

## Ch. 3 Space Time : Gravity in Newtonian Mechanics

1-30-19

What is an inertial reference frame..?

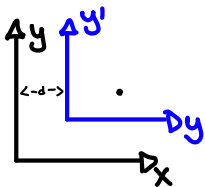
i.e. A free particle obeys

$$\ddot{\vec{r}} = \frac{d^2 \vec{r}}{dt^2} = 0 \Rightarrow \dot{\vec{v}} = \text{constant vector}$$

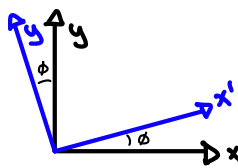
$$(*) \left[ \frac{d^2 x}{dt^2} = \frac{d^2 y}{dt^2} = \frac{d^2 z}{dt^2} = 0 \right]$$

• Laws of Newtonian mechanics take their simplest form in inertial reference frame.

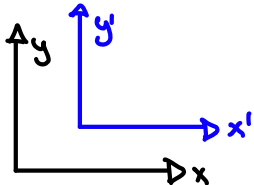
To define the rest

1.)   $x' = x - d$   
 $y' = y$   
 $z' = z$

2.) Rotation by  $\phi$

  $x' = (\cos\phi)x + (\sin\phi)y$   
 $y' = (-\sin\phi)x + (\cos\phi)y$   
 $z' = z$

3.) Uniform motion by velocity  $v$

  $x' = x - vt$   
 $y' = y$   
 $z' = z$

Central assumption of Newtonian mechanics is that there is absolute time,  $t = t'$

Galilean Transformations  $\rightarrow$  Not valid at speeds close to  $c$ .

$$\begin{aligned} x' &= x - vt \\ y' &= y \\ z' &= z \\ t' &= t \end{aligned}$$

### The Principle of Relativity

Identical experiments carried out in different inertial reference frames

If (\*), then

$$\left[ \frac{d^2 x'}{dt^2} = \frac{d^2 y'}{dt^2} = \frac{d^2 z'}{dt^2} = 0 \right] \rightarrow \text{Invariant under Galilean Transformations}$$

- Principle of Relativity is possible only because the geometry of Euclidean Space shares these Symmetries.
- One can verify these Symmetries of Euclidean geometry by examining,

$$ds^2 = dx^2 + dy^2 + dz^2$$

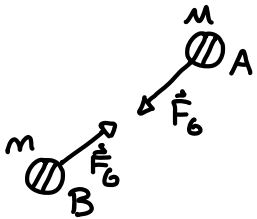
i.e. Rotation

$$\begin{aligned} x &= (\cos\phi)x' - (\sin\phi)y' \\ y &= (\sin\phi)x' + (\cos\phi)y' \\ z &= z' \end{aligned}$$

$$\begin{aligned} dx &= (\cos\phi)dx' - (\sin\phi)dy' \\ dy &= (\sin\phi)dx' + (\cos\phi)dy' \\ dz &= dz' \end{aligned}$$

$$\begin{aligned} ds^2 &= (\cos\phi dx' - \sin\phi dy')^2 + (\sin\phi dx' + \cos\phi dy')^2 + dz'^2 \\ ds^2 &= \cos^2\phi dx'^2 - 2\cos\phi\sin\phi dx'dy' + \sin^2\phi dy'^2 + \sin^2\phi dx'^2 + 2\sin\phi\cos\phi dx'dy' + \cos^2\phi dy'^2 + dz'^2 \\ &= dx'^2(\cos^2\phi + \sin^2\phi) + dy'^2(\cos^2\phi + \sin^2\phi) + dz'^2 = dx'^2 + dy'^2 + dz'^2 \end{aligned}$$

### Newtonian Gravity



$$\left[ \vec{F}_B = -\frac{GmM}{r^2} \hat{e}_r \right] \text{ — Newton's universal law of gravitation}$$

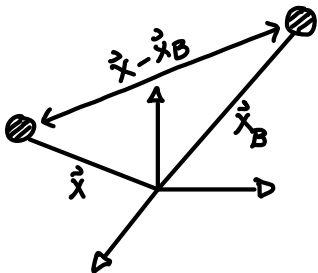
We can write this in terms of a Gravitational Potential

$$\vec{F}_B = -m \vec{\nabla} \Phi(x_B) \text{ Similar for Electrostatics: } \vec{F}_E = -q \vec{\nabla} \phi = qE$$

$\Phi(\vec{x}_B)$  is the gravitational potential produced by A at B's position  $\vec{x}_B$

$$\Phi(\vec{x}) = -\frac{GM}{r} = -\frac{GM}{|\vec{x} - \vec{x}_A|}$$

2-1-19



- Gravitational potential @ B masses  $M_A$ ,  $A=1,2,3,\dots$

$$\Phi(\vec{x}) = - \sum_A \frac{GM_A}{|\vec{x} - \vec{x}_A|}$$

For a Continuous Distribution of Mass Density  $\mu(\vec{x}')$

$$\sum_A M_A \rightarrow \int \mu(\vec{x}') d^3x'$$

$$\left[ \Phi(\vec{x}) = - \int d^3x' \cdot \frac{G\mu(\vec{x}')}{|\vec{x} - \vec{x}'|} \right] \quad \left( \text{Like } \Phi(r) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(r') dV'}{r} \text{ in E: } \mathcal{M} \right)$$

Define The Newtonian Gravitational Field  $\vec{g}(\vec{x})$  (Like  $\vec{E}(\vec{x})$  in E:  $\mathcal{M}$ )

$$\left[ \vec{g}(\vec{x}) \equiv -\vec{\nabla}\Phi(\vec{x}) \right] \quad \therefore \vec{F}_{\text{grav}} = m\vec{g}$$

The Differential Form of (\*) ....

$$\left[ \vec{\nabla} \cdot \vec{g}(\vec{x}) = -4\pi G\mu(\vec{x}) \right] \quad \left( \text{Like } \vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} \right)$$

or

$$\left[ \nabla^2 \Phi(\vec{x}) = 4\pi G\mu(\vec{x}) \right] \quad \text{since } \vec{\nabla} \cdot \vec{\nabla}\Phi = \nabla^2 \Phi \quad \left( \text{Like } \nabla^2 \Phi = -\frac{\rho}{\epsilon_0} \right)$$

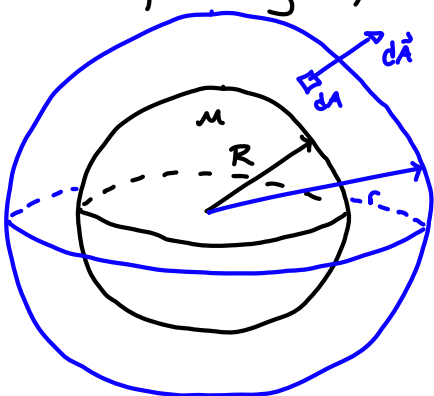
$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \text{ is the Laplacian}$$

To arrive @ the integral form, integrate over the volume ....

$$\int_V \vec{\nabla} \cdot \vec{g}(\vec{x}) d^3x = -4\pi G \int_V \mu(\vec{x}) d^3x$$

$$\int_V \vec{\nabla} \cdot \vec{g}(\vec{x}) d^3x = \oint_S \vec{g}(\vec{x}) \cdot d\vec{A} \quad \rightarrow \quad \oint_S \vec{g} \cdot d\vec{A} = -4\pi G\mu \quad \left( \oint_S \vec{E} \cdot d\vec{A} = \frac{Q_{\text{enc}}}{\epsilon_0} \right)$$

For a spherically-symmetric mass .....



$$\oint_S \vec{g} \cdot d\vec{A} \quad \vec{g} \cdot d\vec{A} = g dA \cos\theta \Big|_{\theta=0}^{\theta=180^\circ} = -g dA, \quad \vec{g} = -g \hat{e}_r$$

$$\oint_S \vec{g} \cdot d\vec{A} = -\oint_S g dA = -g \oint_S dA = -g \cdot 4\pi r^2$$

$$-g \cdot 4\pi r^2 = -4\pi G\mu$$

$$g = \frac{G\mu}{r^2}, \quad \vec{g} = -g \hat{e}_r = -\frac{G\mu}{r^2} \hat{e}_r$$

$$\vec{\nabla}\Phi = \frac{G\mu}{r^2} \hat{e}_r = \frac{\partial\Phi}{\partial r} \hat{e}_r, \quad \frac{d\Phi}{dr} = \frac{G\mu}{r^2}, \quad \Phi = -\frac{G\mu}{r}$$

Gravitational : Inertial Mass

$$\vec{F}_{\text{grav}} = -m_g \vec{\nabla}\Phi(\vec{x}) = m_g \vec{g}$$

Gravitational mass - measures the strength of gravitational force between bodies

• Newton's 2<sup>nd</sup> Law of motion is .....

$$\vec{F} = m_I \vec{a}$$

↑  
Inertial mass - Governs the inertial properties of the body

Newton's 2<sup>nd</sup> Law for a body in a gravitational field ....

$$-m_G \vec{\nabla} \Phi = m_G \vec{g} = m_I \vec{a} \quad \text{Since } [m_I = m_G]$$

$$\vec{a} = -\vec{\nabla} \Phi = \vec{g} - \text{All bodies fall with the same acceleration in a gravitational field}$$