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Surprisal of a **quantum** state: Dynamics, compact representation and coherence effects



Why surprisal?

$$\hat{I} = -\ln \hat{\rho}$$

- Effective computational method*
- High degree of the data compaction
- Novel insights

*Extensively applied in different fields:

-> to analyse the energy disposal in reaction

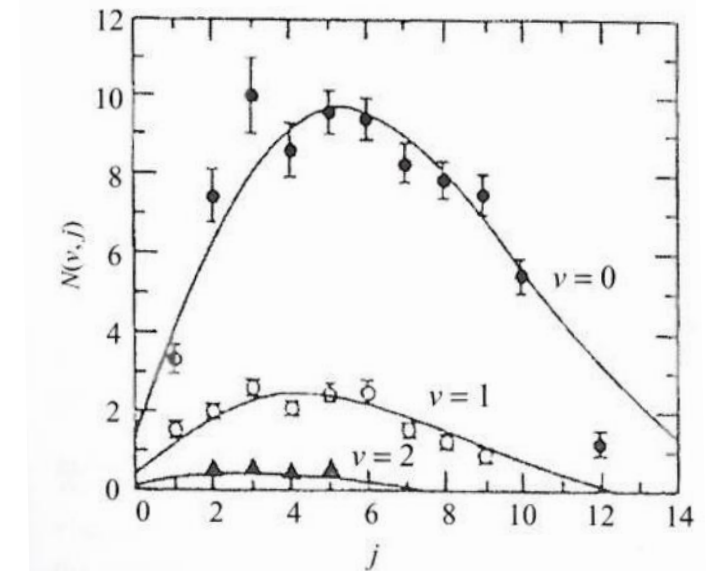
dynamics

R. C. Bernstein and R. D. Levine, "Entropy and Chemical Change. I. Characterization of Product (and Reactant) Energy Distributions in Reactive Molecular Collisions: Information and Entropy Deficiency," J. Chem. Phys. **57**, 434-449 (1972)

-> in system biology to characterize

dominant behavior patterns

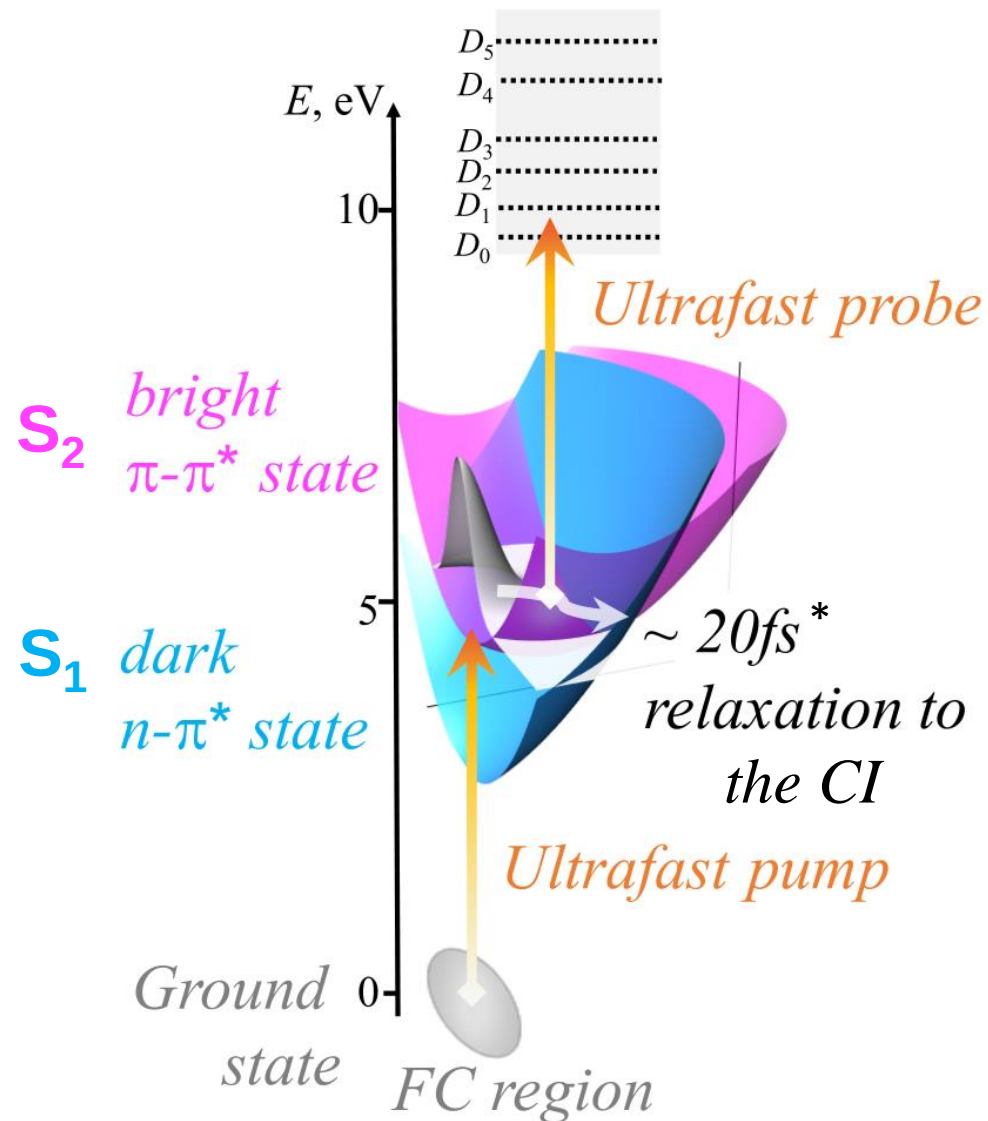
Domini et al. Surprisal analysis characterizes the free energy time course of cancer cells undergoing epithelial-to-mesenchymal transition. PNAS, 111:13235–13240 (2014)



Experimental HD state distributions measured for H+D₂ collision

Ultrafast non-radiative relaxation

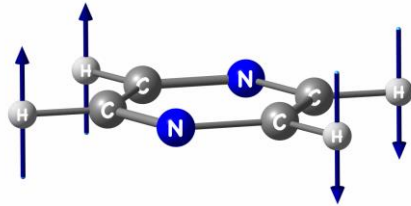
Dynamics in excited electronic states of pyrazine



*Real-time resolved pump-probe
photoelectron spectroscopy experiment:
Y.-I. Suzuki et al, *J Chem Phys*, 2010, 132, 174302

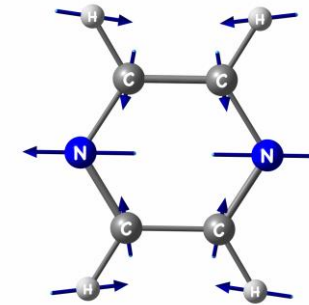
Ultrafast non-radiative relaxation

Dynamics in excited electronic states of pyrazine



Coupling mode Q_1^*

->Diabatic coupling between S_2 and S_1 : κQ_1



Tuning mode Q_2

-> Nuclear motion towards CI

MCTDH: 24-modes diabatic picture for the

* R. Schneider, W. Domcke, Chem Phys Lett, 1988, 150, 235

** I. Burghardt et al, Phys. Scr. 2006, 73, C42

Non-adiabatic quantum dynamics on multiple electronic states

Challenge

Can we have a **compact** representation
for the **quantum treatment** of the dynamics?

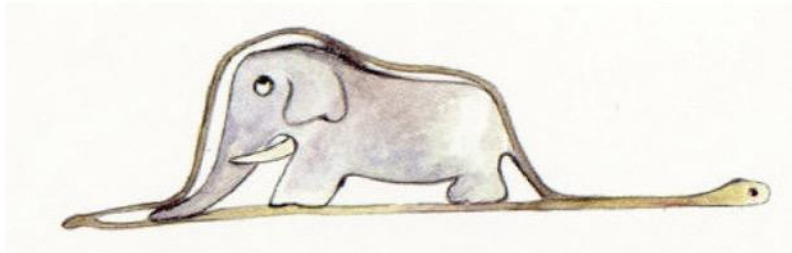
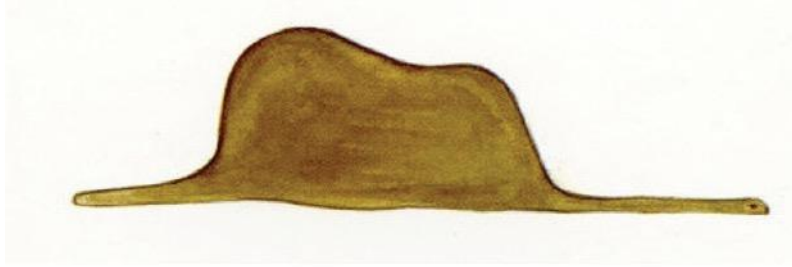
'Quantum' features

- **Branching** of the wave packets
in configurational space
- **Coherence effects** on the
vibrational and/or electronic dynamics
- **Electronic energy redistribution**
for the coherent dynamics on multiple electronic states

**Let us try to address the problem
from a fresh perspective!**



Algebraic approach in chemical dynamics



The state of the system, the density matrix:

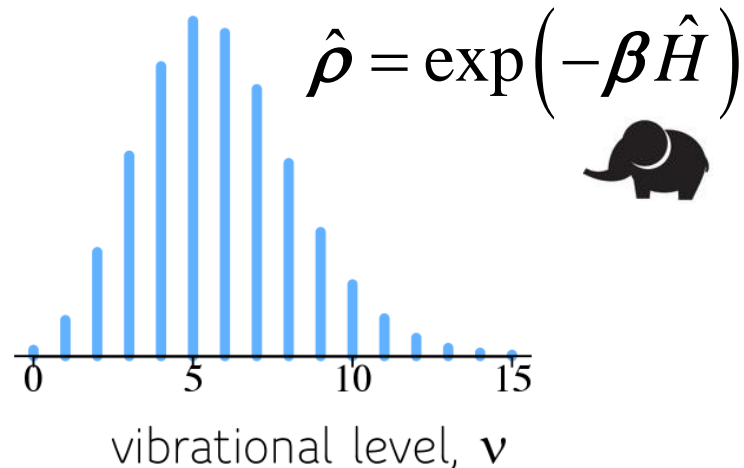
$$\hat{\rho} = \exp \left(- \sum_k \lambda_k \hat{A}_k \right)$$

$$\hat{A}_k : \hat{R}, \hat{P}, \hat{H}, \dots$$

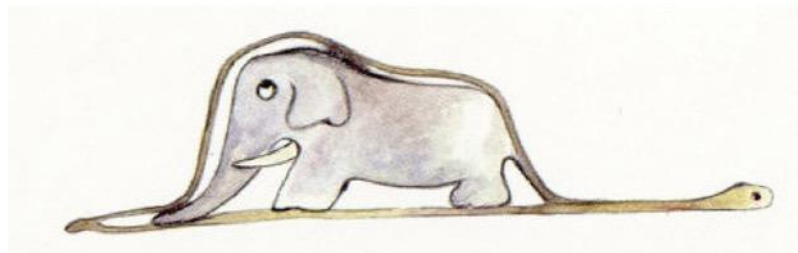
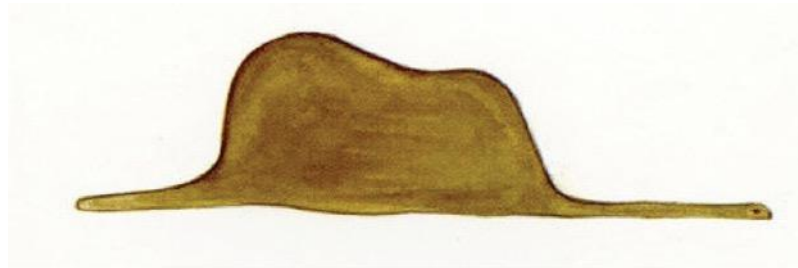
$$\hat{a}, \hat{a}^\dagger, \hat{a}^\dagger \hat{a}, \dots$$

- constraints*

*More examples of different algebras:
Dynamical symmetry, C.E. Wulfman (2010)



Algebraic approach in chemical dynamics



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$$\hat{\rho} = \exp \left(- \sum_k \lambda_k \hat{A}_k \right)$$

$$\hat{A}_k : \hat{R}, \hat{P}, \hat{H}, \dots$$

- constraints*

$$\hat{a}, \hat{a}^\dagger, \hat{a}^\dagger \hat{a}, \dots$$

$$\hat{\rho}(t=0) = \exp \left(- \sum_k \lambda_k(t=0) \hat{A}_k \right)$$

$$\hat{U}(t) = \exp \left\{ - \frac{i}{\hbar} \hat{H} t \right\}$$

$$\hat{\rho}(t) = \hat{U}(t) \hat{\rho}(t=0) \hat{U}^\dagger(t)$$

$$= \exp \left[- \sum_k \lambda_k(t=0) \left\{ \hat{U}(t) \hat{A}_k \hat{U}^\dagger(t) \right\} \right]$$

Time evolution
of the **density**
via the evolution
of the
constraints

Density and its surprisal

$$\hat{\rho} = \exp\left(-\sum_k \lambda_k \hat{A}_k\right)$$

$$\hat{I} = -\ln \hat{\rho}$$

Surprisal

$$\hat{I} = \sum_k \lambda_k \hat{A}_k$$



$$\hat{I}(t=0) = \sum_k \lambda_k(t=0) \hat{A}_k$$

When time evolution of the constraints is closed:

$$i\hbar \frac{d\hat{I}}{dt} = [\hat{H}, \hat{I}] \quad [\hat{H}, \hat{I}] = \sum_k \lambda_k [\hat{H}, \hat{A}_k] \quad [\hat{H}, \hat{A}_k] = \sum_s g_{ks} \hat{A}_s$$

$$\hat{I}(t) = \sum_k \lambda_k(t) \hat{A}_k$$

Information
about the time
evolution is
compressed to
just a **vector**

$\lambda(t)$

Surprisal of a quantum state in the general case

In a given finite basis representation:

$$I_{nm}(t=0) = \langle \varphi_n | \hat{I}(t=0) | \varphi_m \rangle = \sum_k \lambda_k(t=0) A_{nm}^k$$

Equation of motion for the surprisal:

$$i\hbar \frac{d\mathbf{I}}{dt} = (\mathbf{H} \cdot \mathbf{I} - \mathbf{I} \cdot \mathbf{H})$$

Compact representation
via **dominant**

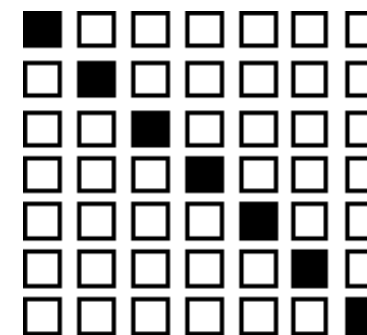
constraints:

$$\mathbf{I}(t) \approx \sum_k \lambda_k(t) \hat{A}_k$$

Straightforward transformation between them via the spectral decomposition:

$$\hat{I} = -\ln \hat{\rho} = \sum_{s=0}^N \mu_s |s\rangle \langle s|$$

$$\hat{\rho} = \sum_{s=0}^N \exp(-\mu_s) |s\rangle \langle s|$$

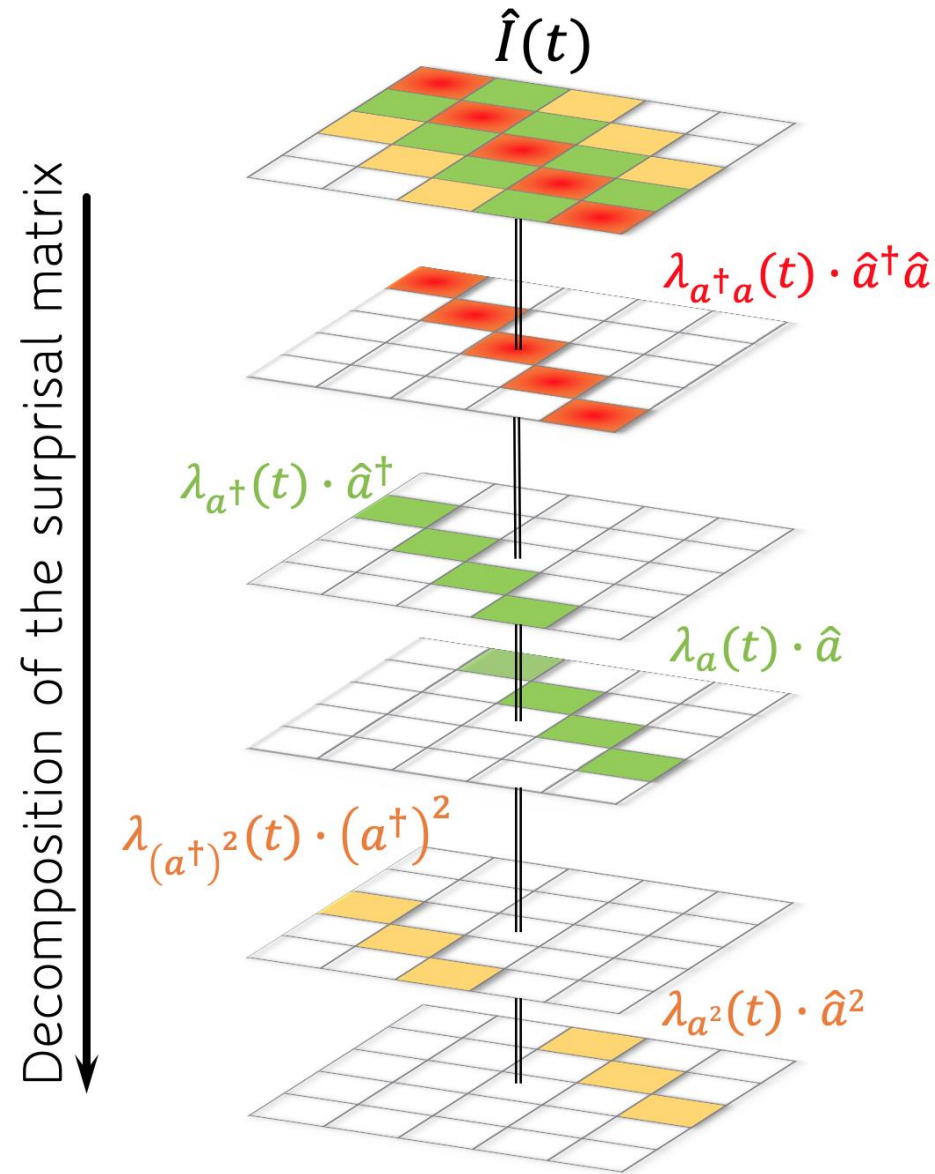


Dominant constraints in the vibrational basis

$$i\hbar \frac{d\mathbf{I}}{dt} = (\mathbf{H} \cdot \mathbf{I} - \mathbf{I} \cdot \mathbf{H})$$

$$\hat{I}(t) \approx \sum_k \lambda_k(t) \hat{A}_k$$

$$\hat{A}_k : \hat{a}, \hat{a}^\dagger, \hat{a}^\dagger \hat{a}, \dots$$



Compact representation of the non-adiabatic transfer in pyrazine

$$\{I_N, \hat{a}, \hat{a}^\dagger, \hat{a}^\dagger \hat{a}\} \otimes \{I_e, \hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z\}$$

Population in S_2

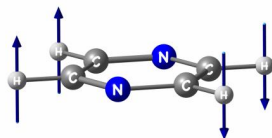
$t = 0$ fs



Population in S_1



Coupling mode



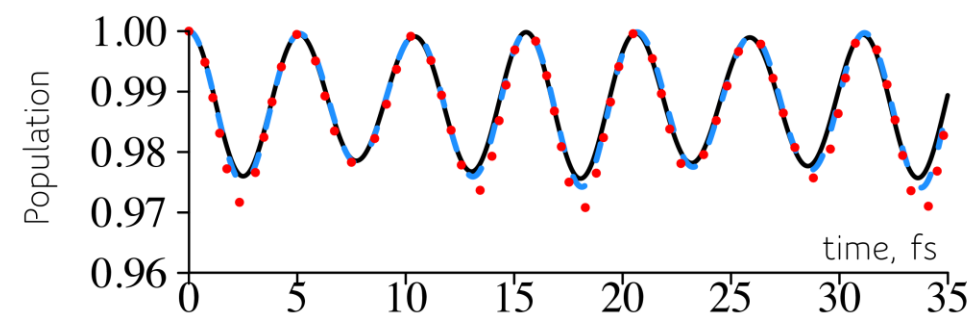
$$\hat{H} = \hat{H}_0 + \hat{V}$$

$$\hat{H}_0 = \hbar\omega \cdot I_e \otimes \left(\hat{a}^\dagger \hat{a} + \frac{1}{2} \right) + V_0 \cdot \hat{\sigma}_z \otimes I_N$$

$$\hat{V} = \hat{\sigma}_x \otimes \kappa (\hat{a} + \hat{a}^\dagger)$$

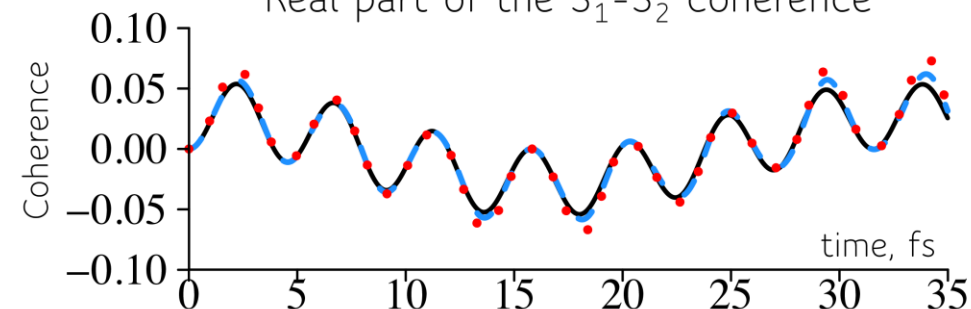
(a)

S_2 population



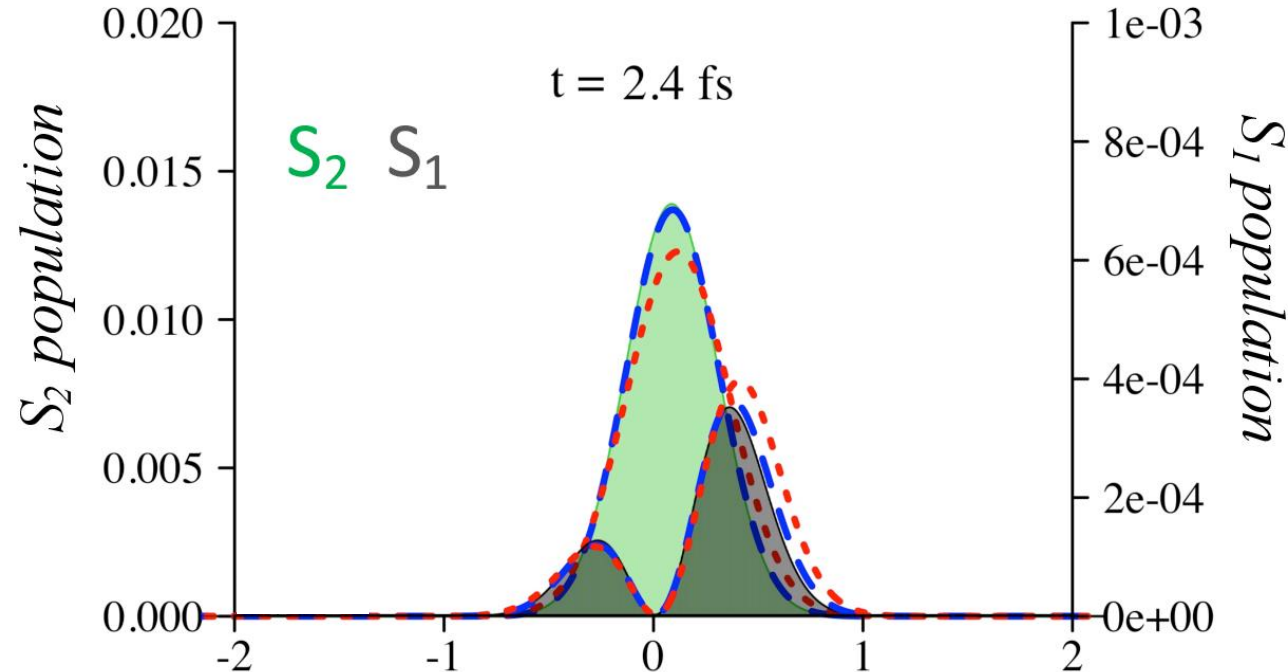
(b)

Real part of the S_1 - S_2 coherence



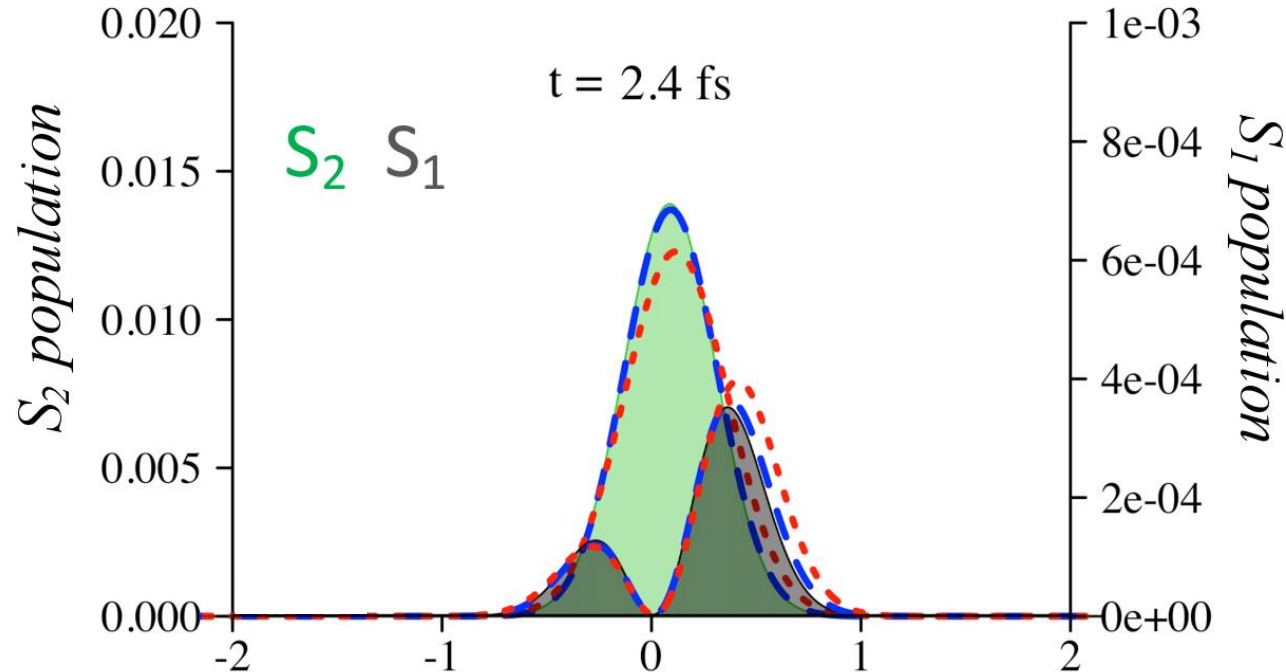
Dominant constraints

	Gelfand basis	Dominant constraints $\hat{I}(t) \approx \sum_k \lambda_k(t) \hat{A}_k$
Small initial shift	400	16 $\left\{ I_N, \hat{a}, \hat{a}^\dagger, \hat{a}^\dagger \hat{a} \right\} \otimes \left\{ I_e, \hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z \right\}$
Large initial shift	1600	24 $\left\{ I_N, \hat{a}, \hat{a}^\dagger, \hat{a}^2, (\hat{a}^\dagger)^2 \hat{a}^\dagger \hat{a} \right\} \otimes \left\{ I_e, \hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z \right\}$



Dominant constraints

	Gelfand basis	Dominant constraints
Small initial shift	400	16 $\{I_N, \hat{a}, \hat{a}^\dagger, \hat{a}^\dagger \hat{a}\} \otimes \{I_e, \hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z\}$
Large initial shift	1600	24 $\{I_N, \hat{a}, \hat{a}^\dagger, \hat{a}^2, (\hat{a}^\dagger)^2 \hat{a}^\dagger \hat{a}\} \otimes \{I_e, \hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z\}$



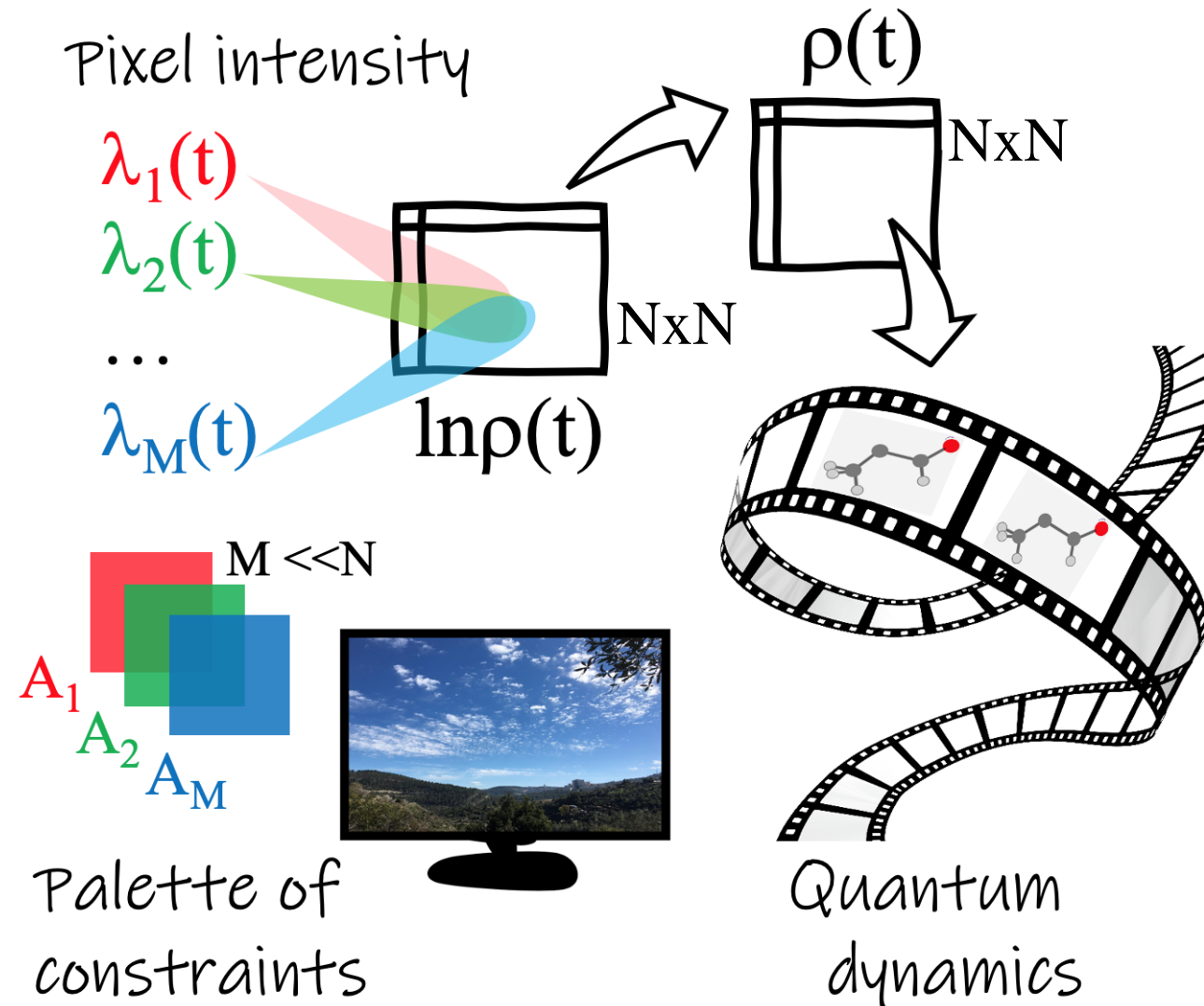
Constraints responsible for the bi-gaussian shape of the S_1 wave packet

$$|1\rangle \hat{a} \langle 2|, |2\rangle \hat{a} \langle 1|$$

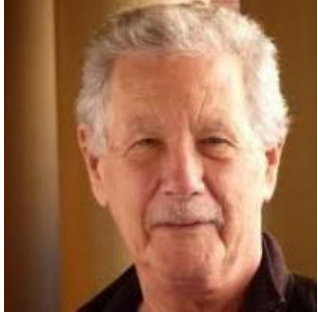
$$|1\rangle \hat{a}^\dagger \langle 2|, |2\rangle \hat{a}^\dagger \langle 1|$$



Surprisal based quantum dynamics



Acknowledgements



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Hebrew University



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Purdue University

COPAC
Coherent Optical PARallel Computing





Thank you for your attention!

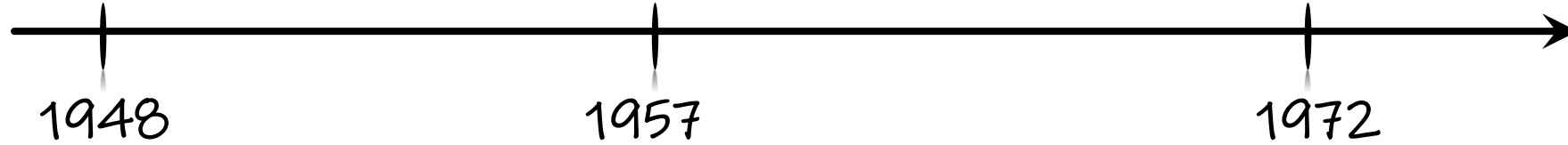
Dominant constraints in the vibrational basis

Set of 16 **time-independent constraints** for the direct product Hilbert space of electronic and nuclear degrees of freedom

	$\hat{\mathbb{I}}_N$	\hat{a}	\hat{a}^\dagger	$\hat{a}^\dagger \hat{a}$
$ 1\rangle\langle 1 $	$ 1\rangle\hat{\mathbb{I}}_N\langle 1 $	$ 1\rangle\hat{a}\langle 1 $	$ 1\rangle\hat{a}^\dagger\langle 1 $	$ 1\rangle\hat{a}^\dagger\hat{a}\langle 1 $
$ 2\rangle\langle 2 $	$ 2\rangle\hat{\mathbb{I}}_N\langle 2 $	$ 2\rangle\hat{a}\langle 2 $	$ 2\rangle\hat{a}^\dagger\langle 2 $	$ 2\rangle\hat{a}^\dagger\hat{a}\langle 2 $
$ 1\rangle\langle 2 $	$ 1\rangle\hat{\mathbb{I}}_N\langle 2 $	$ 1\rangle\hat{a}\langle 2 $	$ 1\rangle\hat{a}^\dagger\langle 2 $	$ 1\rangle\hat{a}^\dagger\hat{a}\langle 2 $
$ 2\rangle\langle 1 $	$ 2\rangle\hat{\mathbb{I}}_N\langle 1 $	$ 2\rangle\hat{a}\langle 1 $	$ 2\rangle\hat{a}^\dagger\langle 1 $	$ 2\rangle\hat{a}^\dagger\hat{a}\langle 1 $

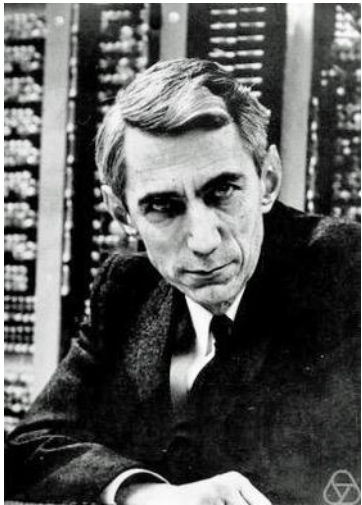
$$\hat{I}(t) \approx \sum_k \lambda_k(t) \hat{A}_k$$

Surprisal analysis and maximal entropy formalism



E.C. Kemble
A. Katz
E.H. Wichman
J.L. Kinsey
And many other..

Information Theory



Claude Shannon

Information Theory
and
**Statistical
Mechanics**



Edwin Jaynes

Information Theory
and
**Chemical
Change**



Richard Bernstein & Raphael Levine

Density and its surprisal

$$\hat{\rho} = \exp\left(-\sum_k \lambda_k \hat{A}_k\right)$$

$$\hat{I} = -\ln \hat{\rho}$$

Surprisal

$$\hat{I} = \sum_k \lambda_k \hat{A}_k$$



$$\hat{\rho}(t) = \hat{U}(t) \hat{\rho}(t=0) \hat{U}^\dagger(t)$$

$$\hat{\rho}^2(t) = \hat{\rho}(t) \cdot \hat{\rho}(t) = \hat{U}(t) \hat{\rho}(t=0) \hat{U}^\dagger(t) \cdot \hat{U}(t) \hat{\rho}(t=0) \hat{U}^\dagger(t) = \hat{U}(t) \hat{\rho}^2(t=0) \hat{U}^\dagger(t)$$

Evolution of the surprisal

$$\hat{I}(t) = -\ln \hat{\rho}(t) = \hat{U}(t) [-\ln \hat{\rho}(t=0)] \hat{U}^\dagger(t)$$

$$i\hbar \frac{d\hat{I}(t)}{dt} = [\hat{H}, \hat{I}(t)]$$