



COLLEGE OF ARTS AND SCIENCES  
HOMER L. DODGE  
DEPARTMENT OF PHYSICS AND ASTRONOMY  
*The* UNIVERSITY *of* OKLAHOMA

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## Math Methods in Physics

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PHYS 5013 HOMEWORK ASSIGNMENT #2

PROBLEMS: {1, 2, 3, 4, 5, 6, 7, 8}

Due: September 13, 2021 By 10:30 AM

STUDENT  
Taylor Larrechea

PROFESSOR  
Dr. Kieran Mullen



**Problem 1:**

Use a Lagrange multiplier approach to find the extremal points  $(x, y)$  of the function

$$f(x, y) = 2x^2 + \frac{1}{2}y^2 - xy$$

subject to the constraint

$$g(x, y) = 4x^2 + y^2 - 4 = 0.$$

$$L(x, y, \lambda) = f(x, y) - \lambda g(x, y)$$

$$L = 2x^2 + \frac{1}{2}y^2 - xy - \lambda(4x^2 + y^2 - 4)$$

$$\frac{\partial L}{\partial x} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} = 0, \quad \frac{\partial L}{\partial y} - \frac{d}{dt} \frac{\partial L}{\partial \dot{y}} = 0$$

$$\frac{\partial L}{\partial x} = 0, \quad \frac{\partial L}{\partial \dot{x}} = 4x - y - 8x\lambda, \quad y = 4x(1 - 2\lambda)$$

$$\frac{\partial L}{\partial y} = 0, \quad \frac{\partial L}{\partial \dot{y}} = y - x - 2y\lambda, \quad x = y(1 - 2\lambda)$$

$$\text{w/ } x = y(1 - 2\lambda) : y = 4 \cdot y(1 - 2\lambda)^2 : 1/4 = (1 - 2\lambda)^2 : \pm \frac{1}{2} = 1 - 2\lambda : \lambda = \frac{1}{2}(1 \pm \frac{1}{2})$$

$$\text{w/ } y = 4x(1 - 2\lambda) : x = 4x(1 - 2\lambda)^2 : 1/4 = (1 - 2\lambda)^2 : \pm \frac{1}{2} = 1 - 2\lambda : \lambda = \frac{1}{2}(1 \pm \frac{1}{2})$$

$$4x^2 + 16x^2(1 - 2\lambda)^2 - 4 = 0 : x^2 + 4x^2(1 - 2\lambda)^2 - 1 = 0 : x = \pm \sqrt{\frac{1}{(1 + 4(1 - 2\lambda)^2)}}$$

$$\lambda = \frac{1}{4}, \quad y = 4x(1 - \frac{1}{2}) = 2x : \lambda = \frac{3}{4}, \quad y = 4x(1 - \frac{3}{2}) = -2x$$

$\lambda = \frac{1}{4}, (1/\sqrt{2}, \sqrt{2})$
$\lambda = \frac{1}{4}, (-1/\sqrt{2}, -\sqrt{2})$
$\lambda = \frac{3}{4}, (1/\sqrt{2}, \sqrt{2})$
$\lambda = \frac{3}{4}, (-1/\sqrt{2}, -\sqrt{2})$

## Problem 1: Review

### Procedure:

- Use the Lagrange multiplier approach:  $h = f - \lambda g$ .
- Use the Euler Lagrange equation to find an equation for  $y$  and for  $x$ .
- Solve for the values of  $\lambda$ .
- Insert values of  $\lambda$  into equation for  $y(x)$ .
- Determine  $(x, y)$  for different values of  $\lambda$ .

### Key Concepts:

- Lagrange multipliers can be used to extremize two functions.
- Values of  $\lambda$  will determine extremal points of both functions.

### Variations:

- $F$  and  $G$  can be changed.
  - This would change the final answer but not the overall procedure.

**Problem 2:**

A bug moves on a surface with cylindrical symmetry. The surface is given by  $z = \ln \rho$  where  $\rho$  is the distance of the surface from the  $z$ -axis. As an intelligent bug, it wants to move the shortest distance between two points.

- (a) Calculate the differential equation of such a curve, between any two points on the surface. (You can solve this using a local constraint, but that's the hard way.)

$$z = \ln(\rho) : ds^2 = d\rho^2 + \rho^2 d\theta^2 + dz^2 : dz = \frac{d\rho}{\rho} : ds^2 = d\rho^2 + \rho^2 d\theta^2 + \frac{d\rho^2}{\rho^2}$$

we wish to minimize this line element:  $ds^2 = d\rho^2 + \rho^2 d\theta^2 + \frac{d\rho^2}{\rho^2}$

$$K[\rho(\alpha)] = \int ds = \int \sqrt{\left(1 + \frac{1}{\rho^2}\right)\rho'^2 + \rho^2} d\alpha$$

$$f = \sqrt{\left(1 + \frac{1}{\rho^2}\right)\rho'^2 + \rho^2}, \quad \alpha = \left(1 + \frac{1}{\rho^2}\right) : \quad \rho' \frac{\partial f}{\partial \rho'} - f = c : \quad \frac{\partial f}{\partial \rho'} = \frac{\alpha \rho'}{f}, \quad \rho' \frac{\partial f}{\partial \rho'} = \frac{\alpha \rho'^2}{f}$$

$$\rho' \frac{\partial f}{\partial \rho'} - f = c : \quad \frac{\alpha \rho'^2}{f} - f = c : \quad \frac{\alpha \rho'^2 - f^2}{f} = c : \quad \frac{\alpha \rho'^2 - \alpha \rho'^2 - \rho^2}{f} = c : \quad \frac{-\rho^2}{f} = c$$

$$\frac{-\rho^2}{c} = f : \quad \frac{\rho^4}{c^2} = \left(1 + \frac{1}{\rho^2}\right)\rho'^2 + \rho^2 \quad \therefore \quad \rho'^2 = \left(\frac{\rho^4}{c^2} - \rho^2\right) / \left(1 + \frac{1}{\rho^2}\right)$$

$$\rho' = \sqrt{\left(\frac{\rho^6}{c^2} - \rho^4\right) / (\rho^2 + 1)} : \quad \frac{dp}{d\alpha} = \sqrt{\left(\frac{\rho^6}{c^2} - \rho^4\right) / (\rho^2 + 1)}$$

$$\boxed{\frac{dp}{d\alpha} = \sqrt{\left(\frac{\rho^6}{c^2} - \rho^4\right) / (\rho^2 + 1)}}$$

- (b) Solve for either the function  $\rho(\theta)$  or  $\theta(\rho)$ , whichever you find easier. (You may need to consult a good integral table).

$$\Theta = \int \sqrt{\frac{(\rho^2 + 1)}{\rho^4 \left(\frac{\rho^2}{c^2} - 1\right)}} d\rho = \frac{1}{c} \left\{ E(\arcsin(c/\rho) - \alpha) - E(\arcsin(c/\rho_0) - \alpha) \right\}$$

$$\boxed{\Theta = \frac{E}{c} \left\{ \arcsin(c/\rho) - \arcsin(c/\rho_0) \right\}}$$

## Problem 2: Review

### Procedure:

- Extremize the functional :  $ds = \sqrt{d\rho^2 + \rho^2 d\theta^2 + dz^2}$ .
- Differentiate  $z$  with respect to  $\rho$  and substitute into  $ds$ .
- Factor our  $d\theta^2$  from the functional.
- Extremize the functional with respect to  $\rho'$  where  $\rho' \equiv d\rho/d\theta$ .
- Use the first derivative Euler Lagrange since the function is not explicitly dependent upon the independent variable.
- Solve for  $d\rho/d\theta$ .
- Rearrange  $d\rho/d\theta$  where  $\theta = [.....]d\rho$  and integrate numerically.

### Key Concepts:

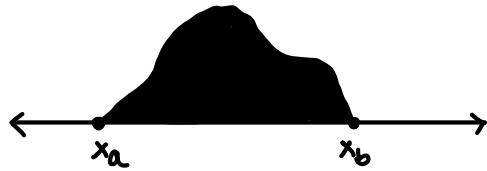
- Extremizing the  $ds$  functional will create a shortest path for our bug.
- Use cylindrical co-ordinates because they are the easiest to work with in this problem.
- Since there is now explicit independent variable dependence, we can use the first order Euler Lagrange equation.

### Variations:

- The surface can change thus changing the functional and the  $ds$  element.
  - If the function changes, then we would use the same procedure but with a different function.

**Problem 3:**

Consider a curve  $y(x)$  where  $y(x_a) = 0$ , and the total length of the curve is  $L$ . Find the curve that gives you the maximum area enclosed between  $y(x)$  and the  $x$ -axis. (Assume that  $L < \pi(x_b - x_a)/2$ ). There is a simple reason why the problem changes when  $L$  is too larger. Can you see why?



$$f(x, y) = \frac{1}{2} \int_{x_a}^{x_b} (xy' - x'y) dt : g(x, y) = \int_{x_a}^{x_b} \sqrt{x'^2 + y'^2} dt \text{ w/ } y(x_a) = y(x_b) = 0$$

$$L(x, y, \lambda) = f(x, y) - \lambda g(x, y) : L(x, y, \lambda) = \frac{1}{2}(xy' - x'y) - \lambda(\sqrt{x'^2 + y'^2})$$

$$\frac{d}{dt} \frac{\partial L}{\partial y'} - \frac{\partial L}{\partial y} : \frac{\partial L}{\partial y} = \frac{x}{2} - \frac{\lambda \dot{y}}{\sqrt{\dot{x}^2 + \dot{y}^2}}, \quad \frac{d}{dt} \frac{\partial L}{\partial y'} = \frac{\dot{x}}{2} - \frac{d}{dt} \frac{\lambda \dot{y}}{\sqrt{\dot{x}^2 + \dot{y}^2}}, \quad \frac{\partial L}{\partial y} = -\frac{\dot{x}}{2}$$

$$\frac{\dot{x}}{2} - \frac{d}{dt} \frac{\lambda \dot{y}}{\sqrt{\dot{x}^2 + \dot{y}^2}} + \frac{\dot{x}}{2} = 0, \quad \dot{x} = \lambda \frac{d}{dt} \frac{\dot{y}}{\sqrt{\dot{x}^2 + \dot{y}^2}}, \quad \dot{y} = \lambda \frac{d}{dt} \frac{\dot{x}}{\sqrt{\dot{x}^2 + \dot{y}^2}}$$

$$(x - x_0)^2 + (y - y_0)^2 = \lambda^2 \frac{\dot{x}^2 + \dot{y}^2}{\dot{x}^2 + \dot{y}^2} : (x - x_0)^2 + (y - y_0)^2 = \lambda^2$$

$$(x - x_0)^2 + (y - y_0)^2 = \lambda^2$$

If  $L$  becomes too large, the function becomes double valued while the radius of the semi circle will explode to infinity.

## Problem 3: Review

### Procedure:

- Set  $f$  to be the maximum area function :  $\int_{x_a}^{x_b} (xy' - x'y) dt$ .
- Set  $g$  to be the  $ds$  element :  $\sqrt{x'^2 + y'^2} dt$ .
- Use the method of Lagrange multipliers :  $h = f - \lambda g$ .
- Use the Euler Lagrange equation on the Lagrange multiplier equation.
- Use the professors trick to show

$$\dot{x} = \lambda \frac{d}{dt} \frac{\dot{y}}{\sqrt{\dot{x}^2 + \dot{y}^2}} \quad \text{and} \quad \dot{y} = \lambda \frac{d}{dt} \frac{\dot{x}}{\sqrt{\dot{x}^2 + \dot{y}^2}}.$$

- Reduce and solve.

### Key Concepts:

- We use the method of Lagrange multipliers to extremize this surface that is to be constrained under the function.
- Use the dependent  $f$  and  $g$  to create a Lagrange multiplier equation.
- Use the Euler Lagrange equation to extremize the functional.
- The professors trick allows us to bypass difficult differentiation.

### Variations:

- The end points of the function,  $x_a, x_b$  can yield different values when plugged into the equation for  $y(x)$ .
  - We would use the same procedure up to this point.
- The length of  $L$  can change.
  - This would in turn change boundary conditions.

**Problem 4:**

Often constrained problems yield difficult integral problems or differential equations to solve. In each case below, apply the calculus of variations and derive a final expression for the curve specified as an integral, that is:

$$x - x_a = \int_{y_a}^y F(y) dy$$

where  $F(y)$  is a known function of  $y$ . (This is called *reducing the problem to quadrature*.)

- (a) Determine the curve  $y(x)$  that connects the points  $(x_a, y_a)$  and  $(x_b, y_b)$  on a path of fixed length but along which a frictionless particle would slide with the shortest time. (That is, similar to the brachistochrone but on a path of fixed length).

$$f(x, y) = \frac{\sqrt{1+x'^2}}{\sqrt{2gy}} = \frac{ds}{\sqrt{V}}, \quad g(x, y) = \sqrt{1+x'^2} : \quad \dot{L}(x, y) = \frac{\sqrt{1+x'^2}}{\sqrt{2gy}} - \lambda \sqrt{1+x'^2}$$

$$\int \frac{\sqrt{1+x'^2}}{\sqrt{2gy}} - \lambda \sqrt{1+x'^2} dy$$

$$\frac{\partial L}{\partial x} - \frac{d}{dy} \frac{\partial L}{\partial x'} = 0 : \quad \frac{\partial L}{\partial x} = 0, \quad \frac{\partial L}{\partial x'} = \frac{x'}{\sqrt{2gy} \cdot \sqrt{1+x'^2}} - \lambda \frac{x'}{\sqrt{1+x'^2}}$$

$$\frac{d}{dy} \left( \frac{x'}{\sqrt{2gy} \cdot \sqrt{1+x'^2}} - \lambda \frac{x'}{\sqrt{1+x'^2}} \right) = 0 : \quad \frac{x'}{\sqrt{2gy} \cdot \sqrt{1+x'^2}} - \lambda \frac{x'}{\sqrt{1+x'^2}} \rightarrow \text{const}$$

We can argue that the derivative of a const is 0  $\therefore$  the terms in parentheses is a const

$$\frac{x'}{\sqrt{2gy} \cdot \sqrt{1+x'^2}} - \lambda \frac{x'}{\sqrt{1+x'^2}} = C : \quad \frac{x'}{\sqrt{1+x'^2}} \left( \frac{1}{\sqrt{2gy}} - \lambda \right) = C : \quad \frac{1}{\sqrt{2gy}} - \lambda \rightarrow \alpha$$

$$\frac{x'^2}{1+x'^2} \cdot \alpha^2 = C^2 : \quad x'^2 \cdot \alpha^2 = C^2 + C^2 x'^2 : \quad x'^2 \alpha^2 - x'^2 C^2 = C^2 : \quad x'^2 = \frac{C^2}{\alpha^2 - C^2} : \quad x' = \sqrt{\frac{C^2}{\alpha^2 - C^2}}$$

$$\frac{dx}{dy} = \sqrt{\frac{C^2}{\alpha^2 - C^2}} : \quad \int dx = \int \sqrt{\frac{C^2}{\alpha^2 - C^2}} dy : \quad x_b - x_a = \int_{y_a}^{y_b} \sqrt{\frac{C^2}{\alpha^2 - C^2}} dy$$

$$\alpha = \frac{1}{\sqrt{2gy}} - \lambda$$

$$x_b - x_a = \int_{y_a}^{y_b} \sqrt{\frac{C^2}{\alpha^2 - C^2}} dy$$

- (b) Determine the shortest length curve  $y(x)$  that connects the points  $(x_a, y_a)$  and  $(x_b, y_b)$  along which a frictionless particle would slide with a fixed time.

$$f(x, y) = \sqrt{1+x'^2}, \quad g(x, y) = \frac{\sqrt{1+x'^2}}{\sqrt{2gy}} = \frac{ds}{\sqrt{V}} : \quad \dot{L}(x, y) = \sqrt{1+x'^2} - \lambda \frac{\sqrt{1+x'^2}}{\sqrt{2gy}}$$

$$\int \sqrt{1+x'^2} - \lambda \frac{\sqrt{1+x'^2}}{\sqrt{2gy}} dy$$

$$\frac{\partial \dot{L}}{\partial x} - \frac{d}{dy} \frac{\partial \dot{L}}{\partial x'} = 0 : \quad \frac{\partial \dot{L}}{\partial x} = 0, \quad \frac{\partial \dot{L}}{\partial x'} = \frac{x'}{\sqrt{1+x'^2}} - \frac{\lambda x'}{\sqrt{2gy} \sqrt{1+x'^2}} : \quad \frac{d}{dy} \left( \frac{x'}{\sqrt{1+x'^2}} - \frac{\lambda x'}{\sqrt{2gy} \sqrt{1+x'^2}} \right)$$

## Problem 4: Continued

$$\frac{d}{dy} \left( \frac{x'}{\sqrt{1+x'^2}} - \frac{\lambda x'}{\sqrt{2gy}\sqrt{1+x'^2}} \right) = 0 : \frac{x'}{\sqrt{1+x'^2}} - \frac{\lambda x'}{\sqrt{2gy}\sqrt{1+x'^2}} \rightarrow \text{const}$$

We can argue that the derivative of a const is 0  $\therefore$  the terms in parentheses is a const

$$\frac{x'}{\sqrt{1+x'^2}} - \frac{\lambda x'}{\sqrt{2gy}\sqrt{1+x'^2}} = c : \frac{x'}{\sqrt{1+x'^2}} \left( 1 - \frac{\lambda}{\sqrt{2gy}} \right) = c : 1 - \frac{\lambda}{\sqrt{2gy}} \rightarrow \alpha$$

$$\frac{x'^2}{1+x'^2} \alpha^2 = c^2 : x'^2 \alpha^2 = c^2 + c^2 x'^2 : x'^2 \alpha^2 - c^2 x'^2 = c^2 : x'^2 (\alpha^2 - c^2) = c^2$$

$$x'^2 = \frac{c^2}{\alpha^2 - c^2} : x' = \sqrt{\frac{c^2}{\alpha^2 - c^2}} : \frac{dx}{dy} = \sqrt{\frac{c^2}{\alpha^2 - c^2}} : \int_{x_0}^{x_b} dx = \int_{y_a}^{y_b} \sqrt{\frac{c^2}{\alpha^2 - c^2}} dy$$

$$\alpha = 1 - \frac{\lambda}{\sqrt{2gy}}$$

$$x_b - x_a = \int_{y_a}^{y_b} \sqrt{\frac{c^2}{\alpha^2 - c^2}} dy$$

- (c) A solid object can be defined by rotating a curve  $\rho(z)$  about the  $z$ -axis to determine a surface of revolution. The top and bottom surfaces of the object would be disks parallel to the  $x - y$  plane. Determine the curve  $\rho(z)$  connecting the points  $\rho(z_a) = \rho_a$  to  $\rho(z_b) = \rho_b$  so that it minimizes the surface area of the object but has a fixed volume and moment of inertia. Assume that the density of the object has a fixed mass/ volume, denoted by the constant  $\alpha$ .

$$V = \pi \int \rho^2 dz : I = \frac{MR^2}{2}, I = \frac{\alpha \pi}{2} \int \rho^4 dz : A = 2\pi \int \rho \sqrt{1+\rho'^2} dz$$

$$\mathcal{L}(\rho, \rho', z, z') = \rho \sqrt{1+\rho'^2} - \lambda_1 \rho^2 - \lambda_2 \rho^4$$

$$\rho' \frac{\partial \mathcal{L}}{\partial \rho'} - \mathcal{L} = c : \frac{\partial \mathcal{L}}{\partial \rho'} = \frac{\rho \rho'}{\sqrt{1+\rho'^2}} : \rho' \frac{\partial \mathcal{L}}{\partial \rho'} = \frac{\rho \rho'^2}{\sqrt{1+\rho'^2}} \rightarrow \alpha$$

$$\rho' \frac{\partial \mathcal{L}}{\partial \rho'} - \mathcal{L} = \frac{\rho \rho'^2}{\sqrt{1+\rho'^2}} - \rho \sqrt{1+\rho'^2} + \lambda_1 \rho^2 + \lambda_2 \rho^4 = c : \frac{-\rho + (\lambda_1 \rho^2 + \lambda_2 \rho^4)\alpha}{\alpha} = \alpha c$$

using Mathematica

$$\frac{dp}{dz} = \pm \sqrt{\frac{(-((c-p(p(\lambda_2 p^2 + \lambda_1) - 1)))((c-p(1 + \lambda_1 p + m p^3)))}}{(c-p^2(\lambda_1 + \lambda_2 p^2))^2}}$$

## Problem 4: Review

### Procedure:

- Use the Lagrange multiplier approach to extremize a functional which is subject to a constraint.
- We want to extremize time  $\therefore f = t = ds/v$  with  $ds = \sqrt{1 + x'^2} dy$  and  $v = \sqrt{2gy}$ .
- Use  $g = \sqrt{1 + x'^2}$ , plug  $f$  and  $g$  into the Lagrange multiplier equation,  $h = f - \lambda g$ .
- Extremize  $h$  with respect to the correct independent variables.
- Because  $\partial L/\partial x = 0$ ,  $d/dt(\partial L/\partial x')$  yields  $\partial L/\partial x' = c$ .
- Solve for  $x'$  and solve the differential equation.
- Reverse  $f$  and  $g$  for part (b), then repeat steps once again.
- Minimize the surface area  $\therefore f$  represents the surface area, constraints are volume and moment of inertia.
- This function will then have two constraints  $\therefore$  two  $\lambda$ 's.
- Use the first derivative trick and extremize the functional.
- Solve for the differential equation.

### Key Concepts:

- With Lagrange multipliers,  $f \rightarrow$  surface, distance, time etc. to be extremized. Then this means the constraints are then defined as  $g$ .
- $\partial L/\partial q_i = \text{constant}$ .
- Situations with multiple constraints will have multiple  $\lambda$ 's.

### Variations:

- The surface can be changed.
  - This would change some small details but not the overall procedure.
- The object rolling in (c) can be different.
  - This then changes the constraints of the functional.

**Problem 5:**

Consider the variational problem:

$$I = \int (f(x, y_1, y'_1, y_2, y'_2) - \lambda(x)g(y_1, y_2)) dx$$

where  $f$  is the optimization function and  $g$  is a local constraint.

- (a) If  $f$  is just a square root such as  $\sqrt{1+y'^2}$ , it would be easier to minimize  $f^2$ . Does this give you the same curve?

$$\mathcal{L}(x, y_1, y_2, y'_1, y'_2) = f(x, y_1, y'_1, y_2, y'_2) - \lambda(x)g(y_1, y_2)$$

$$\text{w/ } f = \sqrt{1+y'^2} : \frac{\partial \mathcal{L}}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial \mathcal{L}}{\partial y'_{1,2}} = \frac{\partial f}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial F}{\partial y'_{1,2}} - \lambda(x) \frac{\partial g}{\partial y_{1,2}} = 0$$

$$\frac{\partial f}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial F}{\partial y'_{1,2}} - \lambda(x) \frac{\partial g}{\partial y_{1,2}} = 0 \longrightarrow F$$

$$\text{w/ } f^2 = 1+y'^2 : \frac{\partial \mathcal{L}}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial \mathcal{L}}{\partial y'_{1,2}} = \frac{\partial f^2}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial f^2}{\partial y'_{1,2}} - \lambda(x) \frac{\partial g}{\partial y_{1,2}} = 0$$

$$\frac{\partial f \cdot \partial f}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial f \cdot \partial f}{\partial y'_{1,2}} - \frac{d}{dx} \frac{\partial f \cdot \partial f}{\partial y'_{1,2}} - \lambda(x) \frac{\partial g}{\partial y_{1,2}} = 0 \longrightarrow G$$

If we compare  $F$  w/  $G$  we see that there are extra terms in  $G$  that will not simplify down to  $F$ . Therefore these will not yield the same curve.

- (b) If  $g$  is just a square root, such as  $\sqrt{y_1^2 + y_2^2}$ , it would be easier to replace  $g$  with  $g^2$ . Does this give you the same curve?

$$\mathcal{L}(x, y_1, y_2, y'_1, y'_2) = f(x, y_1, y'_1, y_2, y'_2) - \lambda(x)g(y_1, y_2)$$

$$\text{w/ } g = \sqrt{y_1^2 + y_2^2} : \frac{\partial \mathcal{L}}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial \mathcal{L}}{\partial y'_{1,2}} = \frac{\partial f}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial F}{\partial y'_{1,2}} - \lambda(x) \frac{\partial g}{\partial y_{1,2}} = 0$$

$$\frac{\partial f}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial F}{\partial y'_{1,2}} - \lambda(x) \frac{\partial g}{\partial y_{1,2}} = 0 \longrightarrow F$$

$$\text{w/ } g^2 = y_1^2 + y_2^2 : \frac{\partial \mathcal{L}}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial \mathcal{L}}{\partial y'_{1,2}} = \frac{\partial f}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial f}{\partial y'_{1,2}} - \lambda(x) \cdot \partial g \cdot \frac{\partial g}{\partial y_{1,2}} = 0$$

$$\frac{\partial f}{\partial y_{1,2}} - \frac{d}{dx} \frac{\partial f}{\partial y'_{1,2}} - \lambda(x) \cdot \partial g \cdot \frac{\partial g}{\partial y_{1,2}} = 0 \longrightarrow G$$

If we compare  $F$  w/  $G$  we can assume that  $\lambda(x) \cdot \partial g \rightarrow \lambda(x)$  because  $\lambda(x)$  is an arbitrary function and thus we can say  $F$  and  $G$  will yield the same curve.

## Problem 5: Review

### Procedure:

- Use the Euler Lagrange equation on the functional, once with  $F = f$  and once with  $F = f^2$ .
- Proceed to show that these do not yield the same curve.
- Repeat with  $G = g$  and  $G = g^2$ .
- Proceed to show that there is no difference between the two curves.

### Key Concepts:

- When we square the function we are extremizing, this will yield a different curve.
- When we square our constraint equation, this will not yield a different curve.

### Variations:

- The function being extremized can change.
  - This should still yield a different curve.
- The constraint equation can change.
  - This shouldn't affect the overall curve.

**Problem 6:**

Consider a 3D crystal with a surface defined by a function  $z(x, y)$ . We seek to minimize the total surface free energy:

$$F[z] = \int dx dy \left\{ \alpha(x, y) \sqrt{1 + z_x^2 + z_y^2} \right\}$$

where  $\alpha(z_x, z_y)$  is the direction dependent surface tension, and we have introduced the compact but confusing notation,  $z_x \equiv \partial z / \partial x$  and  $z_y \equiv \partial z / \partial y$ . We wish to enforce the constraint of a constant volume to the crystal where the volume is given by:

$$V[z] = \int dx dy z$$

via a Lagrange multiplier,  $\lambda$ .

- (a) What is the Euler-Lagrange equation for the system?

$$\mathcal{L}(x, y, z, \lambda) = F[z] - \lambda V[z] = \alpha(x, y) \sqrt{1 + z_x^2 + z_y^2} - \lambda z$$

$$\frac{\partial \mathcal{L}}{\partial z} - \frac{d}{dx} \frac{\partial \mathcal{L}}{\partial z_x} - \frac{d}{dy} \frac{\partial \mathcal{L}}{\partial z_y} = 0$$

$$\frac{\partial \mathcal{L}}{\partial z} = -\lambda, \quad \frac{\partial \mathcal{L}}{\partial z_x} = \frac{\alpha(x, y) z_x}{(1 + z_x^2 + z_y^2)^{1/2}}, \quad \frac{d}{dx} \frac{\partial \mathcal{L}}{\partial z_x} = \frac{d}{dx} \frac{\alpha(x, y) z_x}{(1 + z_x^2 + z_y^2)^{1/2}} : \quad \frac{d}{dx} \frac{\alpha(x, y) z_x}{(1 + z_x^2 + z_y^2)^{1/2}} = -\lambda$$

$$\frac{\partial \mathcal{L}}{\partial z_y} = \frac{\alpha(x, y) z_y}{(1 + z_x^2 + z_y^2)^{1/2}}, \quad \frac{d}{dy} \frac{\partial \mathcal{L}}{\partial z_y} = \frac{d}{dy} \frac{\alpha(x, y) z_y}{(1 + z_x^2 + z_y^2)^{1/2}} : \quad \frac{d}{dy} \frac{\alpha(x, y) z_y}{(1 + z_x^2 + z_y^2)^{1/2}} = -\lambda$$

$$-\lambda - \frac{d}{dx} \frac{\alpha(x, y) z_x}{(1 + z_x^2 + z_y^2)^{1/2}} - \frac{d}{dy} \frac{\alpha(x, y) z_y}{(1 + z_x^2 + z_y^2)^{1/2}} = 0 : - \frac{d}{dx} \frac{\alpha(x, y) z_x}{(1 + z_x^2 + z_y^2)^{1/2}} - \frac{d}{dy} \frac{\alpha(x, y) z_y}{(1 + z_x^2 + z_y^2)^{1/2}} = \lambda$$

$$-\frac{d}{dx} \frac{\alpha(x, y) z_x}{(1 + z_x^2 + z_y^2)^{1/2}} - \frac{d}{dy} \frac{\alpha(x, y) z_y}{(1 + z_x^2 + z_y^2)^{1/2}} = \lambda$$

- (b) Show that in the isotropic case, where  $\alpha(z_x, z_y) = \alpha_0$ , that the shape that minimizes the surface energy for fixed volume is a sphere. You may do this either by direct substitution to verify the solution

$$z(x, y) = \sqrt{R^2 - x^2 - y^2}$$

or by changing to polar co-ordinates  $z(\rho, \phi)$  in the original functional, and invoking cylindrical symmetry, so that your 2D problem becomes effectively a one dimensional problem for  $z(\rho)$ . In either case, determine the relation between  $\lambda$  and the radius of the sphere.

$$-\frac{d}{dx} \frac{\alpha_0 z_x}{(1 + z_x^2 + z_y^2)^{1/2}} - \frac{d}{dy} \frac{\alpha_0 z_y}{(1 + z_x^2 + z_y^2)^{1/2}} = \lambda$$

$$z_x = \frac{-x}{\sqrt{R^2 - x^2 - y^2}}, \quad z_y = \frac{-y}{\sqrt{R^2 - x^2 - y^2}} : z_x^2 = \frac{x^2}{R^2 - x^2 - y^2}, \quad z_y^2 = \frac{y^2}{R^2 - x^2 - y^2}$$

$$z_x^2 + z_y^2 = \frac{x^2 + y^2}{R^2 - x^2 - y^2} : 1 + z_x^2 + z_y^2 = \frac{R^2 - x^2 - y^2 + x^2 + y^2}{R^2 - x^2 - y^2} = \frac{R^2}{R^2 - x^2 - y^2}$$

## Problem 6: Continued

$$\frac{\frac{Zx}{(1+2x^2+2y^2)^{1/2}} = \frac{-x}{\sqrt{R^2-x^2-y^2}} / \frac{R}{\sqrt{R^2-x^2-y^2}} = -\frac{x}{R} : \frac{\frac{Zy}{(1+2x^2+2y^2)^{1/2}} = \frac{-y}{R}}$$

$$-\frac{d}{dx} \frac{-a_0 x}{R} - \frac{d}{dy} \frac{-a_0 y}{R} = \lambda : \frac{a_0}{R} + \frac{a_0}{R} = \lambda : \frac{2a_0}{R} = \lambda$$

$$R = \frac{2a_0}{\lambda} \quad \text{w/ } R = \sqrt{x^2+y^2+z^2} \longrightarrow \text{This is a sphere!}$$

$$R = \frac{2a_0}{\lambda}$$

## Problem 6: Review

### Procedure:

- Define a functional with constraints → Lagrange multiplier method.
- Add extra terms into the Euler Lagrange equation because of multiple independent variables.
- Use Euler Lagrange equations on the functional to create the Equations of Motion.
- Solve for the constraint in terms of other variables.
- Proceed to use the  $z(x, y)$  equation in the Equation of Motion.
- Show that the equation will simplify down to a sphere.

### Key Concepts:

- Two independent variables just get added together in the Euler Lagrange equation.
- Multiple dependent variables will cause multiple Euler Lagrange equations.
- Constraints can be dependent upon multiple independent variables.
- Use Lagrange multipliers to extremize functions with constraints.

### Variations:

- The constraint equation can change.
  - This can add independent variables to the system, or change the function.
- The function to be extremized can be slightly changed.
  - Multiple dependent variables can be added → multiple Euler Lagrange equations.

**Problem 7:****Variational Derivation of Poisson Equation:**

The energy in an electrostatic field in the presence of charges is:

$$E = \int \left\{ \frac{1}{8\pi} (\nabla \phi(\vec{r}))^2 - \rho(\vec{r})\phi(\vec{r}) \right\} d\vec{r}$$

Show that the minimum energy configuration of the potential  $\phi(\vec{r})$  satisfies the equation:

$$\vec{\nabla}^2(\vec{r}) = -4\pi\rho(\vec{r}).$$

$$\begin{aligned} f &= \frac{1}{8\pi} (\nabla \phi(\vec{r}))^2 - \rho(\vec{r})\phi(\vec{r}) : \frac{\partial f}{\partial \phi(\vec{r})} - \frac{d}{dr} \frac{\partial f}{\partial \nabla \phi(\vec{r})} = 0 \\ \frac{\partial f}{\partial \nabla \phi(\vec{r})} &= \frac{1}{4\pi} \nabla \phi(\vec{r}), \quad \frac{d}{dr} \cdot \frac{\partial f}{\partial \nabla \phi(\vec{r})} = \frac{\nabla^2 \phi(\vec{r})}{4\pi}, \quad \frac{\partial f}{\partial \phi(\vec{r})} = -\rho(\vec{r}) \\ -\rho(\vec{r}) - \frac{\nabla^2 \phi(\vec{r})}{4\pi} &= 0 : \frac{\nabla^2 \phi(\vec{r})}{4\pi} = -\rho(\vec{r}) : \nabla^2 \phi(\vec{r}) = -4\pi\rho(\vec{r}) \end{aligned}$$

$\boxed{\nabla^2 \phi(\vec{r}) = -4\pi\rho(\vec{r})}$

Thus the *minimal energy* configuration for the field is also the one given by the Poisson equation.

The above problem is rather simple. I am assigning it to you for the notes below, so that you see *why* it works. It is a common technique in quantum field theory, so it's worth knowing the background.

**Background:** The expression for the energy above avoids counting the "self-energy", so that a charge does not feel the force of the electric field it creates. To see this assume that we have a set of point charges  $q_i$  each with potential  $\phi_i(\vec{r})$ . Then the total potential is the sum of the individual contributions:

$$\phi(\vec{r}) = \sum_i \phi_i(\vec{r}).$$

The energy of the charges (in some units) can be written as

$$E = \frac{1}{8\pi} \int (\vec{\mathcal{E}}(\vec{r}))^2 d\vec{r}$$

(including the self-energy (Jackson, *Classical Electrodynamics*, Pg. 46)) which we can write as

$$2E = \int \left\{ \frac{1}{8\pi} \sum_{i \neq j} (\nabla \phi_i(\vec{r}) \cdot \nabla \phi_j(\vec{r})) \right\} d\vec{r}$$

where the factor of 2 comes from double counting in the sum over  $i$  and  $j$ . The energy can also be written as:

$$E = \int \left\{ \sum_{i \neq j} q_i \delta(\vec{r} - \vec{r}_i) \phi_j(\vec{r}) \right\} d\vec{r}$$

where  $\delta(\vec{r} - \vec{r}_i)$  is a function sharply peaked at  $\vec{r}_i$ , the location of the point charge, and there is no double counting. If we remove the restriction  $i \neq j$ , then both expressions pick up a self-energy term. However, if we subtract the two expressions, these cancel out, and we are left with simply the total energy:

**Problem 7: Continued**

$$\int \left\{ \frac{1}{8\pi} \sum_{ij} (\nabla \phi_i(\vec{r}) \cdot \nabla \phi_j(\vec{r})) - q_i \delta(\vec{r} - \vec{r}_i) \phi_j(\vec{r}) \right\} = 2E + E_s - E - E_s = E.$$

If we take the continuum limit and replace the discrete distribution by a continuous one, then we obtain the expression at the top of the page.

This “background” discussion does not contain any questions. I am including it merely because the above formulation is routinely invoked in some field theory courses without explanation.

## Problem 7: Review

### Procedure:

- Extremize the functional with the Euler Lagrange equation.
- Independent variable in this case will be  $\vec{r}$ .
- Dependent variable in this case is  $\vec{\nabla}(\phi(\vec{r}))$ .

### Key Concepts:

- Functions can be extremized with the Euler Lagrange equation.

### Variations:

- This problem cannot be changed have it remain the same type of problem since it is a derivation.

**Problem 8:**

Obtain access to MATHEMATICA on some computer. As proof of having access, if  $N$  is your OU ID number, calculate the prime factors of  $N^2 + 1$  by typing in the command:

**FactorInteger**[ $N * N + 1$ ]

and then hold down the “shift” key and the “return” key at the same time to execute the command.

```
In[1]:= FactorInteger[113 556 911 * 113 556 911 + 1]
Out[1]= {{2, 1}, {2113, 1}, {265 621, 1}, {11 487 757, 1}}
```

## Problem 8: Review

**Procedure:**

- Download MATHEMATICA and follow the instructions in the problem statement.

**Key Concepts:**

- MATHEMATICA can be used to find the prime factors of an integer.

**Variations:**

- This problem is too simple to be changed.