

## An extra example: spiralling football.

\* Let's consider the rotational motion of a football thrown by a quarterback.

\* We assume  $\rightarrow$  no applied torque (ignore gravity)

A simple treatment is given by Euler's eqns (defined w.r.t. body-fixed axes):

$$I_1 \dot{\omega}_1 + \frac{I_3 - I_2}{2} \omega_2 \omega_3 = 0$$

$$I_2 \dot{\omega}_2 + (I_1 - I_3) \omega_1 \omega_3 = 0$$

$$I_3 \dot{\omega}_3 + (I_2 - I_1) \omega_1 \omega_2 = 0$$

A football has cylindrical symmetry such that

$$I_1 = I_2 = I_{\perp} \neq I_3$$

Consequence

$$\Rightarrow \dot{\omega}_3 = 0 \rightarrow \omega_3 \text{ is a const.}$$

This enables us to obtain  $\omega_1(t)$  &  $\omega_2(t)$  in a straightforward manner.

ie.  $\omega_1(t) = C \sin(\Omega t + \delta)$

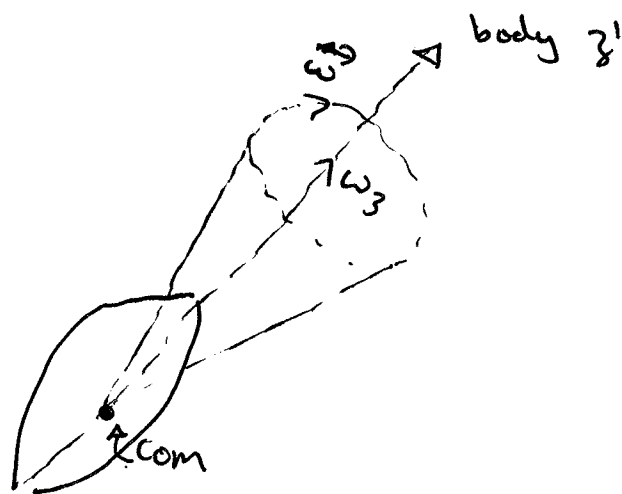
w/  $C = (\omega_{1i}^2 + \omega_{2i}^2)^{1/2}$

$\omega_2(t) = C \cos(\Omega t + \delta)$

$\delta = \arctan\left(\frac{\omega_{1i}}{\omega_{2i}}\right)$

$\omega_3 = \omega_{3i}$

In the body-fixed frame:



$\vec{\omega}$  precesses in  $x'y'$  plane about  $z'$

[In space fixed frame football 'wobbles']

Note also  $\rightarrow$  it is torques acting on football

$\rightarrow$  angular momentum  $\vec{L} = \text{const}$  in space frame. (inertial)

What is the role of precession (wobble, of football?) and it's value?

Proceed by expressing  $\vec{L}$  in body-fixed frame:

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$$\left. \begin{aligned} L_1 &= I_{\perp} \omega_1 = \frac{1}{2} l \sin \theta \sin \psi \\ L_2 &= I_{\perp} \omega_2 = l \sin \theta \cos \psi \\ L_3 &= I_3 \omega_3 = l \cos \theta \end{aligned} \right\} \begin{array}{l} \leftarrow \\ 1) \\ \left( \text{we assume} \right. \\ \left. \vec{L} \text{ is along space} \right. \\ \left. \text{fixed } \hat{z} \right) \end{array}$$

Separately we have instantaneous angular velocity,

$$\vec{\omega} = \begin{pmatrix} \dot{\phi} \sin \theta \sin \psi + \dot{\theta} \cos \psi \\ \dot{\phi} \sin \theta \cos \psi - \dot{\theta} \sin \psi \\ \dot{\phi} \cos \theta + \dot{\psi} \end{pmatrix} \quad \begin{array}{l} 2) \\ \text{Note } \underline{\dot{\theta} = 0} \\ \text{(see diagram!)} \end{array}$$

The rate of precession:



$$\omega_p = \dot{\phi}$$

With  $\dot{\theta} = 0 \rightarrow$  clearly we have:

$$L_1 \Rightarrow I_{\perp} \dot{\phi} \sin \theta \sin \psi = l \sin \theta \sin \psi$$

$$\rightarrow \omega_p = l / I_{\perp}$$

Evaluating  $L$  with initial conditions then yields,

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(\*)

$$\omega_p = \omega_{3i} \left[ \left( \frac{I_3}{I_d} \right)^2 + \frac{(\omega_{1i}^2 + \omega_{2i}^2)}{\omega_{3i}^2} \right]^{1/2}$$

And clearly,

$$\cos \theta = L_3 / L = \left[ 1 + \left( \frac{I_d}{I_3} \right)^2 \frac{(\omega_{1i}^2 + \omega_{2i}^2)}{\omega_{3i}^2} \right]^{-1/2}$$

Important consequences:

i)  $\theta \rightarrow 0$  for  $\omega_{3i} \gg \sqrt{\omega_{1i}^2 + \omega_{2i}^2}$   
 i.e. fast spinning pass

ii)  $\frac{\omega_p}{\omega_{3i}} \rightarrow \frac{I_3}{I_d}$

depends only on properties of ball

Wobble to spin ratio  
 (typically  $\sim 3/5$  for current footballs)

$$I_3 \propto \int dV x^2 + y^2$$

$$I_d \propto \frac{1}{2} \int dV 2r^2 - z^2$$

$\therefore$  want narrow ball!

(\*) compute  $L$  using  $\omega_1(t) + \omega_2(t) + \omega_3(t)$  211e

$$L^2 = I_1^2 \omega_1^2 + I_2^2 \omega_2^2 + I_3^2 \omega_3^2$$

$$= I_{\perp}^2 c^2 + I_3^2 \omega_3^2$$

$$= I_{\perp}^2 \omega_3^2 \left( \left( \frac{I_3}{I_{\perp}} \right)^2 + \frac{c^2}{\omega_3^2} \right)$$

$$\Rightarrow \frac{L}{I_{\perp}} = \cancel{\omega_3} \left( \left( \frac{I_3}{I_{\perp}} \right)^2 + \frac{c^2}{\omega_3^2} \right)^{1/2} \quad \checkmark$$