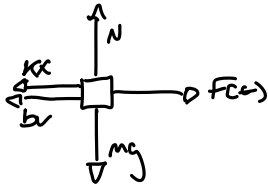
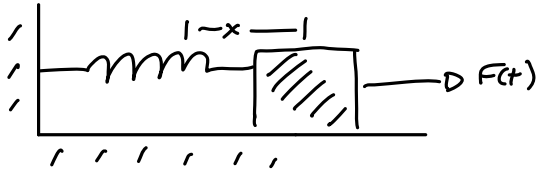


Ch.3 Forced Oscillations

Consider a damped, driven harmonic oscillator



Consider the driving force to be harmonic and of the form

$$F(t) = F_0 \cos(\omega t)$$

$$\gamma = \frac{b}{m} \quad \omega_0 = \sqrt{\frac{k}{m}}$$

Newton's 2nd Law...

$$\begin{aligned} \cdot \vec{F}_{\text{net}} &= m\vec{a}_x \\ \cdot F_0 \cos(\omega t) - kx - bv &= m\ddot{x} \\ \cdot m \frac{d^2 x}{dt^2} + b \frac{dx}{dt} + kx &= F_0 \cos(\omega t) \end{aligned}$$

$$* \left[\frac{d^2 x}{dt^2} + \gamma \frac{dx}{dt} + \omega_0^2 x = \frac{F_0}{m} \cos(\omega t) \right]$$

For \dagger

Try a solution of the form

$$x(t) = A(\omega) \cos(\omega t - \phi)$$

$$\frac{dx}{dt} = -\omega A(\omega) \sin(\omega t - \phi)$$

$$\frac{d^2 x}{dt^2} = -\omega^2 A(\omega) \cos(\omega t - \phi)$$

Notice!

- The amplitude is a function of the frequency
- ϕ is the phase angle between the driving force and the resultant displacement

Plugging into (*)

$$-\omega^2 A(\omega) \cos(\omega t - \phi) - \omega \gamma A(\omega) \sin(\omega t - \phi) + \omega_0^2 A(\omega) \cos(\omega t - \phi) = \frac{F_0}{m} \cos(\omega t)$$

$$\left[(\omega_0^2 - \omega^2) \cos(\omega t - \phi) - \omega \gamma \sin(\omega t - \phi) \right] A(\omega) = \frac{F_0}{m} \cos(\omega t)$$

Remember...

$$\cos(\omega t - \phi) = \cos(\omega t) \cos(\phi) + \sin(\omega t) \sin(\phi)$$

$$\sin(\omega t - \phi) = \sin(\omega t) \cos(\phi) - \cos(\omega t) \sin(\phi)$$

Plugging into *

$$\left[(\omega_0^2 - \omega^2) [\cos(\omega t) \cos(\phi) + \sin(\omega t) \sin(\phi)] \right] - \omega \gamma \left[\sin(\omega t) \cos(\phi) - \cos(\omega t) \sin(\phi) \right] A(\omega) = \frac{F_0}{m} \cos(\omega t)$$

$$\left[[(\omega_0^2 - \omega^2) \cos(\phi) + \omega \gamma \sin(\phi)] \cos(\omega t) + [(\omega_0^2 - \omega^2) \sin(\phi) - \omega \gamma \cos(\phi)] \sin(\omega t) \right] A(\omega) = \frac{F_0}{m} \cos(\omega t)$$

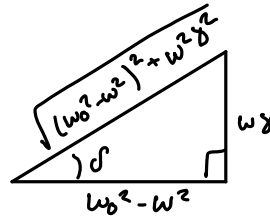
This equation must hold for all time

$$1.) [(\omega_0^2 - \omega^2) \cos(\phi) + \omega \gamma \sin(\phi)] A(\omega) = F_0/m$$

$$2.) [(\omega_0^2 - \omega^2) \sin(\phi) - \omega \gamma \cos(\phi)] A(\omega) = 0$$

For eqn 2)

$$\tan(\phi) = \frac{\omega \gamma}{(\omega_0^2 - \omega^2)} \quad \therefore \phi = \tan^{-1} \left(\frac{\omega \gamma}{(\omega_0^2 - \omega^2)} \right)$$



(*) (*)

$$\sin(\phi) = \frac{\omega \gamma}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}}$$

$$\cos(\phi) = \frac{\omega_0^2 - \omega^2}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}}$$

For (*) (*)

$$\left[\frac{(\omega_0^2 - \omega^2)(\omega_0^2 - \omega^2)}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} + \frac{\omega \gamma \cdot \omega \gamma}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} \right] A(\omega) = \frac{F_0}{m}$$

$$\left[\frac{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} \right] A(\omega) = F_0/m$$

$$\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2} A(\omega) = F_0/m$$

The solution is then . . .

$$[x(t) = A(\omega) \cos(\omega t - \phi)]$$

where

$$\tan(\phi) = \frac{\omega \gamma}{(\omega_0^2 - \omega^2)} \quad \phi \equiv \text{Phase angle}$$

$$\left[A(\omega) = \frac{F_0/m}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} \right]$$

Explore limiting cases

$$CI: \lim_{\omega \rightarrow 0} \tan(\phi) \rightarrow 0 \quad \phi \rightarrow 0 \text{ (or } \pi)$$

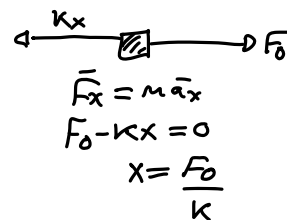
$$\lim_{\omega \rightarrow 0} A(\omega) \rightarrow \frac{F_0/m}{\omega_0^2} = \frac{F_0}{m \omega_0^2} = \frac{F_0}{k}$$

$$\lim_{\omega \rightarrow 0} A(\omega) \rightarrow \frac{F_0/m}{\omega_0^2} = \frac{F_0}{m \omega_0^2} = \frac{F_0}{k}$$

$$x(t) = F_0/k$$

Notice

- For small Frequencies, mass moves with the driving force.



$$CII: \lim_{\omega \rightarrow \infty} \tan(\phi) \rightarrow \frac{\gamma}{-\omega} \rightarrow 0 \quad \therefore \phi = \pi$$

$$\lim_{\omega \rightarrow \infty} A(\omega) \rightarrow 0 \quad \text{Notice}$$

- As ω increases, amplitude decreases, and the mass is in the opposite direction of the driving force
- As $\omega \rightarrow \infty$, mass doesn't move due to inertia

CT III: $\lim_{\omega \rightarrow \omega_0} \tan(\phi) \rightarrow \infty \therefore \phi \rightarrow \pi/2$

$\omega = \omega_0$

$\lim_{\omega \rightarrow \omega_0} A(\omega) = \frac{F_0/m}{\gamma}$

Notice

- When driving frequency is the same as the natural frequency, the displacement lags the driving force by $\pi/2$
- Velocity leads the displacement by $\pi/2$
The mass is moving in the same direction as the driving force

What frequency ω yields $A_{\max} = A(\omega_{\max})$?

$$0 = \frac{dA(\omega)}{d\omega} = \frac{F_0}{m} \frac{d}{d\omega} \left[(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2 \right]^{-1/2} = \frac{-F_0}{2m} \left[(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2 \right]^{-3/2} \left[2(\omega_0^2 - \omega^2) \cdot 2\omega + 2\omega \gamma^2 \right]$$

$$0 = \frac{F_0 \omega}{m} \left[(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2 \right]^{-3/2} \left[2(\omega_0^2 - \omega^2) - \gamma^2 \right]$$

$\underbrace{\quad}_0$ - must be 0

$$2(\omega_0^2 - \omega^2) - \gamma^2 = 0$$

$$\omega_0^2 - \frac{\gamma^2}{2} = \omega^2$$

$$\omega_0^2 \left(1 - \frac{\gamma^2}{2\omega_0^2} \right) = \omega^2$$

$$\omega = \omega_0 \sqrt{1 - \frac{\gamma^2}{2\omega_0^2}}$$

$$A_{\max} = A(\omega_{\max}) = \frac{F_0/m}{\sqrt{\frac{\gamma^4}{4} + \left(\omega_0^2 - \frac{\gamma^2}{2} \right) \gamma^2}} = \frac{F_0/m}{\sqrt{\omega_0^2 \gamma^2 - \frac{\gamma^4}{4}}}$$

$$\left[\omega_{\max} = \omega_0 \sqrt{1 - \frac{\gamma^2}{2\omega_0^2}} \right]$$

↑ maximizes amplitude

$$\left[\begin{array}{l} A_{\max} = \frac{F_0/m}{\gamma \omega_0 \sqrt{1 - \frac{\gamma^2}{4\omega_0^2}}} \\ \omega_{\max} = \omega_0 \sqrt{1 - \frac{\gamma^2}{2\omega_0^2}} \end{array} \right]$$

$$Q = \frac{\omega_0}{\gamma}$$

$$\frac{F_0 Q}{K} (1 - 1/4 Q^2)^{1/2} = \frac{F_0}{K} \frac{Q}{\sqrt{1 - 1/4 Q^2}}$$

For very light damping, $\omega_0 \gg \gamma$ & $Q \gg 1$

$$\lim_{Q \gg 1} A_{\max} \sim Q$$

Harmonic driving force acts as an amplifier

@ $\omega = \omega_{\max}$ Determine $A(\omega) : \rho$

$$\omega_0 = \sqrt{\frac{K}{m}} = \sqrt{\frac{150 \text{ N/m}}{1.5 \text{ kg}}}$$

$$\gamma = \frac{b}{m} = \frac{3.0 \text{ kg/s}}{1.5 \text{ kg}} = 2 \text{ s}^{-1}$$

$$\omega_0 = 10 \text{ rad/s} \quad \gamma = 2 \text{ s}^{-1}$$

FBD

$$K(x - \frac{b}{\gamma})$$

$$F_{\text{net}} = m \ddot{x}$$

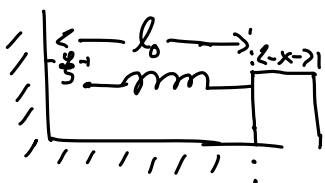
$$-K(x - \frac{b}{\gamma}) - b\dot{x} = m\ddot{x}$$

$$m\ddot{x} + b\dot{x} + Kx = K \frac{b}{\gamma}$$

$$F_0 = K a = (150 \text{ N/m})(5.0 \times 10^{-3} \text{ m}) = 0.750 \text{ N}$$

$$= K a \cos(\omega t)$$

$$A(\omega) = \frac{F_0/m}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} = \frac{0.750 \text{ N} / 1.5 \text{ kg}}{\sqrt{(10 \text{ rad/s}^2 - 6 \text{ rad/s}^2)^2 + (6 \text{ rad/s}^2)^2 (2.0 \text{ s}^{-1})^2}} = 1.9 \times 10^{-3} \text{ m}$$



$$m = 1.5 \text{ kg}$$

$$K = 150 \text{ N/m}$$

$$b = a \cos(\omega t)$$

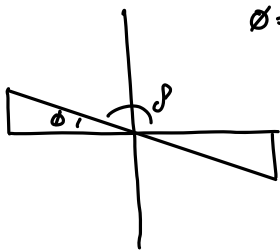
$$a = 5.0 \times 10^{-3} \text{ m}$$

$$\omega = 6 \text{ rad/s}$$

$$b = 3.0 \frac{\text{kg}}{\text{s}}$$

$$\phi = \tan^{-1} \left(\frac{\gamma \omega}{\omega_0^2 - \omega^2} \right)$$

$$= \tan^{-1} \left(\frac{(67 \text{ rad/s})(2.05^{-1})}{(10 \text{ rad/s})^2 - (67 \text{ rad/s})^2} \right) = \tan^{-1}(-0.15)$$



$$\phi = \tan^{-1}(0.15) = 8.5^\circ$$

$$\phi = 171.5^\circ \cdot \frac{\pi \text{ rad}}{180^\circ} = 3.0 \text{ rad}$$

$$\phi = 3.0 \text{ rad}$$

$$\omega_{\max} = \omega_0 \sqrt{1 - \frac{\gamma^2}{4\omega_0^2}} = (10 \text{ rad/s}) \sqrt{1 - \frac{(2.05^{-1})^2}{4(10 \text{ rad/s})^2}} = 9.9 \text{ rad/s}$$

$$A_{\max} = \frac{F_0/m}{\gamma \omega_0} \frac{1}{\sqrt{1 - \gamma^2/4\omega_0^2}} = \frac{0.75 \text{ N}/1.5 \text{ kg}}{(2.05^{-1})(10 \text{ s}^{-1})} \cdot \frac{1}{\sqrt{1 - \frac{(2.05^{-1})^2}{4(10 \text{ s}^{-1})^2}}}$$

$$A_{\max} = 2.5 \times 10^{-2} \text{ m}$$

Transient Phenomena

Claim: Let x_1 be a solution to the damped driven harmonic oscillator

$$\frac{d^2 x_1}{dt^2} + \gamma \frac{dx_1}{dt} + \omega_0^2 x_1 = \frac{F_0}{m} \cos(\omega t)$$

Let x_2 be a solution to the damped harmonic oscillator

$$\frac{d^2 x_2}{dt^2} + \gamma \frac{dx_2}{dt} + \omega_0^2 x_2 = 0$$

Claim: $x = x_1 + x_2$ is a solution to the damped driven harmonic oscillator

Proof

$$\frac{d^2 x}{dt^2} + \gamma \frac{dx}{dt} + \omega_0^2 x = \frac{F_0}{m} \cos(\omega t)$$

$$\frac{d^2 (x_1 + x_2)}{dt^2} + \gamma \frac{d(x_1 + x_2)}{dt} + \omega_0^2 (x_1 + x_2) = \frac{F_0}{m} \cos(\omega t)$$

$$\underbrace{\left(\frac{d^2 x_1}{dt^2} + \gamma \frac{dx_1}{dt} + \omega_0^2 x_1 \right)}_{\frac{F_0}{m} \cos(\omega t)} + \underbrace{\left(\frac{d^2 x_2}{dt^2} + \gamma \frac{dx_2}{dt} + \omega_0^2 x_2 \right)}_0 = \frac{F_0}{m} \cos(\omega t)$$

$$\left[\begin{aligned} x_1(t) &= A(\omega) \cos(\omega t - \phi) - \text{steady state} \\ x_2(t) &= B e^{-\gamma t/2} \cos[\sqrt{\omega_0^2 - (\gamma/2)^2} t + \phi] - \text{transient solution} \end{aligned} \right]$$

The sum of the two is the solution

$$x(t) = x_1(t) + x_2(t)$$

- $A(\omega)$ & ϕ are functions of ω
- B & ϕ are initial conditions