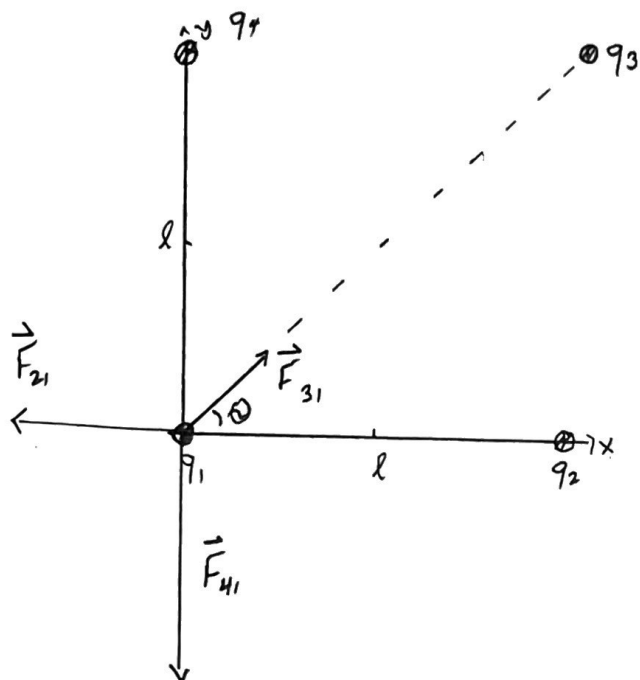


# PHYS 311 HOMEWORK Set 4

1.



$$q_1 = q$$

$$q_2 = 2q$$

$$q_3 = -q$$

$$q_4 = 3q$$

notice:

from trigonometry

$$\tan \theta = \frac{l}{l} = 1 \quad \therefore \theta = 45^\circ$$

$$F_{21} = \frac{1}{4\pi\epsilon_0} \frac{|q_1||q_2|}{l^2} = \frac{1}{4\pi\epsilon_0} \frac{2q^2}{l^2} \quad \text{so} \quad \vec{F}_{21} = -\frac{1}{4\pi\epsilon_0} \frac{2q^2}{l} \hat{x}$$

$$F_{31} = \frac{1}{4\pi\epsilon_0} \frac{|q_1||q_3|}{2l^2} = \frac{1}{4\pi\epsilon_0} \frac{q^2}{2l^2}$$

now

$$F_{31,x} = F_{31} \cos(45^\circ) = \frac{1}{4\pi\epsilon_0} \frac{q^2}{2l^2} \cdot \frac{1}{\sqrt{2}}$$

$$\text{so} \quad \vec{F}_{31} = F_{31,x} \hat{x} + F_{31,y} \hat{y}$$

$$F_{31,y} = F_{31} \sin(45^\circ) = \frac{1}{4\pi\epsilon_0} \frac{q^2}{2l^2} \cdot \frac{1}{\sqrt{2}}$$

$$\vec{F}_{31} = \frac{1}{4\pi\epsilon_0} \frac{q^2}{2l^2} \cdot \frac{1}{\sqrt{2}} (\hat{x} + \hat{y})$$

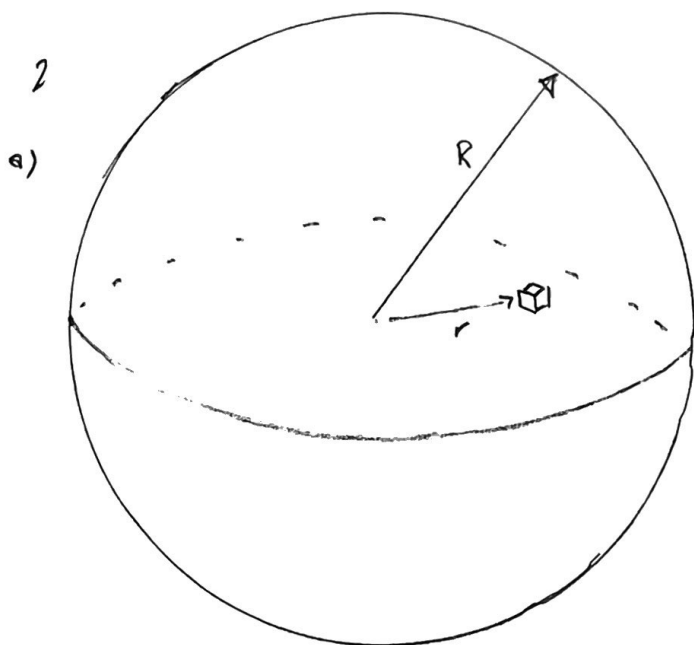
LASTLY...

$$F_{41} = \frac{1}{4\pi\epsilon_0} \frac{|q_1||q_4|}{l^2} = \frac{1}{4\pi\epsilon_0} \frac{3q^2}{l^2}$$

$$\text{so} \quad \vec{F}_{41} = -\frac{1}{4\pi\epsilon_0} \frac{3q^2}{l^2} \hat{y}$$

now

$$\begin{aligned} \vec{F}_1 &= \vec{F}_{21} + \vec{F}_{31} + \vec{F}_{41} = -\frac{1}{4\pi\epsilon_0} \frac{2q^2}{l^2} \hat{x} + \frac{1}{4\pi\epsilon_0} \frac{q^2}{2\sqrt{2}l^2} (\hat{x} + \hat{y}) - \frac{1}{4\pi\epsilon_0} \frac{3q^2}{l^2} \hat{y} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{l^2} \left[ -2\hat{x} + \frac{1}{2\sqrt{2}} (\hat{x} + \hat{y}) - 3\hat{y} \right] \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{l^2} \left[ \left( \frac{1}{2\sqrt{2}} - 2 \right) \hat{x} + \left( \frac{1}{2\sqrt{2}} - 3 \right) \hat{y} \right] \end{aligned}$$



$$\rho(r) = \rho_0 \cos^2 \theta$$

$$dV = r^2 \sin \theta dr d\theta d\phi$$

NOW, THE CHARGE DENSITY IS...

$$\rho = \frac{dq}{dV} \quad \text{so} \quad dq = \rho dV$$

so THE TOTAL CHARGE IS...

$$Q = \int \rho dV = \int_0^R \int_0^\pi \int_0^{2\pi} \rho_0 \cos^2 \theta \cdot r^2 \sin \theta dr d\theta d\phi$$

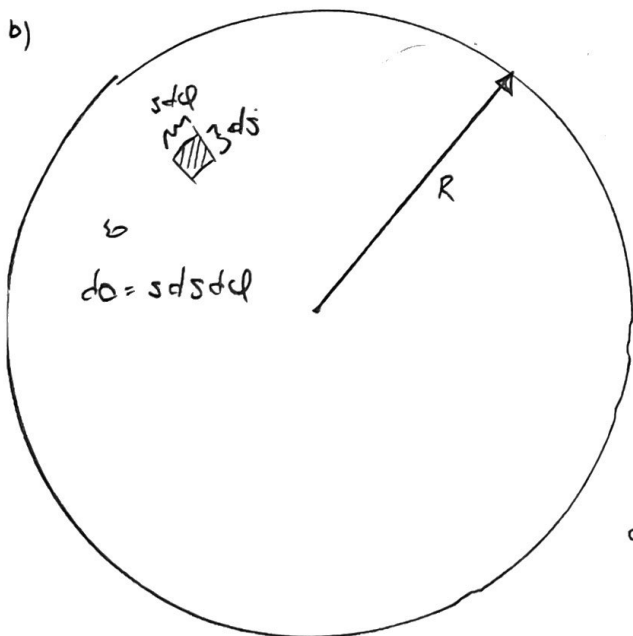
$$Q = \rho_0 \int_0^{2\pi} d\phi \int_0^\pi \cos^2 \theta \sin \theta d\theta \int_0^R r^2 dr = \rho_0 \phi \int_{-1}^1 u^2 du \cdot \frac{r^3}{3} \Big|_0^R = \rho_0 2\pi \frac{u^3}{3} \Big|_{-1}^1 \frac{R^3}{3}$$

$u = \cos \theta$   
 $du = -\sin \theta d\theta$

$$= \rho_0 2\pi \left( \frac{1}{3} - \left( -\frac{1}{3} \right) \right) \frac{R^3}{3} = \int \rho_0 \frac{4\pi R^3}{9} = Q$$



$$dV = dr \cdot r \sin \theta d\theta \cdot r d\phi = r^2 \sin \theta dr d\theta d\phi$$



NOW, THE CHARGE DENSITY IS...

$$\sigma = \frac{dq}{dA} \quad \text{so} \quad dq = \sigma dA$$

so THE TOTAL CHARGE IS...

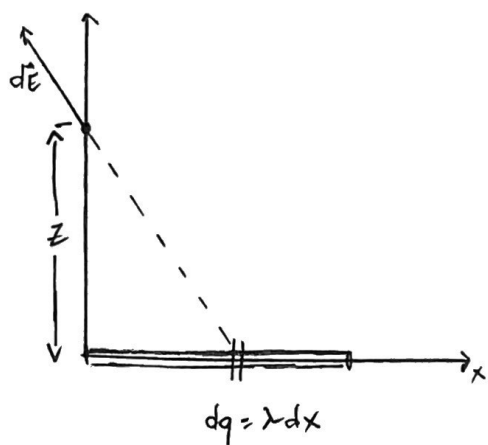
$$Q = \int \sigma dA = \int_0^{2\pi} \int_0^\pi \sigma_0 \frac{R}{s} \cdot s ds d\phi$$

$$Q = \sigma_0 R \int_0^R ds \int_0^{2\pi} d\phi = \sigma_0 R s \Big|_0^R \phi \Big|_0^{2\pi}$$

$$\sigma(s) = \sigma_0 \frac{R}{s}$$

$$[Q = 2\pi \sigma_0 R^2]$$

3



for an infinitesimal chunk of charge...

$$d\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} \hat{n}$$

now

$$\vec{r} = x\hat{x}$$

$$\vec{r} = z\hat{z}$$

$$\text{so } \vec{n} = -x\hat{x} + z\hat{z}$$

$$n^2 = x^2 + z^2 \quad \text{so } n = \sqrt{x^2 + z^2}$$

$$\therefore \hat{n} = \frac{\vec{n}}{n} = \frac{-x\hat{x} + z\hat{z}}{\sqrt{x^2 + z^2}}$$

Putting this all together...

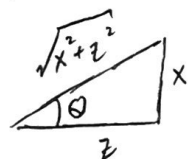
$$d\vec{E} = \frac{1}{4\pi\epsilon_0} \cdot \frac{\lambda dx}{x^2 + z^2} \frac{(-x\hat{x} + z\hat{z})}{\sqrt{x^2 + z^2}} = \frac{\lambda}{4\pi\epsilon_0} \frac{(-x\hat{x} + z\hat{z})}{(x^2 + z^2)^{3/2}} dx$$

So the total  $\vec{E}$ -field is found by integrating over the length  $L$ ..

$$\vec{E} = \int d\vec{E} = \frac{\lambda}{4\pi\epsilon_0} \int_0^L \frac{dx}{(x^2 + z^2)^{3/2}} (-x\hat{x} + z\hat{z}) = \frac{\lambda}{4\pi\epsilon_0} \left[ -\hat{x} \int_0^L \frac{x dx}{(x^2 + z^2)^{3/2}} + z\hat{z} \int_0^L \frac{dx}{(x^2 + z^2)^{3/2}} \right]$$

$$\text{let } u = x^2 + z^2 \\ du = 2x dx$$

$$\text{let } x = z \tan \theta \\ dx = z \sec^2 \theta d\theta$$



$$= \frac{\lambda}{4\pi\epsilon_0} \left[ -\frac{\hat{x}}{2} \int_{z^2}^{L^2 + z^2} \frac{du}{u^{3/2}} + z\hat{z} \int_0^{\tan^{-1}(L/z)} \frac{z \sec^2 \theta d\theta}{z^3 (1 + \tan^2 \theta)^{3/2}} \right]$$

$$= \frac{\lambda}{4\pi\epsilon_0} \left[ -\frac{\hat{x}}{2} \left[ \frac{u^{-1/2}}{-1/2} \right]_{z^2}^{L^2 + z^2} + \frac{\hat{z}}{z} \int_0^{\tan^{-1}(L/z)} \frac{\sec^2 \theta d\theta}{\sec^3 \theta} \right]$$

$$\text{since } \tan^2 \theta + 1 = \sec^2 \theta$$

$$= \frac{\lambda}{4\pi\epsilon_0} \left[ \hat{x} \left[ \frac{1}{\sqrt{L^2 + z^2}} - \frac{1}{z} \right] + \frac{\hat{z}}{z} \int_0^{\tan^{-1}(L/z)} \cos \theta d\theta \right] = \frac{\lambda}{4\pi\epsilon_0} \left[ \left[ \frac{1}{\sqrt{L^2 + z^2}} - \frac{1}{z} \right] \hat{x} + \frac{\hat{z}}{z} \sin \theta \right]_{\theta=0}^{\theta=\tan^{-1}(L/z)}$$

$$\vec{E} = \frac{\lambda}{4\pi\epsilon_0} \left[ \left[ \frac{1}{\sqrt{L^2 + z^2}} - \frac{1}{z} \right] \hat{x} + \frac{L}{\sqrt{L^2 + z^2}} \cdot \frac{\hat{z}}{z} \right] = \frac{\lambda}{4\pi\epsilon_0} \left[ \left[ \frac{1}{z} \left( 1 + \frac{L^2}{z^2} \right)^{-1/2} - \frac{1}{z} \right] \hat{x} + \frac{L}{z^2} \left( 1 + \frac{L^2}{z^2} \right)^{-1/2} \hat{z} \right]$$

now

$$\lim_{z \gg L} \vec{E} = \frac{\lambda}{4\pi\epsilon_0} \left[ \frac{1}{z} \left( 1 - \frac{1}{2} \frac{L^2}{z^2} + O\left(\frac{L^4}{z^4}\right) \right) \hat{x} + \frac{L}{z^2} \left( 1 - \frac{1}{2} \frac{L^2}{z^2} + O\left(\frac{L^4}{z^4}\right) \right) \hat{z} \right] \quad \text{since } (1+x)^n \approx 1 + nx \text{ for small } x$$

$$\text{so } \lim_{z \gg L} \vec{E} \approx \frac{\lambda}{4\pi\epsilon_0} \frac{L}{z^2} \hat{z} = \frac{1}{4\pi\epsilon_0} \frac{q}{z^2} \hat{z} \quad \text{A POINT CHARGE!}$$

4. THIS PROBLEM IS VERY SIMILAR TO THE PREVIOUS ONE, WITH ONLY THE LINEAR CHARGE DENSITY DIFFERING FROM A CONSTANT.

FROM PROBLEM 3

$$d\vec{E} = \frac{\lambda}{4\pi\epsilon_0} \frac{(-x\hat{x} + z\hat{z})}{(x^2 + z^2)^{3/2}} dx \quad \text{WHERE, HOWEVER, } \lambda(x) = \lambda_0 \frac{x^2}{L}$$

SO, THE TOTAL  $\vec{E}$ -FIELD IS OF THE FORM..

$$\vec{E} = \frac{\lambda_0}{4\pi\epsilon_0} \int_0^L \frac{x^2}{L^2} \frac{(-x\hat{x} + z\hat{z})}{(x^2 + z^2)^{3/2}} dx$$

$$= \frac{\lambda_0}{4\pi\epsilon_0 L^2} \left[ -\hat{x} \int_0^L \frac{x^3}{(x^2 + z^2)^{3/2}} dx + z\hat{z} \int_0^L \frac{x^2}{(x^2 + z^2)^{3/2}} dx \right]$$

$$\text{NOW } \cosh^2 \theta - \sinh^2 \theta = 1$$

$$\text{NOW } \int_0^L \frac{x^2 dx}{(x^2 + z^2)^{3/2}} = \int_0^L x \cdot \frac{x dx}{(x^2 + z^2)^{3/2}} = \left. \frac{-x}{\sqrt{x^2 + z^2}} \right|_0^L + \int_0^L \frac{dx}{\sqrt{x^2 + z^2}} = \frac{-L}{\sqrt{L^2 + z^2}} + \int_0^{\sinh^{-1}(L/z)} \frac{z \cosh \theta d\theta}{z \sqrt{1 + \sinh^2 \theta}}$$

let  $x = z \sinh \theta$   
 $dx = z \cosh \theta d\theta$

$$du = dx \quad v = \int \frac{x dx}{(x^2 + z^2)^{3/2}} = \frac{1}{2} \int \frac{dw}{w^{3/2}} = \frac{1}{2} \frac{w^{-1/2}}{-1/2} = -\frac{1}{\sqrt{x^2 + z^2}}$$

let  $w = x^2 + z^2$   
 $dw = 2x dx$

$$\text{SO } \int_0^L \frac{x^2 dx}{(x^2 + z^2)^{3/2}} = \frac{-L}{\sqrt{L^2 + z^2}} + \int_0^{\sinh^{-1}(L/z)} d\theta$$

$$= \frac{-L}{\sqrt{L^2 + z^2}} + \sinh^{-1}\left(\frac{L}{z}\right) \checkmark$$

$$\text{WHEREAS } \int_0^L \frac{x^3 dx}{(x^2 + z^2)^{3/2}} = \int_0^L x \cdot \frac{x^2 dx}{(x^2 + z^2)^{3/2}} = x \left( \frac{-x}{\sqrt{x^2 + z^2}} + \sinh^{-1}\left(\frac{x}{z}\right) \right) \Big|_0^L - \int_0^L \left( \frac{-x}{\sqrt{x^2 + z^2}} + \sinh^{-1}\left(\frac{x}{z}\right) \right) dx$$

let  $u = x$   
 $du = dx$

let  $v = \frac{-x}{\sqrt{x^2 + z^2}} + \sinh^{-1}\left(\frac{x}{z}\right)$

so

$$\int_0^L \frac{x^3 dx}{(x^2+z^2)^{3/2}} = \frac{-L^2}{\sqrt{L^2+z^2}} + L \sinh^{-1}\left(\frac{L}{z}\right) + \int_0^L \frac{x}{\sqrt{x^2+z^2}} dx - \int_0^L \sinh^{-1}\left(\frac{x}{z}\right) dx$$

$u = x^2 + z^2$   
 $du = 2x dx$   
 $\frac{du}{u^{1/2}}$

$w = \frac{x}{z}$   
 $dw = \frac{dx}{z}$

$$= \frac{-L^2}{\sqrt{L^2+z^2}} + L \sinh^{-1}\left(\frac{L}{z}\right) + \frac{1}{2} \int_{z^2}^{L^2+z^2} \frac{du}{u^{1/2}} - z \int_0^{L/z} \sinh^{-1}(w) dw$$

now

$$\int \sinh^{-1}(w) dw = w \sinh^{-1}(w) - \sqrt{w^2+1} \quad (\text{from wolframalpha})$$

so

$$\int_0^L \frac{x^3 dx}{(x^2+z^2)^{3/2}} = \frac{-L^2}{\sqrt{L^2+z^2}} + L \sinh^{-1}\left(\frac{L}{z}\right) + \frac{1}{2} \frac{u^{1/2}}{1/2} \Big|_{z^2}^{L^2+z^2} - z \left[ w \sinh^{-1}(w) - \sqrt{w^2+1} \right]_0^{L/z}$$

$$= \frac{-L^2}{\sqrt{L^2+z^2}} + L \sinh^{-1}\left(\frac{L}{z}\right) + \left[ \sqrt{L^2+z^2} - z \right] - z \left[ \frac{L}{z} \sinh^{-1}\left(\frac{L}{z}\right) - \sqrt{\frac{L^2}{z^2}+1} \right]$$

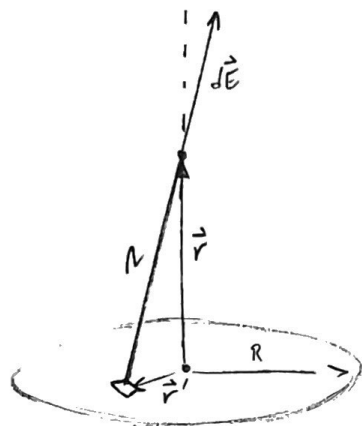
$$= \frac{-L^2}{\sqrt{L^2+z^2}} + \sqrt{L^2+z^2} - z - \sqrt{L^2+z^2}$$

$$= \frac{-L^2}{\sqrt{L^2+z^2}} - z$$

so

$$\vec{E} = \frac{\lambda_0}{4\pi\epsilon_0} \frac{1}{z^2} \left[ \left[ \frac{L^2}{\sqrt{L^2+z^2}} + z \right] \hat{x} + z \left[ -\frac{L}{\sqrt{L^2+z^2}} + \sinh^{-1}\left(\frac{L}{z}\right) \right] \hat{z} \right]$$

5



FOR AN INFINITESIMAL AREA OF CHARGE...

$$d\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} \hat{r}$$

NOW

$$\vec{r}' = s\hat{s} \quad \text{so} \quad \vec{r} = -s\hat{s} + z\hat{z}$$

$$\vec{r} = z\hat{z}$$

$$r^2 = s^2 + z^2 \quad \text{so} \quad r = \sqrt{s^2 + z^2}$$

$$\text{so} \quad \hat{r} = \frac{\vec{r}}{r} = \frac{-s\hat{s} + z\hat{z}}{\sqrt{s^2 + z^2}}$$

thus

$$dq = \sigma da = \sigma s ds d\phi$$

NOW PUTTING THIS ALL TOGETHER...

$$d\vec{E} = \frac{1}{4\pi\epsilon_0} \cdot \frac{\sigma s ds d\phi}{s^2 + z^2} \cdot \frac{(-s\hat{s} + z\hat{z})}{\sqrt{s^2 + z^2}} = \frac{\sigma}{4\pi\epsilon_0} \frac{s(-s\hat{s} + z\hat{z}) ds d\phi}{(s^2 + z^2)^{3/2}}$$

NOW THE TOTAL  $\vec{E}$ -FIELD IS FOUND BY INTEGRATING OVER THE DISK...

$$\vec{E} = \frac{\sigma}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^R \frac{s(-s\hat{s} + z\hat{z}) ds d\phi}{(s^2 + z^2)^{3/2}} = \frac{\sigma}{4\pi\epsilon_0} \left[ -\hat{s} \int_0^{2\pi} d\phi \int_0^R \frac{s^2 ds}{(s^2 + z^2)^{3/2}} + z\hat{z} \int_0^{2\pi} d\phi \int_0^R \frac{s ds}{(s^2 + z^2)^{3/2}} \right]$$

NOTICE!

BECAUSE OF THE CYLINDRICAL-SYMMETRY OF THE PROBLEM, THE  $\vec{E}$ -FIELD MUST POINT ONLY IN THE  $z$ -DIRECTION.

$\therefore$  THE  $\hat{s}$  INTEGRAL MUST BE ZERO, I'LL SHOW THIS NOW...

$$\int_0^R \frac{s^2 ds}{(s^2 + z^2)^{3/2}} = \int_0^R s \cdot \frac{s ds}{(s^2 + z^2)^{3/2}} = -\frac{s}{\sqrt{s^2 + z^2}} \Big|_0^R + \int_0^R \frac{ds}{\sqrt{s^2 + z^2}}$$

$$\text{let } u = s^2 \quad du = 2s ds$$

$$v = \int \frac{s ds}{(s^2 + z^2)^{3/2}} = \frac{1}{2} \int u^{-3/2} du = \frac{1}{2} \frac{u^{-1/2}}{-1/2} = -\frac{1}{\sqrt{s^2 + z^2}}$$

$$\text{let } w = s^2 + z^2 \quad dw = 2s ds$$

$$\begin{aligned} \int_0^R \frac{s^2 ds}{(s^2 + z^2)^{3/2}} &= -\frac{R}{\sqrt{R^2 + z^2}} + \int_0^R \frac{ds}{\sqrt{s^2 + z^2}} = -\frac{R}{\sqrt{R^2 + z^2}} + \int_0^{\sinh^{-1}(R/z)} \frac{z \cosh \theta d\theta}{z \sqrt{1 + \sinh^2 \theta}} = -\frac{R}{\sqrt{R^2 + z^2}} + \int_0^{\sinh^{-1}(R/z)} d\theta \\ &\text{let } s = z \sinh \theta \quad \text{so } \theta = \sinh^{-1}\left(\frac{s}{z}\right) \\ &ds = z \cosh \theta d\theta \end{aligned}$$

$$\int_0^R \frac{s^2 ds}{(s^2 + z^2)^{3/2}} = \frac{-R}{\sqrt{R^2 + z^2}} + \sinh^{-1}\left(\frac{R}{z}\right)$$

so

$$\vec{E} = \frac{Q}{4\pi\epsilon_0} \hat{z} \int_0^\pi d\phi \int_0^R \frac{s ds}{(s^2 + z^2)^{3/2}} = \frac{Q}{4\pi\epsilon_0} \hat{z} \cdot 2\pi \cdot \frac{1}{2} \int_{z^2}^{R^2+z^2} \frac{du}{u^{3/2}} = \frac{Q}{4\epsilon_0} \hat{z} \left[ \frac{u^{-1/2}}{-1/2} \right]_{z^2}^{R^2+z^2}$$

let  $u = s^2 + z^2$   
 $du = 2s ds$

$$= -\frac{Q}{2\epsilon_0} \hat{z} \left[ \frac{1}{\sqrt{R^2+z^2}} - \frac{1}{z} \right] = \frac{Q}{2\epsilon_0} \left[ 1 - \frac{z}{\sqrt{R^2+z^2}} \right] \hat{z}$$

~~~~~

Now

lim  $\vec{E} = \frac{Q}{2\epsilon_0} \hat{z}$  : the  $\vec{E}$  has a finite value!  
 $R \rightarrow \infty$



6. THIS PROBLEM IS VERY SIMILAR TO THE PREVIOUS ONE, WITH ONLY THE SURFACE CHARGE DENSITY DIFFERING FROM THAT OF A CONSTANT.

FROM PROBLEM 5...

$$d\vec{E} = \frac{\sigma}{4\pi\epsilon_0} \frac{s(-s\hat{s} + z\hat{z})}{(s^2 + z^2)^{3/2}} ds d\phi \quad \text{WHERE } \sigma(s) = \frac{\sigma_0 s}{R}$$

$$\text{so}$$

$$d\vec{E} = \frac{\sigma_0}{4\pi\epsilon_0 R} \frac{s^2(-s\hat{s} + z\hat{z})}{(s^2 + z^2)^{3/2}} ds d\phi$$

NOW THE TOTAL  $\vec{E}$ -FIELD IS FOUND BY INTEGRATING OVER THE DISK..

$$\vec{E} = \frac{\sigma_0}{4\pi\epsilon_0 R} \int_0^{2\pi} \int_0^R \frac{s^2(-s\hat{s} + z\hat{z})}{(s^2 + z^2)^{3/2}} ds d\phi = \frac{\sigma_0}{4\pi\epsilon_0 R} \left[ -\hat{s} \int_0^{2\pi} d\phi \int_0^R \frac{s^3 ds}{(s^2 + z^2)^{3/2}} + z\hat{z} \int_0^{2\pi} d\phi \int_0^R \frac{s^2 ds}{(s^2 + z^2)^{3/2}} \right]$$

NOW, IN THE LAST PROBLEM I SHOWED THAT...

$$\int_0^R \frac{s^2 ds}{(s^2 + z^2)^{3/2}} = \frac{-R}{\sqrt{R^2 + z^2}} + \sinh^{-1}(R/z)$$

so

$$\vec{E} = \frac{\sigma_0}{4\pi\epsilon_0 R} \cdot z\hat{z} \cdot 2\pi \cdot \left[ -\frac{R}{\sqrt{R^2 + z^2}} + \sinh^{-1}(R/z) \right]$$

$$= \frac{\sigma_0}{2\epsilon_0} \cdot \frac{z}{R} \left[ -\frac{R}{\sqrt{R^2 + z^2}} + \sinh^{-1}(R/z) \right] \hat{z}$$

so

$$\vec{E} = \frac{\sigma_0}{2\epsilon_0} \left[ -\frac{z}{\sqrt{R^2 + z^2}} + \frac{z}{R} \sinh^{-1}(R/z) \right] \hat{z}$$

$$7. \vec{E} = kr^5 \hat{r} = E_r \hat{r}$$

now

$$a) \vec{\nabla} \cdot \vec{E} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \cdot kr^5) = \frac{k}{r^2} \frac{\partial}{\partial r} (r^7) = \frac{k}{r^2} \cdot 7r^6 = 7kr^4$$

BUT WE ALSO KNOW THAT

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} \quad \text{so} \quad \rho = \epsilon_0 (\vec{\nabla} \cdot \vec{E}) \quad \text{so} \quad \left[ \rho(r) = 7k\epsilon_0 r^4 \right]$$

b) now

$$\rho = \frac{dq}{d\tau} \quad \text{so} \quad dq = \rho d\tau \quad \therefore \quad q = \int \rho d\tau$$

$$\begin{aligned} q &= \int_0^R \int_0^{2\pi} \int_0^\pi 7k\epsilon_0 r^4 \cdot r^2 \sin\theta dr d\theta d\phi = 7k\epsilon_0 \int_0^{2\pi} d\phi \int_0^\pi \sin\theta d\theta \int_0^R r^6 dr \\ &= 7k\epsilon_0 \cdot 2\pi \cdot \cos\theta \Big|_0^\pi \cdot \frac{r^7}{7} \Big|_0^R = 7k\epsilon_0 \cdot 2\pi \cdot 2 \cdot \frac{R^7}{7} = 4\pi\epsilon_0 k R^7 \end{aligned}$$

WE COULD ALSO USE GAUSS' LAW...

$$\oint \vec{E} \cdot d\vec{\sigma} = \frac{Q_{\text{enc}}}{\epsilon_0} \quad \text{so} \quad Q_{\text{enc}} = \epsilon_0 \oint \vec{E} \cdot d\vec{\sigma} = \epsilon_0 \int_0^{2\pi} \int_0^\pi kR^7 \sin\theta d\theta d\phi = \epsilon_0 kR^7 \int_0^{2\pi} d\phi \int_0^\pi \sin\theta d\theta$$

now

$$\begin{aligned} \vec{E} &= kr^5 \hat{r} \Big|_{r=R} \\ d\vec{\sigma} &= R^2 \sin\theta d\theta d\phi \hat{r} \end{aligned} \quad \text{so} \quad \vec{E} \cdot d\vec{\sigma} = kR^7 \sin\theta d\theta d\phi$$

so

$$Q_{\text{enc}} = \epsilon_0 k R^7 \cdot 2\pi \cdot 2 \quad \therefore \quad \left[ Q_{\text{enc}} = 4\pi\epsilon_0 k R^7 \right]$$



NOTICE THAT WE MADE OUR GAUSSIAN SURFACE TO BE A SPHERE OF RADIUS R,

THE EXACT SAME SIZE AS OUR SOURCE OF CHARGE.