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## Math Methods in Physics

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PHYS 5013 HOMEWORK ASSIGNMENT #5

PROBLEMS: {1, 2, 3, 4}

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STUDENT

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PROFESSOR

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# Problem 1:

Consider the parameterization of a position vector in 3D:

$$\vec{r}(\eta, \phi, z) = \cosh \eta \cos \phi \hat{i} + \sinh \eta \sin \phi \hat{j} + z \hat{k}$$

(a) Demonstrate that  $\eta, \phi, z$  constitute a set of orthogonal co-ordinates.

$$\begin{aligned} \frac{\partial \vec{r}}{\partial \eta} \cdot \frac{\partial \vec{r}}{\partial \phi} &= (\sinh(\eta) \cos(\phi) \hat{i} + \cosh(\eta) \sin(\phi) \hat{j}) \cdot (-\cosh(\eta) \sin(\phi) \hat{i} + \sinh(\eta) \cos(\phi) \hat{j}) \\ &= -\sinh(\eta) \cosh(\eta) \cos(\phi) \sin(\phi) + \sinh(\eta) \cosh(\eta) \cos(\phi) \sin(\phi) = 0 \quad \checkmark \end{aligned}$$

$$\frac{\partial \vec{r}}{\partial \eta} \cdot \frac{\partial \vec{r}}{\partial z} = (\sinh(\eta) \cos(\phi) \hat{i} + \cosh(\eta) \sin(\phi) \hat{j} + 0 \hat{k}) \cdot (0 \hat{i} + 0 \hat{j} + 1 \hat{k}) = 0 + 0 + 0 = 0 \quad \checkmark$$

$$\frac{\partial \vec{r}}{\partial \phi} \cdot \frac{\partial \vec{r}}{\partial z} = (-\cosh(\eta) \sin(\phi) \hat{i} + \sinh(\eta) \cos(\phi) \hat{j}) \cdot (0 \hat{i} + 0 \hat{j} + 1 \hat{k}) = 0 + 0 + 0 = 0 \quad \checkmark$$

These are orthogonal co-ordinates

(b) What is the metric tensor?

$$g = \begin{pmatrix} d\eta & d\phi & dz \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} d\eta \\ d\phi \\ dz \end{pmatrix} = \begin{pmatrix} d\eta & 0 & 0 \\ 0 & d\phi & 0 \\ 0 & 0 & dz \end{pmatrix} \begin{pmatrix} d\eta \\ d\phi \\ dz \end{pmatrix} = d\eta^2 + d\phi^2 + dz^2$$

$$dx = \frac{dx}{d\eta} d\eta + \frac{dx}{d\phi} d\phi, \quad dx^2 = \left(\frac{dx}{d\eta}\right)^2 d\eta^2 + 2 \left(\frac{dx}{d\eta}\right) \left(\frac{dx}{d\phi}\right) d\eta d\phi + \left(\frac{dx}{d\phi}\right)^2 d\phi^2$$

$$\frac{dx}{d\eta} = \sinh(\eta) \cos(\phi), \quad \left(\frac{dx}{d\eta}\right)^2 = \sinh^2(\eta) \cos^2(\phi) : \frac{dx}{d\phi} = -\cosh(\eta) \sin(\phi), \quad \left(\frac{dx}{d\phi}\right)^2 = \cosh^2(\eta) \sin^2(\phi)$$

$$\left(\frac{dx}{d\eta}\right)^2 + \left(\frac{dx}{d\phi}\right)^2 = \sinh^2(\eta) \cos^2(\phi) + \cosh^2(\eta) \sin^2(\phi) = \sinh^2(\eta) \cos^2(\phi) + \cosh^2(\eta) (1 - \cos^2(\phi))$$

$$\left(\frac{dx}{d\eta}\right)^2 + \left(\frac{dx}{d\phi}\right)^2 = \sinh^2(\eta) \cos^2(\phi) + \cosh^2(\eta) - \cosh^2(\eta) \cos^2(\phi) = \cosh^2(\eta) - \cos^2(\phi)$$

$$2 \left(\frac{dx}{d\eta}\right) \left(\frac{dx}{d\phi}\right) = -2 \sinh(\eta) \cos(\phi) \cosh(\eta) \sin(\phi) = -\sinh(2\eta) \frac{\sin(2\phi)}{2}$$

$$dx^2 = \cosh^2(\eta) d\phi^2 - \sinh(2\eta) \frac{\sin(2\phi)}{2} d\eta d\phi - \cos^2(\phi) d\eta^2$$

$$dy = \frac{dy}{d\eta} d\eta + \frac{dy}{d\phi} d\phi, \quad dy^2 = \left(\frac{dy}{d\eta}\right)^2 d\eta^2 + 2 \left(\frac{dy}{d\eta}\right) \left(\frac{dy}{d\phi}\right) d\eta d\phi + \left(\frac{dy}{d\phi}\right)^2 d\phi^2$$

$$\frac{dy}{d\eta} = \cosh(\eta) \sin(\phi), \quad \left(\frac{dy}{d\eta}\right)^2 = \cosh^2(\eta) \sin^2(\phi) : \left(\frac{dy}{d\phi}\right) = \sinh(\eta) \cos(\phi), \quad \left(\frac{dy}{d\phi}\right)^2 = \sinh^2(\eta) \cos^2(\phi)$$

$$\left(\frac{dy}{d\eta}\right)^2 + \left(\frac{dy}{d\phi}\right)^2 = \cosh^2(\eta) \sin^2(\phi) + \sinh^2(\eta) \cos^2(\phi) = \cosh^2(\eta) (1 - \cos^2(\phi)) + \sinh^2(\eta) \cos^2(\phi)$$

## Problem 1: Continued

$$\left(\frac{dy}{d\eta}\right)^2 + \left(\frac{dy}{d\varphi}\right)^2 = \cosh^2(\eta) - \cosh^2(\eta)\cos^2(\varphi) + \sinh^2(\eta)\cos^2(\varphi) = \cosh^2(\eta) - \cos^2(\varphi)$$

$$2\left(\frac{dy}{d\eta}\right)\left(\frac{dy}{d\varphi}\right) = 2\sinh(\eta)\cos(\varphi)\cosh(\eta)\sin(\varphi) = \sinh(2\eta)\frac{\sin(2\varphi)}{2}$$

$$dy^2 = -\cos^2(\varphi)d\varphi^2 + \sinh(2\eta)\frac{\sin(2\varphi)}{2}d\eta d\varphi + \cosh^2(\eta)d\varphi^2$$

$$dx^2 + dy^2 = (\cosh^2(\eta) - \cos^2(\varphi))d\varphi^2 + (\cosh^2(\eta) - \cos^2(\varphi))d\eta^2$$

$$dz^2 = 1$$

$$dx^2 + dy^2 + dz^2 = (\cosh^2(\eta) - \cos^2(\varphi))d\eta^2 + (\cosh^2(\eta) - \cos^2(\varphi))d\varphi^2 + dz^2$$

$$g = \begin{pmatrix} \cosh^2(\eta) - \cos^2(\varphi) & 0 & 0 \\ 0 & \cosh^2(\eta) - \cos^2(\varphi) & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

(c) What is  $\vec{\nabla}f$  in these co-ordinates?

$$\vec{\nabla}\varphi = \frac{1}{h_i} \partial_i \varphi \hat{e}_i, \quad h_\eta^2 = \cosh^2(\eta) - \cos^2(\varphi), \quad h_\varphi^2 = \cosh^2(\eta) - \cos^2(\varphi), \quad h_z^2 = 1$$

$$\vec{\nabla}f = \frac{1}{\sqrt{\cosh^2(\eta) - \cos^2(\varphi)}} \left( \frac{\partial f}{\partial \eta} \hat{e}_\eta + \frac{\partial f}{\partial \varphi} \hat{e}_\varphi \right) + \frac{\partial f}{\partial z} \hat{e}_z$$

(d) What is  $\vec{\nabla} \cdot \vec{\mathcal{E}}$  in this co-ordinates, where  $\vec{\mathcal{E}} = \mathcal{E}_\eta \hat{\eta} + \mathcal{E}_\varphi \hat{\varphi} + \mathcal{E}_z \hat{z}$ ?

$$\vec{\nabla} \cdot \vec{\mathcal{E}} = \frac{1}{\cosh^2(\eta) - \cos^2(\varphi)} \left( \frac{\partial}{\partial \eta} \sqrt{\cosh^2(\eta) - \cos^2(\varphi)} \mathcal{E}_\eta + \frac{\partial}{\partial \varphi} \sqrt{\cosh^2(\eta) - \cos^2(\varphi)} \mathcal{E}_\varphi \right) + \frac{\partial \mathcal{E}_z}{\partial z}$$

(e) What is  $\nabla^2 f$  in this co-ordinate system? Where possible, express your answer only in terms of cosines and hyperbolic cosines.

In parts (b)-(e) express your answer in terms of cosines and hyperbolic cosines, where possible.

$$\nabla^2 f = \frac{1}{h_1 h_2 h_3} \left[ \partial_1 \left( \frac{h_2 h_3}{h_1} \partial_1 f \right) + \partial_2 \left( \frac{h_3 h_1}{h_2} \partial_2 f \right) + \partial_3 \left( \frac{h_1 h_2}{h_3} \partial_3 f \right) \right]$$

$$\nabla^2 f = \frac{1}{\cosh^2(\eta) - \cos^2(\varphi)} \left[ \frac{\partial^2 f}{\partial \eta^2} + \frac{\partial^2 f}{\partial \varphi^2} \right] + \frac{\partial^2 f}{\partial z^2}$$

## Problem 1: Review

### Procedure:

- To show that co-ordinates are orthogonal, with co-ordinates labeled as  $q_1, q_2, q_3$ , do

$$\frac{\partial f}{\partial q_1} \cdot \frac{\partial f}{\partial q_2} = 0, \quad \frac{\partial f}{\partial q_1} \cdot \frac{\partial f}{\partial q_3} = 0, \quad \frac{\partial f}{\partial q_2} \cdot \frac{\partial f}{\partial q_3} = 0.$$

- To find the metric tensor for orthogonal co-ordinates, use

$$\hat{g} = \begin{bmatrix} d\eta^2 & 0 & 0 \\ 0 & d\phi^2 & 0 \\ 0 & 0 & dz^2 \end{bmatrix}.$$

- Where to find  $d\eta^2, d\phi^2, dz^2$  we must do the following: Differentiate first component of  $\vec{f}$  with respect to all variables that show up in the equation. In this case it is  $\phi$  and  $\eta$ , then find  $dq_i^2$  by doing the following. (In this case we will examine  $dx^2$ ).

$$dx^2 = \left( \frac{\partial f}{\partial \eta} \right)^2 d\eta^2 + \left( \frac{\partial f}{\partial \eta} \right) \left( \frac{\partial f}{\partial \phi} \right) d\eta d\phi + \left( \frac{\partial f}{\partial \phi} \right)^2 d\phi^2.$$

- Repeat the above procedure for all components, then calculate  $ds^2 = dx^2 + dy^2 + dz^2$ , group the terms together for  $d\eta^2, d\phi^2, dz^2$ . The terms next to the components belong in the diagonals of the metric tensor.
- Use the Gradient and Laplacian equations to compute the rest of the problem.

### Key Concepts:

- For co-ordinates to be orthogonal, they must follow the first rule outlined in the procedure section.
- For orthogonal co-ordinates the metric tensor will be diagonal.
- The Gradient and Laplacian can be calculated easily once we know the values in the diagonal.

### Variations:

- The function in the problem statement can change.
  - Thus leading to a different metric.

## Problem 2:

Consider the function

$$\phi(\mathbf{r}) = \frac{r}{r^2 + \epsilon^2}$$

in three dimensions. Calculate  $\nabla^2 \phi$ .

(a) Show that for any fixed value of  $r = r_0 \neq 0$ ,

$$\lim_{r_0 \rightarrow 0} \left[ \lim_{\epsilon \rightarrow 0} \nabla^2 \phi(r_0) \right] = 0.$$

$$\nabla^2 \phi = \frac{1}{h_1 h_2 h_3} \left( \partial_1 \left( \frac{h_2 h_3}{h_1} \right) \partial_1 \phi \right) + \partial_2 \left( \frac{h_3 h_1}{h_2} \partial_2 \phi \right) + \partial_3 \left( \frac{h_1 h_2}{h_3} \partial_3 \phi \right)$$

$$\nabla^2 \phi = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \phi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \phi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \phi}{\partial \varphi^2}$$

$$\nabla^2 \phi = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \phi}{\partial r} \right) = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 (r^2 + \epsilon^2)^{-1} - 2r^4 (r^2 + \epsilon^2)^{-2} \right) = \frac{1}{r^2} \frac{\partial}{\partial r} \left( \frac{r^2 (\epsilon^2 - r^2)}{(r^2 + \epsilon^2)^2} \right)$$

$$\nabla^2 \phi = \frac{1}{r^2} \left( \frac{2\epsilon^2 r (\epsilon^2 - 3r^2)}{(r^2 + \epsilon^2)^3} \right) = \frac{2\epsilon^2 (\epsilon^2 - 3r^2)}{r (r^2 + \epsilon^2)^3} : \nabla^2 \phi = \frac{2\epsilon^2 (\epsilon^2 - 3r^2)}{r (r^2 + \epsilon^2)^3}$$

$$\lim_{\epsilon \rightarrow 0} \left[ \frac{2\epsilon^2 (\epsilon^2 - 3r_0^2)}{r_0 (r_0^2 + \epsilon^2)^3} \right] = \frac{2 \cdot 0 \cdot (\epsilon^2 - 3r_0^2)}{r_0 (r_0^2 + 0^2)^3} = \frac{0}{r_0^6} = 0 : \lim_{r_0 \rightarrow 0} 0 = 0 \checkmark$$

$$\boxed{\lim_{r_0 \rightarrow 0} \left[ \lim_{\epsilon \rightarrow 0} \frac{2\epsilon^2 (\epsilon^2 - 3r_0^2)}{r_0 (r_0^2 + \epsilon^2)^3} \right] = 0}$$

(b) Show that for a fixed value of  $\epsilon \neq 0$ ,

$$\lim_{\epsilon \rightarrow 0} \left[ \lim_{r_0 \rightarrow 0} \nabla^2 \phi(r_0) \right] = \infty.$$

$$\lim_{r_0 \rightarrow 0} \left[ \frac{2\epsilon^2 (\epsilon^2 - 3r_0^2)}{r_0 (r_0^2 + \epsilon^2)^3} \right] = \frac{2\epsilon^2 (\epsilon^2 - 3(0))}{0 (0^2 + \epsilon^2)^3} = \frac{2\epsilon^4}{0} = \infty : \lim_{\epsilon \rightarrow 0} \infty = \infty \checkmark$$

$$\boxed{\lim_{\epsilon \rightarrow 0} \left[ \lim_{r_0 \rightarrow 0} \frac{2\epsilon^2 (\epsilon^2 - 3r_0^2)}{r_0 (r_0^2 + \epsilon^2)^3} \right] = \infty}$$

(c) Using (a) and (b), construct an argument that  $\nabla^2 \frac{1}{r}$  is zero at all points except at the origin, where it diverges. If  $f(r)$  is a smooth, well behaved function, calculate an estimate for

$$\int f(r) \nabla^2 \left( \frac{1}{r} \right) d^3 r$$

where integration is over all space.

We can Taylor Series expand  $F(r) \rightarrow \frac{F(0) x^0}{1} + \frac{F'(0) x}{1} + \frac{F''(0) x^2}{2} + \dots$

We can keep the first term only,  $F(0)$ , and disregard all of the other terms. Our integral then becomes

## Problem 2: Continued

$$F(0) \int \nabla^2 \left( \frac{1}{r} \right) d^3r$$

From here we have  $\nabla^2 \left( \frac{1}{r} \right) d^3r = -4\pi \delta(r)$ , where our integral now becomes

$-F(0) \cdot 4\pi \int \delta(r) dr$

*Arken Pg. (81) - (1.161)*

Since  $\delta(r)$  is a dirac delta function,  $\int \delta(r) dr = 1 \quad \therefore$

$$\boxed{\int f(r) \nabla^2 \left( \frac{1}{r} \right) dr = -4\pi f(0)}$$

## Problem 2: Review

### Procedure:

- Calculate the Laplacian in spherical co-ordinates.
- Evaluate the limits for the Laplacian.
- Expand the function with a Taylor series, keeping only the first order term.
- Approximate the Laplacian of  $1/r$  and substitute the result.

### Key Concepts:

- The Laplacian of this function will go to 0 or  $\infty$  depending on how the extremum is evaluated.
- We can approximate  $f(r)$  since it is a smooth function.
- We can then use the common integral for the end result.

### Variations:

- The original function can be changed.
  - This in turn would create a different problem. Possibly having to calculate the Laplacian in cylindrical co-ordinates or something similar.

**Problem 3:**

Your textbook (and many others) state that the curl of the gradient of a scalar function is zero. Consider the function  $f(r, \phi) = \phi$  where  $r$  and  $\phi$  are the standard polar co-ordinates in two dimensions.

(a) Calculate  $\vec{g} = \vec{\nabla} f$  in polar co-ordinates.

$$\vec{\nabla} f = \frac{1}{h_i} \partial_i f \hat{e}_i : d\vec{s} = dr + r d\theta + r \sin\theta d\phi$$

$$\text{w/ } \theta = \pi/2 : \vec{\nabla} f = \cancel{\frac{\partial f}{\partial r} \hat{r}} + \frac{1}{r} \cancel{\frac{\partial f}{\partial \theta} \hat{\theta}} + \frac{1}{r} \frac{\partial f}{\partial \phi} \hat{\phi}$$

$$\boxed{\vec{\nabla} f = \frac{1}{r} \hat{\phi} = \vec{g}}$$

(b) Calculate  $\vec{\nabla} \times \vec{\nabla} f = \vec{\nabla} \times \vec{g}$  in polar co-ordinates.

$$\vec{\nabla} \times \vec{\phi} = \frac{1}{h_2 h_3} (\partial_2 (h_3 \phi_3) - \partial_3 (h_2 \phi_2)) \hat{e}_1 + \frac{1}{h_3 h_1} (\partial_3 (h_1 \phi_1) - \partial_1 (h_3 \phi_3)) \hat{e}_2 + \frac{1}{h_1 h_2} (\partial_1 (h_2 \phi_2) - \partial_2 (h_1 \phi_1)) \hat{e}_3$$

$$\vec{\nabla} \times \vec{g} = \frac{1}{r^2 \sin\theta} (\cancel{\partial_\theta (r \sin\theta \cdot 1/r)} - \cancel{\partial_\phi (r \cdot 0)}) \hat{e}_r + \frac{1}{r \sin\theta} (\cancel{\partial_\phi (r \cdot 0)} - \cancel{\partial_r (r \sin\theta \cdot 1/r)}) \hat{e}_\theta + \frac{1}{r} (\cancel{\partial_r (r \cdot 0)} - \cancel{\partial_\theta (0)}) \hat{e}_\phi$$

$$\boxed{\vec{\nabla} \times \vec{g} = 0}$$

(c) From Stoke's Theorem, we know

$$\int (\vec{\nabla} \times \vec{g}) \cdot \vec{n} dA = \oint \vec{g} \cdot d\vec{l}$$

Calculate the last line integral in polar co-ordinates.

$$\oint \frac{1}{r} \hat{\phi} \cdot r d\hat{\phi} = 2\pi$$

$$\boxed{\oint \vec{g} \cdot d\vec{l} = 2\pi}$$

(d) Do your answers agree? Can you explain?

They are not going to agree since there is a discontinuity at  $\phi = 0$ .



## Problem 3: Review

### Procedure:

- Calculate the Gradient of the function in spherical polar co-ordinates.
- Calculate the curl of the Gradient.
- Calculate the line integral of Stokes' theorem.
- Explain why they aren't equal.

### Key Concepts:

- Stokes' theorem is not valid for functions that have discontinuities in them (This is why the line integral is different than the LHS in Stokes' theorem).

### Variations:

- The function that is given to us can change.
  - The curl of the gradient will still be equal to zero but it is possible the line integral will also be zero as long as there are no discontinuities.