

# Introduction to Imaginary-Time Path Integrals

Qijing Zheng

Department of Physics

University of Science and Technology of China



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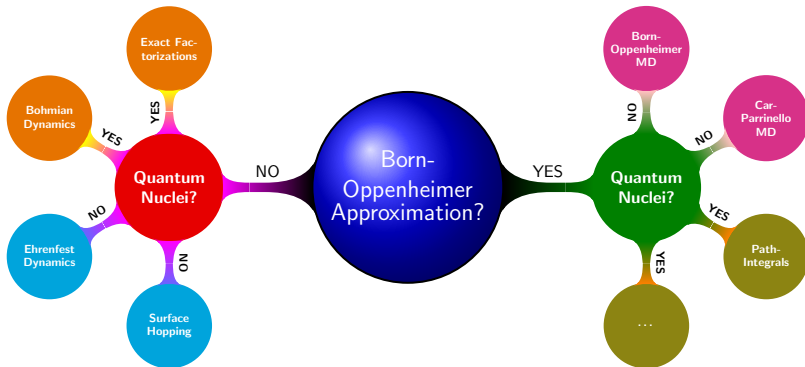
## 1 Introduction

# How to model the dynamics of electrons and nuclei from *ab initio*?

The time-dependent Schrödinger equation:

$$i\hbar \frac{\partial \Psi(\mathbf{r}, \mathbf{R}, t)}{\partial t} = \hat{\mathcal{H}}(\mathbf{r}, \mathbf{R}) \Psi(\mathbf{r}, \mathbf{R}, t)$$

In practice, *approximations* have to be made! <sup>1</sup>

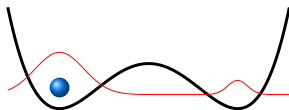


<sup>1</sup> "Ab initio molecular dynamics", Mariana Rossi, DFT Workshop 2017

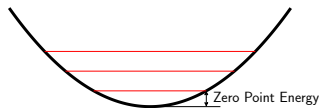
The *thermal* de Broglie wavelength of a particle:

$$\Lambda = \frac{h}{\sqrt{2\pi m k_B T}}$$

for a hydrogen atom at 300 K,  $\Lambda \approx 1.0 \text{ \AA}$ .



Quantum Tunneling



Zero Point Motion

other nuclear quantum effects: exchange effects, quantum coherence. . .

- NQEs are important for any vibrational mode for which  $\hbar\omega/k_B T > 1$ . If  $T = 300 \text{ K}$ ,  $\omega \approx 208 \text{ cm}^{-1}$ .

For a single particle moving in one spatial dimension potential <sup>2</sup>

$$\hat{\mathcal{H}} = \frac{\hat{p}^2}{2m} + \hat{V}(\hat{x}) \equiv \hat{T} + \hat{V}$$

The quantum partition function

$$\begin{aligned} Z &= \text{tr} \left[ e^{-\beta \hat{\mathcal{H}}} \right] = \text{tr} \left[ (e^{-\frac{\beta}{P} \hat{\mathcal{H}}})^P \right] = \text{tr} \left[ (e^{-\beta_P \hat{\mathcal{H}}})^P \right] \\ &= \int dx_1 \langle x_1 | (e^{-\beta_P \hat{\mathcal{H}}})^P | x_1 \rangle \\ &= \int dx_1 \dots \int dx_P \langle x_1 | e^{-\beta_P \hat{\mathcal{H}}} | x_2 \rangle \dots \langle x_P | e^{-\beta_P \hat{\mathcal{H}}} | x_1 \rangle \end{aligned}$$

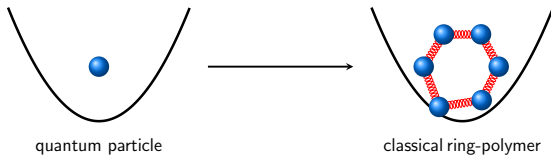
Note the connection between the quantum propagator and the density matrix,

$$\langle x' | e^{-i\hat{\mathcal{H}}t/\hbar} | x \rangle \quad \Rightarrow \quad \langle x' | e^{-\beta \hat{\mathcal{H}}} | x \rangle; \quad t = -i\beta\hbar$$

Hence the name *imaginary-time path integral*. By using the Trotter splitting<sup>3</sup>

$$\begin{aligned} e^{-\beta_P \hat{\mathcal{H}}} &= e^{-\beta_P \hat{V}/2} e^{-\beta_P \hat{T}} e^{-\beta_P \hat{V}/2} + \mathcal{O}(\beta_P^3) \\ \Rightarrow \quad \langle x_i | e^{-\beta_P \hat{\mathcal{H}}} | x_j \rangle &\approx \left( \frac{m}{2\pi\beta_P \hbar^2} \right)^{1/2} e^{-\beta_P \left[ \frac{m}{2(\beta_P \hbar)^2} (x_i - x_j)^2 + \frac{1}{2} (V(x_i) + V(x_j)) \right]} \end{aligned}$$

$$\begin{aligned} Z &= \text{tr} \left[ e^{-\beta \hat{\mathcal{H}}} \right] \\ &= \lim_{P \rightarrow \infty} \int dx_1 \dots \int dx_P \left( \frac{m}{2\pi\beta_P \hbar^2} \right)^{P/2} \exp \left\{ -\beta_P \sum_{i=1}^P \left[ \frac{m}{2(\beta_P \hbar)^2} (x_{i+1} - x_i)^2 + V(x_i) \right] \right\} \Big|_{x_{P+1}=x_1} \end{aligned}$$



- Maps the *quantum partition function* to the *configuration integral* of classical ring-polymers.
- **EXACT** when  $P \rightarrow \infty$ . Reduce to classical partition function when  $P = 1$ .
- The integral can be sampled with Monte Carlo method (PIMC, not covered in this talk), but for MD we need *momenta*!

<sup>2</sup>“Statistical Mechanics: Theory and Molecular Simulation”, Mark E. Tuckerman

<sup>3</sup>There are higher-order splitting techniques.

For a Hermitian operator  $\hat{A}$ , the expectation value follows,

$$\langle \hat{A} \rangle = \frac{1}{Z} \text{tr} [\hat{A} e^{-\beta \hat{\mathcal{H}}}] = \frac{1}{Z} \text{tr} [\hat{A} (e^{-\beta_P \hat{\mathcal{H}}})^P]$$

if  $\hat{A}$  is purely a function of  $\hat{x}$ , i.e.  $\hat{A}(\hat{x}) |x\rangle = a(x) |x\rangle$ , then <sup>4</sup>

$$\langle \hat{A} \rangle = \frac{1}{Z} \lim_{P \rightarrow \infty} \int dx_1 \dots \int dx_P \left( \frac{m}{2\pi\beta_P \hbar^2} \right)^{P/2} \left( \frac{1}{P} \sum_{i=1}^P a(x_i) \right) \exp \left\{ -\beta_P \sum_{i=1}^P \left[ \frac{m}{2(\beta_P \hbar)^2} (x_{i+1} - x_i)^2 + V(x_i) \right] \right\} \Bigg|_{x_{P+1}=x_1}$$

- If  $\hat{A}$  is a function of momentum operator  $\hat{p}$ , then the **cyclic-path** condition  $x_{P+1} = x_1$  is released, in which case one should resort to the so-called **open-path** path integral method.
- The thermodynamic functions, which may depend on both  $\hat{x}$  and  $\hat{p}$ , are exceptional because thermodynamic relations can be used. For example

$$E = \langle \hat{\mathcal{H}} \rangle = \left\langle \frac{\hat{p}^2}{2m} + V(\hat{x}) \right\rangle = -\frac{\partial}{\partial \beta} \ln Z = \frac{1}{Z} \frac{\partial Z}{\partial \beta}$$

<sup>4</sup> "Statistical Mechanics: Theory and Molecular Simulation", Mark E. Tuckerman

Insert  $P$  Gaussian integral into  $Z$ :  $\left(\frac{\beta}{2\pi\tilde{m}}\right)^{1/2} \int d\mathbf{p} e^{-\beta p^2/2\tilde{m}} = 1$

$$Z = \left(\frac{1}{2\pi\hbar}\right)^P \prod_{i=1}^P \sqrt{\frac{m}{\tilde{m}_i}} \lim_{P \rightarrow \infty} \int dx_1 \dots \int dx_P \int d\mathbf{p}_1 \dots \int d\mathbf{p}_P$$

More general:  $\prod_i \tilde{m}_i \Rightarrow \det[\mathbf{M}]$

$$\exp \left\{ -\beta_P \sum_{i=1}^P \left[ \frac{p_i^2}{2\tilde{m}_i} + \frac{m}{2(\beta_P \hbar)^2} (x_{i+1} - x_i)^2 + V(x_i) \right] \right\}$$

real physical mass

$$= C \lim_{P \rightarrow \infty} \int d\mathbf{x}^P \int d\mathbf{p}^P \exp[-\beta_P H_P(\mathbf{x}, \mathbf{p})]$$

$\tilde{m}_i$  need **not** be physical. More general form:  $\frac{1}{2} \mathbf{p}^T \mathbf{M}^{-1} \mathbf{p}$ , where  $\mathbf{M}$  is the mass matrix, and  $\mathbf{p} = (p_1, \dots, p_P)$ .

- PIMD trajectories are obtained by integrating Hamilton's classical equation of motion

$$\frac{d\mathbf{p}_i}{dt} = + \frac{\partial H_P(\mathbf{x}, \mathbf{p})}{\partial \mathbf{x}_i}; \quad \frac{d\mathbf{x}_i}{dt} = - \frac{\partial H_P(\mathbf{x}, \mathbf{p})}{\partial \mathbf{p}_i}$$

- $\beta_P H_P(\mathbf{x}, \mathbf{p}) \Rightarrow \beta (H_P(\mathbf{x}, \mathbf{p})/P)$ , which affects the dynamics but not the statistics.
- Different mass matrices give the same static average, only dynamics will be changed.
- Mathematically, the main difference between PIMD, PACMD and RPMD is the choice of the mass matrix. Physically, they differ dramatically!



# Multiparticle Generalization

For a system of  $N$  *distinguishable* nuclei,<sup>5</sup>

$$\hat{\mathcal{H}}_P = \sum_I^N \frac{\mathbf{p}_I^2}{2m_I} + \hat{V}(\mathbf{x}_1, \dots, \mathbf{x}_N),$$

**Potential energy:** could be from empirical potential or from *ab initio* calculations. *Born-Oppenheimer approximation* implicitly implied.

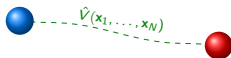
The PIMD Hamiltonian is

$$H_P(\{\mathbf{x}\}, \{\mathbf{p}\}) = \sum_{I=1}^N \sum_{i=1}^P \left[ \frac{(\mathbf{p}_I^i)^2}{2\tilde{m}_I^i} + \frac{1}{2} m_I \omega_P^2 (\mathbf{x}_I^{i+1} - \mathbf{x}_I^i)^2 \right] + \sum_{i=1}^P \hat{V}(\mathbf{x}_1^i, \dots, \mathbf{x}_N^i)$$

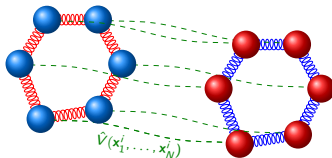
sum over nuclei index

sum over bead index

$$\omega_P = P k_B T / \hbar$$



quantum particle



classical ring-polymers

<sup>5</sup>We are neglecting the exchange effect, which is generally fine for nuclei (unless treating 4 K Helium) but not for electrons.

# Terminologies and Properties of the Ring-Polymer

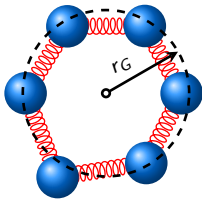
- **Radius of Gyration** – the spread in imaginary time. For a free particle the root mean square radius of gyration is:

$$\langle r_G^2(T) \rangle^{1/2} = \frac{\Lambda(T)}{\sqrt{8\pi}} \quad \Lambda(T) = \frac{h}{\sqrt{2\pi m k_B T}}$$

- **Bead to bead distance** –  
For a free particle the average is:

$$\sqrt{\frac{\beta \hbar^2}{Pm}}$$

Note the distance decreases as  $P$  increases.



- **Centroid** – the center of the Ring-Polymer.

$$x_c = \frac{1}{P} \sum_{i=1}^P x_i$$

- **Bead spring constants**<sup>6</sup> – determined by  $m\omega_P^2$

$$\omega_P = \frac{1}{\beta_P \hbar} = P k_B T / \hbar$$

The overall object is referred to as an *Imaginary Time Path* or a *Ring-Polymer*.

<sup>6</sup>In some textbooks,  $\omega_P = \sqrt{P} k_B T / \hbar$  if the Hamiltonian is  $H_P(\mathbf{x}, \mathbf{p})/P$

# Ring-Polymer Normal Modes

- Normal Modes:** – the ring-polymer potential can be diagonalized

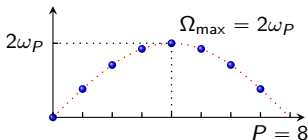
$$\sum_{i=1}^P \frac{1}{2} m \omega_P^2 (x_{i+1} - x_i)^2 \iff \begin{aligned} \tilde{\mathbf{x}} &= \mathbf{U} \cdot \mathbf{x} \\ \mathbf{x} &= \mathbf{U}^T \cdot \tilde{\mathbf{x}} \end{aligned} \iff \sum_{j=1}^P \frac{1}{2} m \Omega_j^2 \tilde{x}_j^2;$$

where  $\mathbf{x}$  is the bead **Cartesian coordinate** and  $\tilde{\mathbf{x}}$  is the **normal mode** coordinate.

- The transformation  $\tilde{\mathbf{x}} \Leftrightarrow \mathbf{x}$  can be done with FFT.<sup>7</sup>
- The normal mode frequencies<sup>8</sup>

$$\Omega_j = 2\omega_P \sin\left(\frac{(j-1)\pi}{P}\right)$$

e.g.  $P = 8$



- The unitary transformation matrix  $\mathbf{U}$  for even  $P$

$$U_{jk} = \begin{cases} \sqrt{1/P}, & k = 1 \\ \sqrt{2/P} \cos(2\pi jk/P), & 2 \leq k \leq P/2 \\ \sqrt{1/P} (-1)^j, & k = P/2 + 1 \\ \sqrt{2/P} \sin(2\pi jk/P), & P/2 + 2 \leq k \leq P \end{cases}$$

- Note that in many books

$$\begin{aligned} \tilde{\mathbf{x}} &= \frac{1}{\sqrt{P}} \mathbf{U} \cdot \mathbf{x} \\ \mathbf{x} &= \sqrt{P} \mathbf{U}^T \cdot \tilde{\mathbf{x}} \end{aligned}$$

then  $\tilde{x}_1 = (1/P) \sum_{i=1}^P x_i$  corresponds to the centroid, which is useful in CMD.

<sup>7</sup> *J. Chem. Phys.*, **104**, 2028(1996).

<sup>8</sup> This reminds me of the phonon dispersion of 1-D atomic chain.

# Initializing PIMD — Positions and Momenta

$$H_P(\{\mathbf{x}\}, \{\mathbf{p}\}) = \sum_{I=1}^N \sum_{i=1}^P \left[ \frac{(\mathbf{p}_I^i)^2}{2\tilde{m}_I^i} + \frac{1}{2} m_I \omega_P^2 (\mathbf{x}_I^{i+1} - \mathbf{x}_I^i)^2 \right] + \sum_{i=1}^P \hat{V}(\mathbf{x}_1^i, \dots, \mathbf{x}_N^i)$$

- To initialize PIMD requires specification of  $3NP$  positions and  $3NP$  momenta.

## MOMENTA

- We inserted  $P$  Gaussians when introducing momenta

$$\left( \frac{\beta}{2\pi\tilde{m}} \right)^{1/2} \int d\mathbf{p} e^{-\beta \mathbf{p}^2 / 2\tilde{m}}$$

- We can then sample the momentum of each bead from a Gaussian distribution with

$$\bar{\mathbf{p}} = 0$$
$$\sigma_P = \sqrt{\frac{\tilde{m}}{\beta_P}} = \sqrt{\tilde{m} k_B P T}$$

## POSITIONS

- 1 Start all beads at the same positions and equilibrate.
  - RP will expand under NVE,  $T$  will drop.
  - Strong Thermostatting is need!
- 2 Sample from the **free** ring-polymer distribution.
  - The only potential is the harmonic springs.
  - In normal mode coordinates

$$\sum_{j=1}^P \frac{1}{2} m \Omega_j^2 \tilde{x}_j^2 \Rightarrow \int d\tilde{x}_j e^{-\beta_P \left( \frac{1}{2} m \Omega_j^2 \tilde{x}_j^2 \right)}$$

- $\bar{\tilde{x}}_j = 0$  and  $\sigma_{\tilde{x}_j} = \sqrt{\frac{1}{\beta_P m \Omega_j^2}}$ .
- In practice, RP slightly too extended.

# Convergence of Standard Path Integral MD

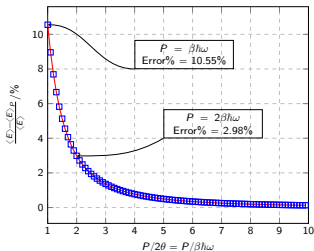
For the 1D harmonic oscillator:  $V(x) = \frac{1}{2}m\omega x^2$ , the RP potential in the normal-mode coordinates

$$\sum_{i=1}^P \left[ \frac{m}{2(\beta_P \hbar)^2} (x_{i+1} - x_i)^2 + V(x_i) \right] \Rightarrow \sum_{i=1}^P \frac{1}{2} m (\Omega_i^2 + \omega^2) \tilde{x}_i^2$$

then  $\langle \tilde{x}_i^2 \rangle = 1/\beta_P m (\Omega_i^2 + \omega^2)$ , the energy of the harmonic oscillator  $\langle E \rangle_P$

$$\begin{aligned} \langle E \rangle_P &= 2 \cdot \frac{1}{2} m \omega^2 \langle x^2 \rangle = \frac{m \omega^2}{P} \sum_{i=1}^P \langle x_i^2 \rangle = \frac{m \omega^2}{P} \sum_{i=1}^P \langle \tilde{x}_i^2 \rangle \\ &= \frac{\theta}{\beta} \sum_{i=1}^P \frac{\theta}{P^2 \sin^2 \left( \frac{(i-1)\pi}{P} \right) + \theta^2} \quad (\theta = \beta \hbar \omega / 2) \end{aligned}$$

the exact solution is  $\langle E \rangle = \frac{\theta}{\beta} \coth(\theta)$



- Commonly used convergence criteria:  
 $P > \hbar\omega_{\max}/k_B T$
- Standard PIMD converges as  $1/P^2$ .
- Some properties converge faster, e.g. RDF converges fast but heat capacity very slowly.
- Methods exist to accelerate the convergence, e.g. PI + GLE.

Thank you!