

Homework # 1

(A)

①

a)

$$j_1 = 3/2 \implies -j_1 \leq m_1 \leq j_1$$

$$j_2 = 1/2 \implies -j_2 \leq m_2 \leq j_2$$

$$|j_1 - j_2| \leq J \leq j_1 + j_2$$

$$\therefore J_1 = 1, 2, \text{ with } -J \leq m \leq J.$$

Defining

$$|3/2, 1/2; Jn\rangle \equiv |Jn\rangle$$

$$|3/2, 1/2; m_1, m_2\rangle \equiv |m_1, m_2\rangle$$

then

$$|22\rangle = |3/2, 1/2\rangle \quad (1)$$

which is the maximally polarized state.

Ⓟ

To find the other states for $J=2$, we apply J_-

$$J_- |J, m\rangle = \hbar \sqrt{(J+m)(J-m+1)} |J, m-1\rangle$$

\therefore

$$J_- |2, 2\rangle = 2\hbar |2, 1\rangle$$

$$= (J_{1-} + J_{2-}) |3/2, 1/2\rangle$$

$$= \sqrt{3}\hbar |1/2, 1/2\rangle + \hbar |1/2, -1/2\rangle$$

$$\Rightarrow |2, 1\rangle = \frac{\sqrt{3}}{2} |1/2, 1/2\rangle + \frac{1}{2} |3/2, -1/2\rangle \quad (2)$$

Also,

$$J_- |2, 1\rangle = \hbar\sqrt{6} |2, 0\rangle$$

(C)

$$= (J_{1-} + J_{2-}) \left(\frac{\sqrt{3}}{2} \left| \frac{1}{2} \frac{1}{2} \right\rangle + \frac{1}{2} \left| \frac{3}{2} -\frac{1}{2} \right\rangle \right)$$

$$= \hbar \sqrt{3} \left| -\frac{1}{2} \frac{1}{2} \right\rangle + \frac{\sqrt{3}}{2} \hbar \left| \frac{1}{2} -\frac{1}{2} \right\rangle$$

$$+ \hbar \frac{\sqrt{3}}{2} \left| \frac{1}{2} -\frac{1}{2} \right\rangle$$

\therefore

$$|20\rangle = \frac{1}{\sqrt{2}} \left| -\frac{1}{2} \frac{1}{2} \right\rangle + \frac{1}{\sqrt{2}} \left| \frac{1}{2} -\frac{1}{2} \right\rangle \quad (3)$$

$$J_- |20\rangle = \hbar \sqrt{6} |2, -1\rangle$$

$$= (J_{1-} + J_{2-}) \left(\frac{1}{\sqrt{2}} \left| -\frac{1}{2} \frac{1}{2} \right\rangle + \frac{1}{\sqrt{2}} \left| \frac{1}{2} -\frac{1}{2} \right\rangle \right)$$

$$= \hbar \frac{\sqrt{3}}{\sqrt{2}} \left| -\frac{3}{2} \frac{1}{2} \right\rangle + \frac{\hbar}{\sqrt{2}} \left| -\frac{1}{2} -\frac{1}{2} \right\rangle$$

$$+ \hbar \sqrt{2} \left| -\frac{1}{2} -\frac{1}{2} \right\rangle$$

(D)

$$|2, -1\rangle = \frac{1}{2} |-\frac{3}{2}, \frac{1}{2}\rangle + \frac{\sqrt{3}}{2} |-\frac{1}{2}, -\frac{1}{2}\rangle \quad (4)$$

and finally:

$$|2, -2\rangle = |-\frac{3}{2}, -\frac{1}{2}\rangle \quad (5)$$

For the states with $J=1$, we start with the state:

$$|1, 1\rangle$$

Imposing the condition that:

$$\langle 1, 1 | 2, 1 \rangle = 0$$

\therefore

$$|1, 1\rangle = \frac{1}{2} |\frac{1}{2}, \frac{1}{2}\rangle - \frac{\sqrt{3}}{2} |\frac{3}{2}, -\frac{1}{2}\rangle \quad (6)$$

(E)

which is well defined up to a global phase. Applying J_- on this state,

$$\begin{aligned}
 J_- |11\rangle &= \hbar\sqrt{2} |10\rangle \\
 &= (J_{1-} + J_{2-}) \left(\frac{1}{2} |1/2, 1/2\rangle - \frac{\sqrt{3}}{2} |3/2, -1/2\rangle \right) \\
 &= \hbar | -1/2, 1/2\rangle + \frac{\hbar}{2} |1/2, -1/2\rangle - \frac{3\hbar}{2} |1/2, -1/2\rangle
 \end{aligned}$$

\Rightarrow

$$|10\rangle = \frac{1}{\sqrt{2}} | -1/2, 1/2\rangle - \frac{1}{\sqrt{2}} |1/2, -1/2\rangle \quad (7)$$

Finally,

$$|1, -1\rangle = \frac{\sqrt{3}}{2} | -3/2, 1/2\rangle - \frac{1}{2} | -1/2, -1/2\rangle$$

by applying J_- on $|10\rangle$.

⑦

b)

$$4|S, m\rangle = \frac{1}{2} \alpha \vec{S}_1 \cdot \vec{S}_2 |S, m\rangle$$

$$= \frac{\alpha \hbar^2}{2} \left[S(S+1) - \underbrace{\frac{3}{2} \left(\frac{S}{2} \right) - \frac{1}{2} \left(\frac{S}{2} \right)}_{-9/2} \right] |S, m\rangle$$

For $S=1$,

$$\pm_1 = \frac{-5\alpha\hbar^2}{4}$$

For $S=2$,

$$\pm_2 = \frac{3\alpha\hbar^2}{4}$$

From the C.G. coefficients,

$$|21\rangle = \frac{\sqrt{3}}{2} |1/2, 1/2\rangle + \frac{1}{2} |3/2, -1/2\rangle$$

$$|11\rangle = \frac{1}{2} |1/2, 1/2\rangle - \frac{\sqrt{3}}{2} |3/2, -1/2\rangle$$

(6)

Inverting,

$$\begin{cases} |1/2, 1/2\rangle = \frac{\sqrt{3}}{2} |2, 1\rangle + \frac{1}{2} |1, 1\rangle \\ |3/2, -1/2\rangle = \frac{1}{2} |2, 1\rangle - \frac{\sqrt{3}}{2} |1, 1\rangle \end{cases}$$

If $\psi(0) = |1/2, 1/2\rangle$, then:

$$\psi(t) = \frac{\sqrt{3}}{2} e^{-iE_2 t/\hbar} |2, 1\rangle + \frac{1}{2} e^{-iE_1 t/\hbar} |1, 1\rangle$$

\therefore the probability of finding it in the state $|3/2, -1/2\rangle$ is:

$$P = |\langle 3/2, -1/2 | \psi(t) \rangle|^2$$

$$= \frac{3}{16} \left| e^{-iE_2 t/\hbar} - e^{-iE_1 t/\hbar} \right|^2$$

(H)

$$= \frac{3}{8} \left\{ 1 - \cos \left[(\pm_1 - \pm_2) \frac{t}{t_h} \right] \right\}$$

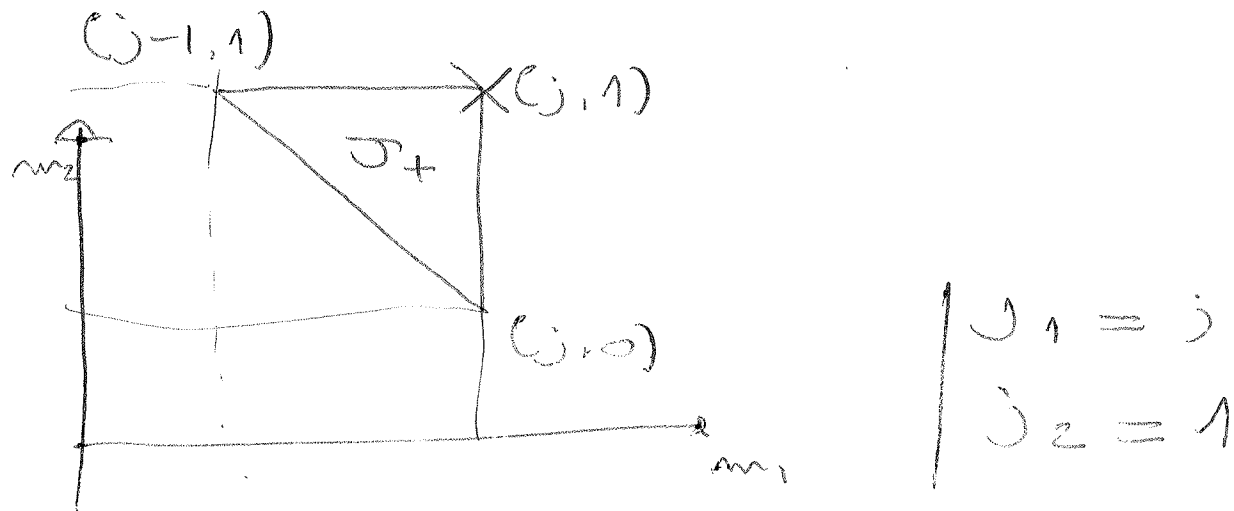
$$= \frac{3}{8} \left[1 - \cos(\alpha t_h t) \right]$$

$$= \frac{3}{4} \sin^2(\alpha t_h t).$$

(1)

(2)

Using the recursion relation J_+



where $\left. \begin{array}{l} m_1 = j \\ m_2 = 1 \end{array} \right\} \begin{array}{l} m_1 + m_2 = m + 1 \\ \therefore m = j \end{array}$

$$\begin{aligned} & \sqrt{(j-m)(j+m+1)} \langle j_1 j_2, m_1 m_2 | j_1 j_2, j m+1 \rangle \\ &= \sqrt{(j_1 - m_1 + 1)(j_1 + m_1)} \langle j_1 j_2, m_1 - 1 m_2 | j m \rangle \\ &+ \sqrt{(j_2 - m_2 + 1)(j_2 + m_2)} \langle j_1 j_2, m_1 m_2 - 1 | j m \rangle \end{aligned}$$

(2)

$$0 = \sqrt{2j} \langle j, j-1, 1 | j, j \rangle \\ + \sqrt{2} \langle j, j, 0 | j, j \rangle$$

Using the normalization relation:

$$\sum_{m_1 m_2} |\langle j_1 j_2 m_1 m_2 | j, j, j \rangle|^2 = 1$$

where $m_1 + m_2 = j = j$

$$\therefore |\langle j, j-1, 1 | j, j, j \rangle|^2 + |\langle j, j, 0 | j, j, j \rangle|^2 = 1$$

$$\Rightarrow \langle j, j, 0 | j, j, j \rangle = \sqrt{\frac{j}{j+1}}$$

①

③

$$Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos\theta = \sqrt{\frac{3}{4\pi}} m_l$$

$$Y_1^{\pm 1} = \mp \sqrt{\frac{3}{8\pi}} \sin\theta e^{\pm i\phi} = \mp \sqrt{\frac{3}{8\pi}} (m_x \pm im_y)$$

a)

$$\text{Mapping} \quad \sqrt{\frac{4\pi}{3}} Y_1^m \rightarrow T_m^{(1)} \\ \hat{m} \rightarrow \vec{p}$$

Then:

$$\begin{cases} \phi_z = T_0^1 \\ \mp \frac{\phi_x \pm i\phi_y}{\sqrt{2}} = T_{\pm 1}^1 \end{cases}$$

$$\therefore \phi_y = -\frac{1}{\sqrt{2}i} (T_{+1}^1 + T_{-1}^1)$$

$$\phi_x = \frac{1}{\sqrt{2}} (-T_{+1}^1 + T_{-1}^1)$$

(2)

Following the Wigner-Eckart theorem,

$$\langle \alpha, j', m' | T_q^{(1)} | \alpha, j, m \rangle$$

$$= \langle j' 1; m_q | j' 1; j' m' \rangle$$

$$\times \underbrace{\frac{\langle \alpha, j' || T^{(1)} || \alpha, j \rangle}{\sqrt{2j+1}}}_{\text{Const.}}$$

The selection rules are:

$$m' = m + q$$

$$|j-j'| \leq 1 \leq j+j'$$

Also, since $T^{(1)}$ is odd under parity,

$$j+j' = 2n+1, n \in \mathbb{Z}, \therefore j' = |j \pm 1|$$

For $\dot{0} = 2$
 $\dot{0} = 1$

$$\begin{aligned}
 \langle \alpha, 2, m | T_0^{(1)} | \alpha, 1, m \rangle \\
 &= \langle 11; m0 | 11; 2m \rangle \times c \\
 &= \begin{cases} c/\sqrt{2}, & m=1 \\ \sqrt{2}c/\sqrt{3}, & m=0 \\ c/\sqrt{2}, & m=-1 \end{cases}
 \end{aligned}$$

Also,

$$\begin{aligned}
 \langle \alpha, 2, m+1 | T_1^{(1)} | \alpha, 1, m \rangle \\
 &= \langle 11; m1 | 11; 2, m+1 \rangle \times c \\
 &= \begin{cases} c, & m=1 \\ c/\sqrt{2}, & m=0 \\ \cancel{c/\sqrt{6}}, & m=-1 \end{cases}
 \end{aligned}$$

(4)

$$\langle \alpha, 2, m-1 | T_{-1}^{(1)} | \alpha, 1, m \rangle$$

$$= \langle 11; m, -1 | 11; 2, m-1 \rangle \times C$$

$$= \begin{cases} C/\sqrt{6}, & m=1 \\ C/\sqrt{2}, & m=0 \\ C, & m=-1 \end{cases}$$

Hence,

$$\langle \alpha, 2, m | P_z | \alpha, 2, m \rangle$$

$$= \begin{cases} 1/\sqrt{2}, & m=1 \\ \sqrt{2}/\sqrt{3}, & m=0 \\ 1/\sqrt{2}, & m=-1 \end{cases} \times \underbrace{\frac{\langle \alpha, 2 | P_z | \alpha, 1 \rangle}{\sqrt{3}}}_C$$

(5)

$$\langle \alpha, 2, m' | \mathcal{P}_x | \alpha, 1, m \rangle$$

$$= - \langle \alpha, 2, m' | \frac{T_{+1}^{(1)} - T_{-1}^{(1)}}{\sqrt{2}} | \alpha, 1, m \rangle$$

$$= \frac{1}{\sqrt{2}} \times \left\{ \begin{array}{l} c/\sqrt{6}, \quad m=1, \quad m'=0 \\ c/\sqrt{2}, \quad m=0, \quad m'=-1 \\ c, \quad m=-1, \quad m'=-2 \\ -c, \quad m=1, \quad m'=2 \\ -c/\sqrt{2}, \quad m=0, \quad m'=1 \\ -c/\sqrt{6}, \quad m=-1, \quad m'=0 \end{array} \right.$$

Finally,

(6)

$$\langle \alpha, 2, m' | P_y | \alpha, 1, m \rangle$$

$$= - \langle \alpha, 2, m' | \frac{T_{+1}^{(1)} + T_{-1}^{(1)}}{\sqrt{2}i} | \alpha, 1, m \rangle$$

$$= - \frac{\hbar}{\sqrt{2}i} \times \left\{ \begin{array}{l} \frac{1}{\sqrt{6}}, m=1, m'=0 \\ \frac{1}{\sqrt{2}}, m=0, m'=-1 \\ 1, m=-1, m'=-2 \\ 1, m=1, m'=2 \\ \frac{1}{\sqrt{2}}, m=0, m'=1 \\ \frac{1}{\sqrt{6}}, m=-1, m'=0 \end{array} \right.$$

$$b) T_{\gamma}^{(2)} = \sum_{\gamma_1, \gamma_2} T_{\gamma_1}^{(1)} T_{\gamma_2}^{(1)} \langle 11; \gamma_1 \gamma_2 | 11; 2\gamma \rangle$$

where $\gamma_1 + \gamma_2 = \gamma$.

$$\begin{aligned} T_0^{(2)} &= T_0^{(1)} T_0^{(1)} \overbrace{\langle 11; 00 | 11, 20 \rangle}^{2/\sqrt{6}} \\ &\quad + 2 T_{-1}^{(1)} T_{+1}^{(1)} \underbrace{\langle 11; 1-1 | 11, 20 \rangle}_{1/\sqrt{6}} \\ &= \phi_z^2 \frac{2}{\sqrt{6}} - \frac{\phi_x^2 + \phi_y^2}{\sqrt{6}} \end{aligned}$$

$$\begin{aligned} T_1^{(2)} &= 2 T_0^{(1)} T_1^{(1)} \underbrace{\langle 11; 01 | 11, 21 \rangle}_{1/\sqrt{2}} \\ &= -\phi_z (\phi_x + i\phi_y) \end{aligned}$$

$$\begin{aligned} T_{-1}^{(2)} &= 2 T_{-1}^{(1)} T_0^{(1)} \underbrace{\langle 11; -10 | 11, 2-1 \rangle}_{1/\sqrt{2}} \\ &= \phi_z (\phi_x - i\phi_y) \end{aligned}$$

$$T_{2}^{(2)} = T_{1}^{(1)} T_{1}^{(1)} \underbrace{\langle 11; 11 | 11, 22 \rangle}_1$$

$$= \frac{1}{2} (p_x + i p_y)^2$$

$$T_{-2}^{(2)} = T_{-1}^{(1)} T_{-1}^{(1)} \underbrace{\langle 11; -1 -1 | 11, 2, -2 \rangle}_1$$

$$= \frac{1}{2} (p_x - i p_y)^2$$