

(\*working with the eigen energies of the 1D harmonic oscillator\*)

In[30]:=

$$\text{explicitPartitionfunctionH0}[\text{beta\_}] := \sum_{n=0}^{\infty} \text{Exp}[-\text{beta} (n + 1 / 2)]$$

In[33]:=

FullSimplify[explicitPartitionfunctionH0[beta]]

$$\text{Out[33]} = \frac{1}{2} \text{Csch}\left[\frac{\text{beta}}{2}\right]$$

In[34]:=

explicitEnergyH0[beta\_] := - $\partial_{\text{beta}}$  Log[explicitPartitionfunctionH0[beta]]

In[35]:=

FullSimplify[explicitEnergyH0[beta]]

$$\text{Out[35]} = \frac{1}{2} \text{Coth}\left[\frac{\text{beta}}{2}\right]$$

(\*these are the 1D harmonic oscillator eigen states\*)

$$\text{In[46]} := \text{EigenstatesH0}[n_, x_, \text{beta\_}] := \frac{1}{\sqrt{2^n (n!)} \sqrt{\pi}} \text{Exp}\left[-\frac{x^2}{2}\right] \text{HermiteH}[n, x]$$

(\*to get the density,

we need to square the eigen states and multiply by the wieght factor;

of course, we also need to normalize by the partition function\*)

In[47]:=

$$\text{explicitDensityH0}[x_, \text{beta\_}] := \frac{\sum_{n=0}^{\infty} (\text{EigenstatesH0}[n, x, \text{beta\_}])^2 \text{Exp}[-\text{beta} (n + 1 / 2)]}{\text{FullSimplify}[\text{explicitPartitionfunctionH0}[\text{beta}]]}$$

(\* we can try to see if Mathemtica can evaluate the infinite sum...

turns out that Mathemtica does not know how to do this...\*)

In[48]:=

FullSimplify[explicitDensityH0[x, beta]]

$$\text{Out[48]} = 2 \text{Sinh}\left[\frac{\text{beta}}{2}\right] \sum_{n=0}^{\infty} \frac{2^{-n} e^{-\text{beta} \left(\frac{1}{2}+n\right)-x^2} \text{HermiteH}[n, x]^2}{\sqrt{\pi} n!}$$

(\*if we want to use the explicit expression,  
we need to change the sum from infinity to a sufficiently large value;  
here, I am using 100 (if the temperature is really high,  
then we need to include even more states in the sum)\*)

In[49]:=

```
explicitDensityCutoffH0[x_, beta_] :=
  
$$\frac{\sum_{n=0}^{100} (\text{EigenstatesH0}[n, x, \text{beta}])^2 \text{Exp}[-\text{beta} (n + 1 / 2)]}{\text{FullSimplify}[\text{explicitPartitionfunctionH0}[\text{beta}]]}$$

```

(\*implementing the analytical expressions  
for the 1D harmonic oscillator from the notes\*)

In[1]:=

```
densityH0[x_, beta_] := 
$$\sqrt{\frac{\text{Tanh}[\text{beta} / 2]}{\pi}} \text{Exp}[-x^2 \text{Tanh}[\text{beta} / 2]]$$

```

In[18]:=

```
partitionfunctionH0[beta_] := 
$$\frac{\text{Exp}[-\text{beta} / 2]}{1 - \text{Exp}[-\text{beta}]}$$

```

In[17]:=

```
energyH0[beta_] := 
$$\text{Coth}[\text{beta} / 2] / 2$$

```

In[24]:=

```
Series[energyH0[beta], {beta, 0, 2}]
```

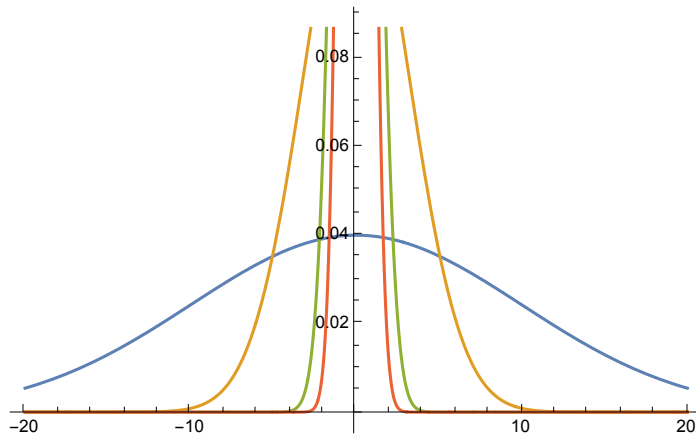
```
Out[24]= 
$$\frac{1}{\text{beta}} + \frac{\text{beta}}{12} + O[\text{beta}]^3$$

```

In[51]:=

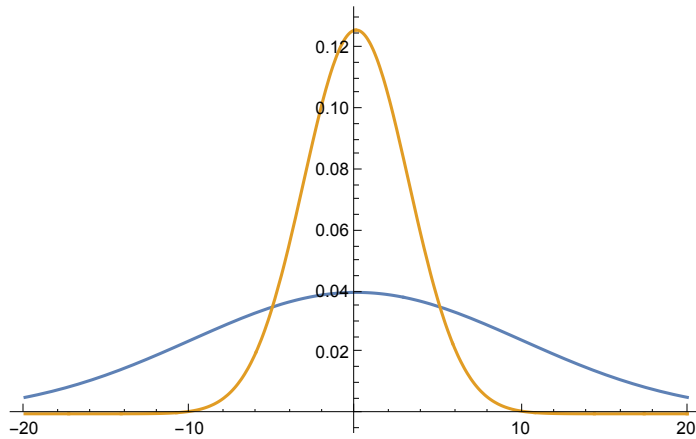
```
Plot[{densityH0[x, 0.01], densityH0[x, 0.1],
      densityH0[x, 1], densityH0[x, 10]}, {x, -20, 20}]
```

Out[51]=



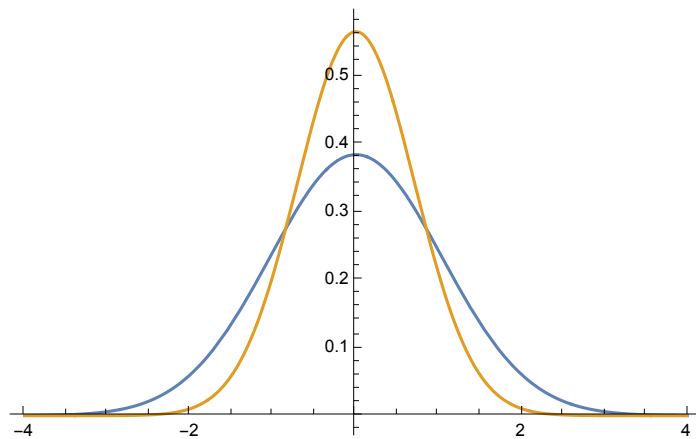
In[52]:= Plot[{densityH0[x, 0.01], densityH0[x, 0.1]}, {x, -20, 20}]

Out[52]=



In[40]:= Plot[{densityH0[x, 1], densityH0[x, 100]}, {x, -4, 4}]

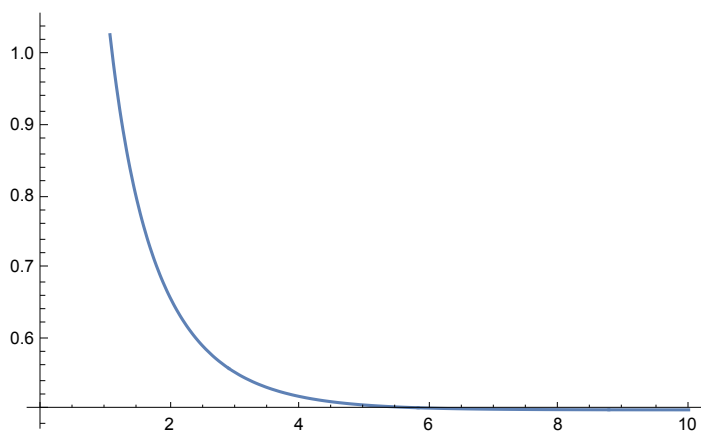
Out[40]=



In[19]:=

**Plot[energyH0[beta], {beta, 0, 10}]**

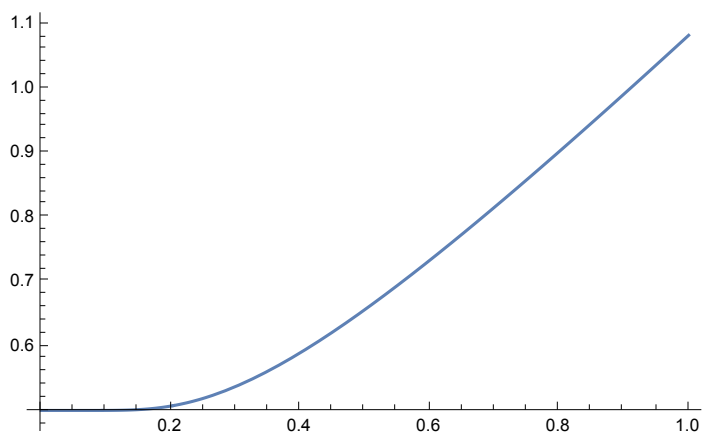
Out[19]=



In[20]:=

**Plot[energyH0[beta] /. beta -> 1 / temp, {temp, 0, 1}]**

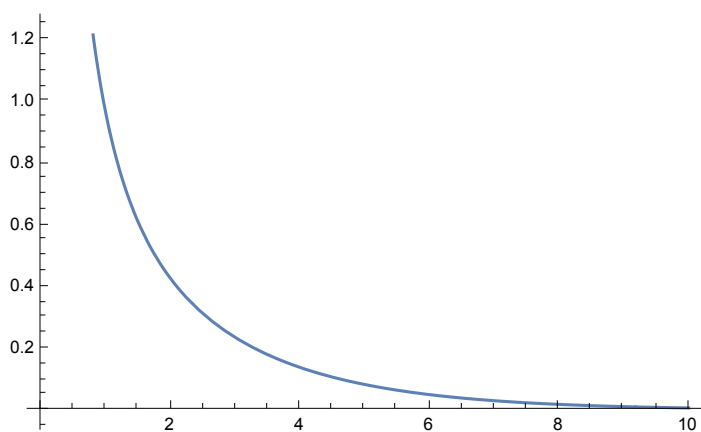
Out[20]=



In[22]:=

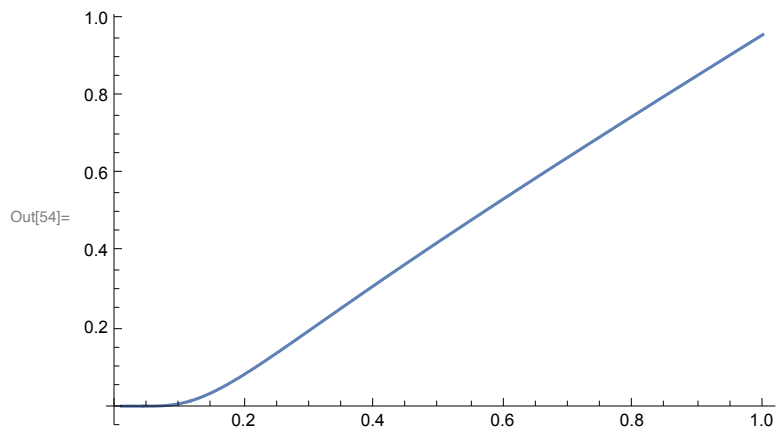
**Plot[partitionfunctionH0[beta], {beta, 0, 10}]**

Out[22]=



In[54]:=

```
Plot[partitionfunctionH0[beta] /. beta -> 1 / temp, {temp, 0.01, 1}]
```



(\*the following graph compares the analytical expression for the density from class with the "brute-force" implementation above (i.e., the implementation where I introduced a finite cutoff for the sum): as expected, the expressions give the same result!\*)

In[55]:=

```
Plot[{explicitDensityCutoffH0[x, 0.9], densityH0[x, 0.9]}, {x, -5, 5}]
```

