

Problem Set 10 Solutions

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Question 1: Expressing $|2; \ell, m\rangle$ in the $|n_1, n_2, n_3\rangle$ Basis

We have two ways to label the $n = 2$ energy eigenstates of the 3D isotropic harmonic oscillator:

- **Cartesian basis:** $|n_1, n_2, n_3\rangle$, where $n_1 + n_2 + n_3 = 2$. There are exactly 6 such states:

$$|2, 0, 0\rangle, |0, 2, 0\rangle, |0, 0, 2\rangle, |1, 1, 0\rangle, |1, 0, 1\rangle, |0, 1, 1\rangle.$$

- **Spherical basis:** $|n; \ell, m\rangle$ with $n = 2$. Then $n = 2$ can split into (n_r, ℓ) satisfying $2n_r + \ell = 2$. We get:

- $\ell = 0, n_r = 1$ (1 state: $|2; 0, 0\rangle$),
- $\ell = 2, n_r = 0$ (5 states: $|2; 2, m\rangle, m = -2 \cdots +2$).

So total of 6 states again.

Below we show how each $|2; \ell, m\rangle$ is written as a linear combination of $|n_1, n_2, n_3\rangle$.

1. The $\ell = 2$ Multiplet (Five States)

Since $\ell = 2$, m runs from -2 to $+2$. We have 5 states: $m = -2, -1, 0, 1, 2$.

(a) $|2; 2, \pm 2\rangle$. It is given (and can be verified by symmetry arguments) that:

$$|2; 2, \pm 2\rangle = \frac{1}{2} \left[|2, 0, 0\rangle - |0, 2, 0\rangle \pm i\sqrt{2}|1, 1, 0\rangle \right].$$

The ± 2 indicates that we attach the $\pm i$ factor times $\sqrt{2}$. This ensures these states behave like the $m = \pm 2$ spherical harmonics.

(b) $|2; 2, 0\rangle$. One finds

$$|2; 2, 0\rangle = \frac{1}{\sqrt{6}} \left[|2, 0, 0\rangle + |0, 2, 0\rangle - 2 |0, 0, 2\rangle \right].$$

Here we see a combination that is symmetric in (x, y) but subtracts out the z direction in a certain proportion. This matches the $\ell = 2, m = 0$ spherical harmonic in the HO Fock basis.

(c) $|2; 2, \pm 1\rangle$. Often written as linear combinations involving $|1, 0, 1\rangle$ and $|0, 1, 1\rangle$, plus some portion of $|2, 0, 0\rangle$ and $|0, 2, 0\rangle$, with appropriate phases. A typical set (with some sign conventions) is:

$$\begin{aligned} |2; 2, +1\rangle &= \frac{1}{2} \left[|2, 0, 0\rangle + |0, 2, 0\rangle - \sqrt{2} (|1, 0, 1\rangle + i |0, 1, 1\rangle) \right], \\ |2; 2, -1\rangle &= \frac{1}{2} \left[|2, 0, 0\rangle + |0, 2, 0\rangle + \sqrt{2} (|1, 0, 1\rangle - i |0, 1, 1\rangle) \right]. \end{aligned}$$

Exact sign details can vary based on phase conventions.

2. The $\ell = 0$ State (One State)

For $n = 2$ and $\ell = 0$, we must have $n_r = 1$. The single state is $|2; 0, 0\rangle$. By orthogonality with the $\ell = 2$ subspace, we find:

$$|2; 0, 0\rangle = \frac{1}{\sqrt{3}} \left[|1, 1, 0\rangle + |1, 0, 1\rangle + |0, 1, 1\rangle \right].$$

That combination is fully symmetric among the (x, y, z) directions and carries no angular momentum ($\ell = 0$).

Conclusion

These six states $|2; \ell, m\rangle$ in spherical coordinates map to linear combinations of the six $|n_1, n_2, n_3\rangle$ with $n_1 + n_2 + n_3 = 2$. As examples:

- $|2; 2, \pm 2\rangle = \frac{1}{2} (|2, 0, 0\rangle - |0, 2, 0\rangle \pm i\sqrt{2} |1, 1, 0\rangle)$,
- $|2; 0, 0\rangle = \frac{1}{\sqrt{3}} (|1, 1, 0\rangle + |1, 0, 1\rangle + |0, 1, 1\rangle)$,
- $|2; 2, \pm 1\rangle$ have slightly more complicated combos involving $|1, 0, 1\rangle$ and $|0, 1, 1\rangle$ as well.

These superpositions exhaust the full $n = 2$ manifold in both bases.

Question 2: Δx and Δp_x in the Hydrogen Ground State

We consider the hydrogenic ground state, $\psi_{100}(r)$, given by:

$$\psi_{100}(r) = \frac{1}{\sqrt{\pi a_0^3}} e^{-r/a_0}, \quad (1)$$

where a_0 is the Bohr radius. We want Δx and Δp_x .

1. Δx

Since the ground state is spherically symmetric, $\langle x \rangle = 0$, and $\langle x^2 \rangle = \frac{1}{3} \langle r^2 \rangle$. A standard result (or from integrals) is $\langle r^2 \rangle = 3a_0^2$ for $n = 1$ hydrogen. Hence:

$$\begin{aligned} \langle x^2 \rangle &= \frac{1}{3} \times 3a_0^2 = a_0^2, \\ \Delta x &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = a_0. \end{aligned}$$

2. Δp_x

In momentum space, one can compute the Fourier transform of ψ_{100} , or use that $\langle \mathbf{p}^2 \rangle = \hbar^2/a_0^2$ in the hydrogen ground state. Again by symmetry, $\langle p_x^2 \rangle = \frac{1}{3} \langle \mathbf{p}^2 \rangle = \frac{\hbar^2}{3a_0^2}$. Thus:

$$\Delta p_x = \sqrt{\langle p_x^2 \rangle} = \sqrt{\frac{\hbar^2}{3a_0^2}} = \frac{\hbar}{\sqrt{3} a_0}.$$

Final Results

$\Delta x = a_0, \quad \Delta p_x = \frac{\hbar}{\sqrt{3} a_0}.$

(2)

Question 3:

Question 4: