

Lecture Set 07

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Angular Momentum – *Algebra*

The algebra that specifies angular momentum, either spin or orbital, is defined by the commutation relations.

- $[\tilde{\mathbf{J}}_i, \tilde{\mathbf{J}}_j] = i\hbar\epsilon_{ijk}\tilde{\mathbf{J}}_k$
- Define $\tilde{\mathbf{J}}^2 = \tilde{\mathbf{J}}_x^2 + \tilde{\mathbf{J}}_y^2 + \tilde{\mathbf{J}}_z^2$
- $[\tilde{\mathbf{J}}^2, \tilde{\mathbf{J}}_i] = 0$
- Define $\tilde{\mathbf{J}}_{\pm} = \tilde{\mathbf{J}}_x \pm i\tilde{\mathbf{J}}_y$
- $[\tilde{\mathbf{J}}^2, \tilde{\mathbf{J}}_{\pm}] = 0 \quad [\tilde{\mathbf{J}}_z, \tilde{\mathbf{J}}_{\pm}] = \pm\hbar\tilde{\mathbf{J}}_{\pm} \quad [\tilde{\mathbf{J}}_+, \tilde{\mathbf{J}}_-] = 2\hbar\tilde{\mathbf{J}}_z$
- $\tilde{\mathcal{D}}(\hat{\mathbf{n}}, \phi) = e^{-i\tilde{\mathbf{J}} \cdot \hat{\mathbf{n}}\phi/\hbar} \approx \mathbf{1} - i\tilde{\mathbf{J}} \cdot \hat{\mathbf{n}}\phi/\hbar$
- $\tilde{\mathbf{J}}^2 |j, m\rangle = j(j+1)\hbar^2 |j, m\rangle \quad \tilde{\mathbf{J}}_z |j, m\rangle = m\hbar |j, m\rangle$
 $m = -j, \dots, j$ integer steps.

- Translation

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$$\langle x| (\tilde{\mathbf{1}} - i\tilde{\mathbf{p}}_x\delta x) |\alpha\rangle = \int dx' \langle x|x'\rangle \langle x' - \delta x|\alpha\rangle$$

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$$\Rightarrow \langle x|\tilde{\mathbf{1}} - i\tilde{\mathbf{p}}_x\delta x|\alpha\rangle = \langle x - \delta x|\alpha\rangle$$

$$\left[\tilde{\mathbf{1}} - \frac{i\tilde{\mathbf{p}}_y}{\hbar}(\tilde{\mathbf{x}}\delta\phi) + \frac{i\tilde{\mathbf{p}}_x}{\hbar}(\tilde{\mathbf{y}}\delta\phi) \right] |x', y', z'\rangle = |x' - y'\delta\phi, y' + x'\delta\phi, z'\rangle$$

$$\left[\tilde{\mathbf{1}} - \frac{i\tilde{\mathbf{p}}_y}{\hbar}(\tilde{\mathbf{x}}\delta\phi) + \frac{i\tilde{\mathbf{p}}_x}{\hbar}(\tilde{\mathbf{y}}\delta\phi) \right] |x', y', z'\rangle = |x' - y'\delta\phi, y' + x'\delta\phi, z'\rangle$$

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} \cos \phi & -\sin \phi & 0 \\ \sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix}$$

$$\approx \begin{pmatrix} 1 & -\delta\phi & 0 \\ \delta\phi & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} x' - y'\delta\phi \\ y' + x'\delta\phi \\ z' \end{pmatrix}$$

$$\left[\tilde{\mathbf{1}} - \frac{i\tilde{\mathbf{P}}_y}{\hbar}(\tilde{\mathbf{x}}\delta\phi) + \frac{i\tilde{\mathbf{P}}_x}{\hbar}(\tilde{\mathbf{y}}\delta\phi) \right] |x', y', z'\rangle = |x' - y'\delta\phi, y' + x'\delta\phi, z'\rangle$$

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$$\left\langle x', y', z' \left| \tilde{\mathbf{1}} - \frac{i\tilde{\mathbf{L}}_z\delta\phi}{\hbar} \right| \alpha \right\rangle = \langle x' + y'\delta\phi, y' - x'\delta\phi, z' | \alpha \rangle$$

$$\langle x', y', z' | \alpha \rangle \rightarrow \langle r, \theta, \phi | \alpha \rangle$$

$$\Rightarrow \left\langle r, \theta, \phi \left| \tilde{\mathbf{1}} - \frac{i\tilde{\mathbf{L}}_z \delta\phi}{\hbar} \right| \alpha \right\rangle = \langle r, \theta, \phi - \delta\phi | \alpha \rangle$$

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$$\left. \begin{aligned} \langle r, \theta, \phi - \delta\phi | \alpha \rangle &= \langle r, \theta, \phi | \alpha \rangle - \delta\phi \frac{\partial}{\partial\phi} \langle r, \theta, \phi | \alpha \rangle \\ \left\langle r, \theta, \phi \left| \tilde{\mathbf{1}} - \frac{i\tilde{\mathbf{L}}_z \delta\phi}{\hbar} \right| \alpha \right\rangle &= \langle r, \theta, \phi | \alpha \rangle - \frac{i\delta\phi}{\hbar} \langle r, \theta, \phi | \tilde{\mathbf{L}}_z | \alpha \rangle \end{aligned} \right\}$$

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$$\Rightarrow \langle r, \theta, \phi | \tilde{\mathbf{L}}_z | \alpha \rangle = -i\hbar \frac{\partial}{\partial\phi} \langle r, \theta, \phi | \alpha \rangle.$$

Angular Momentum Operators – *Coordinate Basis*

- $\tilde{\mathbf{L}}_z \doteq -i\hbar \frac{\partial}{\partial \phi}$
- $\tilde{\mathbf{L}}_x \doteq i\hbar \left(\sin \phi \frac{\partial}{\partial \theta} + \cot \theta \cos \phi \frac{\partial}{\partial \phi} \right)$
- $\tilde{\mathbf{L}}_y \doteq -i\hbar \left(\cos \phi \frac{\partial}{\partial \theta} - \cot \theta \sin \phi \frac{\partial}{\partial \phi} \right)$
- $\tilde{\mathbf{L}}_{\pm} = \tilde{\mathbf{L}}_x \pm i\tilde{\mathbf{L}}_y \doteq -i\hbar e^{\pm i\phi} \left(\pm i \frac{\partial}{\partial \theta} - \cot \theta \frac{\partial}{\partial \phi} \right)$
- $\tilde{\mathbf{L}}^2 = \tilde{\mathbf{L}}_z^2 + \left(\tilde{\mathbf{L}}_+ \tilde{\mathbf{L}}_- + \tilde{\mathbf{L}}_- \tilde{\mathbf{L}}_+ \right) / 2 \doteq$
 $-\hbar^2 \left[\frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) \right]$
 - Used $[\tilde{\mathbf{L}}_+, \tilde{\mathbf{L}}_-] = 0$

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Spherical Harmonics – *Derivation*

- Azimuthal coordinate $\tilde{\mathbf{L}}_z |\ell, m\rangle = m\hbar |\ell, m\rangle$
 - $-i\hbar \frac{\partial}{\partial \phi} \Phi(\phi) = m\hbar \Phi(\phi) \Rightarrow \Phi(\phi) \propto e^{im\phi}$
 - Depends only on ϕ .
- Use ladder operators to derive the spectrum
 - $\tilde{\mathbf{L}}_+ |\ell, \ell\rangle = 0$ or $\tilde{\mathbf{L}}_- |\ell, -\ell\rangle = 0$
 - $\left(i\frac{\partial}{\partial \theta} - \cot \theta \frac{\partial}{\partial \phi}\right) \Theta(\theta) \Phi(\phi) = 0$
 - $\left(\frac{\partial}{\partial \theta} - l \cot \theta\right) \Theta(\theta) = 0$
 - $\Theta(\theta) \propto \sin^\ell \theta$
- $Y_\ell^{m=\ell}(\theta, \phi) = c_\ell e^{i\ell\phi} \sin^\ell \theta$
- Normalization $\iint Y_{\ell'}^{m'*}(\theta, \phi) Y_\ell^m(\theta, \phi) d\Omega = \delta_{m'm} \delta_{\ell'\ell}$

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 - $\left(\frac{\partial}{\partial \theta} - l \cot \theta\right) \Theta(\theta) = 0$
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Spherical Harmonics

$$Y_0^0(\theta, \phi) = \left(\frac{1}{4\pi}\right)^{1/2}$$

$$Y_2^{\pm 2}(\theta, \phi) = \left(\frac{15}{32\pi}\right)^{1/2} \sin^2 \theta e^{\pm i 2\phi}$$

$$Y_1^0(\theta, \phi) = \left(\frac{3}{4\pi}\right)^{1/2} \cos \theta$$

$$Y_3^0(\theta, \phi) = \left(\frac{7}{16\pi}\right)^{1/2} (5 \cos^3 \theta - 3 \cos \theta)$$

$$Y_1^{\pm 1}(\theta, \phi) = \mp \left(\frac{3}{8\pi}\right)^{1/2} \sin \theta e^{\pm i \phi}$$

$$Y_3^{\pm 1}(\theta, \phi) = \mp \left(\frac{21}{64\pi}\right)^{1/2} \sin \theta (5 \cos^2 \theta - 1) e^{\pm i \phi}$$

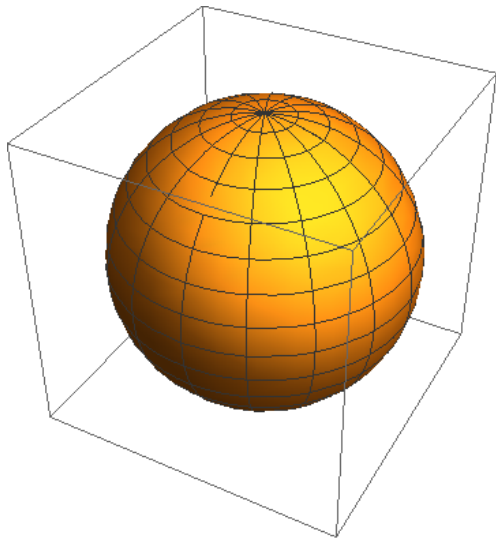
$$Y_2^0(\theta, \phi) = \left(\frac{5}{16\pi}\right)^{1/2} (3 \cos^2 \theta - 1)$$

$$Y_3^{\pm 2}(\theta, \phi) = \left(\frac{105}{32\pi}\right)^{1/2} \sin^2 \theta \cos \theta e^{\pm i 2\phi}$$

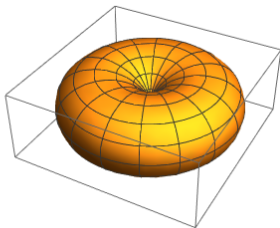
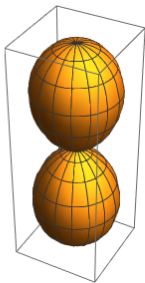
$$Y_2^{\pm 1}(\theta, \phi) = \mp \left(\frac{15}{8\pi}\right)^{1/2} \sin \theta \cos \theta e^{\pm i \phi}$$

$$Y_3^{\pm 3}(\theta, \phi) = \mp \left(\frac{35}{64\pi}\right)^{1/2} \sin^3 \theta e^{\pm i 3\phi}$$

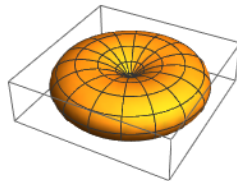
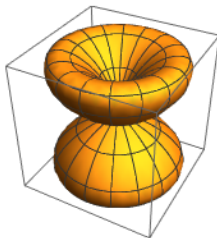
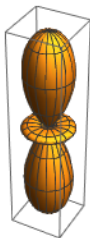
$\ell = 0$ Wavefunction



$\ell = 1$ Wavefunction



$\ell = 2$ Wavefunction



Rotation Operator

- $|\hat{\mathbf{n}}\rangle = \tilde{\mathcal{D}}(R) |\hat{\mathbf{z}}\rangle$
 - A rotation from the z axis to a direction along n axis is given by θ and ϕ .
 - $|\hat{\mathbf{n}}\rangle = \tilde{\mathcal{D}}(\alpha = \phi, \beta = \theta, \gamma = 0) |\hat{\mathbf{z}}\rangle$.
 - $|\hat{\mathbf{n}}\rangle = \sum_{l,m} \tilde{\mathcal{D}}(\alpha = \phi, \beta = \theta, \gamma = 0) |l, m\rangle \langle l, m | \hat{\mathbf{z}}\rangle$
 - $\langle l, m' | \hat{\mathbf{n}}\rangle = \sum_m \langle l, m' | \tilde{\mathcal{D}}(\alpha = \phi, \beta = \theta, \gamma = 0) | l, m\rangle \langle l, m | \hat{\mathbf{z}}\rangle = \sum_m \tilde{\mathcal{D}}_{m'm}^{(l)}(R) \langle l, m | \hat{\mathbf{z}}\rangle$
 - $\langle l, m' | \hat{\mathbf{n}}\rangle = Y_l^{m'*}(\theta, \phi)$
 - $\langle \ell, m | \hat{\mathbf{z}}\rangle = Y_\ell^m(\theta = 0, \phi) = Y_\ell^0(0, \phi) = \sqrt{\frac{2\ell+1}{4\pi}}$
 - $\langle l, m' | \hat{\mathbf{n}}\rangle = \tilde{\mathcal{D}}_{m'0}^{(l)}(\theta, \phi, 0) \sqrt{\frac{2l+1}{4\pi}}$
 - $\tilde{\mathcal{D}}_{m0}^{(l)}(\theta, \phi, 0) = \sqrt{\frac{4\pi}{2l+1}} Y_l^{m*}(\theta, \phi)$

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 - $|\hat{\mathbf{n}}\rangle = \sum_{l,m} \tilde{\mathcal{D}}(\alpha = \phi, \beta = \theta, \gamma = 0) |l, m\rangle \langle l, m | \hat{\mathbf{z}}\rangle$
 - $\langle l, m' | \hat{\mathbf{n}}\rangle = \sum_m \langle l, m' | \tilde{\mathcal{D}}(\alpha = \phi, \beta = \theta, \gamma = 0) | l, m\rangle \langle l, m | \hat{\mathbf{z}}\rangle = \sum_m \tilde{\mathcal{D}}_{m'm}^{(l)}(R) \langle l, m | \hat{\mathbf{z}}\rangle$
 - $\langle l, m' | \hat{\mathbf{n}}\rangle = Y_l^{m'*}(\theta, \phi)$
 - $\langle \ell, m | \hat{\mathbf{z}}\rangle = Y_\ell^m(\theta = 0, \phi) = Y_\ell^0(0, \phi) = \sqrt{\frac{2\ell+1}{4\pi}}$
 - $\langle l, m' | \hat{\mathbf{n}}\rangle = \tilde{\mathcal{D}}_{m'0}^{(l)}(\theta, \phi, 0) \sqrt{\frac{2l+1}{4\pi}}$
 - $\tilde{\mathcal{D}}_{m0}^{(l)}(\theta, \phi, 0) = \sqrt{\frac{4\pi}{2l+1}} Y_l^{m*}(\theta, \phi)$

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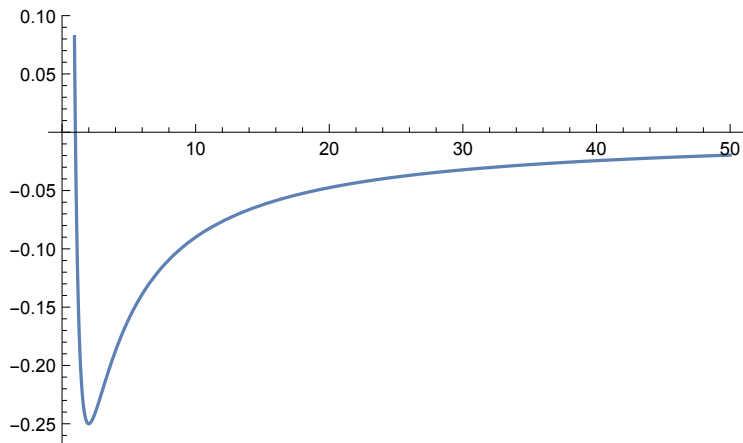
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Effective Potential – Coulomb



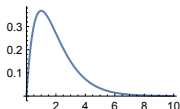
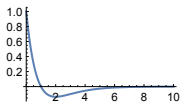
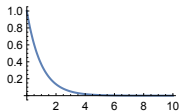
Confluent Hypergeometric Function — Expansion

$$\begin{aligned} {}_1F_1(a; c; x) = & 1 + \frac{ax}{c} + \frac{a(a+1)x^2}{2c(c+1)} + \frac{a(a+1)(a+2)x^3}{6c(c+1)(c+2)} \\ & + \frac{a(a+1)(a+2)(a+3)x^4}{24c(c+1)(c+2)(c+3)} \\ & + \frac{a(a+1)(a+2)(a+3)(a+4)x^5}{120c(c+1)(c+2)(c+3)(c+4)} + \dots \end{aligned}$$

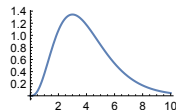
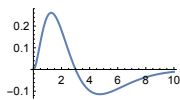
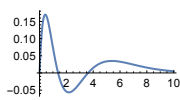
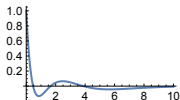
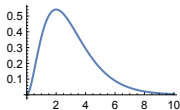
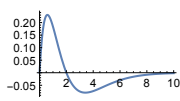
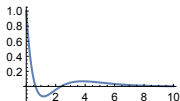
Radial Wavefunction

$$R_{n\ell}(r) = \frac{1}{(2\ell+1)!} \left(\frac{2Zr}{na_0} \right)^\ell e^{-Zr/na_0} \left[\left(\frac{2Z}{na_0} \right)^3 \frac{(n+\ell)!}{2n(n-\ell-1)!} \right]^{1/2} \\ \times {}_1F_1(-n+\ell+1; 2\ell+2; 2Zr/na_0)$$

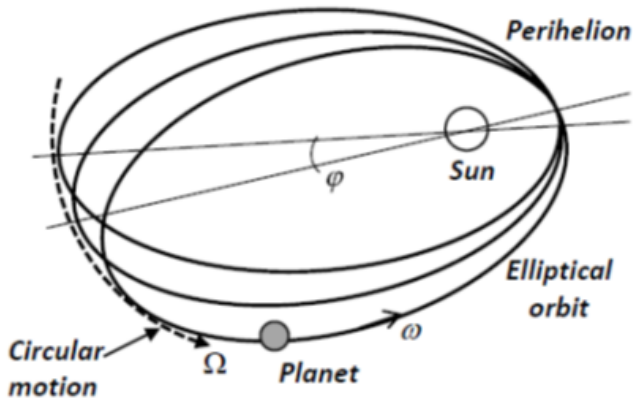
Radial Wavefunction



n	N, ℓ	N, ℓ	N, ℓ	N, ℓ
1	0, 0	-	-	-
2	1, 0	0, 1	-	-
3	2, 0	1, 1	0, 2	-
4	3, 0	2, 1	1, 2	0, 3



Orbit Precession



Laplace-Runge-Lenz Vector

