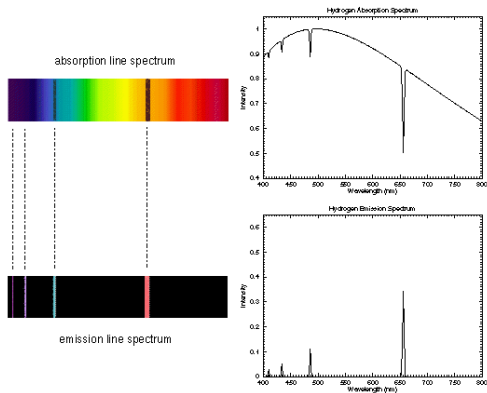


Lectures 01

P. Gutierrez

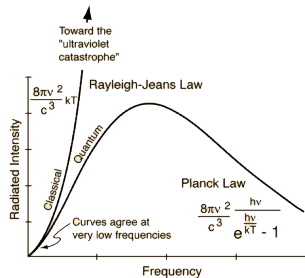
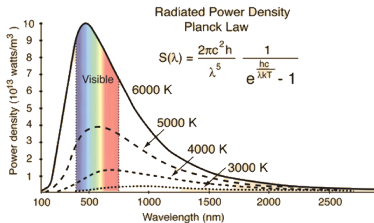
Department of Physics & Astronomy
University of Oklahoma

Absorption Emission Spectra

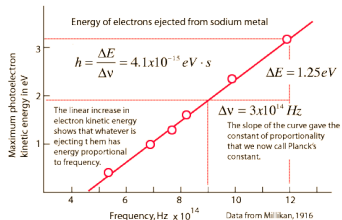
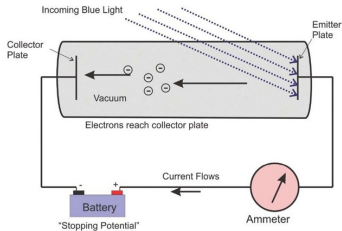


Two ways of showing the same spectra: on the **left** are pictures of the dispersed light and on the **right** are plots of the Intensity vs. wavelength. Notice that the pattern of spectral lines in the absorption and emission line spectra are the **same** since the gas is the same.

Black-Body Radiation



Photoelectric Effect



Electron Diffraction

3.2 Electron diffraction

Davisson-Germer experiment

- Figure 10.1 shows a tube for demonstrating electron diffraction by Davisson and Germer.

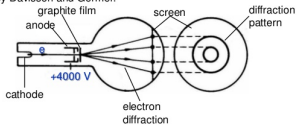


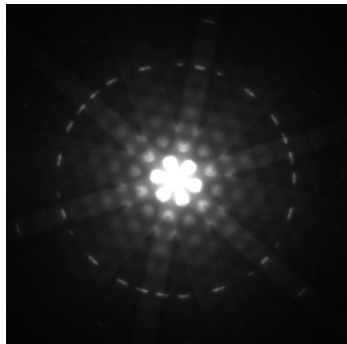
Figure 10.1: electron diffraction tube

- A beam of accelerated electrons strikes on a layer of graphite which is extremely thin and a diffraction pattern consisting of rings is seen on the tube face.



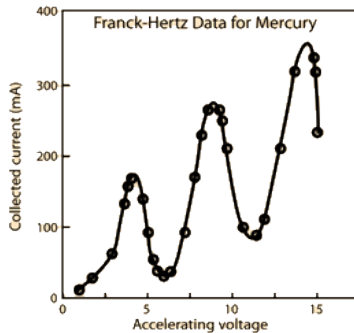
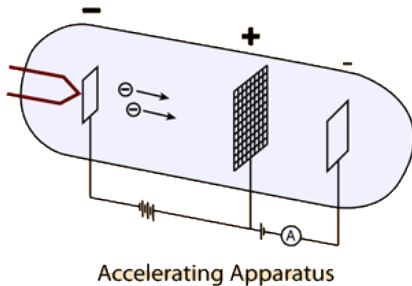
PHY3212 - Photoelectric Effect

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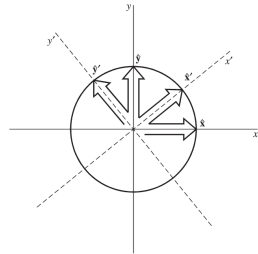
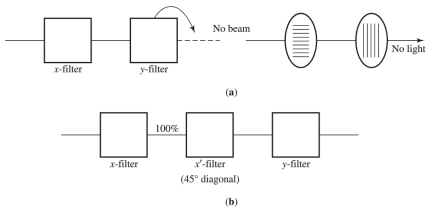


Davisson and Germer

Franck Hertz Experiment

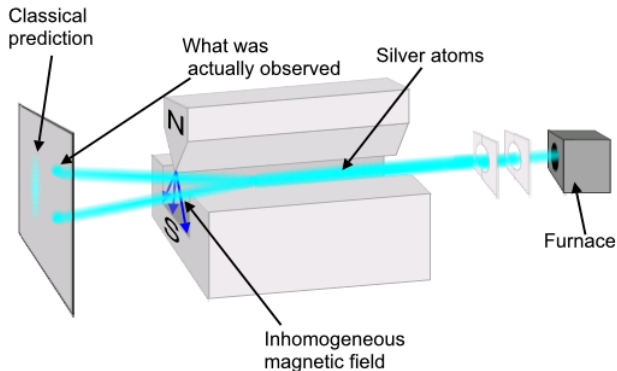


Light Polarization

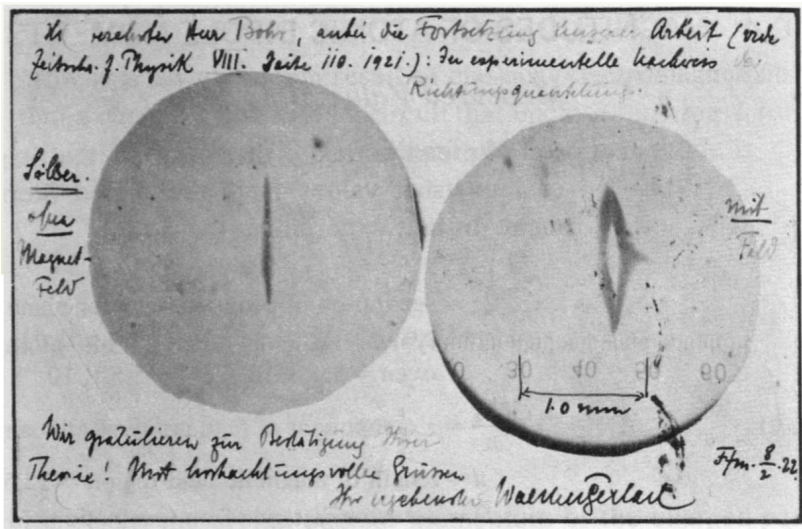


Stern Gerlach Experiment

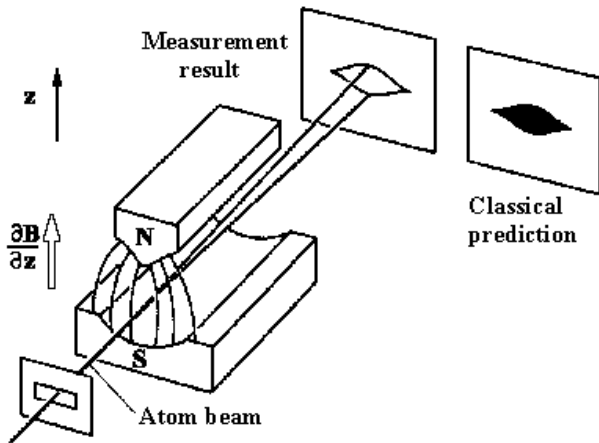
Apparatus

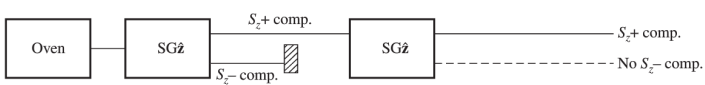


Result

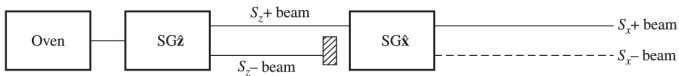


Apparatus

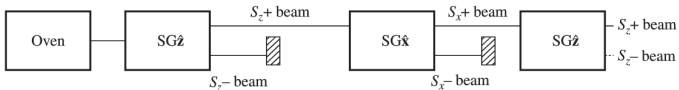




(a)



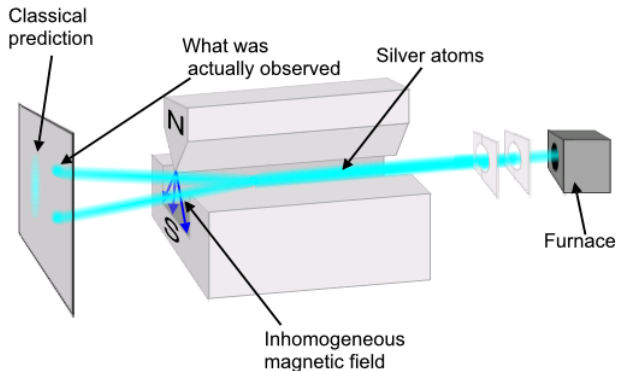
(b)



(c)

Stern Gerlach Experiment

Apparatus



$$|S_z; +\rangle$$

$$|S_z; -\rangle$$

$$|S_x; +\rangle = \frac{1}{\sqrt{2}} [|S_z; -\rangle + |S_z; +\rangle]$$

$$|S_x; -\rangle = \frac{1}{\sqrt{2}} [|S_z; -\rangle - |S_z; +\rangle]$$

$$|S_y; +\rangle = \frac{1}{\sqrt{2}} [|S_z; +\rangle + i |S_z; -\rangle]$$

$$|S_y; -\rangle = \frac{1}{\sqrt{2}} [|S_z; +\rangle - i |S_z; -\rangle]$$

Stern-Gerlach States

$$|S_z; +\rangle$$

$$|S_z; -\rangle$$

$$|S_x; +\rangle = \frac{1}{\sqrt{2}} [|S_z; -\rangle + |S_z; +\rangle]$$

$$|S_x; -\rangle = \frac{1}{\sqrt{2}} [|S_z; -\rangle - |S_z; +\rangle]$$

$$|S_y; +\rangle = \frac{1}{\sqrt{2}} [|S_z; +\rangle + i |S_z; -\rangle]$$

$$|S_y; -\rangle = \frac{1}{\sqrt{2}} [|S_z; +\rangle - i |S_z; -\rangle]$$

$$|S_z; +\rangle$$

$$|S_z; -\rangle$$

$$|S_x; +\rangle = \frac{1}{\sqrt{2}} [|S_z; -\rangle + |S_z; +\rangle]$$

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$$|S_y; -\rangle = \frac{1}{\sqrt{2}} [|S_z; +\rangle - i |S_z; -\rangle]$$

Linear Vector Space

Definition 1: A linear vector space \mathbb{V} is a collection of objects of the form $|V\rangle$, called vectors, for which there exists:

- 1 A definite rule for forming the vector sum, denoted $|V\rangle + |W\rangle$
- 2 A definite rule for multiplication by scalars a denoted $a|V\rangle$ with the following features:
 - The results of these operations is another element of the space, a feature called closure:
 $|V\rangle + |W\rangle \in \mathbb{V}$ and $a|V\rangle \in \mathbb{V}$.
 - Scalar multiplication is distributive in the vectors:
 $a(|V\rangle + |W\rangle) = a|V\rangle + a|W\rangle$.
 - Addition is commutative:
 $|V\rangle + |W\rangle = |W\rangle + |V\rangle$.
 - Addition is associative:
 $|V\rangle + (|W\rangle + |Z\rangle) = (|V\rangle + |W\rangle) + |Z\rangle$
 - There exists a null vector $|0\rangle$ obeying
 $|V\rangle + |0\rangle = |V\rangle$
 - For every vector $|V\rangle$ there exists an inverse under addition, $|-V\rangle$, such that $|V\rangle + |-V\rangle = |0\rangle$

Linear Vector Space

Definition 2: The numbers a are called the field over which the vector space is defined.

If the field consists of all real numbers, we have a real vector space, if they are complex, we have a complex vector space. Note, the vectors are neither real nor complex, the adjective applies only to the scalars.

Dual Vector Space

- $\langle \alpha | \xleftrightarrow{\text{DC}} | \alpha \rangle$
- $c^* \langle \alpha | \xleftrightarrow{\text{DC}} c | \alpha \rangle$
- Inner product $\langle \alpha | \beta \rangle = \text{Complex Number}$
- Orthogonal $\langle \alpha | \beta \rangle = 0$
- Normalization $\langle \alpha | \alpha \rangle = 1$
- Postulates:
 - $\langle \alpha | \beta \rangle = \langle \beta | \alpha \rangle^*$
 - $\langle \alpha | \alpha \rangle \geq 0$
- Operators (*in general*) $\tilde{X} | \alpha \rangle \neq \langle \alpha | \tilde{X}$
 - $\tilde{X} | \alpha \rangle = | \beta \rangle$
 - $\langle \alpha | \tilde{X} = \langle \gamma |$
 - $\langle \alpha | \tilde{X}^\dagger \xleftrightarrow{\text{DC}} \tilde{X} | \alpha \rangle$
 - In general $\tilde{X} \tilde{Y} \neq \tilde{Y} \tilde{X}$
 - $\tilde{X} = | \alpha \rangle \langle \beta |$

Dual Vector Space

- $\langle \alpha | \stackrel{\text{DC}}{\Longleftrightarrow} | \alpha \rangle$
- $c^* \langle \alpha | \stackrel{\text{DC}}{\Longleftrightarrow} c | \alpha \rangle$
- Inner product $\langle \alpha | \beta \rangle = \text{Complex Number}$
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 - $\langle \alpha | \tilde{\mathbf{X}}^\dagger \xleftrightarrow{\text{DC}} \tilde{\mathbf{X}} | \alpha \rangle$
 - In general $\tilde{\mathbf{X}} \tilde{\mathbf{Y}} \neq \tilde{\mathbf{Y}} \tilde{\mathbf{X}}$
 - $\tilde{\mathbf{X}} = | \alpha \rangle \langle \beta |$

Matrix Representation

- $\tilde{\mathbf{A}} = \sum_{ij} |a_i\rangle\langle a_i| \tilde{\mathbf{A}} |a_j\rangle\langle a_j| = \sum_{ij} \langle a_i | \tilde{\mathbf{A}} | a_j \rangle |a_i\rangle\langle a_j|$
- Matrix elements $\langle a_i | \tilde{\mathbf{A}} | a_j \rangle = \langle \text{row} | \tilde{\mathbf{A}} | \text{column} \rangle$
 - In it's own basis matrix is diagonal $\langle a_i | \tilde{\mathbf{A}} | a_j \rangle = a_j \delta_{ij}$
- State can be expressed as a column vector;
 $|\alpha\rangle = \sum_i |a_i\rangle\langle a_i| |\alpha\rangle$
- Matrix elements $\langle a_i | \alpha \rangle$
- Matrix multiplication

$$\begin{aligned}\tilde{\mathbf{A}} |\alpha\rangle &= \sum_{ij} |a_i\rangle\langle a_i| \tilde{\mathbf{A}} |a_j\rangle\langle a_j| |\alpha\rangle \\ &= \sum_{ij} |a_i\rangle \langle a_i | \tilde{\mathbf{A}} | a_j \rangle \langle a_j | \alpha \rangle\end{aligned}$$

Uncertainty Principle

- Incompatible observable $[\tilde{\mathbf{A}}, \tilde{\mathbf{B}}] \neq 0$
 - Measure $\tilde{\mathbf{A}}$ in the state $|a_2\rangle$
 - Measure $\tilde{\mathbf{B}}$ in the state $|a_2\rangle = \sum_j |b_j\rangle \langle b_j | a_2 \rangle$; only the probability of the value is known.
- Uncertainty principle
 - Schwartz inequality $\langle \alpha | \alpha \rangle \langle \beta | \beta \rangle \geq |\langle \alpha | \beta \rangle|^2$
 - Dispersion $\Delta \tilde{\mathbf{A}} = \tilde{\mathbf{A}} - \langle \tilde{\mathbf{A}} \rangle \mathbf{1}$
 - $\langle (\Delta \tilde{\mathbf{A}})^2 \rangle = \langle \tilde{\mathbf{A}}^2 - 2\tilde{\mathbf{A}} \langle \tilde{\mathbf{A}} \rangle + \langle \tilde{\mathbf{A}} \rangle^2 \rangle = \langle \tilde{\mathbf{A}}^2 \rangle - \langle \tilde{\mathbf{A}} \rangle^2$
 - Hermitian operator $\{\tilde{\mathbf{A}}, \tilde{\mathbf{B}}\}$ has real eigenvalues.
 - Anti-Hermitian operator $[\tilde{\mathbf{A}}, \tilde{\mathbf{B}}]$ has imaginary eigenvalues.

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 - Hermitian operator $\{\tilde{\mathbf{A}}, \tilde{\mathbf{B}}\}$ has real eigenvalues.
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Uncertainty Principle – *continued*

- Schwartz inequality

- $\left\langle (\Delta \tilde{\mathbf{A}})^2 \right\rangle \left\langle (\Delta \tilde{\mathbf{B}})^2 \right\rangle \geq \left| \left\langle (\Delta \tilde{\mathbf{A}} \Delta \tilde{\mathbf{B}}) \right\rangle \right|^2$

- $\Delta \tilde{\mathbf{A}} \Delta \tilde{\mathbf{B}} = \frac{1}{2} [\Delta \tilde{\mathbf{A}}, \Delta \tilde{\mathbf{B}}] + \frac{1}{2} \{ \Delta \tilde{\mathbf{A}}, \Delta \tilde{\mathbf{B}} \}$

- $\Delta \tilde{\mathbf{A}} \Delta \tilde{\mathbf{B}} = \frac{1}{2} [\tilde{\mathbf{A}}, \tilde{\mathbf{B}}] + \frac{1}{2} \{ \Delta \tilde{\mathbf{A}}, \Delta \tilde{\mathbf{B}} \}$

- $\left\langle (\Delta \tilde{\mathbf{A}})^2 \right\rangle \left\langle (\Delta \tilde{\mathbf{B}})^2 \right\rangle \geq \frac{1}{4} \left\langle [\tilde{\mathbf{A}}, \tilde{\mathbf{B}}] \right\rangle^2 + \frac{1}{4} \left\langle \{ \Delta \tilde{\mathbf{A}}, \Delta \tilde{\mathbf{B}} \} \right\rangle^2$

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Uncertainty Principle – *continued*

- Schwartz inequality

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Uncertainty Principle – *continued*

- Schwartz inequality

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- $\Delta \tilde{\mathbf{A}} \Delta \tilde{\mathbf{B}} = \frac{1}{2} \left[\Delta \tilde{\mathbf{A}}, \Delta \tilde{\mathbf{B}} \right] + \frac{1}{2} \left\{ \Delta \tilde{\mathbf{A}}, \Delta \tilde{\mathbf{B}} \right\}$

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- $\left\langle \left(\Delta \tilde{\mathbf{A}} \right)^2 \right\rangle \left\langle \left(\Delta \tilde{\mathbf{B}} \right)^2 \right\rangle \geq \frac{1}{4} \left\langle \left[\tilde{\mathbf{A}}, \tilde{\mathbf{B}} \right] \right\rangle^2 + \frac{1}{4} \left\langle \left\{ \Delta \tilde{\mathbf{A}}, \Delta \tilde{\mathbf{B}} \right\} \right\rangle^2$

- $\left\langle \left(\Delta \tilde{\mathbf{A}} \right)^2 \right\rangle \left\langle \left(\Delta \tilde{\mathbf{B}} \right)^2 \right\rangle \geq \frac{1}{4} \left\langle \left[\tilde{\mathbf{A}}, \tilde{\mathbf{B}} \right] \right\rangle^2$

Uncertainty Principle – *continued*

- Schwartz inequality

- $\left\langle (\Delta \tilde{\mathbf{A}})^2 \right\rangle \left\langle (\Delta \tilde{\mathbf{B}})^2 \right\rangle \geq \left| \left\langle (\Delta \tilde{\mathbf{A}} \Delta \tilde{\mathbf{B}}) \right\rangle \right|^2$

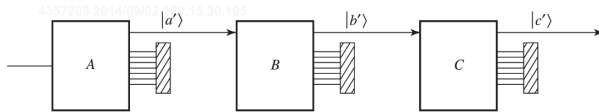
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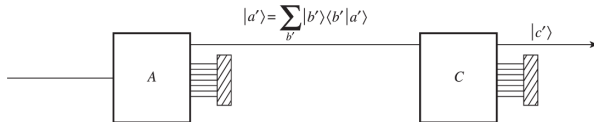
- $\left\langle (\Delta \tilde{\mathbf{A}})^2 \right\rangle \left\langle (\Delta \tilde{\mathbf{B}})^2 \right\rangle \geq \frac{1}{4} \left\langle [\tilde{\mathbf{A}}, \tilde{\mathbf{B}}] \right\rangle^2 + \frac{1}{4} \left\langle \{ \Delta \tilde{\mathbf{A}}, \Delta \tilde{\mathbf{B}} \} \right\rangle^2$

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Apparatus



(a)

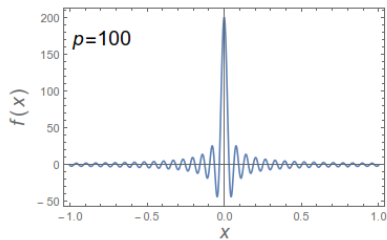
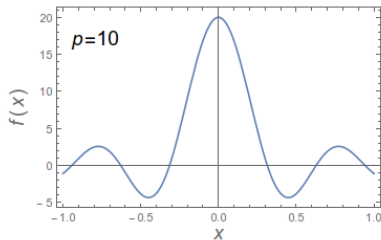


(b)

Lectures 02

P. Gutierrez

Department of Physics & Astronomy
University of Oklahoma



$$f(x) = \frac{2 \sin(px)}{x} \Rightarrow \begin{cases} \lim_{x \rightarrow \infty} f(x) \rightarrow 0 \\ \lim_{x \rightarrow 0} f(x) \rightarrow 2p \end{cases}$$

Translation Operator

- $\mathcal{J}(dx) |x\rangle = |x + dx\rangle$
- Propose $\mathcal{J}(d\vec{x}) = 1 - i\tilde{\mathbf{K}} \cdot d\vec{x}$
- Unitary operator (*maintain normalization*):

$$\begin{aligned}\langle\alpha|\alpha\rangle &= \left\langle\alpha\left|\mathcal{J}^\dagger(d\vec{x}')\mathcal{J}(d\vec{x}')\right|\alpha\right\rangle \Rightarrow \mathcal{J}^\dagger(d\vec{x}')\mathcal{J}(d\vec{x}') = 1 \\ &\Rightarrow \tilde{\mathbf{K}} = \tilde{\mathbf{K}}^\dagger\end{aligned}$$

- Successive translations:

$$\mathcal{J}(d\vec{x}'')\mathcal{J}(d\vec{x}') = \mathcal{J}(d\vec{x}'' + d\vec{x}').$$

- Inverse translation:

$$\mathcal{J}(-d\vec{x}') = \mathcal{J}^{-1}(d\vec{x}').$$

- No translation:

$$\lim_{d\vec{x}' \rightarrow 0} \mathcal{J}(d\vec{x}') = \mathcal{J}(0) = 1.$$

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- No translation:

$$\lim_{d\vec{x}' \rightarrow 0} \mathcal{J}(d\vec{x}') = \mathcal{J}(0) = 1.$$

Translation Operator

- $\mathcal{J}(dx) |x\rangle = |x + dx\rangle$
- Propose $\mathcal{J}(d\vec{x}) = 1 - i\tilde{\mathbf{K}} \cdot d\vec{x}$
- Unitary operator (*maintain normalization*):

$$\begin{aligned}\langle\alpha|\alpha\rangle &= \left\langle\alpha\left|\mathcal{J}^\dagger(d\vec{x}')\mathcal{J}(d\vec{x}')\right|\alpha\right\rangle \Rightarrow \mathcal{J}^\dagger(d\vec{x}')\mathcal{J}(d\vec{x}') = 1 \\ &\Rightarrow \tilde{\mathbf{K}} = \tilde{\mathbf{K}}^\dagger\end{aligned}$$

- Successive translations:

$$\mathcal{J}(d\vec{x}'')\mathcal{J}(d\vec{x}') = \mathcal{J}(d\vec{x}'' + d\vec{x}').$$

- Inverse translation:

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