



COLLEGE OF ARTS AND SCIENCES

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The UNIVERSITY *of* OKLAHOMA

Electrodynamics 1

CH. 5 MAGNETOSTATICS IN VACUUM LECTURE NOTES

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Faradays Law

Faradays Law mathematically is

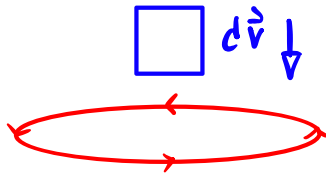
$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Where we have,

$$\oint \vec{E} \cdot d\vec{\lambda} = -\frac{d}{dt} \iint \vec{B} \cdot \hat{n} ds$$

Faradays Law says that an electro/magnetic Field can be created

We look at a simple scenario



Where the same happens if the loop is moving and the block is stationary

We look at the moving block,

$$\oint \vec{E} \cdot d\vec{\lambda} = -\frac{\partial}{\partial t} \int \vec{B} \cdot \hat{n} ds + \int \vec{\nabla} \times (\vec{r} \times \vec{B}) \cdot \hat{n} ds$$

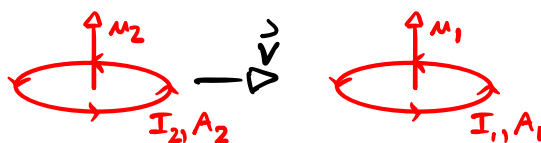
We then have

$$\oint (\vec{E} - \vec{v} \times \vec{B}) \cdot d\vec{\lambda} = -\frac{\partial}{\partial t} \int \vec{B} \cdot \hat{n} ds \Rightarrow \oint \vec{E} \cdot d\vec{\lambda} = -\frac{\partial}{\partial t} \int \vec{B} \cdot \hat{n} ds$$

We have our electric potential energy as

$$U = \frac{1}{2} \sum q \phi(\vec{r}_i) = \frac{1}{2} \int \rho(\vec{r}) \phi(\vec{r}) d^3r = \frac{\epsilon_0}{2} \int E^2 d^3r$$

Lets look at two separate loops



We then have

$$dW_e = -I_2 \Delta t \mathcal{E}, \quad \text{w/} \quad \mathcal{E}_1 = -A_2 \frac{\partial B_{12}}{\partial x} \frac{\Delta x}{\Delta t} = -A_2 \frac{\partial B_{12}}{\partial x} v_x$$

Where further we see

$$dW_e = I_2 A_2 \frac{\partial B_{12}}{\partial x} V_x \Delta t = \mu_2 \frac{\partial B_{12}}{\partial x} V_x \Delta t$$

The force is then

$$\vec{F} = -\vec{\nabla}(\vec{m}_2, \vec{B}_1) = -\vec{\nabla}(m_2 B_{12})$$

We then finally have

$$dW_m = \vec{F} \cdot \Delta \vec{x} = -m_2 \frac{\partial B_{12}}{\partial x} V \Delta t$$

For the hoop that is static,

$$dW_T = dW_{ez} + dW_{mz} + dW_{e1} = -dW_{m1}$$

Looking at the math we see

$$dW = \vec{m}_1 \cdot \vec{B}_2, \quad U_m = I_1 \int \vec{B}_2 \cdot \hat{n} \, ds = I_1 \int (\vec{\nabla} \times \vec{A}_2) \cdot \hat{n} \, ds = I_1 \oint \vec{A}_2 \cdot d\vec{\ell}$$

If we examine the current along the loop,

$$I_1 d\vec{\ell} = \vec{J} \cdot \hat{n} \, ds \, d\ell$$

The electric potential energy is

$$U_m = \int \vec{J} \cdot \vec{A} \, d^3r \quad \text{w/} \quad \vec{J} = \frac{1}{\mu_0} (\vec{\nabla} \times \vec{B}) \quad \therefore \quad U_m = \frac{1}{2\mu_0} \int (\vec{\nabla} \times \vec{B}) \cdot \vec{A} \, d^3r$$

We then proceed to calculate

$$\begin{aligned} (\vec{\nabla} \times \vec{B}) \cdot \vec{A} &= \epsilon_{ijk} (\partial_i B_j) A_k = \epsilon_{ijk} (\partial_j B_k A_i) - B_k (\partial_j A_i) \\ &= \partial_j \epsilon_{jki} (B_k A_i) - B_k \epsilon_{kij} (\partial_j A_i) = \vec{\nabla} \cdot (\vec{B} \times \vec{A}) + \vec{B} \cdot (\vec{\nabla} \times \vec{A}) \end{aligned}$$

This means our potential energy will be

$$U_m = \frac{1}{2\mu_0} \int \vec{\nabla} \cdot (\vec{B} \times \vec{A}) \, d^3r + \int \vec{B} \cdot (\vec{\nabla} \times \vec{A}) \, d^3r = \frac{1}{2\mu_0} \int B^2 \, d^3r$$

$\hookrightarrow = 0 \rightarrow \text{Gauss' Law}$

So, finally we have

$$U_m = \frac{1}{2} \int \vec{J} \cdot \vec{A} \, d^3r = \frac{1}{2\mu_0} \int \frac{\vec{J}(\vec{r}) \cdot \vec{J}(\vec{r}')}{|\vec{r} - \vec{r}'|} \, d^3r \, d^3r'$$

We now move on to looking at an example

Example: Solenoid



We can say

$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 I_{enc}, \quad \vec{B} \cdot \vec{L} = \mu_0 N I \quad \therefore B = \mu_0 \frac{N}{L} I$$

The potential energy is then

$$U = \frac{1}{2\mu_0} \int B^2 d^3r = \frac{1}{2\mu_0} \left(\mu_0 \frac{N}{L} I \right)^2 A \cdot L = \frac{\mu_0}{2} \frac{N^2}{L^2} I^2 A$$

We can then say

$$dU = I \Delta t \cdot N \frac{dB}{\Delta t} = N I \Delta t A \left(\frac{\mu_0 N}{L} \right) \frac{\Delta I}{\Delta t} = \frac{\mu_0}{L} A N^2 I dI$$

We then say

$$U = \frac{\mu_0}{L} A N^2 \int_0^{I_1} I dI = \frac{\mu_0}{2} \frac{A N^2 I^2}{L}$$