

I. 基本电磁理论.

$$\left\{ \begin{array}{l} \nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t} \\ \nabla \times \vec{H} = \frac{\partial \vec{D}}{\partial t} + \vec{J} \\ \nabla \cdot \vec{D} = \rho \\ \nabla \cdot \vec{B} = 0 \end{array} \right.$$

$$\vec{B} = \nabla \times \vec{A} \quad \rightarrow \quad \nabla \times (\vec{E} + \frac{\partial \vec{A}}{\partial t}) = 0$$

无旋场，可以用标势中的梯度来表示。

$$\vec{E} = -\nabla \psi - \frac{\partial \vec{A}}{\partial t}$$

Current J_A^μ

- δA_μ :
- Electromagnetism background gauge field. NOT DYNAMICAL, Only parameter.
 - NOT include original background magnetic field (produce QHE)
 - Only perturbations around Hall state (turn on an electric or magnetic field)

J_μ

- Current, DYNAMICAL degrees of freedom.
- Always couple to gauge field

$$S_A = \int d^3x J^\mu A_\mu$$

Perturbation - Response pair

(2+1)d

$$\vec{T}_\mu = \sigma^{xy} \vec{J}_E \times \vec{z}$$

$$\vec{J}_H = \sigma^{xy} \delta \vec{E} \times \hat{z}$$

$$\delta \vec{E} = -\partial_t \delta \vec{A} - \nabla \delta A_0$$

Continuity equation: $\partial_t P_H + \nabla \cdot \vec{J}_H = 0$

$$\Rightarrow \partial_t P_H = -\nabla \cdot \vec{J}_H = -\sigma^{xy} \hat{z} \cdot (\nabla \times \vec{E})$$

$$= \sigma^{xy} \hat{z} \cdot (\partial_t \nabla \times \delta \vec{A})$$

$$= \sigma^{xy} \partial_t \cdot \delta B_z$$

$$\Rightarrow P_H = \sigma^{xy} \cdot \delta B_z$$

$$J_H^\mu = \sigma^{xy} \partial_\nu \delta A_\lambda \epsilon^{\mu\nu\lambda} \iff \begin{cases} P_H = \sigma^{xy} \delta B_z \\ \vec{J}_H = \sigma^{xy} \delta \vec{E} \times \hat{z} \end{cases}$$

$$J_H^\mu = (P_H, J_H^x, J_H^y)$$



$$P_H = \sigma^{xy} (\partial_x \delta A_y - \partial_y \delta A_x) = \sigma^{xy} \cdot \delta B_z$$

$$J_H^x = \sigma^{xy} (\partial_y \delta A_0 - \partial_0 \delta A_y) = \sigma^{xy} \cdot \delta E_y$$

$$J_H^y = \sigma^{xy} (\partial_0 \delta A_x - \partial_x \delta A_0) = \sigma^{xy} \cdot \delta E_x$$

$$J_H^\mu \stackrel{\text{def.}}{=} \frac{\delta S[\delta A]}{\delta [\delta A^\mu]}$$

推測 Action $S[\delta A]$

- Say nothing about microscopic: electrons, Landau level
- General assumption: at low energies, there are no degrees of freedom that can affect the physics when the system is perturbed.
 - Gap
 - GSD (FQHE)

Partition function:

$$Z[A_\mu] = \int D(\phi) e^{iS[\phi, A_\mu]/\hbar}$$

A_μ : NOT dynamic, source.

ϕ : field, dynamic degrees of freedom,
be integrated out.

$$Z[A_\mu] = e^{iS_{\text{eff}}[A_\mu]/\hbar}$$

$S_{\text{eff}}[A_\mu]$ depend on parameter! rather than dynamic field!

NO dynamics in A_μ .

$$J^\mu(x) = \frac{\delta S_{\text{eff}}[A]}{\delta A_\mu(x)}$$

Assumption of $S_{\text{eff}}[A_\mu]$:

1. gauge invariance.

2. symmetry. (rotation, translation)

3. Local. $S_{\text{eff}} = \int d^3x \dots$

a gap ΔE , non-locality arise at distance $\sim \frac{\sqrt{\hbar}}{\Delta E}$

⋮
⋮

The Chern-Simons Term

$$S_{\text{CS}}[A] = \frac{k}{4\pi} \int d^3x \epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho$$

k : level of the Chern-Simons term.

gauge invariance

$$A_\mu \rightarrow A_\mu + \partial_\mu w$$

$$S_{CS} \rightarrow \frac{k}{4\pi} \int d^3x \epsilon^{\mu\nu\rho} (A_\mu + \partial_\mu w) \partial_\nu (A_\rho + \partial_\rho w)$$

$$= S_{CS} + \frac{k}{4\pi} \int d^3x \epsilon^{\mu\nu\rho} [\partial_\mu w \cdot \partial_\nu A_\rho + A_\mu \cdot \partial_\nu \partial_\rho w + \partial_\mu w \partial_\nu \partial_\rho w]$$

$$= S_{CS} + \frac{k}{4\pi} \int d^3x \partial_\mu (w \epsilon^{\mu\nu\rho} \partial_\nu A_\rho)$$

Second term: throw total derivative. \rightarrow gauge invariance

HOWEVER -

Symmetry

rotation symmetry : \checkmark

parity reversal symmetry: \times

$$x^0 \rightarrow x^0, \quad x^1 \rightarrow -x^1, \quad x^2 \rightarrow x^2$$

correspondingly $A_0 \rightarrow A_0, \quad A_1 \rightarrow -A_1, \quad A_2 \rightarrow A_2$ (线性于 \vec{x})

$$\Rightarrow \epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho \rightarrow -\epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho$$

* the Chern-Simons effective interaction only arise in systems that BREAK PARITY SYMMETRY.

time reversal symmetry: \times

J^{μ}

$$J_i = \frac{\delta S_{CS}[A]}{\delta A_i} = \frac{k}{4\pi} \frac{1}{\delta A_i} \int d^3x \epsilon^{\mu\nu\rho} [\delta A_\mu \partial_\nu A_\rho + A_\mu \partial_\nu \delta A_\rho]$$

$$= -\frac{k}{2\pi} \epsilon_{ij} E_j$$

$$\text{QHE: } \sigma_{xy} = \frac{k}{2\pi} \underset{k=e^2v/h}{=} v \cdot \frac{e^2}{h}$$

$$\begin{aligned} J_0 &= \frac{\delta S_{CS}[A]}{\delta A_0} = \frac{k}{4\pi} \frac{1}{\delta A_0} \cdot \delta \int d^3x \epsilon^{\nu\rho} A_0 \partial_\nu A_\rho \\ &= \frac{k}{4\pi} \epsilon^{\nu\rho} \partial_\nu A_\rho \\ &= \frac{k}{2\pi} B \end{aligned}$$

A_μ is additional gauge field. $\rightarrow J^0$ is the change in the charge density over and above that already present in the ground state.

Quantisation of the Chern-Simons level k

Wick rotation : τ periodic S^1