

Phs 396 Homework Set 3 Solutions

$$1. a) dt' = \gamma(dt - \frac{v}{c^2} dx)$$

$$dx' = \gamma(dx - v dt)$$

$$dy' = dy$$

$$dz' = dz$$

Now, plugging these into the line element yields...

$$\begin{aligned} ds'^2 &= -c^2 dt'^2 + dx'^2 + dy'^2 + dz'^2 \\ &= -c^2 \gamma^2 (dt - \frac{v}{c^2} dx)^2 + \gamma^2 (dx - v dt)^2 + dy^2 + dz^2 \\ &= -c^2 \gamma^2 (dt^2 - 2 \frac{v}{c^2} dx dt + \frac{v^2}{c^4} dx^2) \\ &\quad + \gamma^2 (dx^2 - 2v dx dt + v^2 dt^2) + dy^2 + dz^2 \end{aligned}$$

(Now grouping like terms...

$$\begin{aligned} ds'^2 &= \gamma^2 (-c^2 + v^2) dt^2 + \gamma^2 (1 - \frac{v^2}{c^2}) dx^2 + dy^2 + dz^2 \\ &= -c^2 \gamma^2 (1 - \frac{v^2}{c^2}) dt^2 + \gamma^2 (1 - \frac{v^2}{c^2}) dx^2 + dy^2 + dz^2 \end{aligned}$$

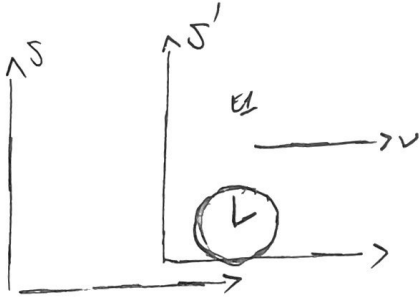
Now

$$\gamma^2 (1 - \frac{v^2}{c^2}) = \frac{1}{(1 - v^2/c^2)} \cdot (1 - v^2/c^2) = 1$$

so

$$ds'^2 = -c^2 dt^2 + dx^2 + dy^2 + dz^2 = ds^2 \checkmark$$

b) Let S' be the rest frame of the clock, that moves w/ speed v relative to the S frame...



Let E_1 be the event describing coordinates of the clock at t_1, t'_1 ,
 Let E_2 be the event describing coordinates of the clock at t_2, t'_2

notice that $\Delta x' = 0 = x'_2 - x'_1$

now

$$\Delta t = \gamma \left(\Delta t' + \frac{v}{c^2} \Delta x' \right)$$

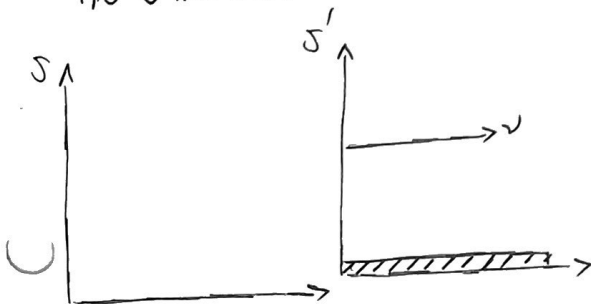
$$\Delta x = \gamma (\Delta x' + v \Delta t')$$

the first eqn yields

$$\Delta t = \gamma \Delta t' \quad \text{so} \quad \left[\Delta t' = \frac{1}{\gamma} \Delta t = \sqrt{1 - v^2/c^2} \Delta t \right]$$

so $\Delta t' < \Delta t$ - the 'moving' clock measures a shorter time interval

∴ Let S' be the rest frame of an object of length L_0 , which moves w/ speed v relative to the S frame...



to measure the length, the S observer notes the left and right ends at the same moment in time, relative to the S frame.

$$\text{so } \Delta t = 0$$

$$\Delta t' = \gamma \left(\Delta t - \frac{v}{c^2} \Delta x \right)$$

$$\Delta x' = \gamma (\Delta x - v \Delta t)$$

THE SECOND EQN YIELDS...

$$\Delta x' = \gamma \Delta x$$

IF THE 2 EVENTS CORRESPOND TO THE LOCATIONS OF THE LEFT AND RIGHT ENDS OF THE OBJECT, THEN

$$\Delta x' \equiv L_* \quad \text{— THE REST LENGTH OF OBJECT}$$

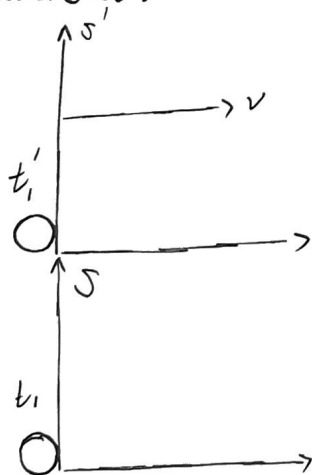
$$\Delta x \equiv L \quad \text{— THE LENGTH OF THE OBJECT, AS MEASURED BY THE MOVING FRAME SO}$$

$$L_* = \gamma L = \frac{L}{\sqrt{1-v^2/c^2}} \quad \therefore \left[L = \sqrt{1-v^2/c^2} L_* \right]$$

NOTICE: $L < L_*$ THE 'MOVING' FRAME MEASURES A SHORTER LENGTH

\therefore MOVING OBJECTS ARE CONTRACTED

4) LET S' BE MOVING W/ SPEED v RELATIVE TO THE S FRAME. EACH IRF HAS A SET OF CLOCKS AT DIFFERENT LOCATIONS, BOTH AT REST RELATIVE TO ONE ANOTHER, AND BOTH SYNCHRONIZED IN THEIR OWN FRAME OF REFERENCE.



AT THE ORIGIN OF EACH COORDINATE SYSTEM LIES A CLOCK. BOTH OBSERVERS SET THEIR ORIGIN CLOCKS EQUAL TO ONE ANOTHER THE MOMENT THEY ALIGN SO

$$t_1 = t_1' = 0$$

LET A SECOND CLOCK IN THE S' FRAME BE TRAILING THE ORIGIN CLOCK BY A REST DISTANCE D (DISTANCE MEASURED IN S' FRAME)

THE S FRAME HAS A CLOCK AT EACH SPATIAL POSITION AND ONE OF THEM, CLOCK 2, ALIGNS W/ CLOCK 2' WHEN CLOCK 1 : 1' ALIGN.

NOW, ACCORDING TO THE S FRAME...

CLOCK 1 ALIGNS W/ 1' AT PRECISELY THE SAME TIME

AS CLOCK 2 ALIGNS W/ 2' \therefore THESE EVENTS ARE SIMULTANEOUS SO

$$\Delta t = t_2 - t_1 = 0$$

$$\Delta t' = \gamma(\Delta t - \frac{v}{c} \Delta x)$$

$$\therefore \Delta t' = -\frac{v}{c} \Delta x$$

$$\Delta x' = \gamma(\Delta x - v \Delta t)$$

$$\Delta x' = \Delta x$$

Plugging the 2nd into the 1st yields..

$$\Delta t' = -\frac{v}{c} \Delta x' = +\frac{v}{c} D \quad \text{so} \quad t_2' - t_1' = t_2' - t_1' = 0$$

$$\Delta x' = x_2' - x_1' = -D$$

since $x_2' < x_1'$

$$t_2' = +\frac{v}{c} D$$

THE POSITIVE SIGN INDICATES THAT CLOCK 2' READS A LATER TIME THAN CLOCK 1' & LEADING CLOCKS LAG BY

$$\left[\Delta t' = \frac{v}{c} D \right] \checkmark$$

2. According to the S frame...

$$\nabla^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \vec{E} = 0 \quad \text{where } \vec{E} = \vec{E}(x, y, z, t)$$

According to the S' frame, which moves at a speed v relative to the S frame, this observer

measures an \vec{E} -field where $\vec{E} = \vec{E}(x', y', z', t')$

where

$$x' = \gamma(x, t)$$

$$t' = \gamma(t - \frac{v}{c^2}x) \quad \text{via the Lorentz boosts which are the form}$$

$$t' = \gamma(t - \frac{v}{c^2}x)$$

$$x' = \gamma(x - vt)$$

We need to transform the derivatives by using the chain rule...

$$\frac{\partial}{\partial t} \vec{E}(x', y', z', t') = \frac{\partial \vec{E}}{\partial x'} \frac{\partial x'}{\partial t} + \frac{\partial \vec{E}}{\partial t'} \frac{\partial t'}{\partial t}$$

$$\text{but } \frac{\partial x'}{\partial t} = -\gamma v \quad ; \quad \frac{\partial t'}{\partial t} = \gamma$$

so

$$\frac{\partial \vec{E}}{\partial t} = -\gamma v \frac{\partial \vec{E}}{\partial x'} + \gamma \frac{\partial \vec{E}}{\partial t'} = \gamma \left(\frac{\partial}{\partial t'} - v \frac{\partial}{\partial x'} \right) \vec{E}$$

A second time-derivative takes the form...

$$\frac{\partial^2 \vec{E}}{\partial t^2} = \frac{\partial}{\partial t} \left(\frac{\partial \vec{E}}{\partial t} \right) = \gamma^2 \left(\frac{\partial}{\partial t'} - v \frac{\partial}{\partial x'} \right) \left(\frac{\partial}{\partial t'} - v \frac{\partial}{\partial x'} \right) \vec{E} = \gamma^2 \left(\frac{\partial^2}{\partial t'^2} - 2v \frac{\partial^2}{\partial x' \partial t'} + v^2 \frac{\partial^2}{\partial x'^2} \right) \vec{E}$$

Now transform the spatial derivatives in an analogous way...

$$\frac{\partial \vec{E}(x', y', z', t')}{\partial x} = \frac{\partial \vec{E}}{\partial x'} \frac{\partial x'}{\partial x} + \frac{\partial \vec{E}}{\partial t'} \frac{\partial t'}{\partial x}$$

$$\text{but } \frac{\partial x'}{\partial x} = \gamma \quad ; \quad \frac{\partial t'}{\partial x} = -\gamma \frac{v}{c^2}$$

$$\frac{\partial \vec{E}}{\partial x} = \gamma \left(\frac{\partial}{\partial x'} - \frac{v}{c^2} \frac{\partial}{\partial t'} \right) \vec{E}$$

A SECOND X-DERIVATIVE TAKES THE FORM...

$$\frac{\partial^2}{\partial x'^2} \vec{E} = \frac{\partial}{\partial x'} \left(\frac{\partial \vec{E}}{\partial x'} \right) = \gamma^2 \left(\frac{\partial}{\partial x'} - \frac{v}{c} \frac{\partial}{\partial t'} \right) \left(\frac{\partial}{\partial x'} - \frac{v}{c} \frac{\partial}{\partial t'} \right) \vec{E}$$

$$= \gamma^2 \left(\frac{\partial^2}{\partial x'^2} - 2 \frac{v}{c} \frac{\partial^2}{\partial x' \partial t'} + \frac{v^2}{c^2} \frac{\partial^2}{\partial t'^2} \right) \vec{E}$$

NOW

$$\frac{\partial}{\partial y} \vec{E} = \frac{\partial \vec{E}}{\partial y'} \frac{\partial y'}{\partial y} = \frac{\partial \vec{E}}{\partial y'} \quad \text{AND} \quad \frac{\partial^2}{\partial y^2} \vec{E} = \frac{\partial^2}{\partial y'^2} \vec{E}$$

LIKewise

$$\frac{\partial}{\partial z} \vec{E} = \frac{\partial \vec{E}}{\partial z'} \frac{\partial z'}{\partial z} = \frac{\partial \vec{E}}{\partial z'} \quad \text{AND} \quad \frac{\partial^2}{\partial z^2} \vec{E} = \frac{\partial^2}{\partial z'^2} \vec{E}$$

NOW, PUTTING IT ALL TOGETHER...

$$\left[\frac{\partial^2}{\partial x'^2} + \frac{\partial^2}{\partial y'^2} + \frac{\partial^2}{\partial z'^2} - \frac{1}{c^2} \frac{\partial^2}{\partial t'^2} \right] \vec{E} = 0$$

BECOMES

$$\gamma^2 \left(\frac{\partial^2}{\partial x'^2} - 2 \frac{v}{c} \frac{\partial^2}{\partial x' \partial t'} + \frac{v^2}{c^2} \frac{\partial^2}{\partial t'^2} \right) \vec{E} + \frac{\partial^2}{\partial y'^2} \vec{E} + \frac{\partial^2}{\partial z'^2} \vec{E} - \frac{1}{c^2} \gamma^2 \left(\frac{\partial^2}{\partial t'^2} - 2v \frac{\partial^2}{\partial x' \partial t'} + v^2 \frac{\partial^2}{\partial x'^2} \right) \vec{E}$$

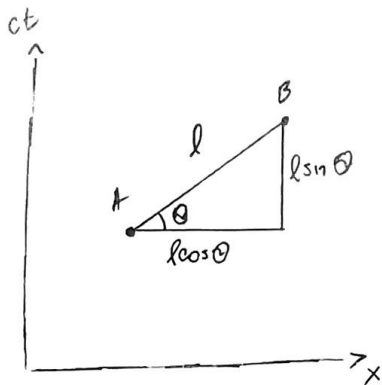
NOW GROUPING LIKE TERMS...

$$\gamma^2 \left(1 - \frac{v^2}{c^2} \right) \frac{\partial^2}{\partial x'^2} \vec{E} - \gamma^2 \left(-2 \frac{v}{c} + \frac{2v}{c} \right) \frac{\partial^2}{\partial x' \partial t'} \vec{E} + \gamma^2 \left(\frac{v^2}{c^2} - \frac{1}{c^2} \right) \frac{\partial^2}{\partial t'^2} \vec{E} + \frac{\partial^2}{\partial y'^2} \vec{E} + \frac{\partial^2}{\partial z'^2} \vec{E} = 0$$

$$\frac{1}{\text{SINCE}} \quad \gamma^2 \left(1 - \frac{v^2}{c^2} \right) = \frac{1}{1 - v^2/c^2} \left(1 - v^2/c^2 \right) = 1$$

BECOMES

$$\frac{\partial^2}{\partial x'^2} \vec{E} + \frac{\partial^2}{\partial y'^2} \vec{E} + \frac{\partial^2}{\partial z'^2} \vec{E} - \frac{1}{c^2} \frac{\partial^2}{\partial t'^2} \vec{E} = \nabla'^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2}{\partial t'^2} \vec{E} = 0 \quad \checkmark$$



$$\begin{aligned}
 b) \Delta S^2 &= -(\Delta ct)^2 + \Delta x^2 \\
 &= -(l^2 \sin^2 \theta) + l^2 \cos^2 \theta \\
 \Delta S^2 &= l^2 (\cos^2 \theta - \sin^2 \theta)
 \end{aligned}$$

c) setting $l=1$,

$$\Delta S^2 = \cos^2 \theta - \sin^2 \theta$$

$$d) -1 \leq \Delta S^2 \leq 1$$

SMALLEST MAGNITUDE : $\theta = \pi/4$ or 45° where $\Delta S^2 = 0$

MAXIMUM VALUE : $\theta = 0$ where $|\Delta S| = 1$

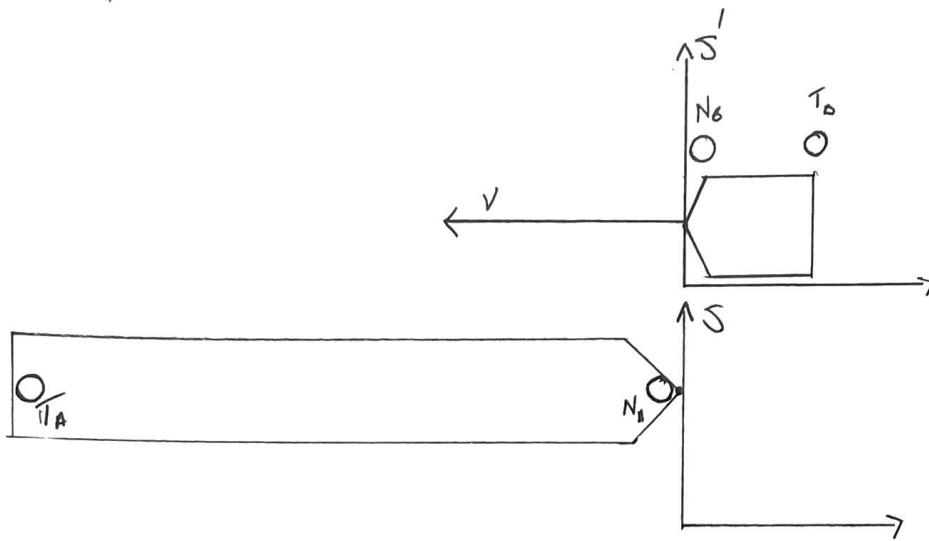
MAXIMUM VALUE : $\theta = \pi/2$ or 90° where $|\Delta S| = 1$

Scale: 100m : 1.5cm

4. $v = -\frac{12}{13} c$

$L_{AB} = 598m$, $L_{AB} = 299m$

a)



As we analyze this problem from the S frame, and as both S frame clocks are synchronized,

both

$$t_{N_A} = t_{T_A} = 0$$

The nose B clock sets is set to zero when passing the nose A clock,

$$t'_{N_B} = 0$$

Now, although the N_B & T_B clocks are synchronized in the S' frame of reference. To get the time on the T_B clock, takes its reading as the 2nd event so that

$$t_2 = t_1, \quad x_1' = 0, \quad x_2' = L_{AB}, \quad v = -\frac{12}{13} c$$

so

$$t_2' = \gamma \left(t_2 - \frac{v}{c^2} x_2 \right)$$

$$x_2' = \gamma (x_2 - v t_2)$$

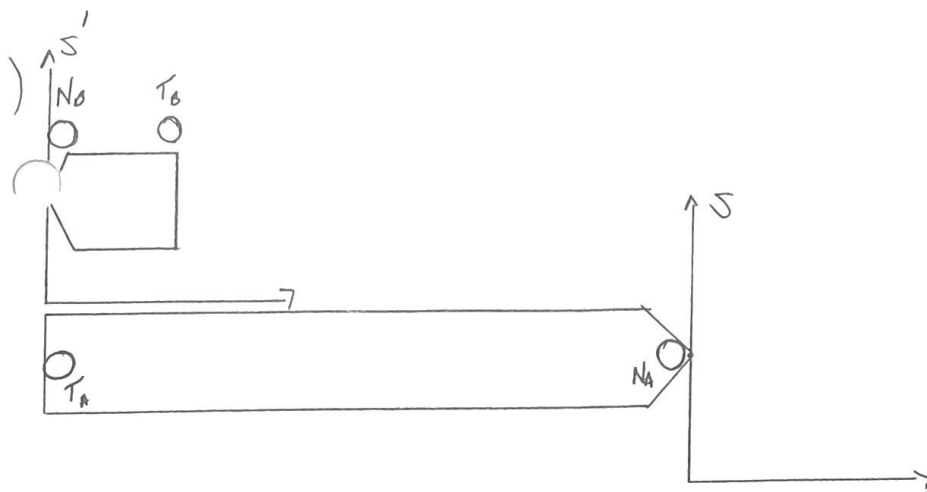
\therefore

$$t_2' = -\gamma \frac{v}{c^2} x_2 \quad \therefore \quad t_2' = -\frac{v}{c^2} x_2' = -\frac{v}{c^2} L_{AB}$$

plugging our values in yield..

$$t_2' = -\left(-\frac{12}{13} c\right) (299m) = \frac{276m}{c} \times \frac{1c}{3.0 \times 10^8 m/s} = 9.2 \times 10^{-7} s$$

$$\therefore \left[t_2' = 276m/c = 9.2 \times 10^{-7} s \right] \checkmark$$



- c) Let E_1 be the event when the nose of B reaches the nose of A.
 Let E_2 be the event when the nose of B reaches the tail of A.

Now $\Delta x = x_2 - x_1 = -L_{AB} - 0 = -L_{AB}$
 $\Delta x' = x'_2 - x'_1 = 0$ (both events occur at the same x' coordinate)

so

$$\Delta t' = \gamma(\Delta t - \frac{v}{c^2} \Delta x)$$

$$\Delta x' = \gamma(\Delta x - v \Delta t) = 0$$

so $\Delta t = \frac{\Delta x}{v} = \frac{-L_{AB}}{v} = -\frac{(598\text{m})}{\frac{-12}{13}c}$
 $= 648\text{m}/c \times \left(\frac{13}{3.0 \times 10^8\text{m/s}}\right)$

$\left[\Delta t = 648\text{m}/c = 2.16 \times 10^{-6}\text{s} \right]$ - the time interval that ticks by on the S frame's watch

⇒ plugging $\Delta x = v \Delta t$ into the first eq yields:

$$\Delta t' = \gamma(\Delta t - \frac{v^2}{c^2} \Delta t) = \frac{1}{\gamma(1 - v^2/c^2)} (1 - v^2/c^2) \Delta t = \sqrt{1 - v^2/c^2} \Delta t$$

so

$$\Delta t' = (648\text{m}/c) \sqrt{1 - (12/13)^2} = (648\text{m}/c) \cdot \frac{5}{13} = 249\text{m}/c \times \left(\frac{13}{3.0 \times 10^8\text{m/s}}\right)$$

$\left[\Delta t' = 249\text{m}/c = 8.31 \times 10^{-7}\text{s} \right]$

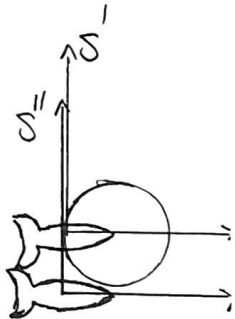
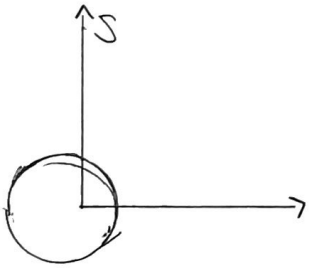
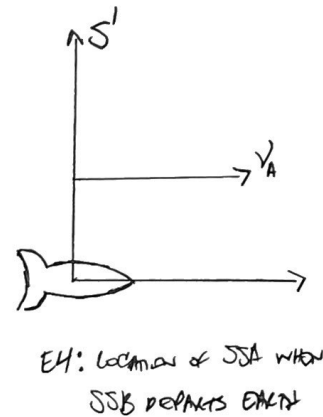
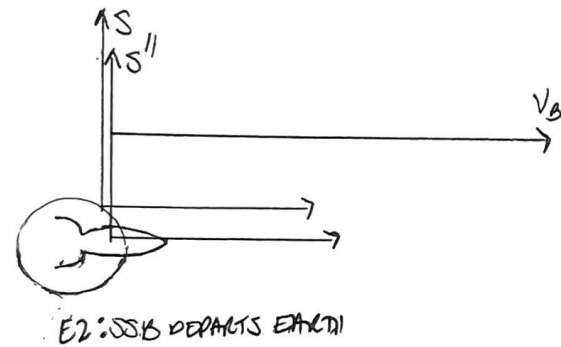
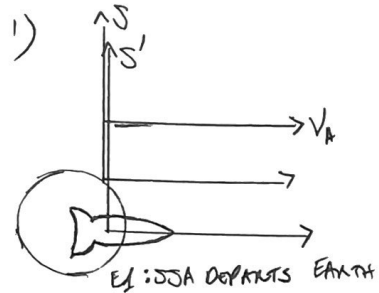
$\left[\begin{aligned} t_{NA2} &= t_{NA1} = 648\text{m}/c \\ t'_{NB1} &= t'_{NB1} + \Delta t' = 249\text{m}/c \\ t'_{NB2} &= t'_{NB1} + \Delta t' = 249\text{m}/c + 249\text{m}/c = 525\text{m}/c \end{aligned} \right]$

5. Let S be the Earth/Star rest frame. Let S' be spaceship A's rest frame, which moves at a speed $V_A = \frac{5}{13}c$. Let S'' be the rest frame of spaceship B, which moves at speed $V_B = \frac{12}{13}c$ relative to S .

Let E_1 be the departure of JSA from Earth

Let E_2 be the departure of JSB from Earth

Let E_3 be the arrival of $JSA \div JSB$ at the distant star



hence:

$$\Delta t_{12} = t_2 - t_1 = 7 \text{ ks}$$

$$\Delta x_{13} = \Delta x_{23} \quad (\text{the rest length between the Earth \& star})$$

$$\Delta t_{13} = \Delta t_{23} + \Delta t_{12} = \Delta t_{23} + 7 \text{ ks}$$

E_3 : $JSA \div JSB$ arriving at distant star

$$\Delta x'_{13} = \Delta x''_{23} = 0 \quad (\text{as both events occur at the same spatial coordinate})$$

Now, write down the Lorentz boosts and inverse Lorentz boosts between $E_1 \div E_2$ and between $E_2 \div E_3$

$$\text{let } \gamma_A \equiv \frac{1}{\sqrt{1 - v_A^2/c^2}}$$

$$\gamma_B \equiv \frac{1}{\sqrt{1 - v_B^2/c^2}}$$

$$\Delta t_{13} = \gamma_A (\Delta t'_{13} + \frac{v_A}{c^2} \Delta x'_{13})$$

$$\Delta x_{13} = \gamma_A (\Delta x'_{13} + v_A \Delta t'_{13})$$

$$\Delta t'_{13} = \gamma_A (\Delta t_{13} - \frac{v_A}{c^2} \Delta x_{13})$$

$$\Delta x'_{13} = \gamma_A (\Delta x_{13} - v_A \Delta t_{13})$$

which simplify to...

$$\Delta t_{13} = \gamma_A \Delta t'_{13}$$

$$\Delta x_{13} = \gamma_A v_A \Delta t'_{13}$$

$$\Delta t'_{13} = \gamma_A (\Delta t_{13} - \frac{v_A}{c^2} \Delta x_{13})$$

$$\Delta x_{13} = v_A \Delta t_{13}$$

$$\Delta t_{23} = \gamma_B (\Delta t''_{23} + \frac{v_B}{c} \Delta x''_{23})$$

$$\Delta x_{23} = \gamma_B (\Delta x''_{23} + v_B \Delta t''_{23})$$

$$\Delta t''_{23} = \gamma_B (\Delta t_{23} - \frac{v_B}{c^2} \Delta x_{23})$$

$$\Delta x''_{23} = \gamma_B (\Delta x_{23} - v_B \Delta t_{23})$$

$$\Delta t_{23} = \gamma_B \Delta t''_{23}$$

$$\Delta x_{23} = \gamma_B v_B \Delta t''_{23}$$

$$\Delta t''_{23} = \gamma_B (\Delta t_{23} - \frac{v_B}{c^2} \Delta x_{23})$$

$$\Delta x_{23} = v_B \Delta t_{23}$$

2) now $\Delta x_{13} = \Delta x_{23}$, so the last set of eqns give..

$$v_A \Delta t_{13} = v_B \Delta t_{23} = v_A (\Delta t_{23} + 7 \mu s)$$

$$\uparrow_{\text{usual}} \Delta t_{13} = \Delta t_{23} + 7 \mu s$$

now solve for Δt_{13} ...

$$(v_B - v_A) \Delta t_{23} = v_A \cdot 7 \mu s$$

$$\Delta t_{23} = \frac{v_A}{(v_B - v_A)} \cdot 7 \mu s = \frac{\frac{5}{13} c}{(\frac{12}{13} c - \frac{5}{13} c)} \cdot 7 \mu s = \frac{5/13}{7/13} \cdot 7 \mu s = 5 \mu s$$

$$\Delta t_{13} = \Delta t_{23} + 7 \mu s = 12 \mu s$$

$$\therefore \left[\Delta t_{13} = 12 \mu s, \Delta t_{23} = 5 \mu s \right]$$

IT TAKES JOA 12 μs TO COMPLETE THE JOURNEY, IT TAKES JOB 5 μs TO COMPLETE THE JOURNEY, RELATIVE TO THE EARTH-STAR FRAME.

$$c) \Delta x_{13} = v_A \Delta t_{13} = \left(\frac{5}{13} c\right) (12 \mu s) = \frac{60}{13} c \cdot \mu s$$

also

$$\Delta x_{23} = v_B \Delta t_{23} = \left(\frac{12}{13} c\right) (5 \mu s) = \frac{60}{13} c \cdot \mu s \quad \checkmark$$

so

$$\left[\Delta x_{13} = \Delta x_{23} = \frac{60}{13} c \cdot \mu s \right]$$

d. now, let's find the spacetime coordinates of SSA when SSB departs Earth

$$\Delta t'_{14} = \gamma_A \left(\Delta t_{14} - \frac{v_A}{c^2} \Delta x_{14} \right)$$

$$\Delta x'_{14} = \gamma_A \left(\Delta x_{14} - v_A \Delta t_{14} \right)$$

$$\Delta t_{14} = \gamma_A \left(\Delta t'_{14} + \frac{v_A}{c^2} \Delta x'_{14} \right)$$

$$\Delta x_{14} = \gamma_A \left(\Delta x'_{14} + v_A \Delta t'_{14} \right)$$

$$\text{now } \Delta x'_{14} = 0 \quad (\text{ca: EH occur at the same } x \text{ coordinate})$$

$$\Delta t_{14} = \Delta t_{12} \quad (\text{same time interval according to the } S \text{ frame})$$

using our 2 boosts the LORENTZ BOOSTS : INVERSE LORENTZ BOOSTS BECOME ...

$$\Delta t'_{14} = \gamma_A \left(\Delta t_{12} - \frac{v_A}{c^2} \Delta x_{14} \right)$$

$$\Delta x_{14} = v_A \Delta t_{12}$$

$$\Delta t_{12} = \gamma_A \Delta t'_{14}$$

$$\Delta x_{14} = \gamma_A v_A \Delta t'_{14}$$

so, the 3rd eqn yields..

$$\Delta t'_{14} = \frac{1}{\gamma_A} \Delta t_{12} = \sqrt{1 - v_A^2/c^2} \Delta t_{12} = \sqrt{1 - (5/13)^2} (7 \mu s) = \frac{12}{13} \cdot 7 \mu s$$

$$\left[\Delta t'_{14} = \frac{84}{13} \mu s \right] -$$

e) WE NEED $\Delta t'_{13}$...

FROM EARLIER...

$$\Delta t'_{13} = \frac{1}{\gamma_A} \Delta t_{13} = \sqrt{1 - v_A^2/c^2} \Delta t_{13} = \sqrt{1 - (5/13)^2} (12 \mu s) = \frac{12}{13} \cdot 12 \mu s$$

$$\left[\Delta t'_{13} = \frac{144}{13} \mu s \right]$$

f) WE NEED $\Delta t''_{23}$...

FROM EARLIER...

$$\Delta t''_{23} = \frac{1}{\gamma_B} \Delta t_{23} = \sqrt{1 - v_B^2/c^2} \Delta t_{23} = \sqrt{1 - (12/13)^2} (5 \mu s) = \frac{5}{13} \cdot 5 \mu s$$

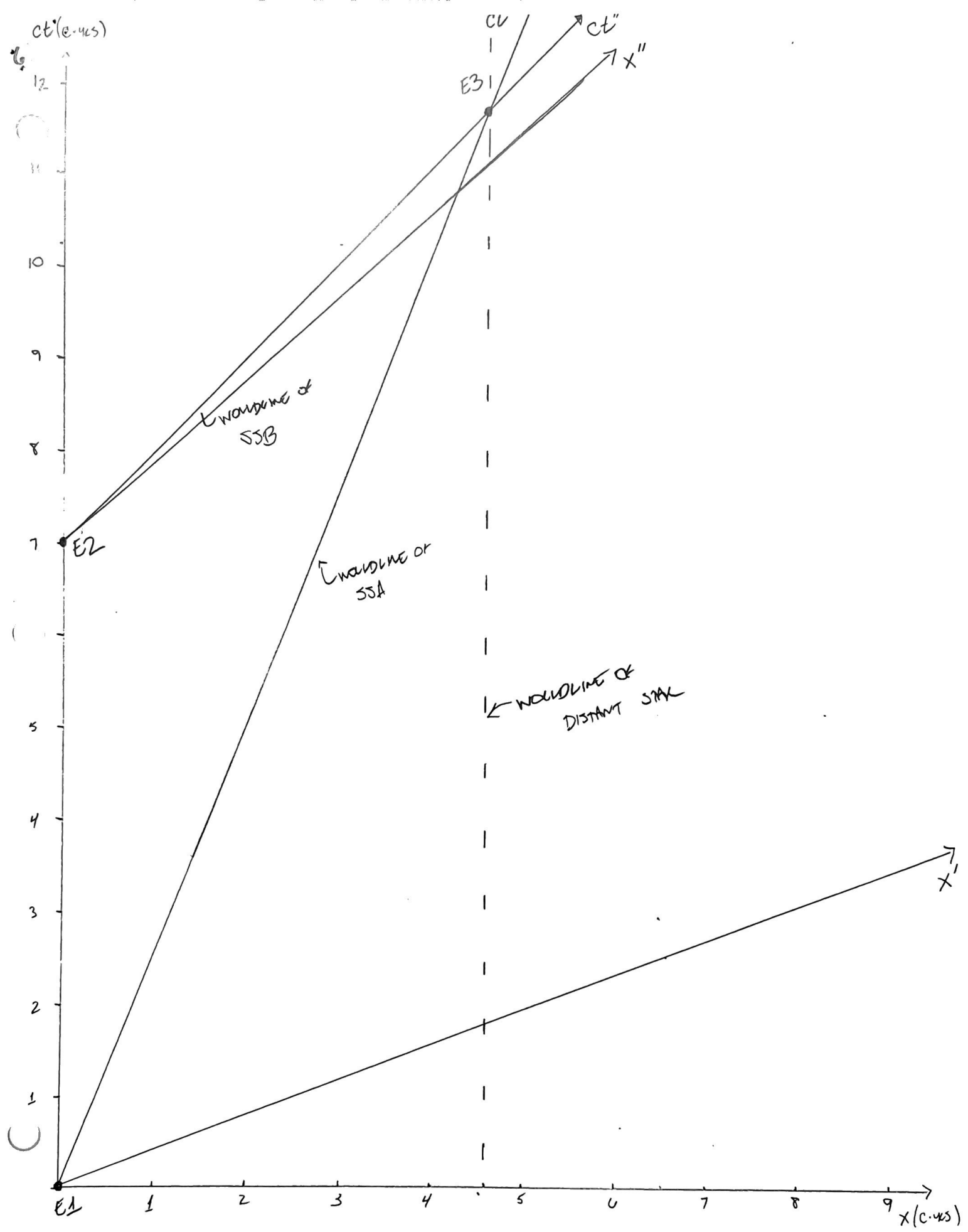
$$\left[\Delta t''_{23} = \frac{25}{13} \mu s \right]$$

$$g) \ell' = v_A \Delta t'_{13} = \left(\frac{5}{13} c \right) \left(\frac{144}{13} \mu s \right) = \frac{720}{169} c \cdot \mu s$$

$$\left[\ell' = \frac{720}{169} c \cdot \mu s \right]$$

$$h) \ell'' = v_B \Delta t''_{23} = \left(\frac{12}{13} c \right) \left(\frac{25}{13} \mu s \right) = \frac{300}{169} c \cdot \mu s$$

$$\left[\ell'' = \frac{300}{169} c \cdot \mu s \right]$$



b) SINCE SPACESHIP A TRAVELS AT A SPEED

$$V_A = \frac{5}{13} c = \frac{\Delta x}{\Delta t} \quad \text{so} \quad \frac{c \Delta t}{\Delta x} = \frac{13}{5}$$

- THE WORLDLINE OF SSA HAS A SLOPE OF $\frac{13}{5}$ RELATIVE TO THE S FRAME.

SINCE SPACESHIP B TRAVELS AT A SPEED

$$V_B = \frac{12}{13} c = \frac{\Delta x}{\Delta t}$$

then

$$\frac{c \Delta t}{\Delta x} = \frac{13}{12}$$

- THE WORLDLINE OF SSB HAS A SLOPE OF $\frac{13}{12}$ RELATIVE TO THE S FRAME, w/ A

CT INTERCEPT AT $ct = 17c \mu s$

THE WORLDLINES OF SSA & SSB CORRESPOND TO THE ct' & ct'' AXES.

• THE x' AXIS HAS A SLOPE OF $\frac{5}{13}$; THE x'' AXIS HAS A SLOPE OF $\frac{12}{13}$