Improving Quantum Parameter Estimation

Taylor Larrechea

What is Quantum Estimation?

Spin-1/2 Farti

What is Estimation

Coin Toss

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Single Particle

Multiple Particles

Density Op

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Improving Quantum Parameter Estimation

Taylor Larrechea

Colorado Mesa University

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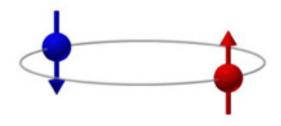
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Spin-1/2 Particles

Quantum estimation is the use of quantum mechanical systems as measuring devices for physical parameters.



In the context of this presentation, these quantum mechanical systems are spin-1/2 particles.

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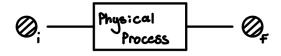
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Spin-1/2 Particles in Magnetic Field

These spin-1/2 particles can be subjected to something like a magnetic field.



Can measure the change in the spin vector.

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Example of Quantum Estimation

An example of spin-1/2 particles and parameters is in MRI machines.



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Estimation

Estimation

The fine art of guessing





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Coin Toss

The simplest example of estimation is a coin toss.



The result of this coin flip is either heads or tails.

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Flip

Probability

Wish to estimate the probability of obtaining heads.

N is the number of times flipped, n_H is how many times heads appears.

Probability estimate of obtaining heads is

$$P_{\rm est} = \frac{n_H}{N}.\tag{1}$$

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Multiple runs of N flips will give different estimates.

Fluctuation of estimates is calculated via

$$\nu = \bar{P}_{est}^2 - (\bar{P}_{est})^2. \tag{2}$$

 \bar{P}^2_{est} and \bar{P}_{est} are calculated via

$$\bar{P}_{est} = \frac{1}{N} \bar{n}_H$$
 and $\bar{P}_{est}^2 = \frac{1}{N^2} \bar{n}_H^2$. (3)

We aim to have the smallest variance possible.

$$\nu = \frac{P_{est}}{N} (1 - P_{est}). \tag{4}$$

The Fisher Information is calculated via

$$F = \sum_{\text{All Possible Outcomes}}^{\infty} \frac{1}{\text{Prob(Outcome)}} \left(\frac{\partial \text{Prob(Outcome)}}{\partial P} \right)^{2}. \tag{5}$$

The Cramer-Rao bound tells us that

$$\nu \ge \frac{1}{F}.\tag{6}$$

where F is calculated in equation (5). The Fisher Information of our coin flip experiment via

$$F_{\text{N Toss}} = N \cdot F_{\text{Single Toss}} = N \cdot \frac{1}{P(1-P)} = \frac{N}{P(1-P)}.$$
 (7)

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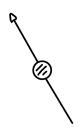
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The state of a single particle mathematically is represented by

$$|+\hat{n}\rangle = \cos(\theta/2)|0\rangle + e^{i\phi}\sin(\theta/2)|1\rangle$$
 (8)

and

$$|-\hat{n}\rangle = \sin(\theta/2)|0\rangle - e^{i\phi}\cos(\theta/2)|1\rangle.$$
 (9)



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States can evolve via physical processes, mathematically this is calculated by

$$|\psi_f\rangle = \hat{U}|\psi_i\rangle. \tag{10}$$



It should be noted that \hat{U} are $2^n \times 2^n$ matrices where n is the number of particles.

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Unitary Evolution

This evolution is calculated by

$$|\psi_f\rangle = \hat{U}(\lambda) |\psi_i\rangle.$$
 (11)

An example of this unitary operator is

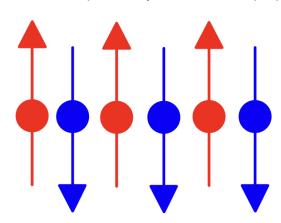
$$\hat{U}(\lambda) = \begin{pmatrix} e^{-i\lambda/2} & 0\\ 0 & e^{i\lambda/2} \end{pmatrix}. \tag{12}$$

$$|\psi_{i}\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \rightarrow |\psi_{f}\rangle = \frac{1}{\sqrt{2}}(e^{-i\lambda/2}|0\rangle + e^{i\lambda/2}|1\rangle)$$
(13)

 λ is the parameter that we wish to estimate.

Multiple Particles

We can now discuss quantum systems with multiple particles.



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Product States

A generic state is

$$|\psi\rangle = c_0 |0\rangle |0\rangle + c_1 |0\rangle |1\rangle + c_2 |1\rangle |0\rangle + c_3 |1\rangle |1\rangle.$$
 (14)

A product state can be written as

$$|\psi\rangle = |\psi_A\rangle |\psi_B\rangle. \tag{15}$$

An example is

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) \cdot \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$
 (16)

where

$$|\psi\rangle = \frac{1}{2}(|0\rangle |0\rangle + |0\rangle |1\rangle - |1\rangle |0\rangle - |1\rangle |1\rangle). \tag{17}$$

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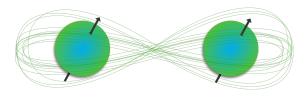
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Entangled States

An example of an entangled state is

$$|\psi\rangle = \frac{1}{2}(|0\rangle|0\rangle + |0\rangle|1\rangle + |1\rangle|0\rangle - |1\rangle|1\rangle)$$
 (18)

Entangled states do not exist in classical physics.



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Density Operators

Exact state unknown. Represented by state $|\psi_1\rangle$ with probability a_1 and $|\psi_2\rangle$ with probability a_2

$$\hat{\rho} = a_1 |\psi_1\rangle \langle \psi_1| + a_2 |\psi_2\rangle \langle \psi_2|. \tag{19}$$

For example

$$|\psi_1\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \text{ with } a_1 = \frac{1+r}{2}$$
 (20)

$$|\psi_2\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) \text{ with } a_2 = \frac{1-r}{2}$$
 (21)

the density operator is then

$$\hat{\rho} = \frac{1}{2} \begin{pmatrix} 1 & r \\ r & 1 \end{pmatrix}. \tag{22}$$

Equation (22) is the density operator that will be used without the rest of the talk.

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Density Operators

Parameter dependent density operators evolve via

$$\hat{\rho}(\lambda) = \hat{U}(\lambda)\hat{\rho}_0\hat{U}^{\dagger}(\lambda) \tag{23}$$

An example

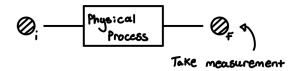
$$\hat{U}(\lambda) = \begin{pmatrix} e^{-i\lambda/2} & 0\\ 0 & e^{i\lambda/2} \end{pmatrix}$$
 and $\hat{\rho}_0 = \frac{1}{2} \begin{pmatrix} 1 & r\\ r & 1 \end{pmatrix}$ (24)

then evolves to

$$\hat{\rho}_f = \frac{1}{2} \begin{pmatrix} 1 & re^{-i\lambda} \\ re^{i\lambda} & 1 \end{pmatrix}. \tag{25}$$

Quantum Estimation

A key component of quantum estimation is called the Quantum Fisher Information.



Relationship between F and H is

$$F \le H. \tag{26}$$

H is not dependent upon measurement choice, only initial state.

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The way this Quantum Fisher Information is calculated is by

$$H = \text{Tr}[\hat{\rho}\hat{\mathcal{L}}^2] \tag{27}$$

where

$$\frac{\partial \hat{\rho}}{\partial \lambda} = \frac{1}{2} [\hat{\rho} \hat{\mathcal{L}} + \hat{\mathcal{L}} \hat{\rho}]. \tag{28}$$

Want to maximize Quantum Fisher Information.

$$\hat{\rho}_{i} = \frac{1}{2} \begin{pmatrix} 1 + r_{zi} & r_{xi} - ir_{yi} \\ r_{xi} + ir_{yi} & 1 - r_{zi} \end{pmatrix}$$
 (29)

Where $\hat{\rho}_f$ is

$$\hat{\rho}_f = \frac{1}{2} \begin{pmatrix} 1 + r_{zf} & r_{xf} - r_{yf} \\ r_{xf} + r_{yf} & 1 - r_{zf} \end{pmatrix}$$
(30)

Where

$$r_{xf} = r_{xi}\cos(\lambda) - r_{yi}\sin(\lambda)$$

$$r_{yf} = r_{xi}\sin(\lambda) + r_{yi}\cos(\lambda)$$

$$r_{zf} = r_{zi}.$$
(31)

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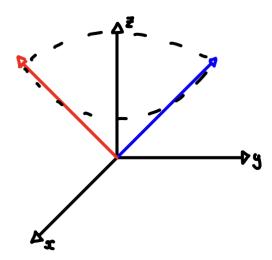
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Visualizing of Phase Shifts

We can picture this Phase Shift in the figure.



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One Particle Phase Shift

We now examine single particle phase shifts.



The Quantum Fisher Information is

$$H = r^2. (32)$$

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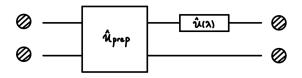
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Two Particle one Channel Phase Shift

We now examine a two particle phase shift.



The \hat{U}_{prep} matrix involves single qubit operations, one interaction is between two particles.

The Quantum Fisher Information is

$$H = \frac{2r^2}{1+r^2}. (33)$$

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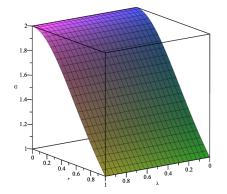
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Gain of Phase Shifts

One can calculate what is called the gain of two scenarios by $G = \frac{H_{both}}{H_{circle}}$. The gain of the last two scenarios is

$$G = \frac{2}{1+r^2}. (34)$$



When G > 1 the correlated scheme is estimating better.

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$$\hat{\rho}_f = (1 - \lambda)\hat{\rho}_i + \lambda\hat{\sigma}_z\hat{\rho}_i\hat{\sigma}_z.$$

$$\hat{\rho}_{i} = \frac{1}{2} \begin{pmatrix} 1 + r_{zi} & r_{xi} - ir_{yi} \\ r_{xi} + ir_{yi} & 1 - r_{zi} \end{pmatrix}$$
(35)

Where $\hat{\rho}_f$ is

$$\hat{\rho}_f = \frac{1}{2} \begin{pmatrix} 1 + r_{zi} & (1 - 2\lambda)(r_{xi} - ir_{yi}) \\ (1 - 2\lambda)(r_{xi} + ir_{yi}) & 1 - r_{zi} \end{pmatrix}. \quad (36)$$

The evolution of each component of these particles is

$$r_{xf} = r_{xi}(1 - 2\lambda)$$

$$r_{yf} = r_{yi}(1 - 2\lambda)$$

$$r_{zf} = r_{zi}.$$
(37)

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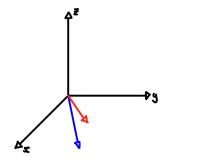
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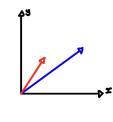
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Visualizing of Phase Flips

We can visualize how this Phase Flip evolves in the figure.





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QFI of Phase Flips

We examine the simplest scenario for a phase flip.



The Quantum Fisher Information is

$$H = \frac{4r^2}{1 - r^2(1 - 2\lambda)^2}. (38)$$

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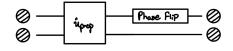
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The next scenario is with two particles, one channel phase flip.

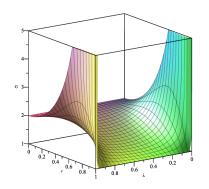


The Quantum Fisher Information is

$$H = \frac{8r^2(1+r^2)}{(1+r^2)^2 - 4r^2(1-2\lambda)^2}.$$
 (39)

The gain of the last two scenarios is

$$G = \frac{2(1+r^2)(1-r^2(1-2\lambda)^2)}{(1+r^2)^2 - 4r^2(1-2\lambda)^2}.$$
 (40)



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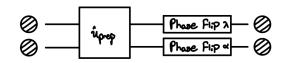
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.

We now examine calculating the Quantum Fisher Information of a measurement with noise in a spectator.



The Quantum Fisher Information is

$$H = \frac{8r^2(1-2\alpha)^2(1+r^2)}{(1+r^2)^2-4r^2(1-2\lambda)^2(1-2\alpha)^2}.$$
 (41)

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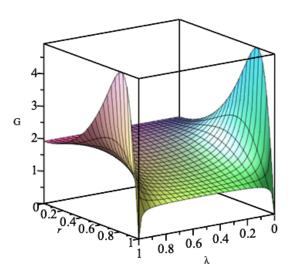
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The gain of the last scenario is

$$G = \frac{2(1-2\alpha)^2(1+r^2)(1-r^2(1-2\lambda)^2)}{(1+r^2)^2 - 4r^2(1-2\lambda)^2(1-2\alpha)^2}.$$
 (42)

We will examine gains of when $\alpha = 0.01, 0.1, 0.3, 0.5$.

Gain of when $\alpha = 0.01$.



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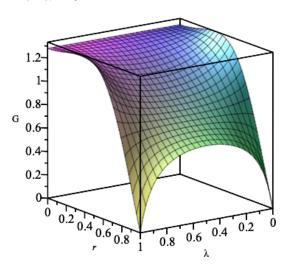
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Gain of when $\alpha = 0.1$.



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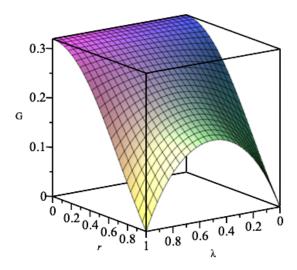
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Gain of when $\alpha = 0.3$.



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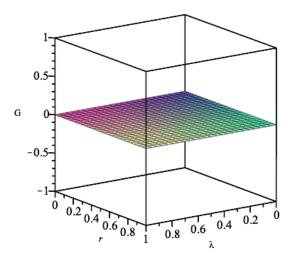
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Gain of when $\alpha = 0.5$.



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Phase Flip Conclusions

When the noise in the spectator is small or large enough (i.e. α is small or large) there is still advantageous gain.

As α approaches 0.5 the gain disappears.

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We now wish to estimate parameter λ with $\hat{\rho}_f = \frac{(1-\lambda)}{2} \text{Tr}[\hat{\rho}_i] \hat{I} + \lambda \hat{\rho}_i$.

$$\hat{\rho}_{i} = \frac{1}{2} \begin{pmatrix} 1 + r_{zi} & r_{xi} - ir_{yi} \\ r_{xi} + ir_{yi} & 1 - r_{zi} \end{pmatrix}. \tag{43}$$

Where $\hat{\rho}_f$ is

$$\hat{\rho}_f = \frac{1}{2} \begin{pmatrix} 1 + \lambda r_{zi} & \lambda (r_{xi} - ir_{yi}) \\ \lambda (r_{xi} + ir_{yi}) & 1 - \lambda r_{zi} \end{pmatrix}$$
(44)

The way the components of these vectors change can be described mathematically as

$$r_{xf} = \lambda r_{xi}$$

$$r_{yf} = \lambda r_{yi}$$

$$r_{zf} = \lambda r_{zi}.$$
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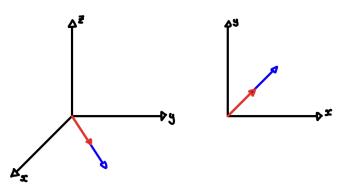
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Visualizing of Depolarizing Channel

We can visualize how this Depolarizing evolution.



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We now examine the simplest depolarizing channel evolution.



The Quantum Fisher Information is

$$H = \frac{r^2}{1 - r^2 \lambda^2}.\tag{46}$$

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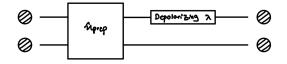
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Depolarizing Channels

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We now examine the next scenario with two particles and one channel depolarizing evolution.

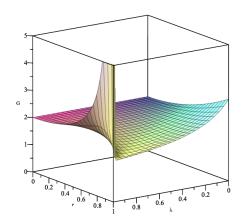


The Quantum Fisher Information is

$$H = \frac{r^2(r^2(1-4\lambda) + r^4\lambda + 2)}{(1+2r\lambda + r^2\lambda)(1-2r\lambda + r^2\lambda)(1-r^2\lambda)}.$$
 (47)

The gain of the last two scenarios is

$$G = \frac{(r^2\lambda^2 - 1)(r^4\lambda - 4r^2\lambda + r^2 + 2)}{(1 + 2r\lambda + r^2\lambda)(1 - 2r\lambda + r^2\lambda)(r^2\lambda - 1)}.$$
 (48)



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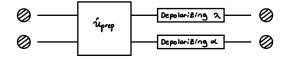
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We now examine a two particle two channel depolarizing evolution.



The Quantum Fisher Information is

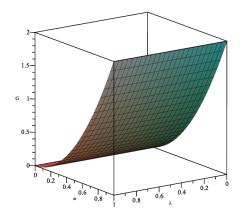
$$H = \frac{r^2 \alpha^2 (\alpha \lambda r^4 - 4r^2 \lambda \alpha + r^2 + 2)}{(1 - r^2 \lambda \alpha)(1 + r\alpha(r+2)\lambda)(1 + r\alpha(r-2)\lambda)}.$$
 (49)

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The gain of the last scenario is

$$G = \frac{\alpha(\alpha\lambda r^4 - 4\lambda r^2\alpha + r^2 + 2)^2(r^2\lambda^2 - 1)}{(\lambda r^2\alpha - 1)(1 + r\alpha(r+2)\lambda)(1 + r\alpha(r-2)\lambda)}.$$
 (50)

With r = 0.01.



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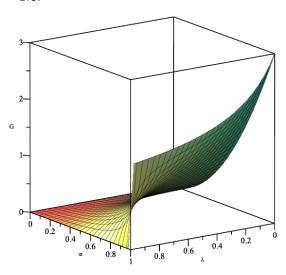
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With r = 1.0.



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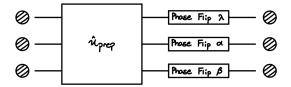
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We now examine a three particle measurement.



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3 Channel Phase Flip Conclusion 47 / 54

The Quantum Fisher Information of this measurement comes out to be

$$H = r^{2}(1 - 2\alpha)^{2}(1 - 2\beta)^{2} \cdot \left(\frac{(r+3)^{2}(3r^{2} + 1)(\eta) + 3(1 - r^{2})^{3}(\xi)}{(\eta)(\xi)}\right)$$
(51)

where
$$\eta = (1 - r^2)^2 - \gamma^2 r^2 (1 - r^2)^2$$
, $\xi = (3r^2 + 1)^2 - r^2 (r + 3)^2 \gamma^2$, and $\gamma = (1 - 2\lambda)(1 - 2\alpha)(1 - 2\beta)$.

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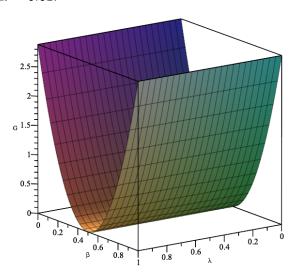
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The gain of the three particle scenario comes out to be

$$G = \frac{1 - r^2 (1 - 2\lambda)^2 (1 - 2\alpha)^2 (1 - 2\beta)^2}{4} \cdot \left(\frac{(r+3)^2 (3r^2 + 1)(\eta) + 3(1 - r^2)^3(\xi)}{(\eta)(\xi)}\right).$$
 (52)

With the next gain scenarios r = 0.01.

With $\alpha = 0.01$.



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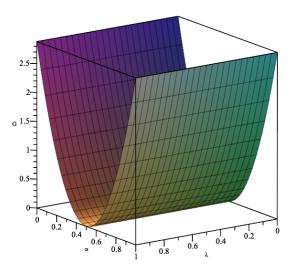
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With $\beta = 0.01$.



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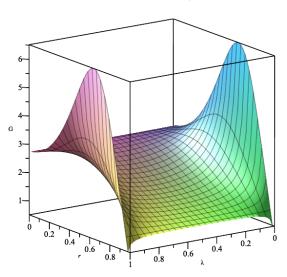
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With this last scenario, $\alpha = 0.01$ and $\beta = 0.01$.



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Conclusion

In the presence of Phase Shifts, Phase Flips, and Depolarizing Channels the Quantum Fisher Information will vary depending on initial conditions.

Pure states are difficult to obtain, so in practicality noisy states are commonly used for estimation.

Typically the use of multiple particles in a measurement yield better outcomes with some exceptions.

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Questions?

Any Questions?

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