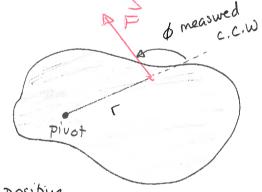
Thes: Discussion / 9 mit
Chiz Conc @ 1,8

Prob 4,7,13,20,21

Torque:

The rotational effects of forces are determined via associated quantities called torque. With reference to the diagram, the torque produced by a force is:

and this has units of N·m. In this formula:



- 1) Ford F are magnitudes and are positive
- z) ϕ can be such that torque is:

 positive = D counterclockwise

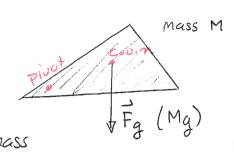
 negative = 0 clockwise,

In general many forces might act on an object. Then the net torque

$$T_{\text{net}} = T_1 + T_2 + \dots = \sum_{\text{all fores}} T_i$$

Additionally one can show that.

The torque produced by the gravitational forces acting on various parts of on object is exactly the same as the torque produce by a downward force with magnitude Mg acting at the center of mass



Example: A 2.0m long long rod can pivot about one end. A rope pulls horizontally with a 2000 tension at the other end, Determine

- a) torque produced by rope
- " gravity
- " pivet
- d) net torque

Answer: a) T = r.F.sin &

= 2.0mx 200N sin 1500

(Trope = 200 N·m

Tg= rFsing = 1.0m × 98N sin 240° 7g = -85 N.M

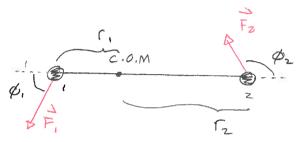
30°

 $\phi = 180^{\circ} - 30^{\circ} = 150^{\circ}$ φ=180° +60° =240° = 10kg × 968m/s2 = 98N

Torque and dynamics

How do torques determine the dynamics of any object? As an example, consider two masses connected by a rigid rod with negligible mass. The rod can rotate about an

axle through its center of mass. We compute the angular acceleration about the c.o.m, via the tangential acceleration. For 1



where $F_{i,t}$ is the tangential component of the force: $F_{i,t} = F_{i,t} \sin \phi_{i,t}$. Then $a_{i,t} = r_{i,t} \times gives$

 $F_1 \sin \phi_1 = M_1 \Gamma_1 \propto = 0$ $\Gamma_1 F_1 \sin \phi_1 = M_1 \Gamma_1^2 \propto = 0$ $\Gamma_1 = M_1 \Gamma_1^2 \propto$ Similarly $\Gamma_2 = M_2 \Gamma_2^2 \propto$ and combining gives:

$$T_{\text{net}} = T_1 + T_2 = \left[M_1 \Gamma_1^2 + M_2 \Gamma_2^2 \right] \alpha$$

This shows that not only mass, but also the location of mass eletermines angular acceleration. This motivates the definition of the moment of inertia.

The moment of inertia of a collection of pointlike objects about a point O is $T = m_1 r_1^2 + m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$ $T = \sum_{m \in \Gamma_1} m_2 r_2^2 + m_3 r_3^2 + \dots$

Units kg-m2

Thus the moment of inertia is determined by:

- 1) mass present
- 2) location of the mass.

Worm Upl

The derivation presented is an example of a general rule:

The net torque on an object determines its angular acceleration via $T_{net} = I \times I$

Worm Up 2

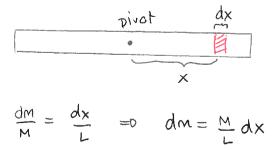
Demo Red/blue sticks

Moment of inertia: continuous mass distributions

The previous derivation only considered point masses. But the calculation can be extended to continuous mass distributions. The following example illustrates this:

Example: Determine the moment of inertia of a uniform rod with length L and total mass M. (about its center)

Answer: 1) Decompose object into small pieces. Consider the piece of width dx at location x. The mass of this piece is dm. Ratios give



2) Determine moment of inertia of each piece:

$$dI = \Gamma^2 dm = x^2 \frac{M}{L} dx$$

3) Add all such moments for all segments

For the rod, this means adding all contributions from $x=-\frac{1}{z}$ to $x=\frac{1}{z}$.

$$I = \sum dI = \int dI$$
 all segments

$$= \int_{-\frac{L}{z}}^{\frac{L}{z}} \frac{M}{L} dx$$

$$= \frac{M}{L} \int_{-L/2}^{L/2} x^2 dx = \frac{M}{L} \frac{1}{3} x^3 \Big|_{-L/2}^{L/2}$$

$$= \frac{M}{L} \frac{1}{3} \left\{ \left(\frac{L}{2}\right)^3 - \left(\frac{-L}{2}\right)^3 \right\}$$

$$= D I = \frac{M}{3L} \left\{ \frac{L^3}{8} + \frac{L^3}{8} \right\} = \frac{ML^2}{12}$$

$$= 0 \qquad I = \frac{1}{12} M L^2$$

M

Similar calculations can be done for various symmetrical extended objects