aligned}$$

Interatomic force constants

$$\begin{aligned}\Phi(\mathbf{R} - \mathbf{R}') &= -\frac{\partial^2 E}{\partial \mathbf{u}_R \partial \mathbf{u}_{R'}} && \text{Short ranged +} \\ &= \frac{\Omega}{(2\pi)^3} \int e^{i\mathbf{q} \cdot (\mathbf{R} - \mathbf{R}')} D(\mathbf{q}) d\mathbf{q} && \text{dipole-dipole}\end{aligned}$$

- Do FFT's

\mathbf{q} 's = # \mathbf{R} 's

remove singularities in $D(\mathbf{q})$

Interatomic force constants

$$\begin{aligned}\Phi(\mathbf{R} - \mathbf{R}') &= -\frac{\partial^2 E}{\partial \mathbf{u}_R \partial \mathbf{u}_{R'}} && \text{Short ranged +} \\ &= \frac{\Omega}{(2\pi)^3} \int e^{i\mathbf{q} \cdot (\mathbf{R} - \mathbf{R}')} D(\mathbf{q}) d\mathbf{q} && \text{dipole-dipole}\end{aligned}$$

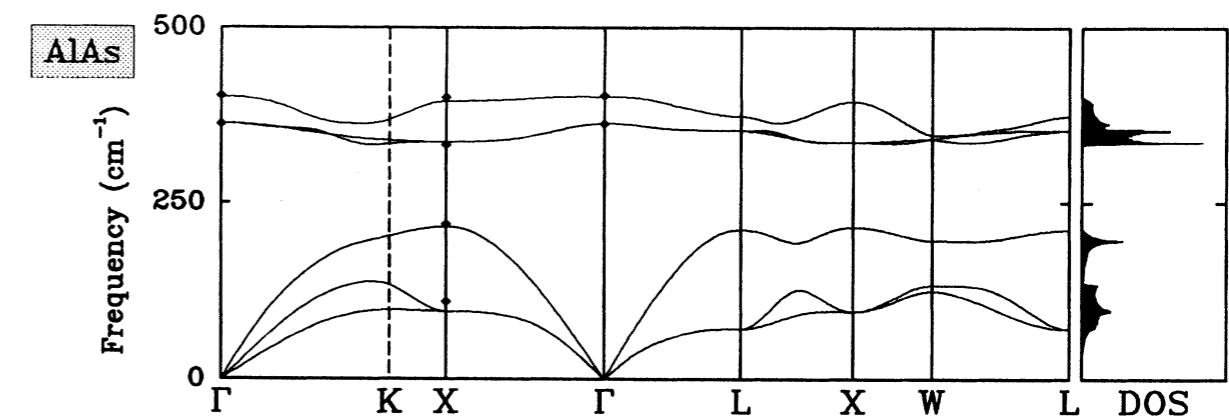
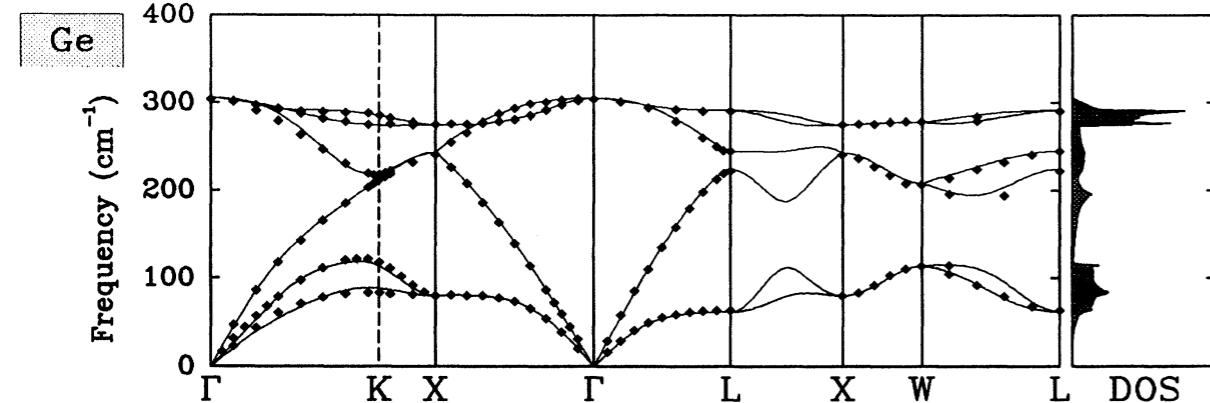
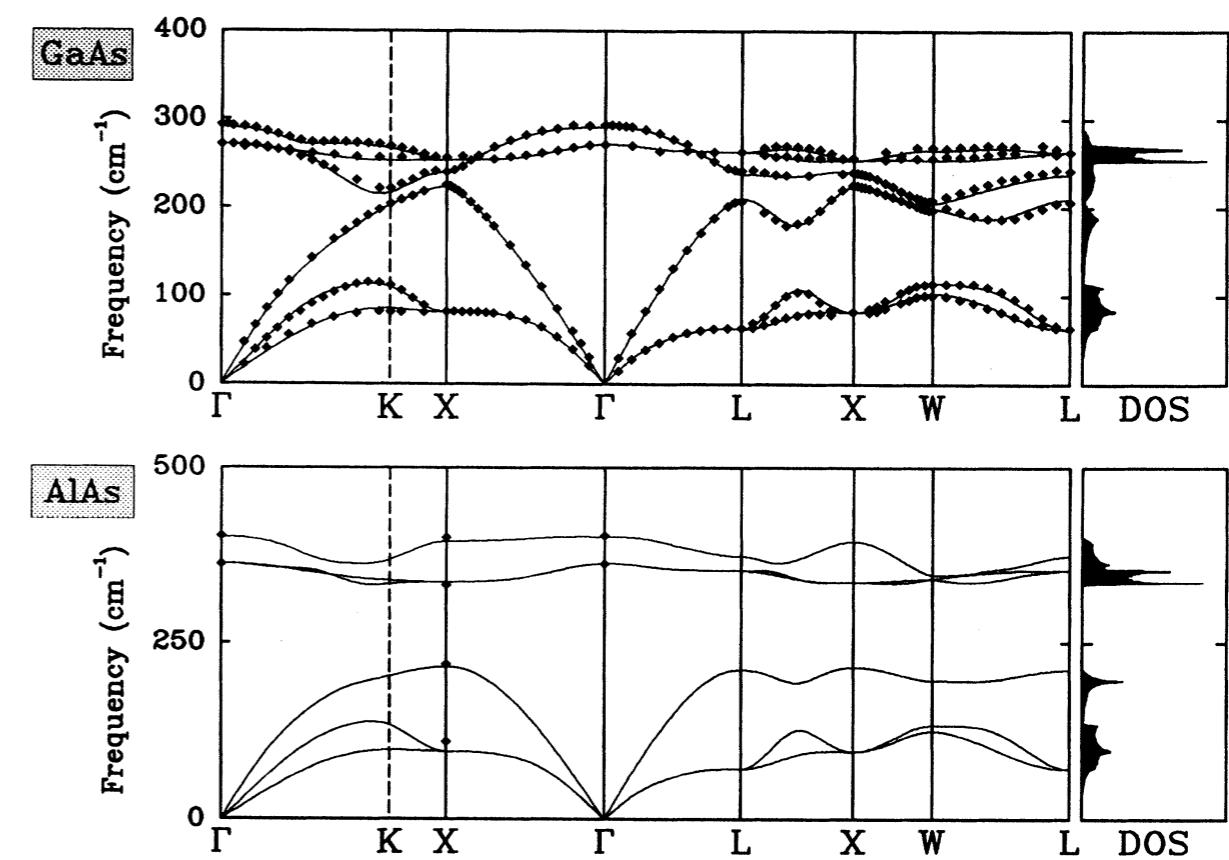
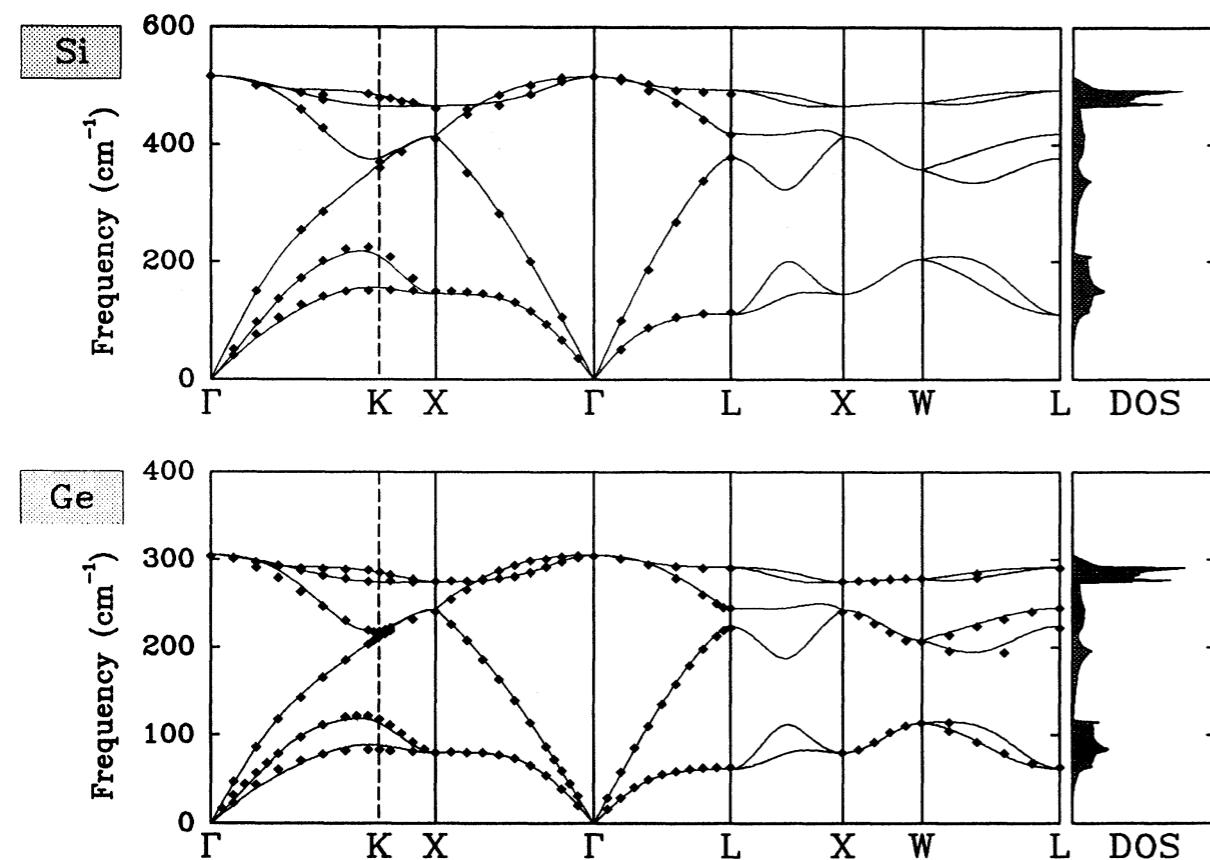
- Do FFT's
 - # \mathbf{q} 's = # \mathbf{R} 's
 - remove singularities in $D(\mathbf{q})$
- Store information

Interatomic force constants

$$\begin{aligned}\Phi(\mathbf{R} - \mathbf{R}') &= -\frac{\partial^2 E}{\partial \mathbf{u}_R \partial \mathbf{u}_{R'}} && \text{Short ranged +} \\ &= \frac{\Omega}{(2\pi)^3} \int e^{i\mathbf{q} \cdot (\mathbf{R} - \mathbf{R}')} D(\mathbf{q}) d\mathbf{q} && \text{dipole-dipole}\end{aligned}$$

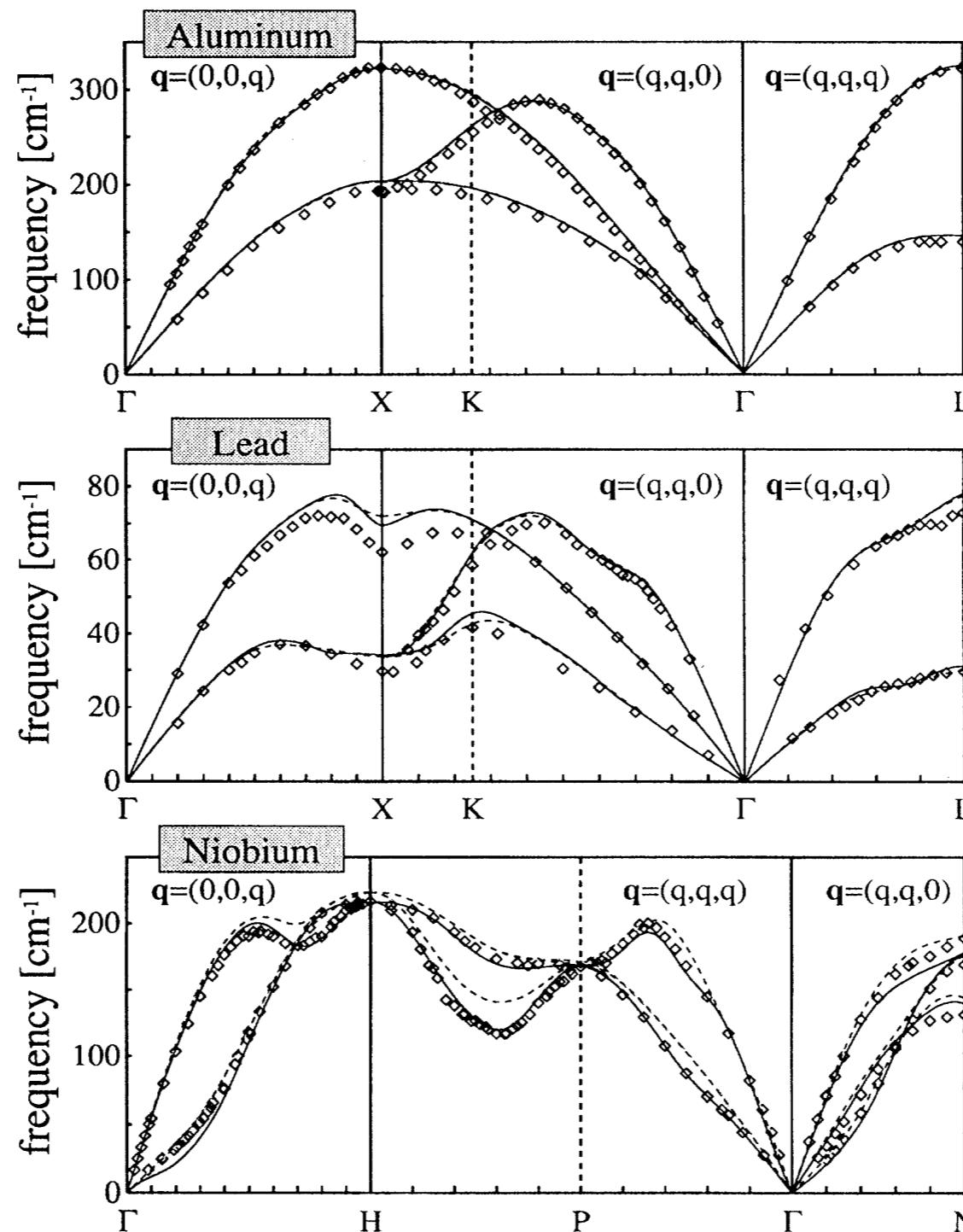
- Do FFT's
 - # \mathbf{q} 's = # \mathbf{R} 's
 - remove singularities in $D(\mathbf{q})$
- Store information
- Interpolate phonon bands

Phonons from DFPT



P. Giannozzi, S. de Gironcoli, P. Pavone, and S. Baroni, Phys. Rev. B **43**, 7231 (1991)

Phonons from DFPT



S. de Gironcoli, Phys. Rev. B **51**, 6773 (1995)

Applications done so far

- Dielectric properties
- Piezoelectric properties
- Elastic properties
- Phonon in crystals and alloys
- Phonon at surfaces,
interfaces, super-lattices, and
nano-structures
- Raman and infrared activities
- Anharmonic couplings and
vibrational line widths
- Mode softening and structural
transitions
- Electron-phonon interaction
and superconductivity
- Thermal expansion
- Isotopic effects on structural
and dynamical properties
- Thermo-elasticity and other
thermal properties of
minerals
- ...

S. Baroni, A. Dal Corso, S. de Gironcoli, and P. Giannozzi,
*Phonons and related crystal properties from density-functional
perturbation theory*, Rev. Mod. Phys. **73**, 515 (2001)

www.quantum-espresso.org