

# On the dynamical origins of nonstatisticality

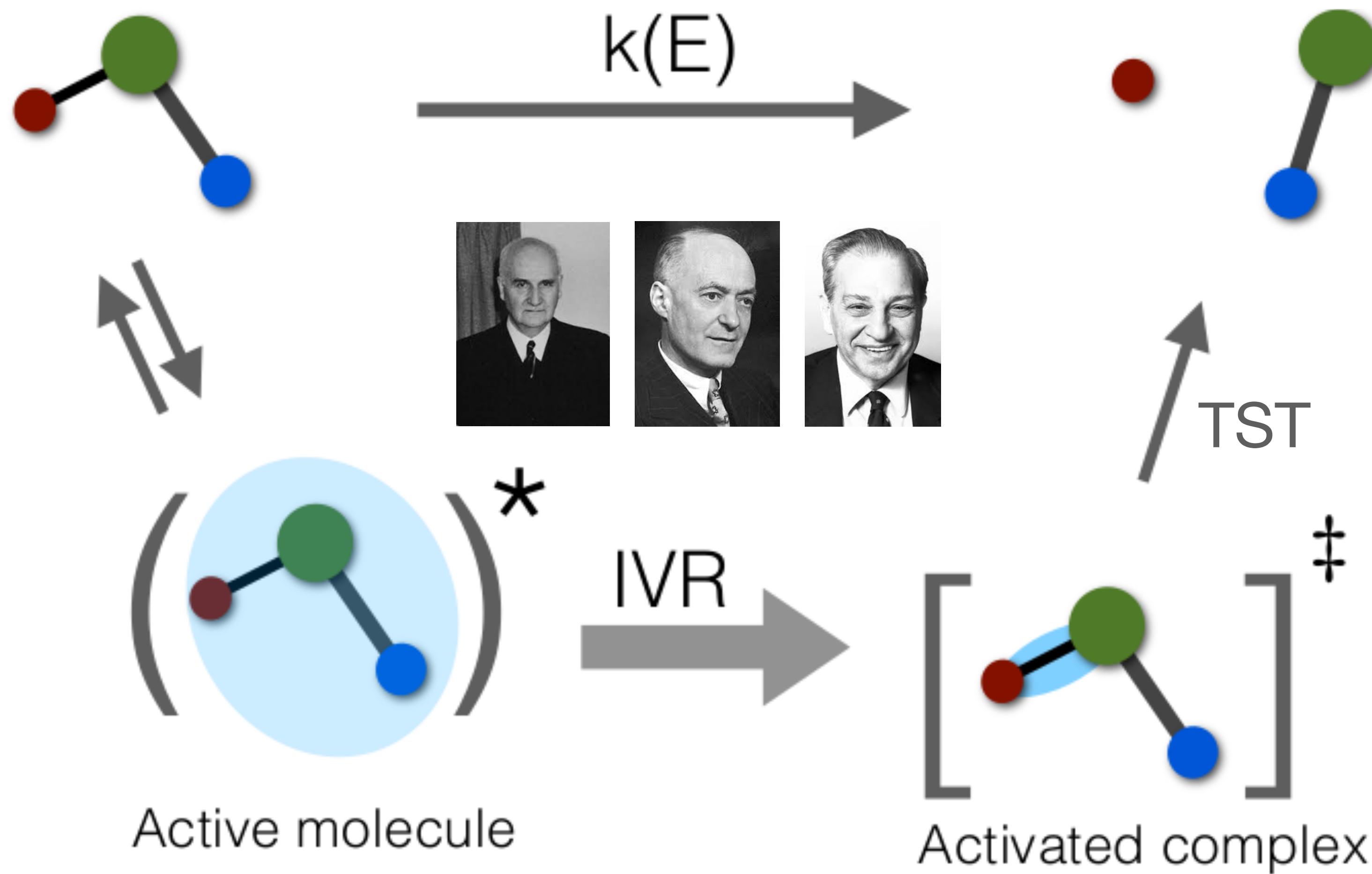
To flow or not to flow...

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# Unimolecular reactions ~ 1950s

Hinshelwood-Lindemann-Marcus



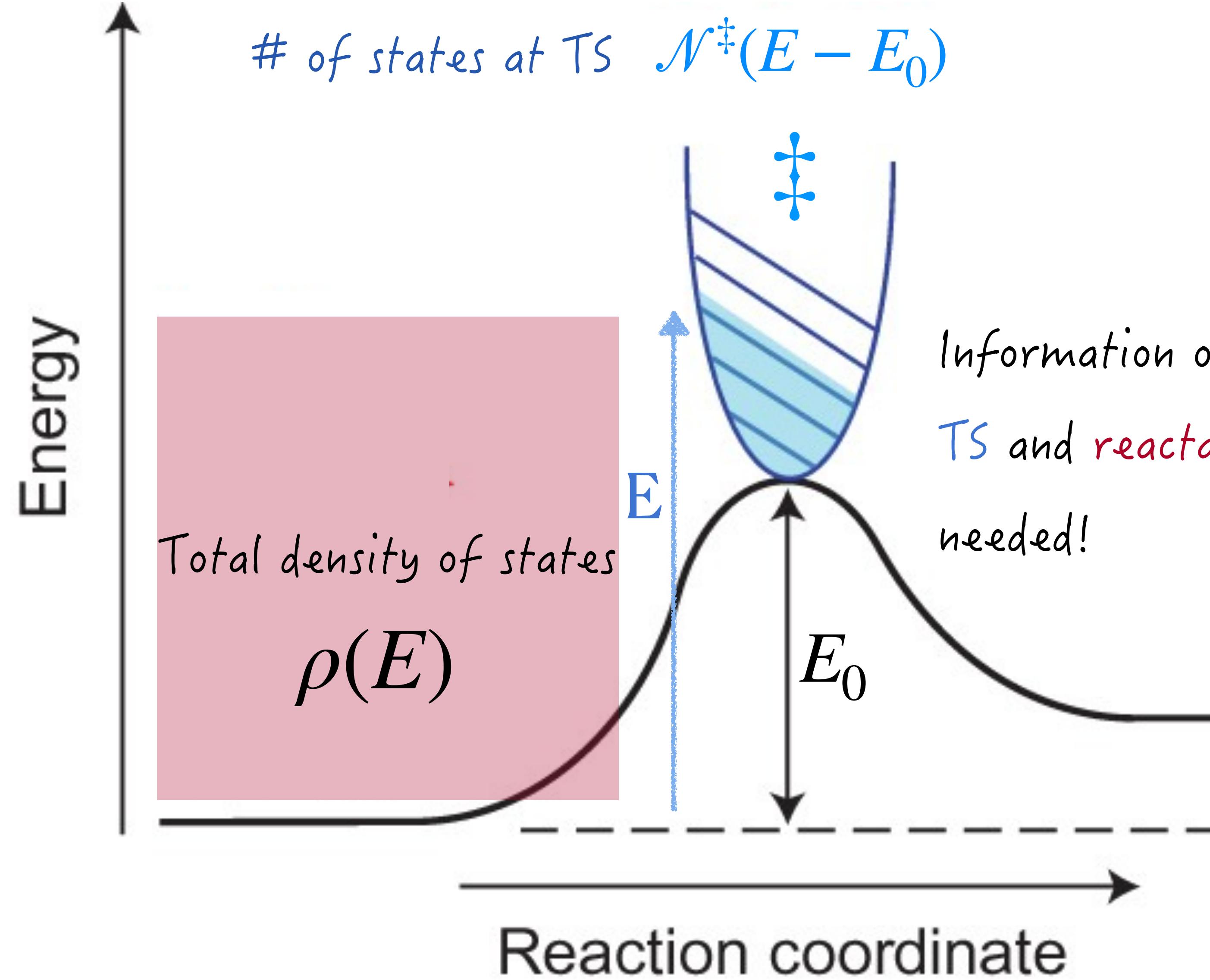
$$\tau_{rxn}^{\ddagger} \gg \tau_{IVR}$$

RRKM

(FPUT)

Dynamics

$$k(E) \approx \frac{\mathcal{N}^{\ddagger}(E - E_0)}{2\pi\hbar\rho(E)}$$



If RRKM then:

$\frac{N(t)}{N(0)} = e^{-k(E)t}$

✓ ↗ ✗ ↙

Exponential survival probability

Subtleties...

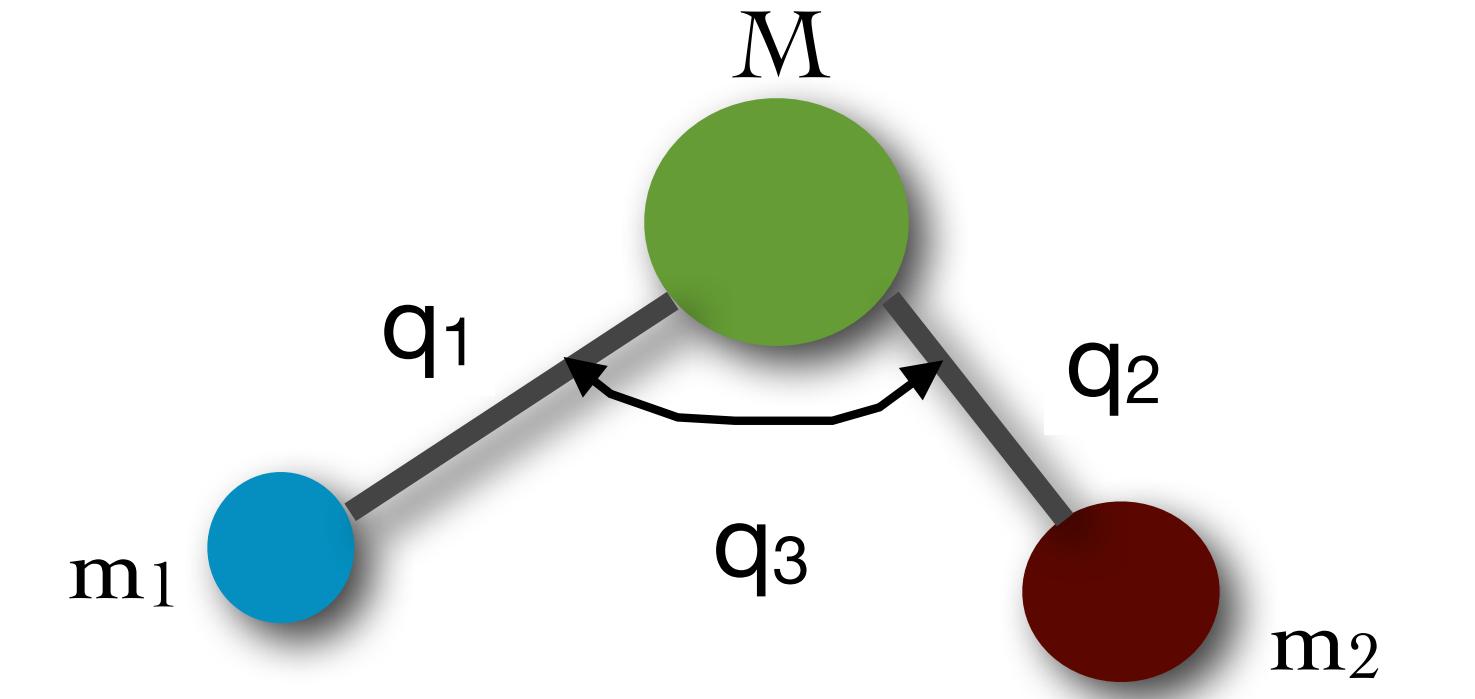
Ignore !?

# Model Hamiltonian

D L Bunker, J. Chem. Phys. (1962)

$$H(\mathbf{p}, \mathbf{q}) = \sum_{k=1}^3 \left[ \frac{1}{2} G_{kk}^{(0)} p_k^2 + V_k(q_k) \right] + \sum_{k < l=1}^3 G_{kl}^{(0)} p_k p_l$$

Zeroth-order Hamiltonian      Coupling  
Integrable       $\epsilon \neq 0$  ; Non-integrable



$$V_k(q_k) = D_k \left[ 1 - e^{-\alpha_k(q_k - q_k^0)} \right]^2$$

$$V_3(q_3) \sim \omega_3(q_3 - q_3^0)^2$$

Modes 1 and 2 (stretches) anharmonic (nonlinear); dissociation for  $E > D_{\text{low}}$

Mode 3 (bending) is harmonic ; cannot dissociate

NB: Inspired by Bunker's model Hamiltonian       $G_{ij}(\mathbf{q}) \rightarrow G_{ij}^{(0)}$

# Question...

(necessary & sufficient)

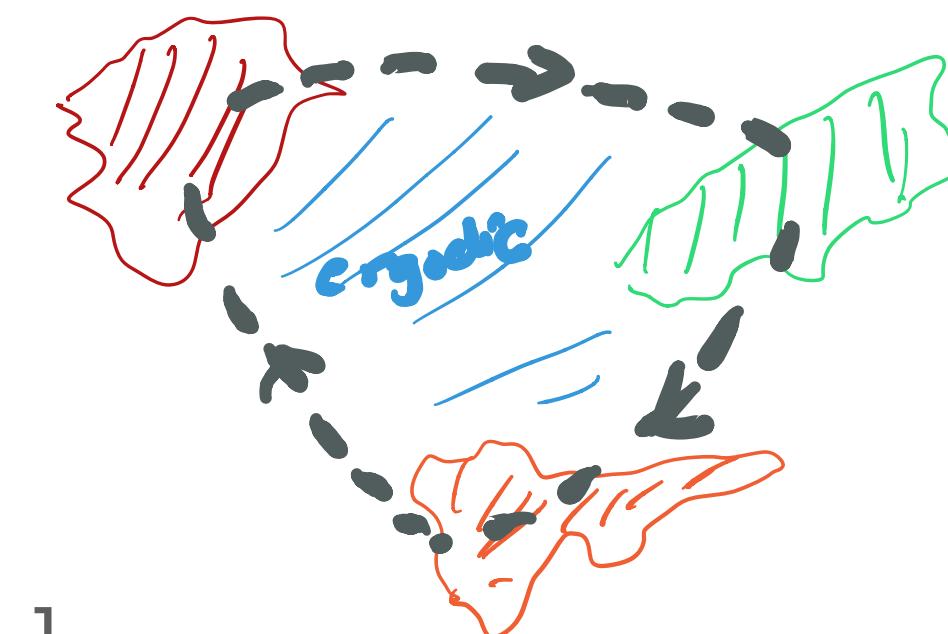
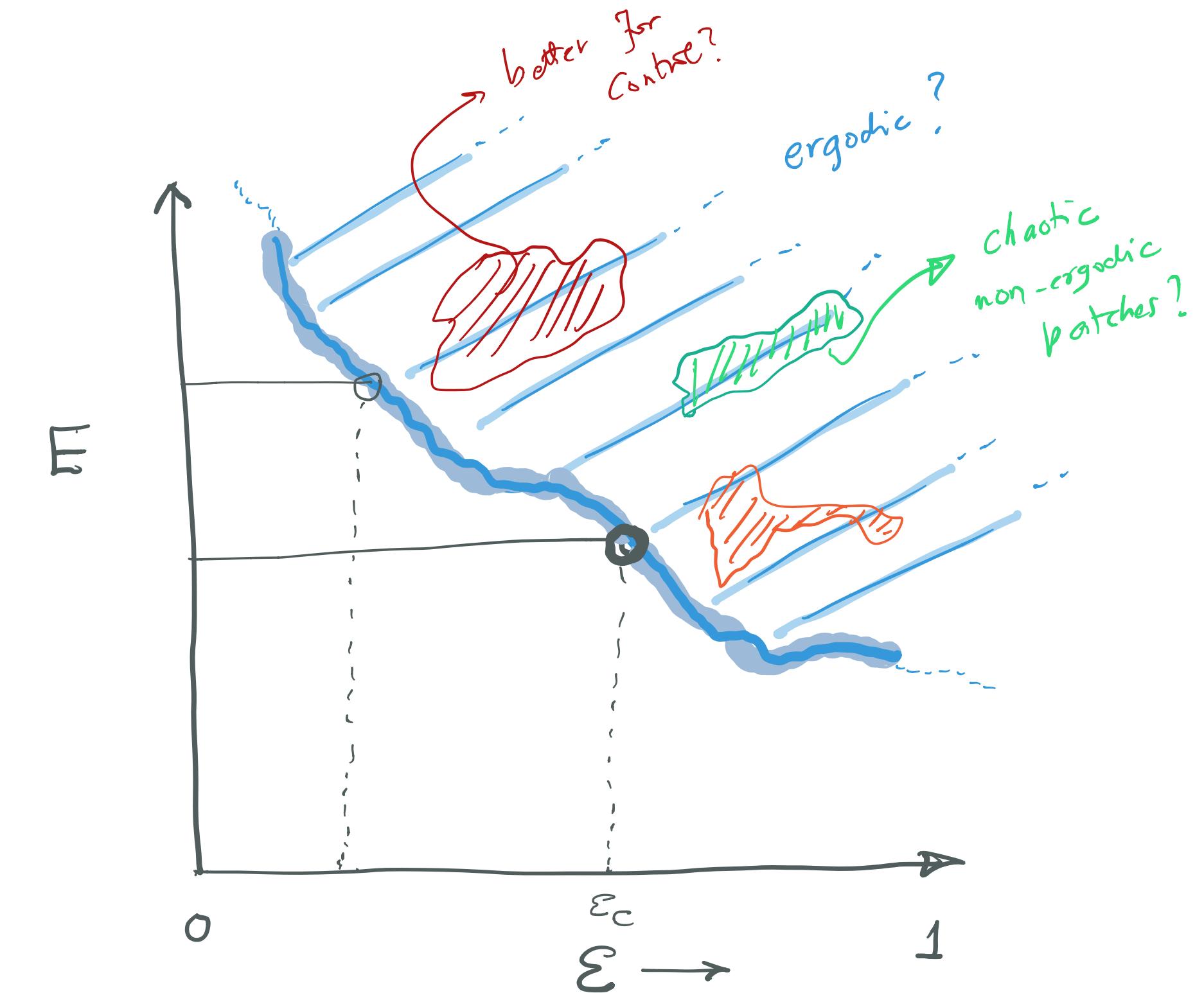
$$H(\mathbf{p}, \mathbf{q}) = \sum_{k=1}^3 \left[ \frac{1}{2} G_{kk}^{(0)} p_k^2 + V_k(q_k) \right] + \epsilon \sum_{k < l=1}^3 G_{kl}^{(0)} p_k p_l$$

Given Hamiltonian parameters,  
can one predict/estimate  $\epsilon_{\text{crit}}$  for which RRKM valid?

Classical - Oxtoby and Rice, J. Chem. Phys. 65, 1676 (1976)

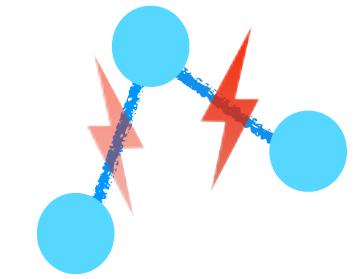
Quantum - Local Random Matrix Theory, Bose Statistics Triangle Rule:  
Logan, Wolynes, Leitner, Gruebele, ...

$\epsilon_{\text{crit}}^{\text{QM}} \sim \epsilon_{\text{crit}}^{\text{CM}}$  ? [Tunnelling, coherence, entanglement,...]

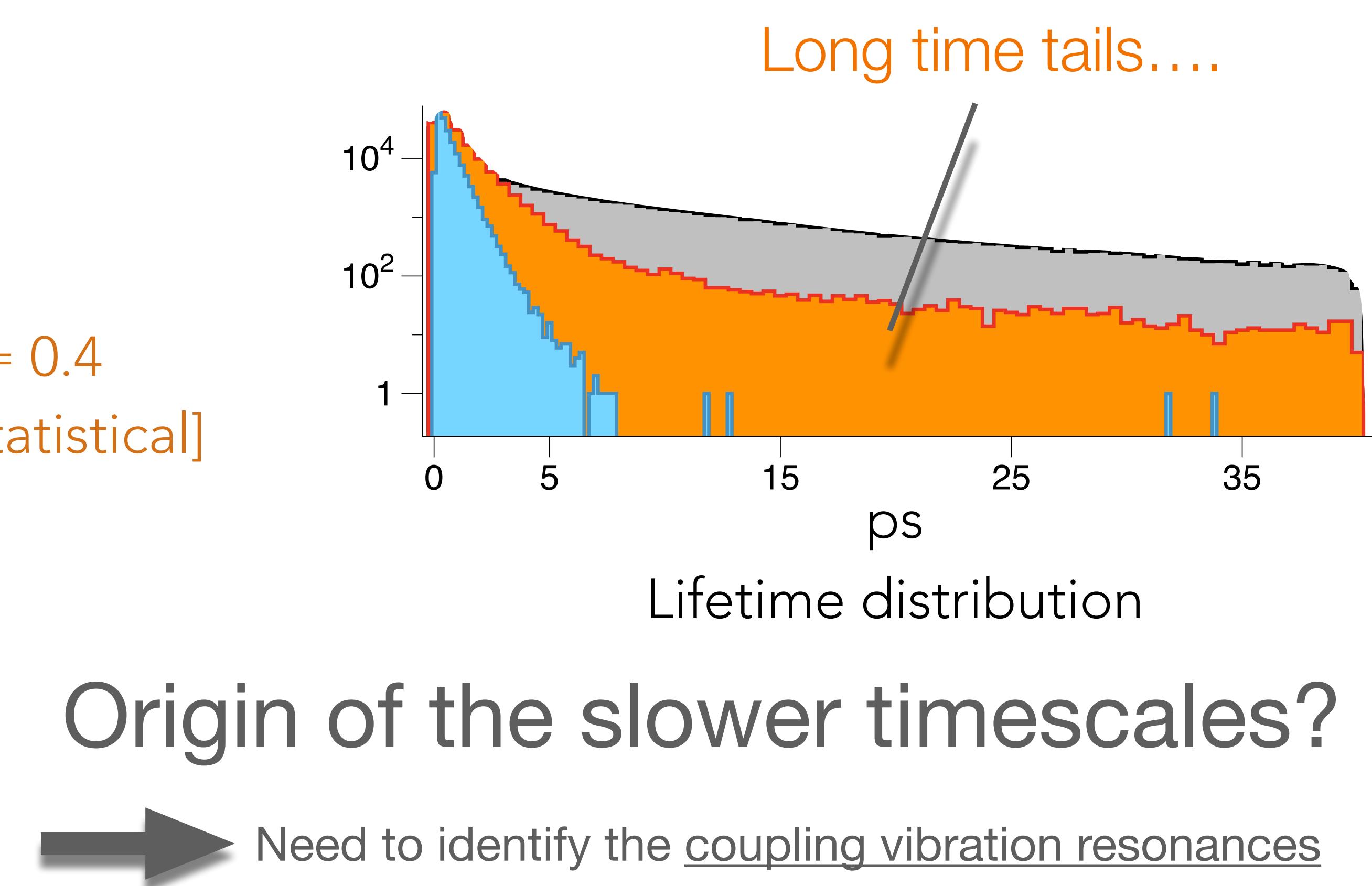
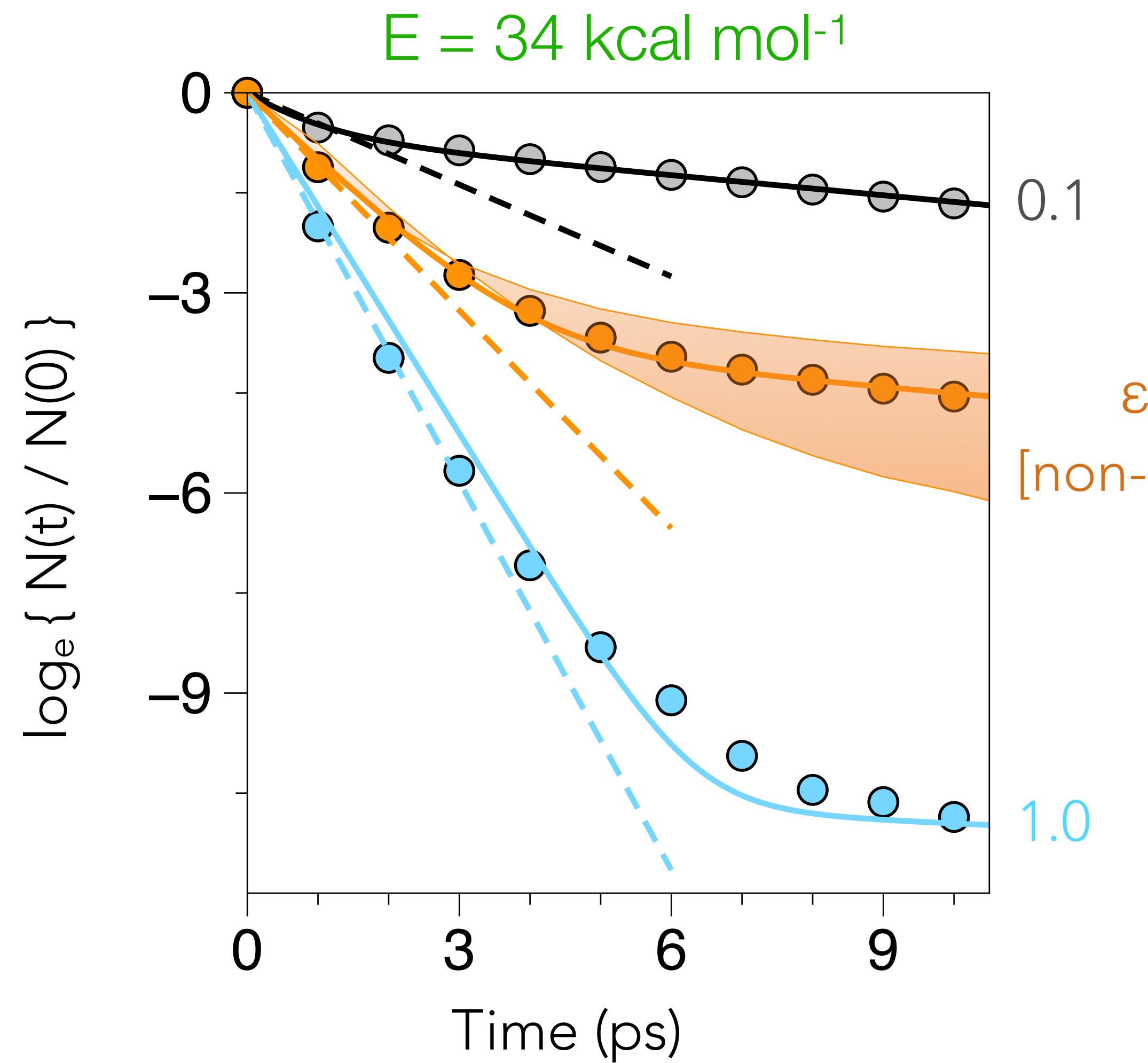


# Specific Bunker model

Model #6 ~ ozone  $D = 24 \text{ kcal mol}^{-1}$



[~ 200000 trajectories, time ~ 40 ps]



Momentum  
coupling in  
bond coordinates

$Q_1 Q_2 Q_3^2$

QM  
ladder r  
obs.

$a \leftrightarrow \sqrt{J} e^{i\theta}$   
Heisenberg

CM  
action-angle

For e.g.  $a_1^+ a_2^+ a_3^2 + \text{H.C.} \xrightarrow{\text{Heisenberg}} J_3 \sqrt{J_1 J_2} \cos(\theta_1 + \theta_2 - 2\theta_3)$

## Fermi Resonances

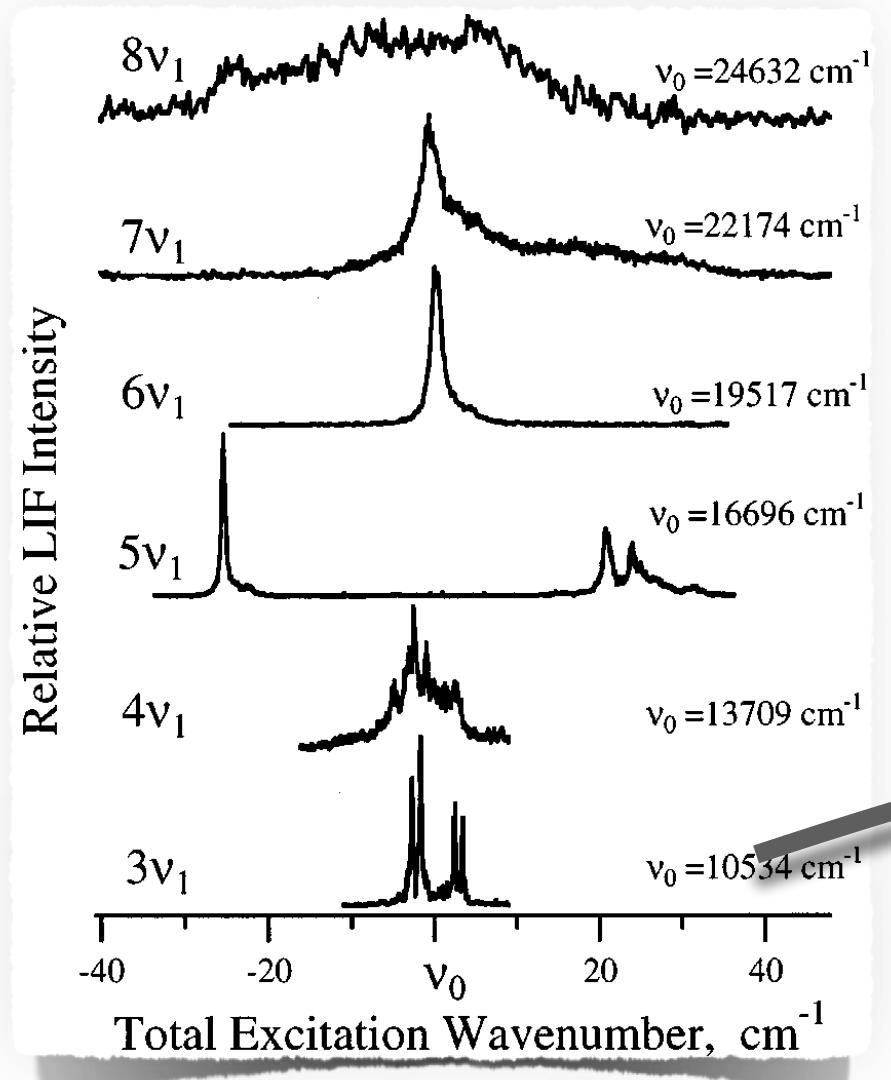
$$|v_1, v_2, v_3\rangle \leftrightarrow |v_1 \pm 1, v_2 \mp 1, v_3 \mp 2\rangle$$

$$|\underline{v}\rangle \quad |\underline{v}'\rangle$$

degenerate

$$E_{\underline{v}}^{\text{HO}} \approx E_{\underline{v}'}^{\text{HO}}$$

$$\Rightarrow \omega_1 + \omega_2 - 2\omega_3 \approx 0$$



anharmonicity :  $f_{\text{QM}}(\underline{v}, \underline{v}') \approx 0$

$$\frac{d}{dt}(\theta_1 + \theta_2 - 2\theta_3) \approx 0$$

"locked, in phase"

$$\Omega_1 + \Omega_2 - 2\Omega_3 \approx 0$$

$$f_{\text{CM}}(\underline{J}, \underline{J}') \approx 0$$

## Nonlinear Resonances

Slow:  
cannot be  
averaged away!

$J_k \longleftrightarrow \left(v_k + \frac{1}{2}\mu_k\right)\hbar$

(correlation!)

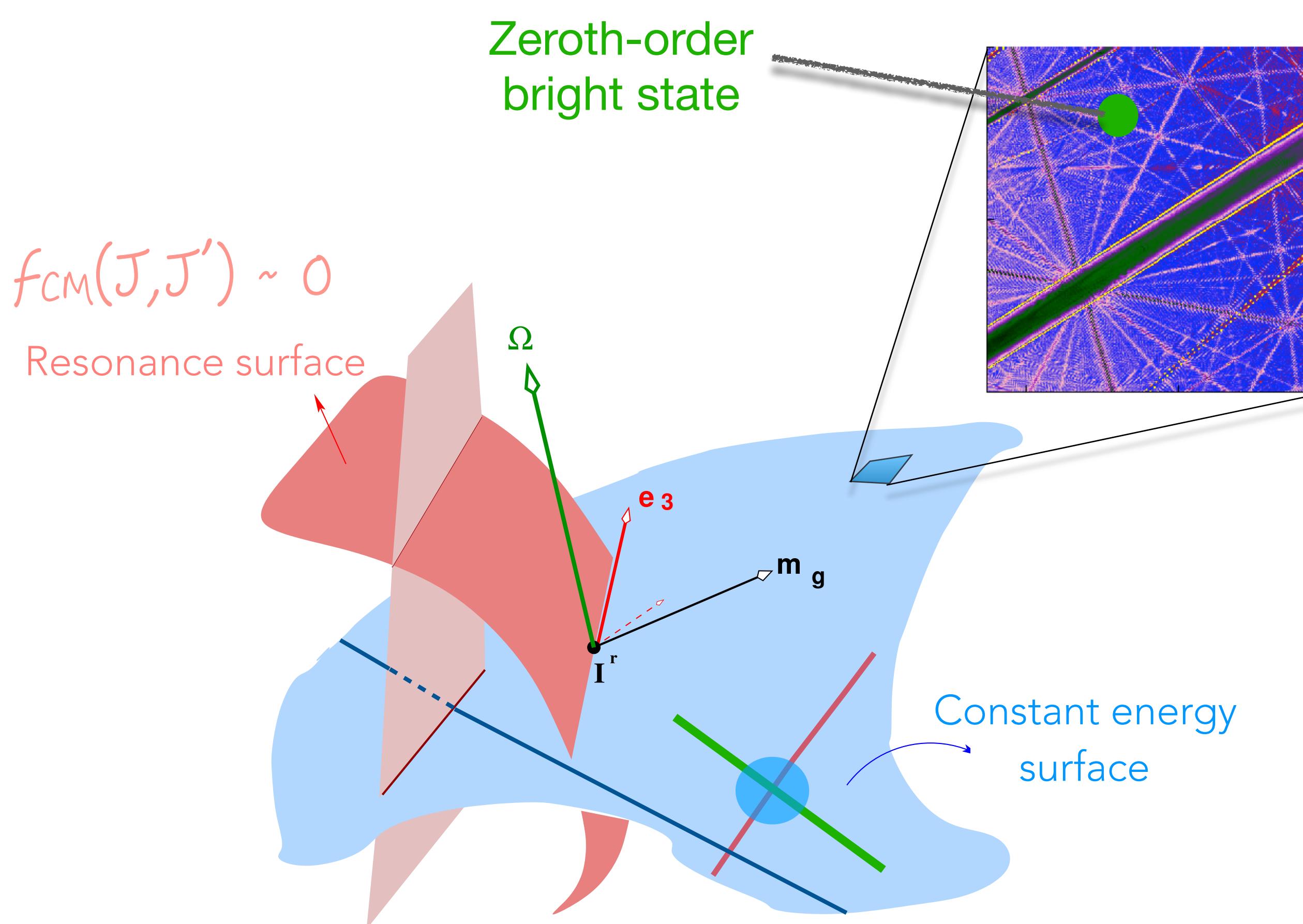
Depend on  $E$   
"nonlinear"

# The Arnold web

“IVR traffic routes”...

Rich “Geography” of the Arnold web ~ IVR pathways

**only** for > 2 degrees of freedom!

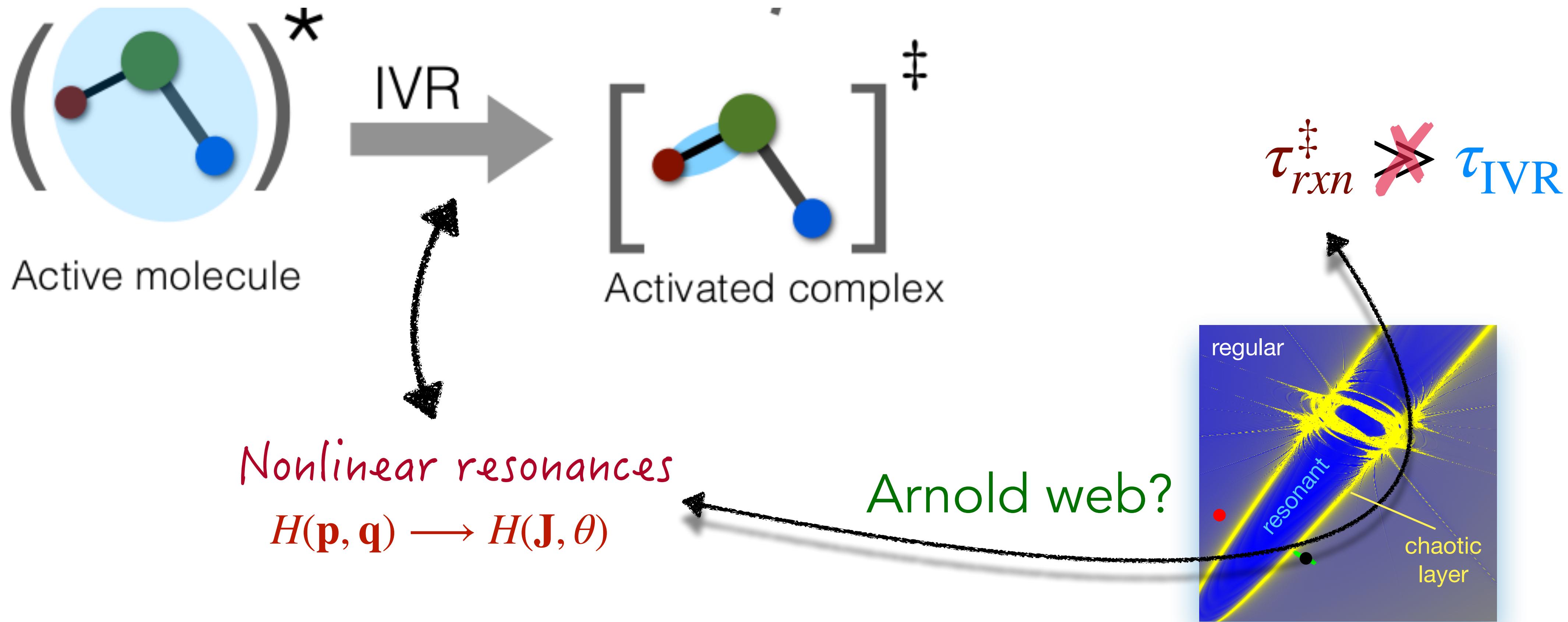


Oxtoby and Rice, J. Chem. Phys. 65, 1676 (1976)

For two interacting oscillators (the case we have discussed so far), the zeroth order energy surface is one dimensional. This case is simple in that there is only one path from a given point to any other, so that narrow resonances do not overlap. For  $n > 2$  oscillators, however, the energy hypersurface is of dimension  $n - 1$ , and centers of resonances are not points but  $n - 2$  dimensional surfaces. For a given energy, one can study the positions and extent of overlap of the resonances on the energy hypersurface; the degree to which the energy hypersurface is covered by overlapping resonances determines the extent to which the system may be described stochastically. However, such a many-dimensional picture rapidly becomes very complicated even for quite small molecules, so the question arises of whether it can be replaced by an approximate one-dimensional picture.

# Slower timescales...

Where are the bottlenecks?



**Fast Lyapunov Indicator (FLI); Froeschlé, Guzzo & Lega, Science 2000**

# Transformations

Action-angle variables

$$H(\mathbf{p}, \mathbf{q}) = \sum_{k=1}^3 \left[ \frac{1}{2} G_{kk}^{(0)} p_k^2 + V_k(q_k) \right] + \epsilon \sum_{k < l=1}^3 G_{kl}^{(0)} p_k p_l$$

Action - Angle transformation

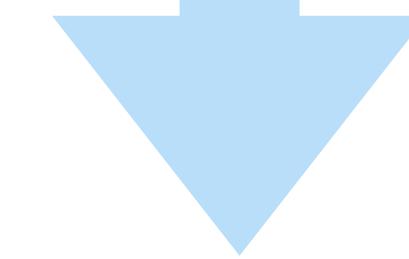
$$H(\mathbf{J}, \theta) = H_0(\mathbf{J}) + \boxed{\epsilon V(\mathbf{J}, \theta)}$$

$$f_{lmn}(\mathbf{J}) \cos(l\dot{\theta}_1 + m\dot{\theta}_2 + n\dot{\theta}_3)$$

“frequencies lock”

Nonlinear resonances

$$l\dot{\theta}_1 + m\dot{\theta}_2 + n\dot{\theta}_3 \approx 0$$



IVR

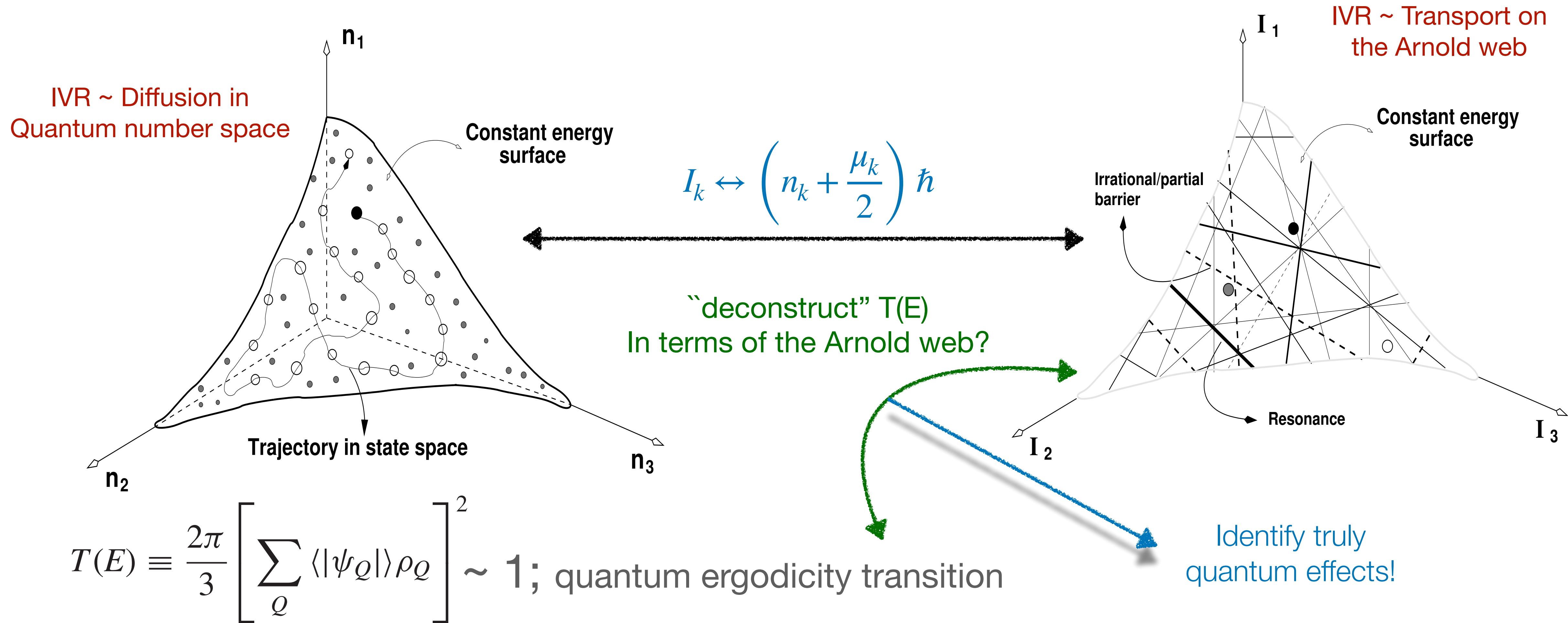
$\mathbf{H} = \mathbf{E}$  is 5-dimensional

Pick a  $(\theta_1, \theta_2, \theta_3)$  angle slice

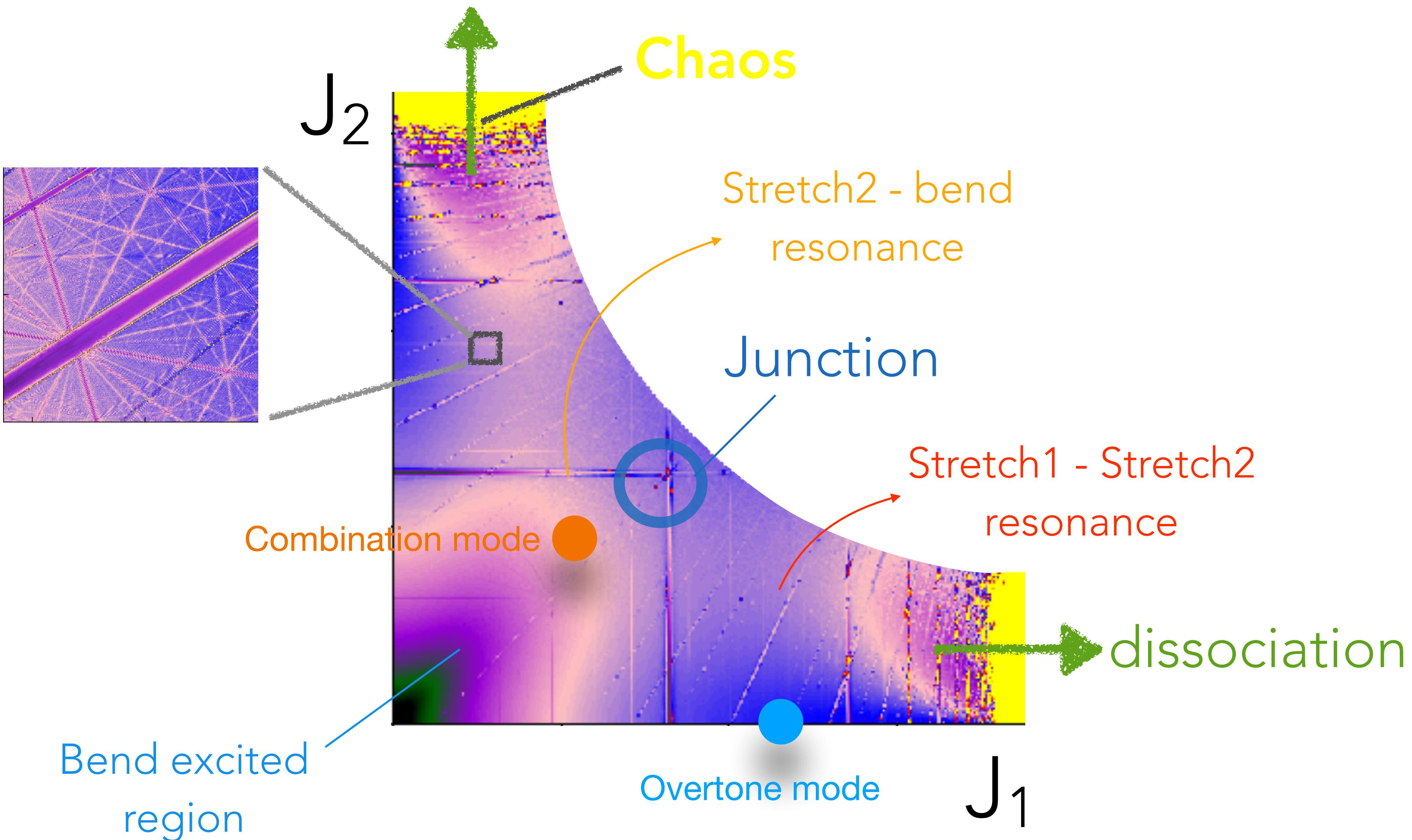
Map the Arnold web in  $(J_1, J_2, J_3)$  action space

# Advantage of using action space

Classical-Quantum correspondence...

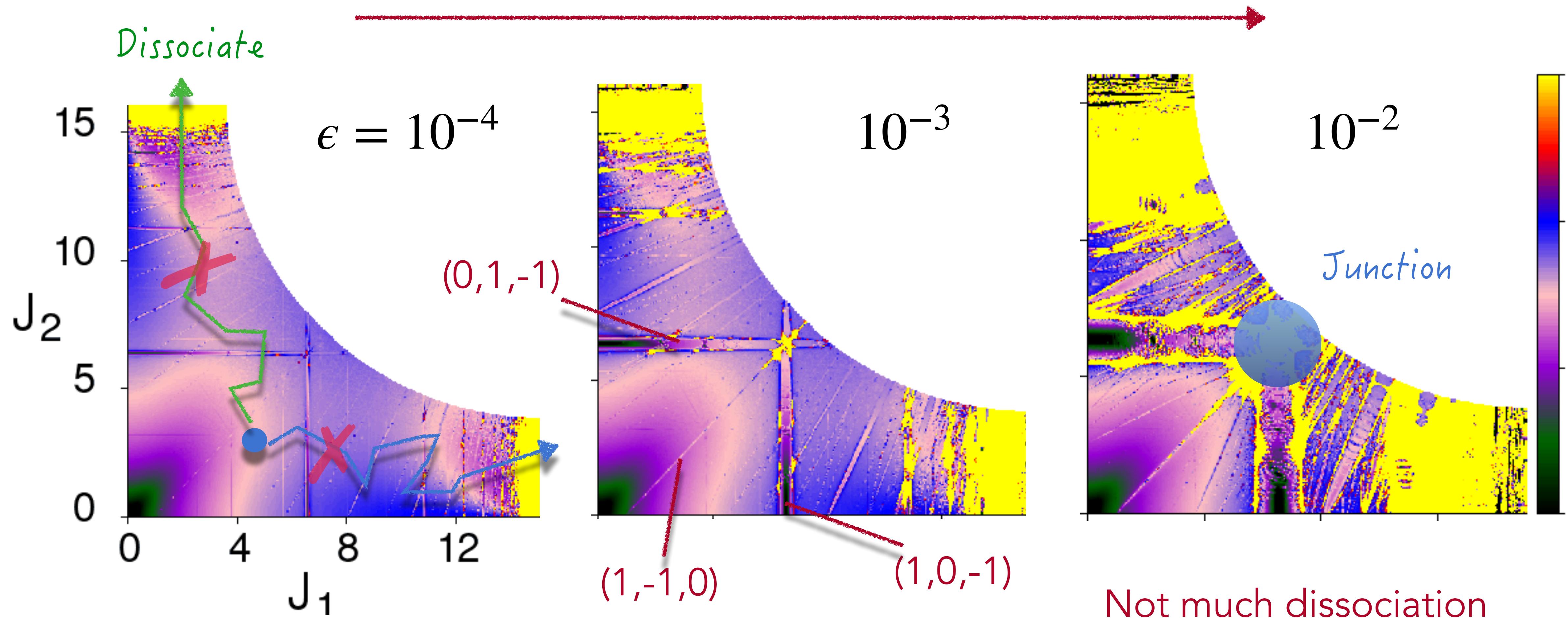


# WHAT TO EXPECT?



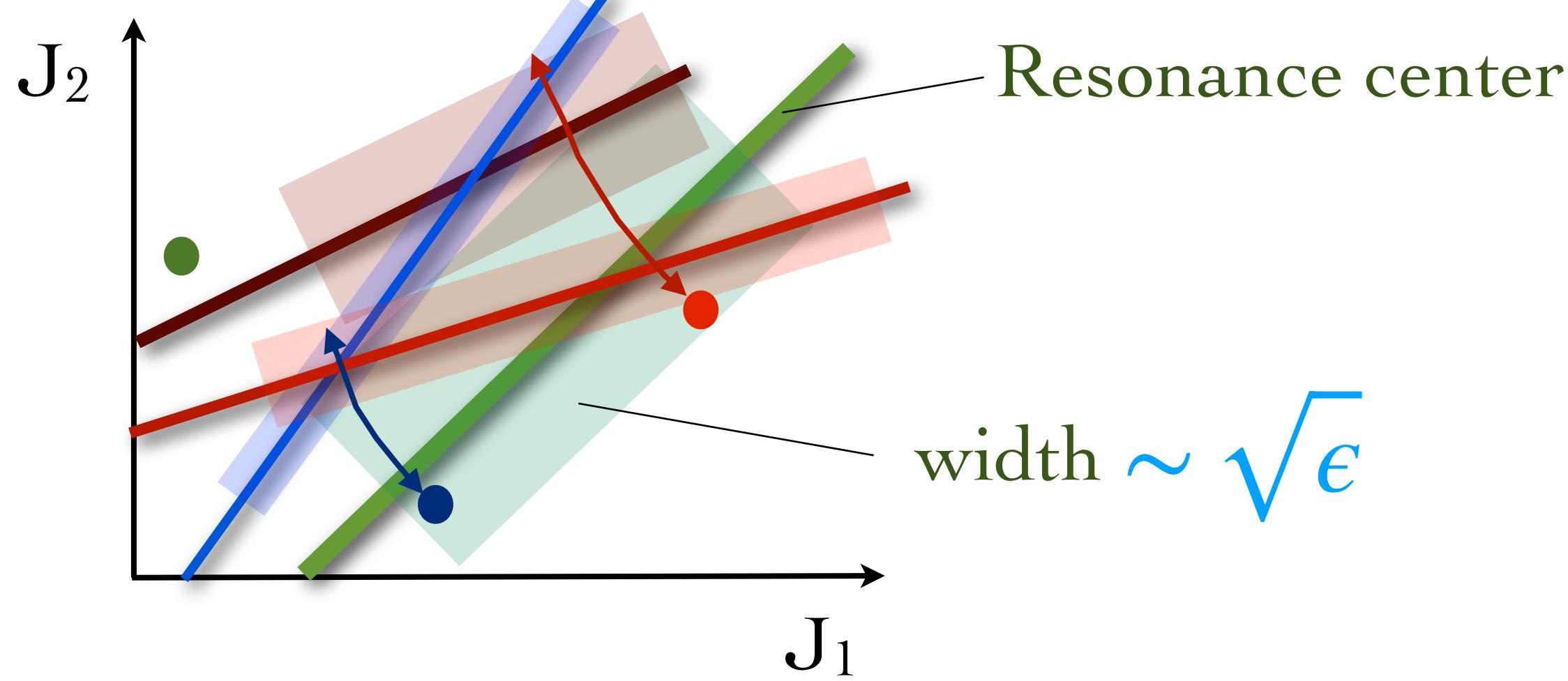
# Low coupling strengths

KAM - Nekhoroshev - Chirikov ...



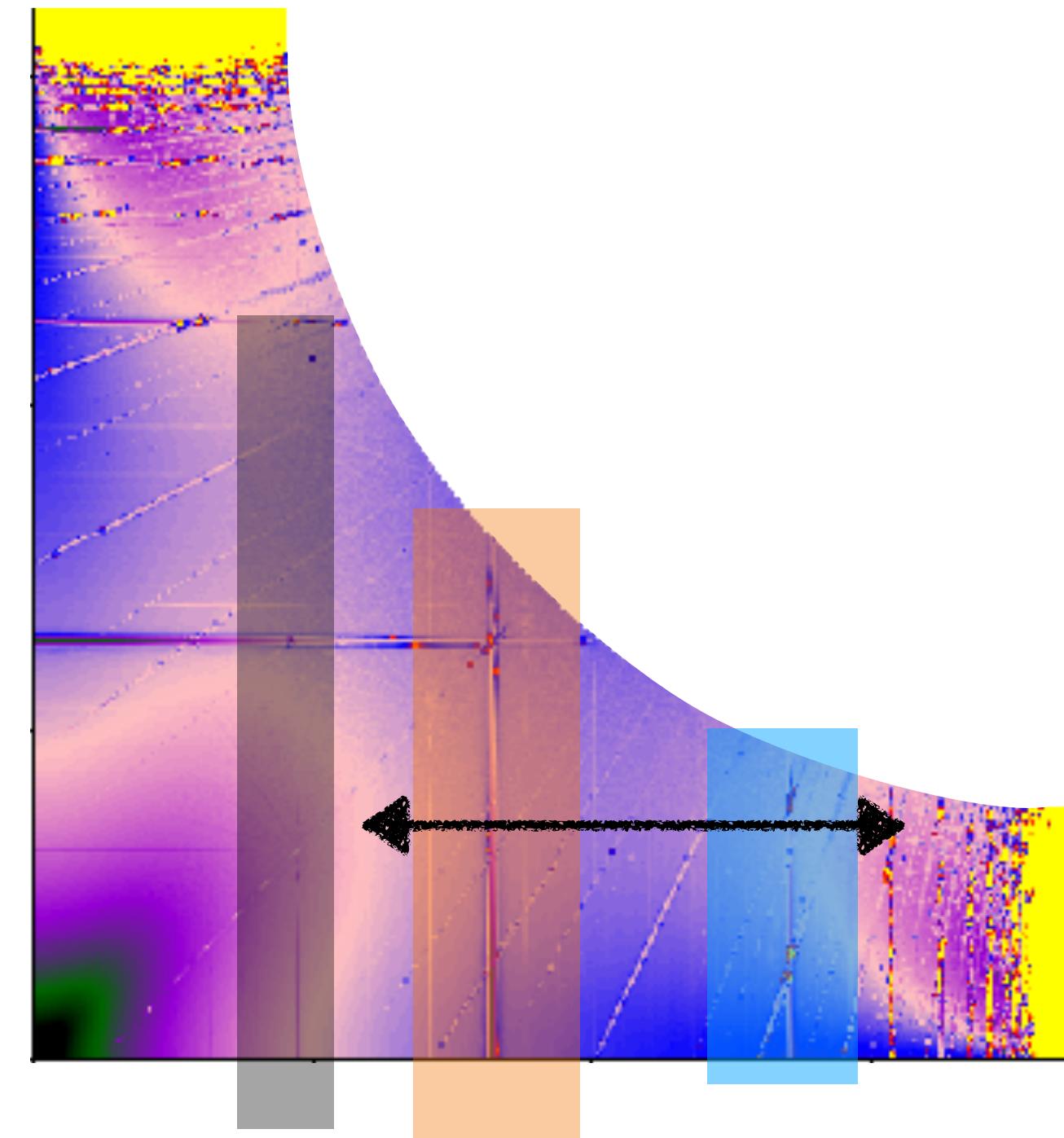
# Transition to statisticality?

Large scale resonance overlap...[“Chirikov” criterion]



2 degrees of freedom  
Oxtoby and Rice, J. Chem. Phys. 65, 1676 (1976)

> 2 degrees of freedom?

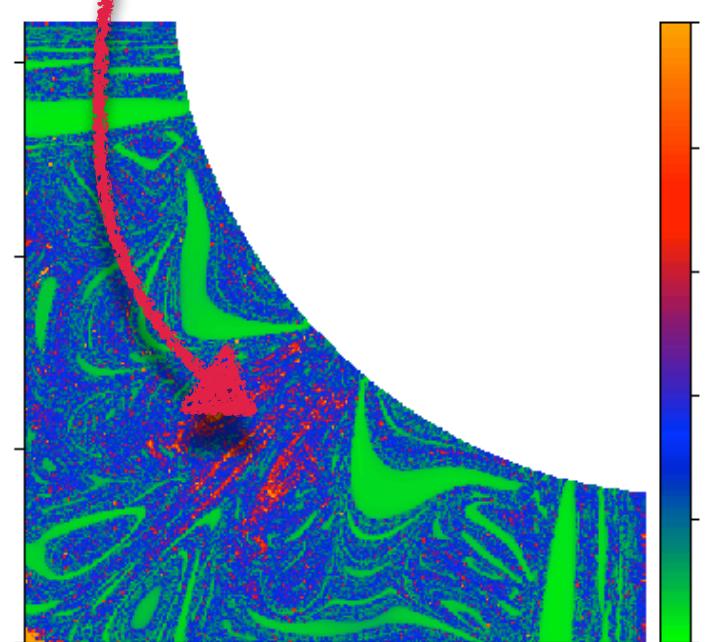


$$\epsilon_c \sim 0.1 - 0.2$$

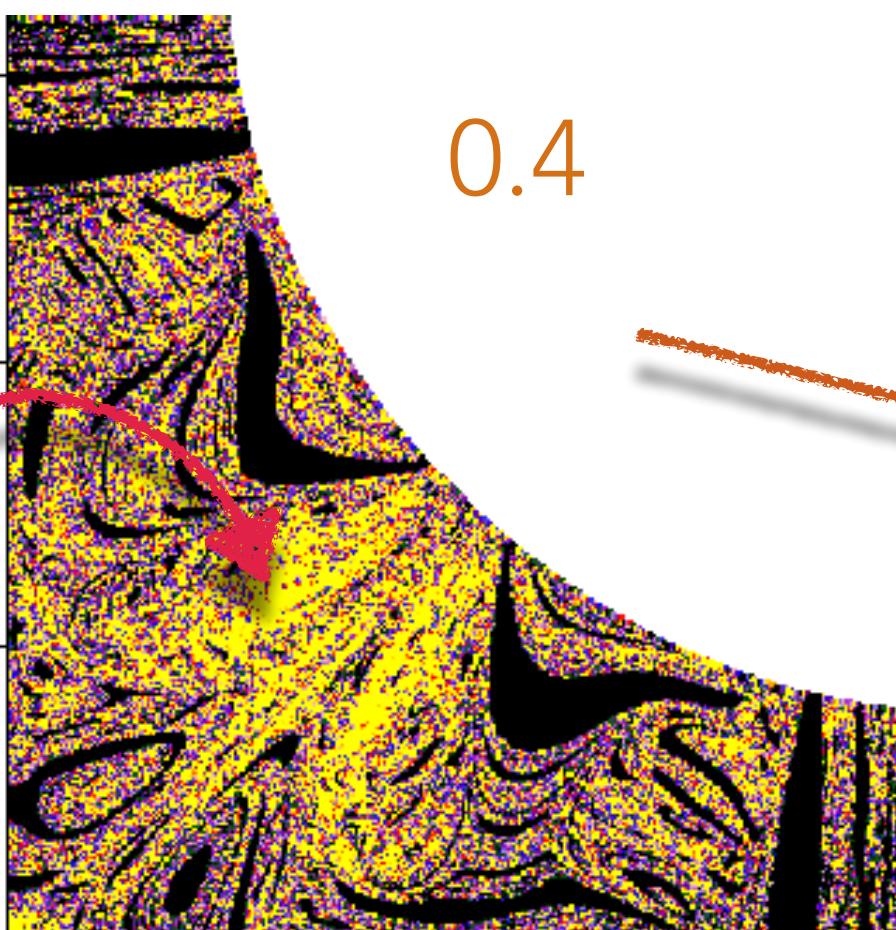
# However....

For  $\varepsilon \sim 0.4$

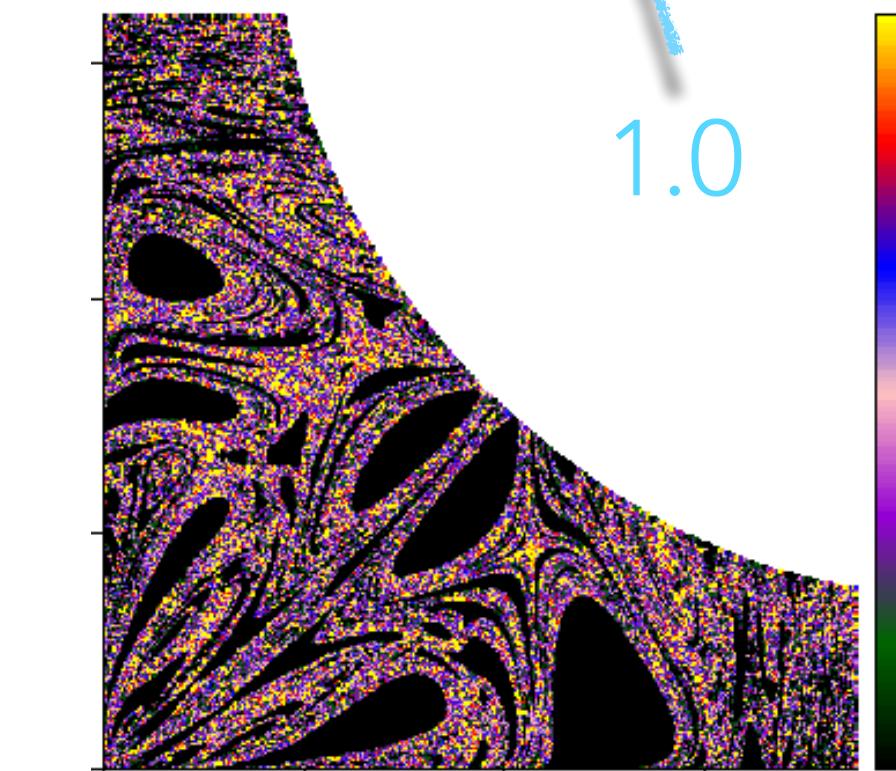
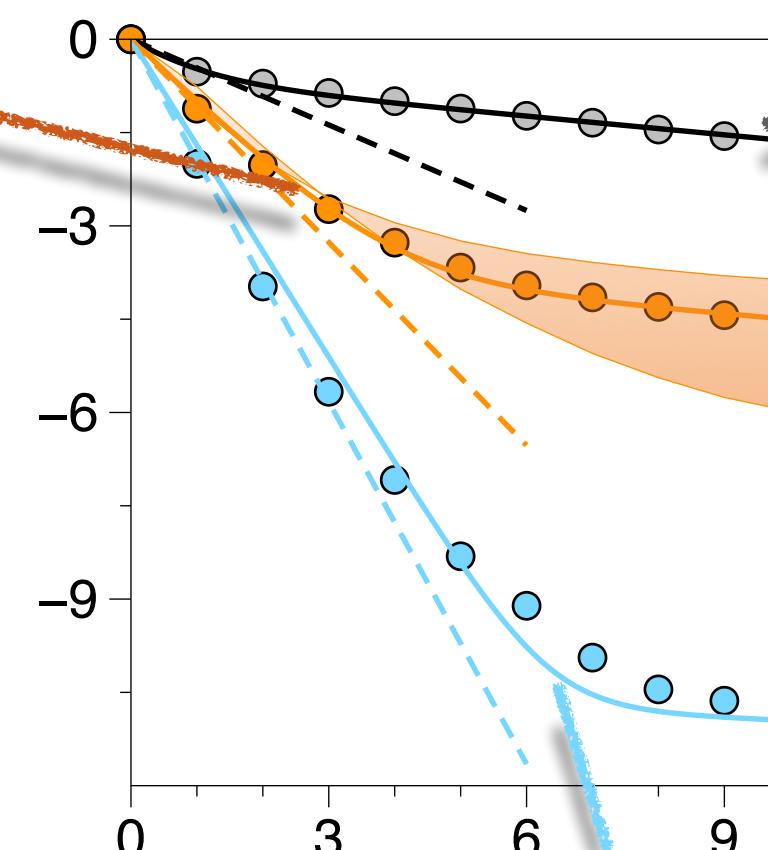
not much of  
a web!



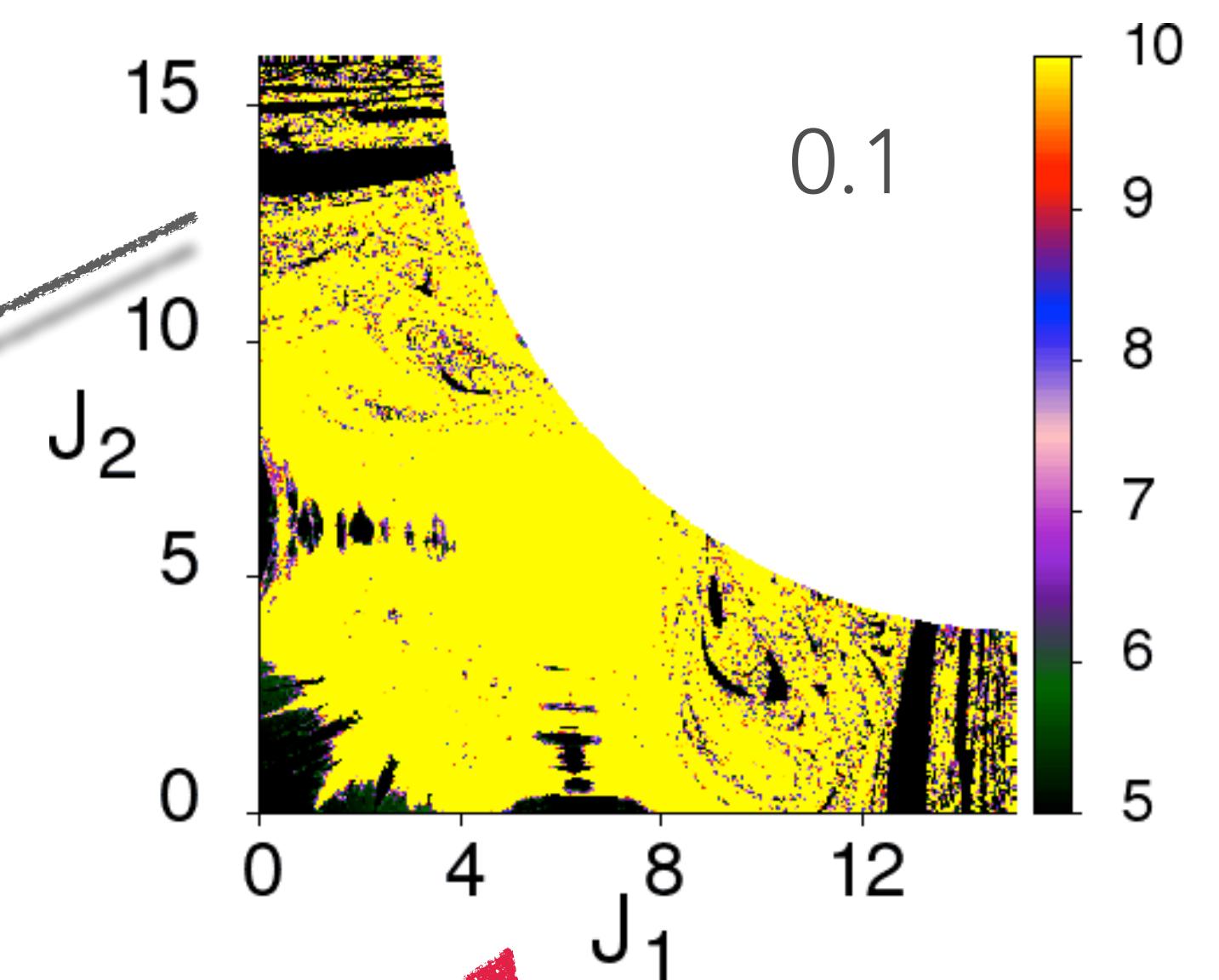
Lifetime ( $\text{ps}$ )



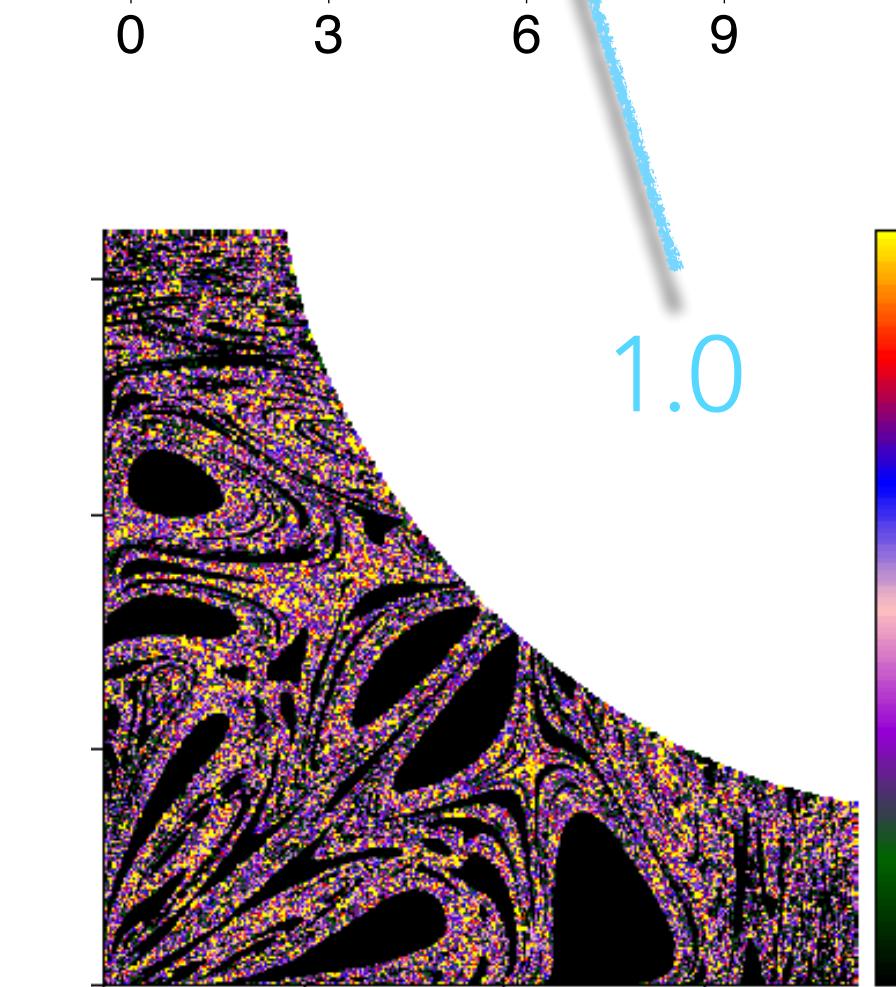
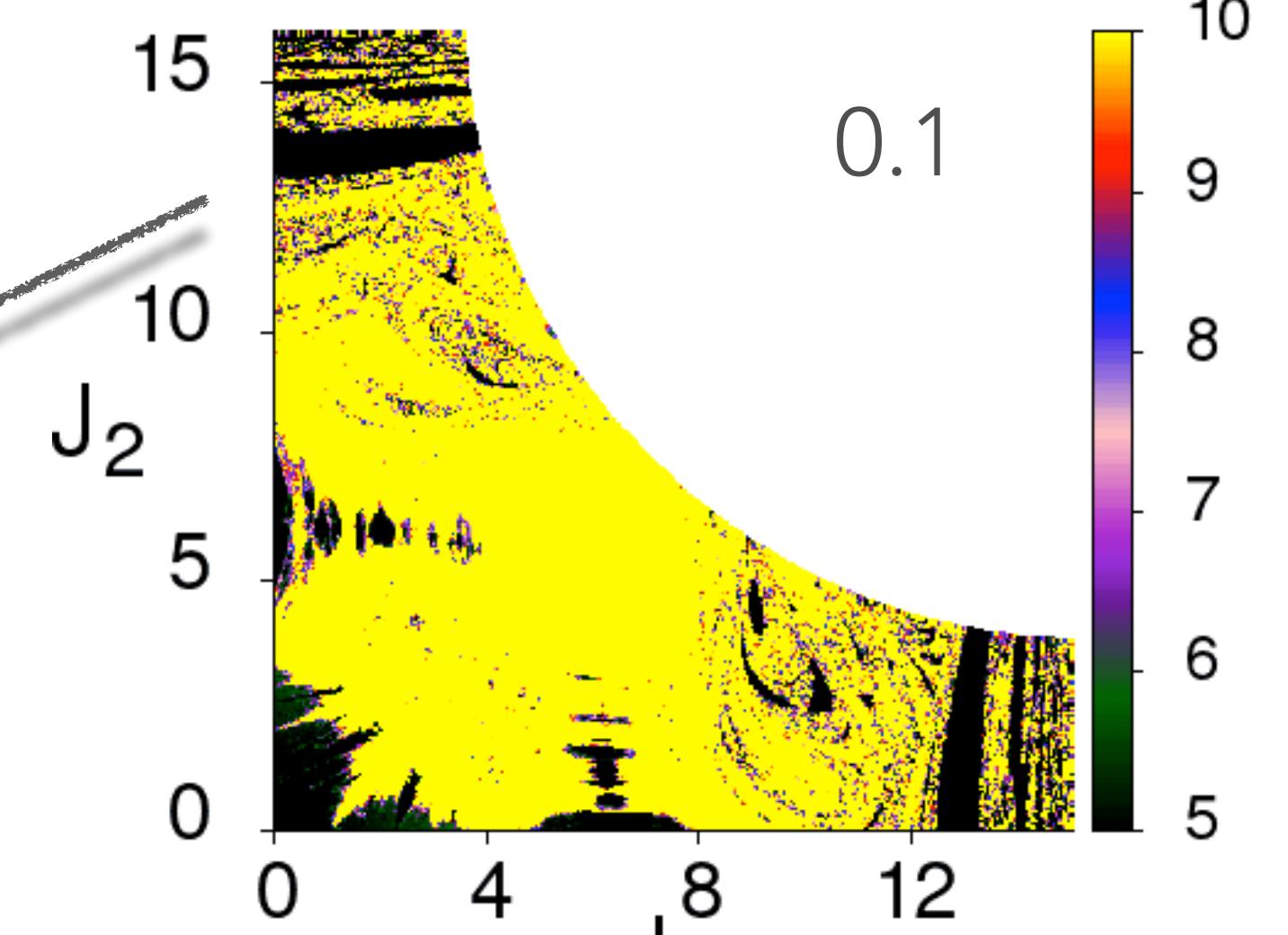
0.4



1.0



0.1

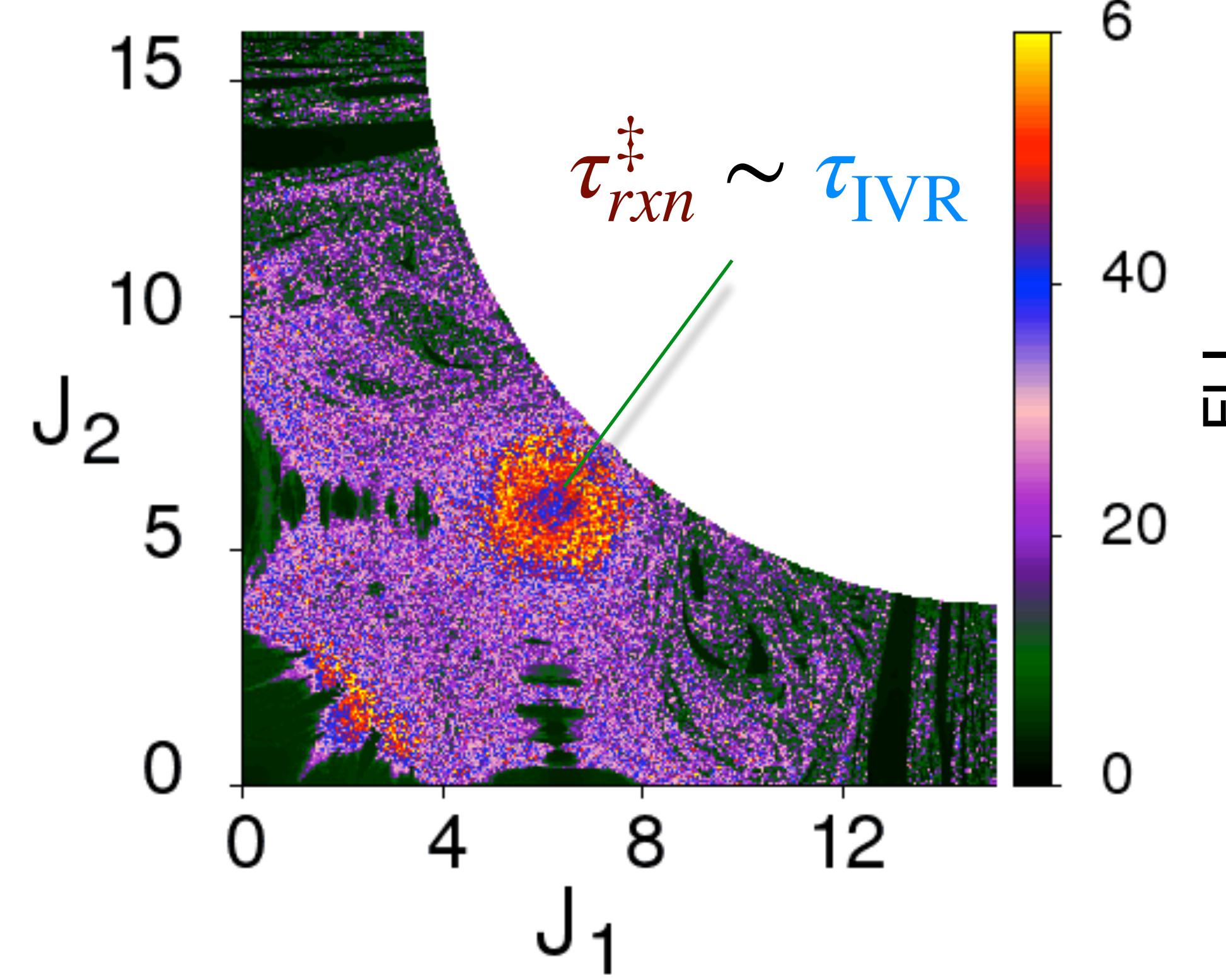


Lifetime ( $\text{ps}$ )

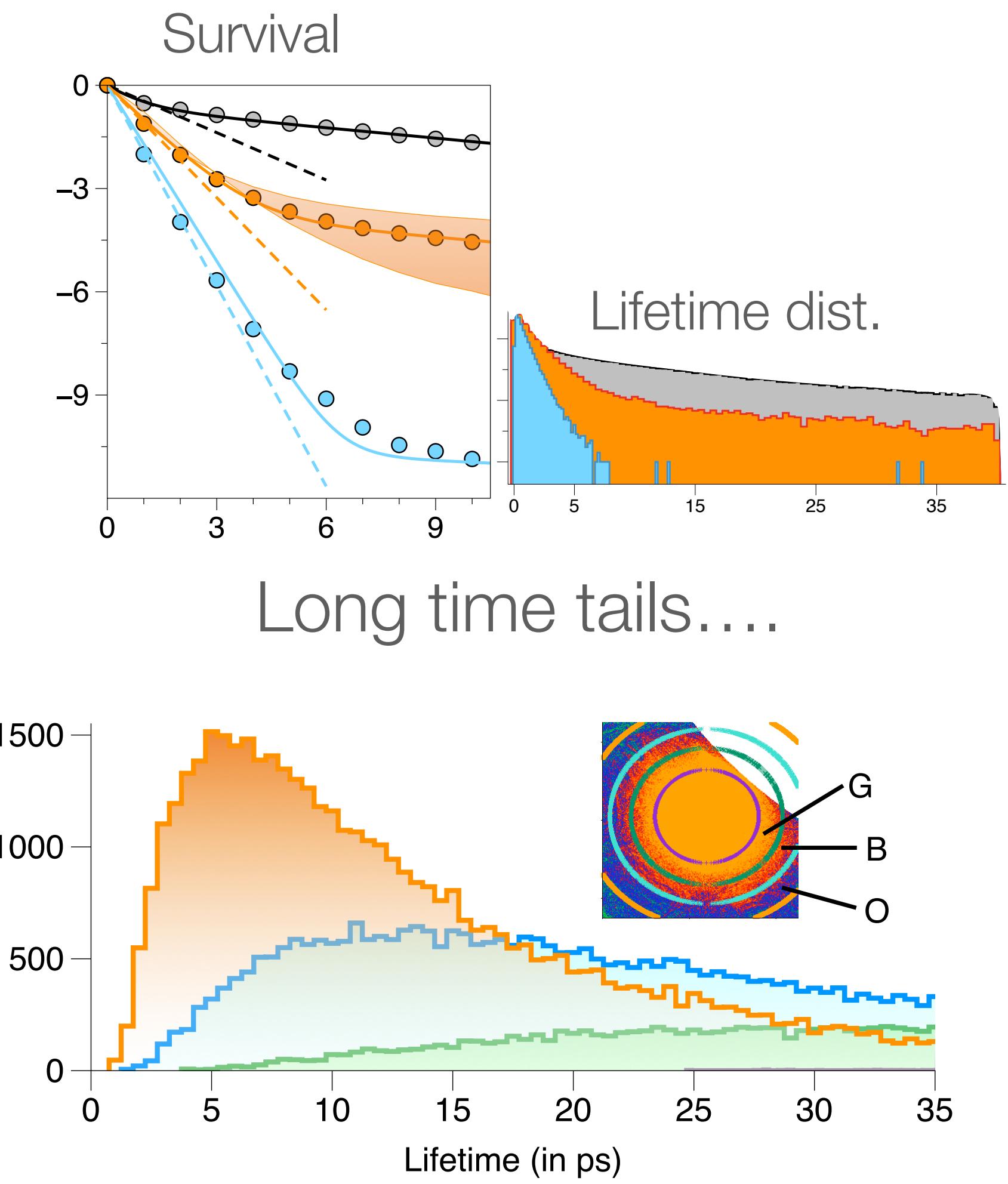
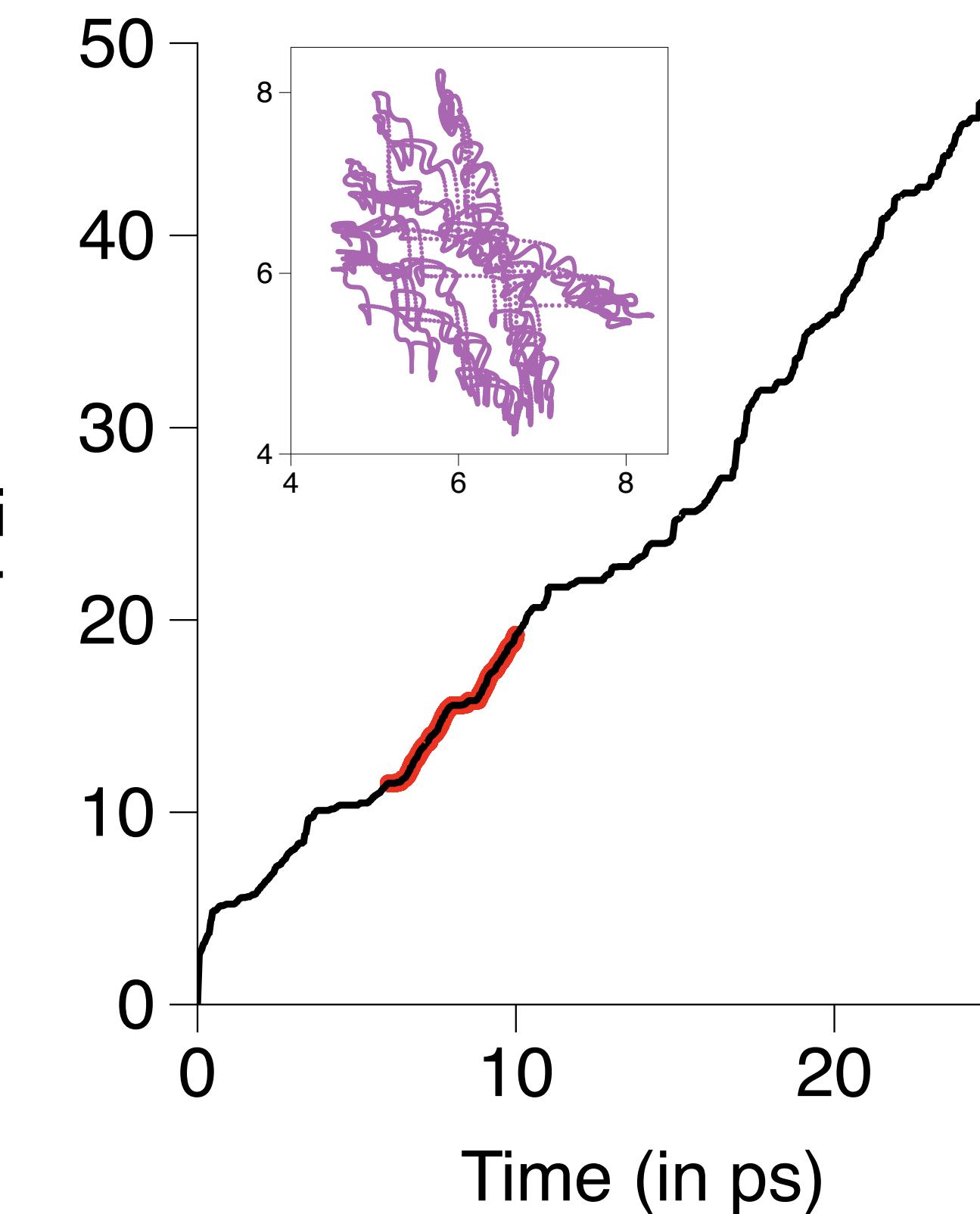
# Junctions are key!

Local dynamical traps...stable chaos...

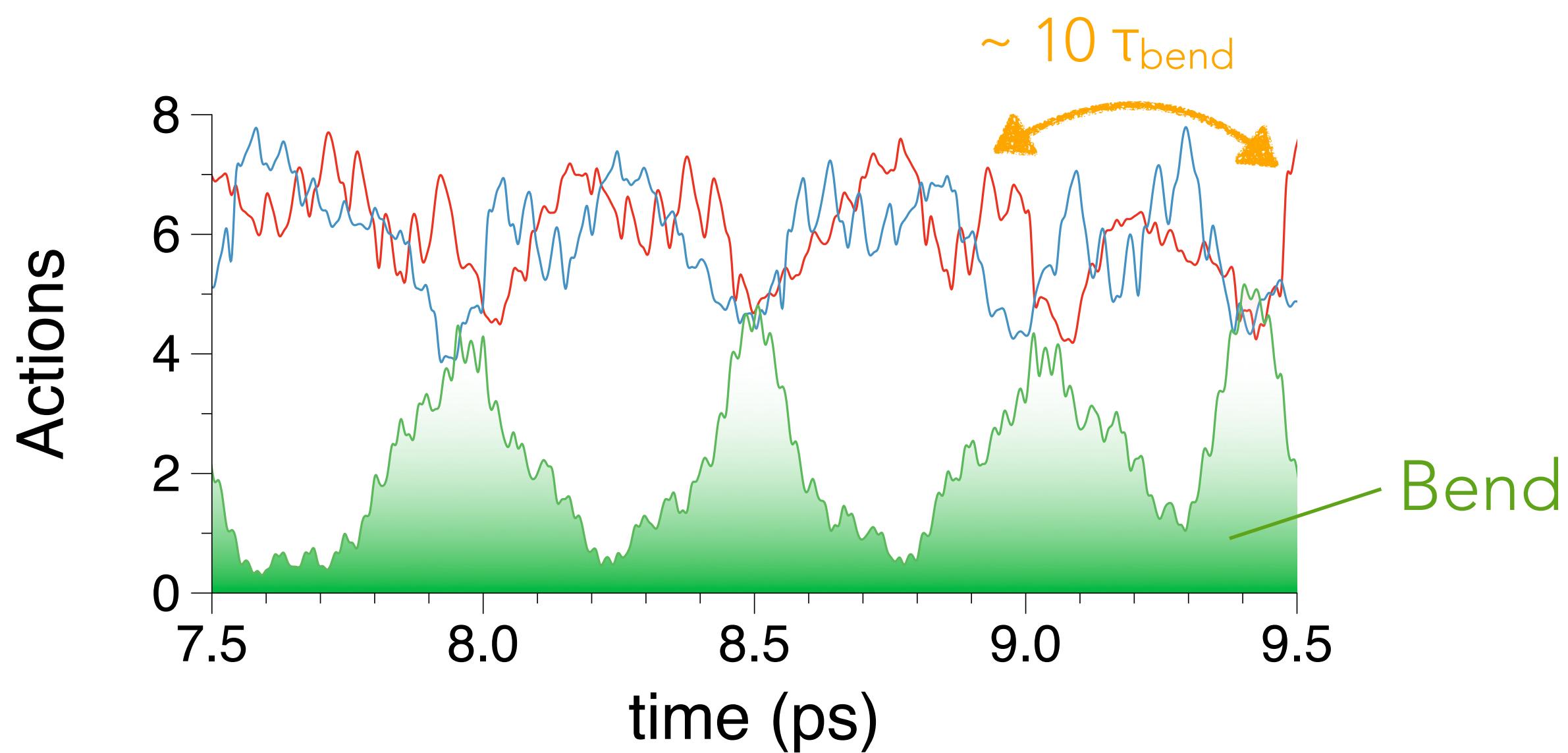
$$\varepsilon = 0.1$$



Expanded FLI scale...



# Correlated dynamics...

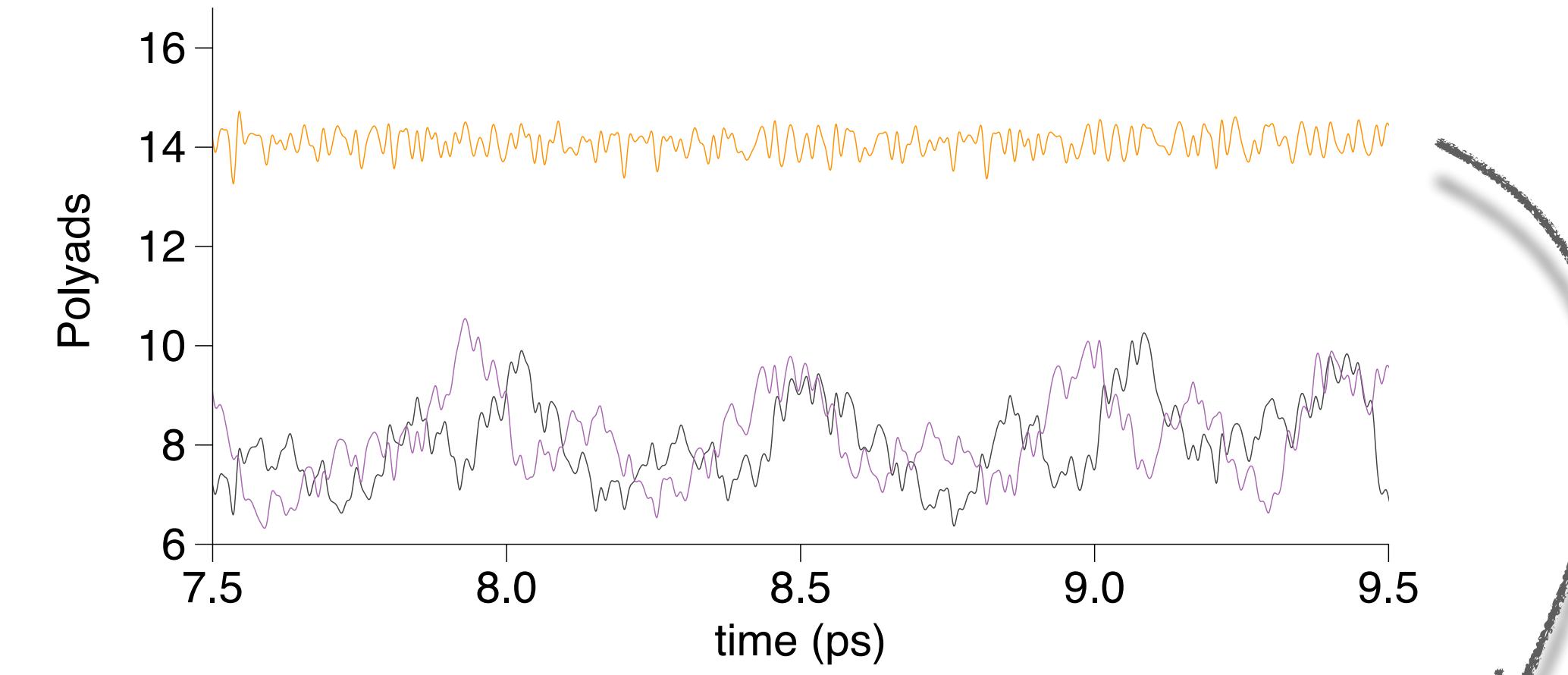


OCS molecule

Martens, Davis, Ezra , Chem. Phys. Lett. (1987)

Paškauskas, Chandre, Uzer, J. Chem. Phys (2009)

Vela-Arevalo, Wiggins (2002)



$$P \equiv J_1 + J_2 + J_3 ; [H, P] \approx 0$$

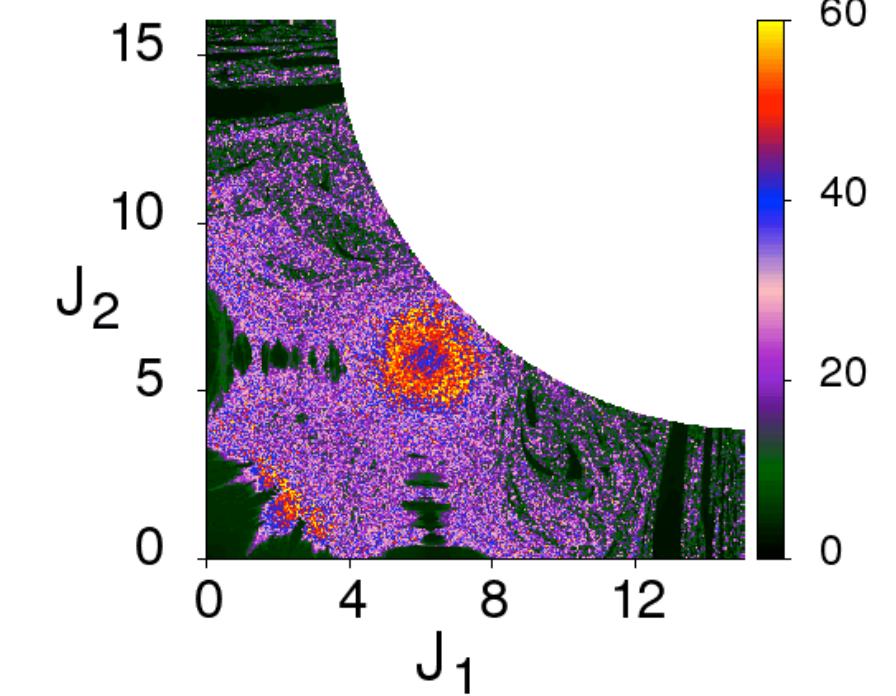
Local integrals of motion



Holme, Levine, Chem. Phys. (1989)

Atkins & Logan, J. Chem. Phys. (1992)

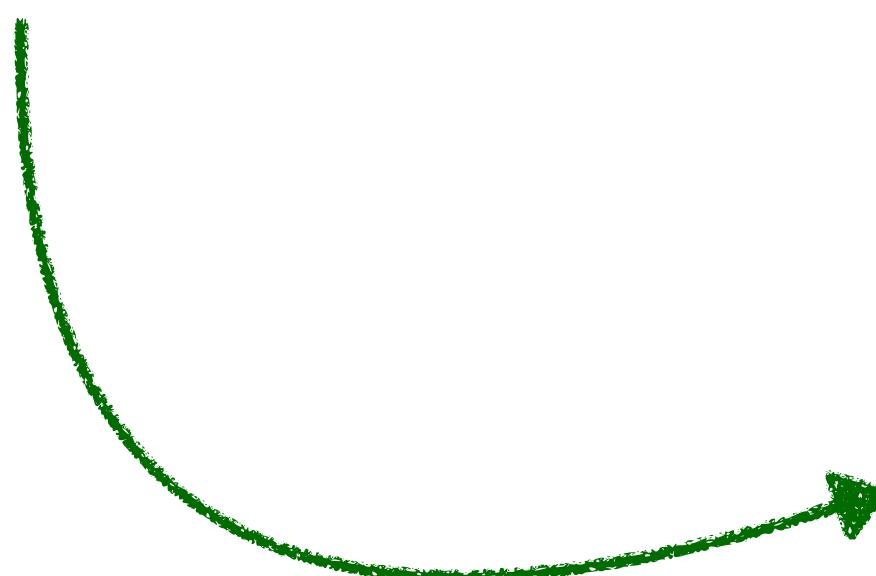
Reinhardt, Hutchinson, Sibert, Hynes,:::



# Connections ... only a glimpse

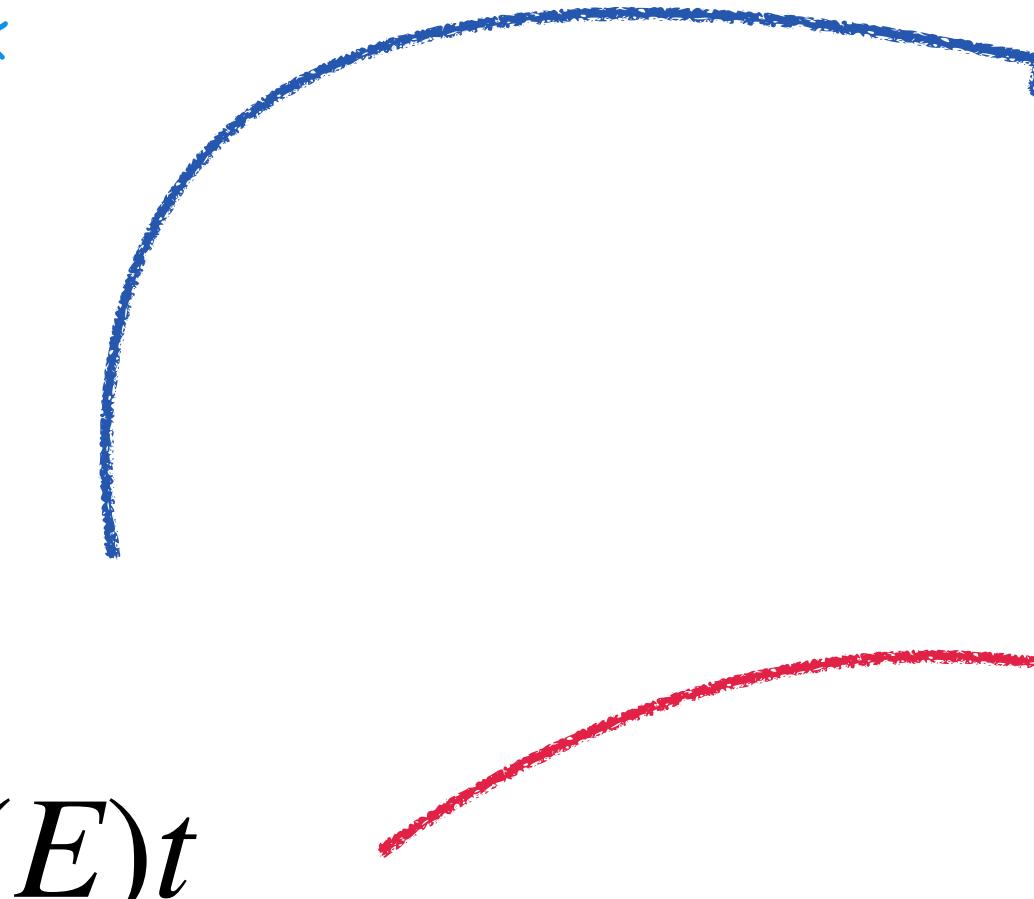
Lange, Bäcker & Ketzmerick  
EPL 116, 30002 (2016)

$$\frac{N(t)}{N(0)} \neq e^{-k(E)t}$$



Exponential/power law  
[Shoiguchi, Li, Komatsuzaki, Toda  
2008]

Shirts & Reinhardt  
JCP 77, 5204 (1982)



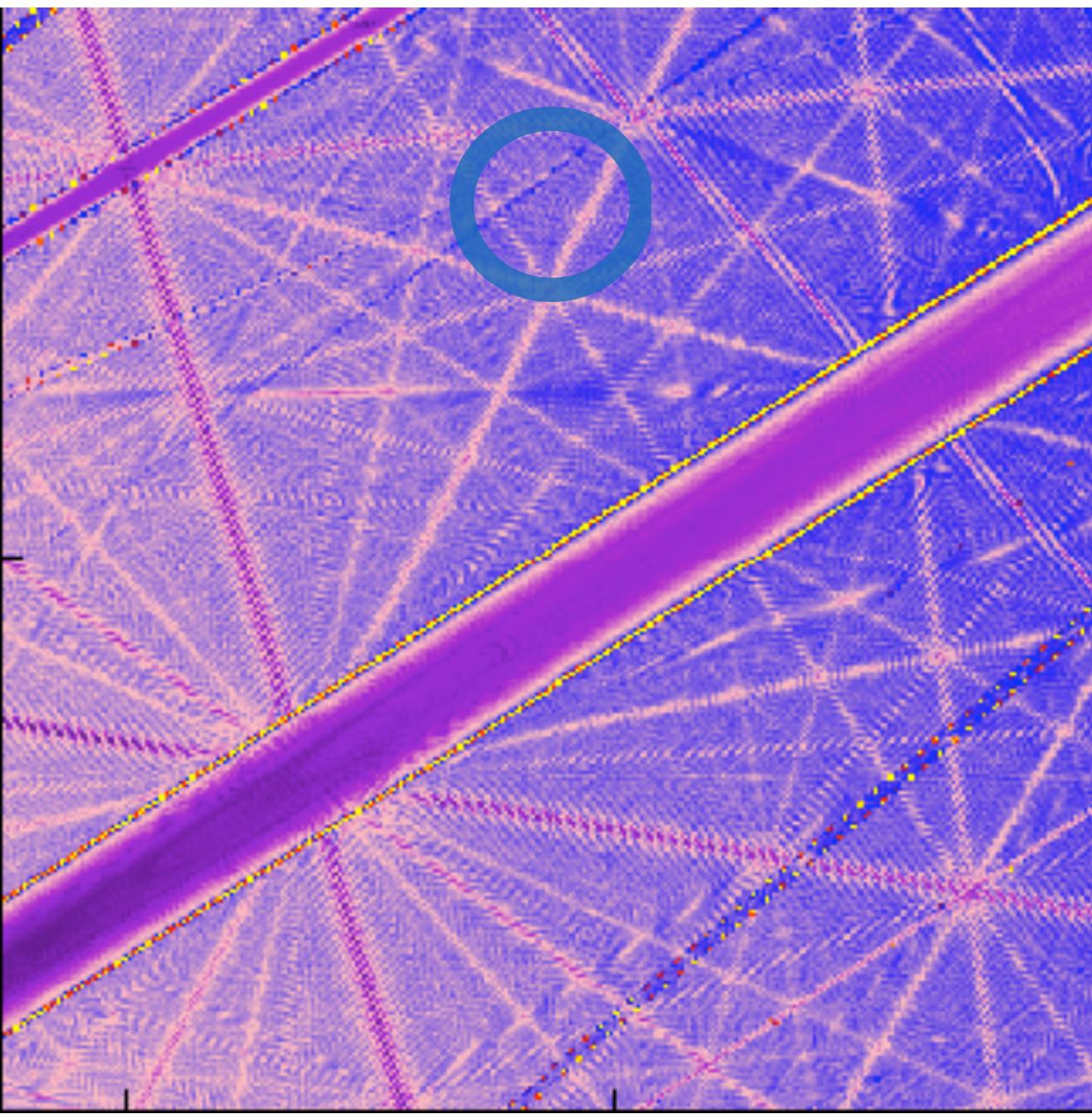
Multi-exponential  
[Marcus, Hase, Swamy  
1984]

Paskauskas, Chandre, Uzer  
JCP 130, 164105 (2009)

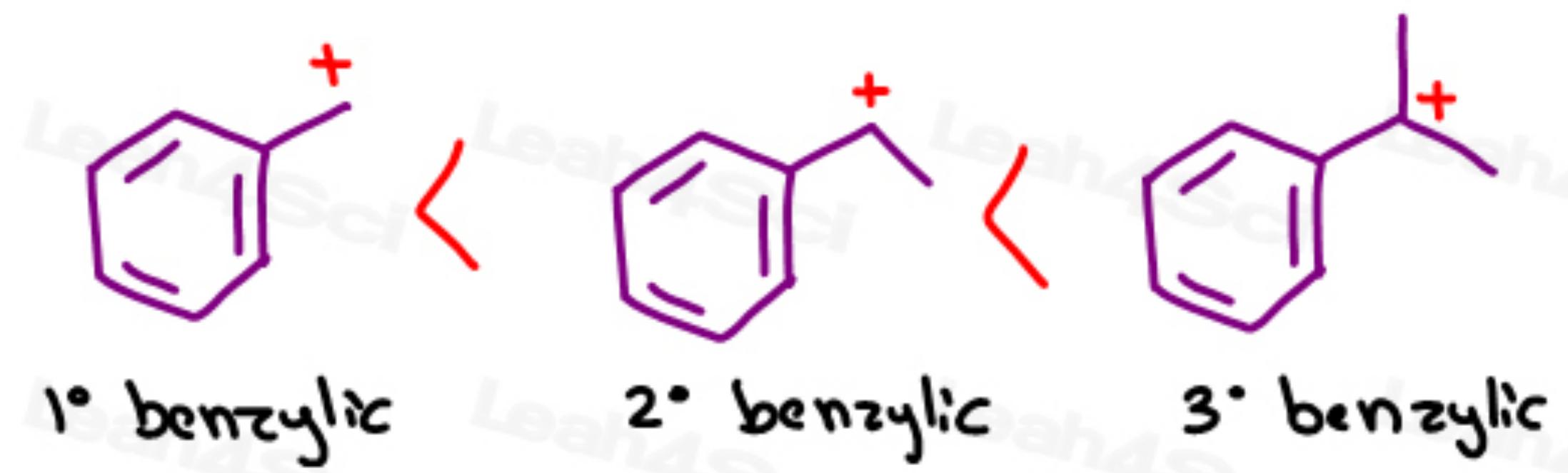
Stretched exponential/power law  
[Ezra, Waalkens, Wiggins 2009]

Leitner & Wolynes  
JCP 105, 11226 (1996)

# An analogy of sorts ...



At a junction, independent  
resonances delocalise energy  
“Dynamical stability”



More resonant structures  
delocalise charge  
“Energetic stability”

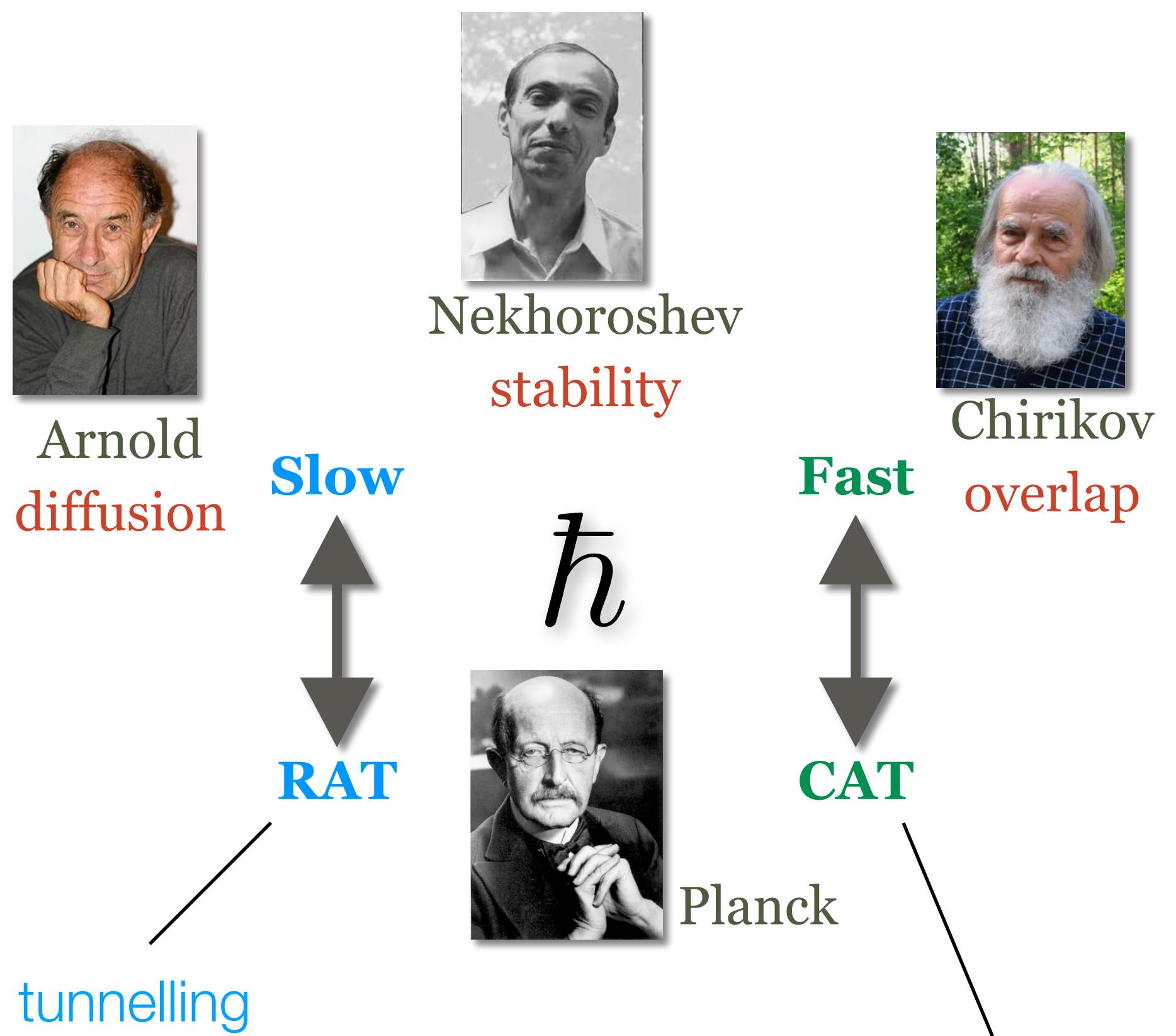
# 1. Is quantum sensitive to the junction?

~ Yes, modulo dynamical tunnelling

Quantum ergodicity transition:

$$T(E) \equiv \frac{2\pi}{3} \left[ \sum_Q \langle |\psi_Q| \rangle \rho_Q \right]^2 \sim 1$$

Conjecture: Influence of junctions minimal?



$\text{SCCl}_2$  IVR dynamics: Manikandan and Keshavamurthy, PNAS (2014)

Martens' model & connections to LRMT: Karmakar and Keshavamurthy, PCCP (2020)

# 2. Stable chaos

## Nekhoroshev's theorem

### An example of stable chaos in the Solar System

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5 place Jules Janssen, F-92195 Meudon, France



MANY planets have been shown to have chaotic instabilities in their orbital motions, but the long-term significance of this is not fully understood<sup>1</sup>. The eccentricity of Mercury, for example, changes by about 25% of its value over 40 times the Lyapunov time<sup>2,3</sup> (the e-folding time for divergence of nearby orbits), but the orbit of Pluto, in an integration lasting 50 Lyapunov times<sup>4</sup>, shows no significant change. Here we show that the orbit of the near-Jupiter asteroid 522 Helga is chaotic, with an unusually short Lyapunov time of 6,900 yr. We integrate its motion, including perturbations from the outer giant planets, over a period 1,000 times longer than this, and find no significant instability. Chaos in the orbit of 522 Helga is caused by a 7:12 resonance with the orbit of Jupiter, but the size of the chaotic region in phase space is small; stability is ensured because the eccentricity and precession of the orbit are such that it avoids close encounters with Jupiter. Asteroid orbits with larger proper eccentricity would, we suggest, be genuinely unstable, consistent with the sparse asteroid population near Helga. Although Helga is the first clear-cut example of a stable chaotic orbit, we argue that ‘stable chaos’ may be a rather common feature of Solar System dynamics.

Given the set up as above, the celebrated Nekhoroshev's theorem provides results regarding the stability of various types of initial conditions on the Arnold web. Roughly speaking<sup>§</sup>, for  $\varepsilon < 1$ , an initial condition  $(\mathbf{J}(0), \boldsymbol{\theta}(0))$  satisfies

$$\|\mathbf{J}(t) - \mathbf{J}(0)\| \leq \varepsilon^{1/2(f-m)} \quad (25)$$

for times

$$|t| \leq \exp(c\varepsilon^{-1/2(f-m)}) \quad (26)$$

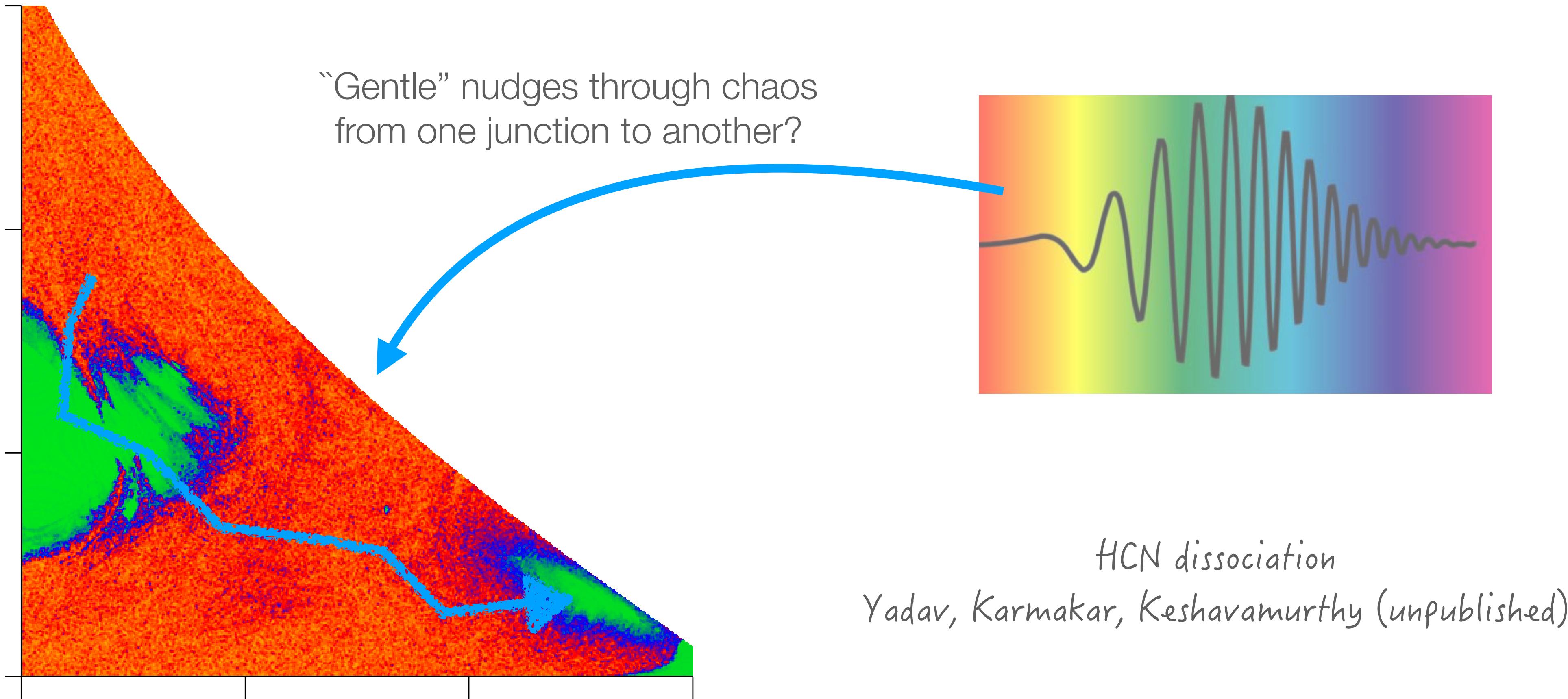
f: degrees of freedom  $\sim 3N-6$ , m: multiplicity of the junction  
(How many independent resonances intersect?)

For large molecules junctions of  
various multiplicities....non-RRKM?

Will quantum effects “de-trap”?

### 3. Use the junctions to control?

“Field chirping” for polyatomic molecules...



# Quo vadis?

- Which junctions important? Quantum, Classical....
- Mapping the web for large systems?
- Junctions unimportant in the presence of a “bath”?
- Nonadiabatic dynamics....



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