

Quantum Computing

Problem sheet 1

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2022-04-09

Exercise 1.1 (Complex vector spaces)

a)

First we simplify:

$$\begin{aligned}z_1 &= 3 - i + (2 - i)(-1 + i) \\&= 3 - i + (-2 + 2i) + (i + 1) \\&= 3 - i + (-1 + 3i) \\&= 3 + (-1 + 2i) \\&= 2 + 2i\end{aligned}$$

This gives us the real and imaginary part $\Re(z_1) = \Im(z_1) = 2$, as well as the complex conjugate $z_1^* = 2 - 2i$ and by pythagoras theorem the absolute value $\sqrt{2^2 + 2^2} = \sqrt{8}$.

b)

$$\begin{aligned}& \prod_{k=-4}^4 \sqrt{36-k^2} + ik \\&= (\sqrt{36} + 0i) \cdot \prod_{\substack{k=-4 \\ k \neq 0}}^4 \sqrt{36-k^2} + ik \\&= 6 \cdot \prod_{k=1}^4 \underbrace{(\sqrt{36-k^2} + ik)(\sqrt{36-k^2} - ik)}_{(36-k^2) - \cancel{ik(36-k^2)} + \cancel{ik(36-k^2)} - \underbrace{i^2 k^2}_{+k^2}} \\&= 6 \cdot (36-1)(36-2)(36-3)(36-4)\end{aligned}$$

c)

Given that $\{|0\rangle, |1\rangle, \dots, |d-1\rangle\}$ is an ONB, we also know that $\{\langle 0|, \langle 1|, \dots, \langle d-1|\}$ is an ONB of the dual-space \mathbb{C}^{d*} with $\langle i|j\rangle = \delta_{ij}$.

Therefore given a ket $|\psi\rangle$ which has a decomposition into basis vectors $\sum_{k=0}^{d-1} \psi_k |k\rangle$, because \mathbb{C}^{d*} is a vector space, we can find the ψ_i by

$$\begin{aligned}\langle i|\psi\rangle &= \langle i|\sum_{k=0}^{d-1} \phi_k |k\rangle \\ &= \delta_{0i}\psi_0 + \dots + \delta_{ii}\psi_i + \dots + \delta_{i(d-1)}\psi_{d-1} \\ &= \psi_i\end{aligned}$$

d)

A Hilbert space H is a real or complex inner product space that is also a complete metric space with respect to the distance function induced by the inner product.

We already know that \mathbb{C}^d is a complete vector space under the euclidean norm. Moreover it is also a complete vector space under the induced distance function of the inner product, defined by:

$$\langle x, y \rangle = y^\dagger x = \bar{y}^t x$$

we therefore only need to check that $\langle \cdot, \cdot \rangle$ is an inner product.

1. We have *conjugate symmetry*:

$$\begin{aligned}\overline{\langle y, x \rangle} &= \overline{x^t y} \\ &= x^t \bar{y} \\ &= \bar{y}^t x \\ &= y^\dagger x \\ &= \langle x, y \rangle\end{aligned}$$

2. We have *linearity in the first argument*:

$$\begin{aligned}
\langle \lambda x_1 + \mu x_2, y \rangle &= \bar{y}^t (\lambda x_1 + \mu x_2) \\
&= \begin{pmatrix} \bar{y}_1 & \dots & \bar{y}_d \end{pmatrix} \begin{pmatrix} \lambda x_{1_1} + \mu x_{2_1} \\ \vdots \\ \lambda x_{1_d} + \mu x_{2_d} \end{pmatrix} \\
&= \begin{pmatrix} \bar{y}_1 & \dots & \bar{y}_d \end{pmatrix} \left(\begin{pmatrix} \lambda x_{1_1} \\ \vdots \\ \lambda x_{1_d} \end{pmatrix} + \begin{pmatrix} \mu x_{2_1} \\ \vdots \\ \mu x_{2_d} \end{pmatrix} \right) \\
&= \begin{pmatrix} \bar{y}_1 & \dots & \bar{y}_d \end{pmatrix} \begin{pmatrix} \lambda x_{1_1} \\ \vdots \\ \lambda x_{1_d} \end{pmatrix} + \begin{pmatrix} \bar{y}_1 & \dots & \bar{y}_d \end{pmatrix} \begin{pmatrix} \mu x_{2_1} \\ \vdots \\ \mu x_{2_d} \end{pmatrix} \\
&= \lambda \begin{pmatrix} \bar{y}_1 & \dots & \bar{y}_d \end{pmatrix} \begin{pmatrix} x_{1_1} \\ \vdots \\ x_{1_d} \end{pmatrix} + \mu \begin{pmatrix} \bar{y}_1 & \dots & \bar{y}_d \end{pmatrix} \begin{pmatrix} x_{2_1} \\ \vdots \\ x_{2_d} \end{pmatrix} \\
&= \lambda \langle x_1, y \rangle + \mu \langle x_2, y \rangle
\end{aligned}$$

3. We have that the inner product of an element with itself is *positive-definite*:

$$\begin{aligned}
\langle x, x \rangle &= \bar{x}^t x \\
&= \begin{pmatrix} \bar{x}_1 & \dots & \bar{x}_d \end{pmatrix} \begin{pmatrix} x_1 \\ \vdots \\ x_d \end{pmatrix} \\
&= \underbrace{\underbrace{\bar{x}_1 x_1}_{=0 \Leftrightarrow x_1=0} + \dots + \underbrace{\bar{x}_d x_d}_{=0 \Leftrightarrow x_d=0}}_{=0 \Leftrightarrow x=0} \\
&\quad \quad \quad >0 \forall x
\end{aligned}$$

Exercise 1.2 (Postulates: Quantum states and their transformations)

a)

$$\begin{aligned}
Y &= \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix} \\
Y^\dagger &= \begin{pmatrix} \bar{1} & \bar{-i} \\ \bar{i} & \bar{1} \end{pmatrix}^t = \begin{pmatrix} 1 & i \\ -i & 1 \end{pmatrix}^t = \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix} = Y
\end{aligned}$$

b)

$$M \in U(\mathbb{C}^{3 \times 3}) \Leftrightarrow M^\dagger M = \mathbb{1}_{\mathbb{C}^3}$$

$$\begin{aligned}
\det(M) &= \cos(\theta)^2 e^{i\phi} - i \sin(\theta) i \sin(\theta) e^{i\phi} \\
&= e^{i\phi} (\cos(\theta)^2 + \sin(\theta) \sin(\theta)) \\
&= e^{i\phi} (\cos(\theta)^2 + \sin(\theta)^2) \\
&= e^{i\phi}
\end{aligned}$$

Therefore $|\det(M)| \neq 1$ for all $\phi, \theta \in \mathbb{R}$

$$\begin{aligned}
\begin{pmatrix} \cos(\theta) & 0 & -i \sin(\theta) \\ 0 & e^{i\phi} & 0 \\ -i \sin(\theta) & 0 & \cos(\theta) \end{pmatrix}^\dagger &= \begin{pmatrix} \cos(\theta) & 0 & i \sin(\theta) \\ 0 & e^{-i\phi} & 0 \\ i \sin(\theta) & 0 & \cos(\theta) \end{pmatrix} \\
&= \begin{pmatrix} \cos(\theta) & 0 & -i \sin(\theta) \\ 0 & e^{i\phi} & 0 \\ -i \sin(\theta) & 0 & \cos(\theta) \end{pmatrix} \begin{pmatrix} \cos(\theta) & 0 & i \sin(\theta) \\ 0 & e^{-i\phi} & 0 \\ i \sin(\theta) & 0 & \cos(\theta) \end{pmatrix} \\
&= \begin{pmatrix} \cos(\theta)^2 + \sin(\theta)^2 & 0 & i \cos(\theta) \sin(\theta) - i \sin(\theta) \cos(\theta) \\ 0 & e^{i\phi} e^{-i\phi} & 0 \\ -i \sin(\theta) \cos(\theta) + \cos(\theta) i \sin(\theta) & 0 & -i^2 \sin(\theta)^2 + \cos(\theta)^2 \end{pmatrix} \\
&= \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\phi - i\phi} & 0 \\ 0 & 0 & 1 \end{pmatrix} \\
&= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}
\end{aligned}$$

Finally we can conclude that for all $\phi, \theta \in \mathbb{C}$ the matrix M is unitary.

c)

$$H|0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$H|1\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

There cannot be another matrix that describes the same linear transformation. I am unsure what is meant with: > any other unitary matrices that describe the same transformation of states? Does this mean the same linear transformation?

1.3 (Observables)

a)

We need $M^\dagger = M$ or equivalently $a_{ij} = \bar{a}_{ji}$ for a matrix to be hermitian. For the special case

$$\begin{bmatrix} a & b \\ c & d \end{bmatrix}$$

we have $a = \bar{a}$, $b = \bar{c}$ and $d = \bar{d}$

b)

Assume M is self-adjoint and given a choice of orthonormal basis vectors $(e_i)_{1 \leq i \leq d}$ we choose the following representation

$$|\phi\rangle = \begin{bmatrix} \psi_1 \\ \vdots \\ \psi_d \end{bmatrix} \quad \langle\phi| = [\bar{\psi}_1 \dots \bar{\psi}_d] \quad M = \begin{bmatrix} m_{11} & m_{12} & \dots & m_{1d} \\ \overline{m_{12}} & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ \overline{m_{1d}} & \dots & \dots & m_{dd} \end{bmatrix}$$

We can now calculate the following using the fact that M is a linear operator¹ :

$$\begin{aligned} \langle\psi| M |\psi\rangle &= \langle\psi| (M |\psi\rangle) = (\langle\psi| M) |\psi\rangle \\ &= \psi^\dagger (M\psi) \\ &= [\bar{\psi}_1 \dots \bar{\psi}_d] \begin{bmatrix} m_{11} & m_{12} & \dots & m_{1d} \\ \overline{m_{12}} & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ \overline{m_{1d}} & \dots & \dots & m_{dd} \end{bmatrix} \begin{bmatrix} \psi_1 \\ \vdots \\ \psi_d \end{bmatrix} \\ &= [\bar{\psi}_1 \dots \bar{\psi}_d] \begin{bmatrix} m_{11}\psi_1 + \dots + m_{1d}\psi_d \\ \vdots \\ \overline{m_{1d}}\psi_1 + \dots + m_{dd}\psi_d \end{bmatrix} \\ &= \sum_{k=1}^d \bar{\psi}_k \left(\sum_{l=1}^d m_{kl}\psi_l \right) \\ &= \underbrace{\sum_{k=1}^d m_{kk}\bar{\psi}_k\psi_k}_{\in \mathbb{R}} + \underbrace{\sum_{k=2}^d \sum_{l=1}^{k-1} m_{kl}\bar{\psi}_k\psi_l + \overline{m_{kl}}\bar{\psi}_l\psi_k}_{\in \mathbb{R}} \end{aligned}$$

The right underbrace can be checked by writing $m_{kl} = (a + bi)$, $\psi_k = (c + di)$, ... and then multiplying out.

¹https://en.wikipedia.org/wiki/Bra%E2%80%93ket_notation#Linear_operators

c)

$$\begin{aligned}\langle 1 | \sigma_x | 1 \rangle &= \begin{bmatrix} 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \begin{bmatrix} 0 & 1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= 0\end{aligned}$$

$$\begin{aligned}\langle 1 | \sigma_y | 1 \rangle &= \begin{bmatrix} 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \begin{bmatrix} 0 & 1 \end{bmatrix} \begin{bmatrix} -i \\ 0 \end{bmatrix} \\ &= 0\end{aligned}$$

$$\begin{aligned}\langle 1 | \sigma_z | 1 \rangle &= \begin{bmatrix} 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \begin{bmatrix} 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ -1 \end{bmatrix} \\ &= -1\end{aligned}$$

$$\begin{aligned}\langle 0 | \sigma_x | 0 \rangle &= \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= 0\end{aligned}$$

$$\begin{aligned}\langle 0 | \sigma_y | 0 \rangle &= \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ i \end{bmatrix} \\ &= 0\end{aligned}$$

$$\begin{aligned}\langle 0 | \sigma_z | 0 \rangle &= \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= 1\end{aligned}$$

d)

Due to the associativity of matrix-multiplication, we can calculate $\langle v | (XY) | v \rangle$ instead of $\langle v | X(Y | v \rangle)$. Therefore applying the Pauli observables to $H | 0 \rangle$ and $H | 1 \rangle$ yields:

$$\begin{aligned} \langle 1 | (\sigma_x H) | 1 \rangle &= [0 \quad 1] \left(\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \right) \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} [0 \quad 1] \begin{bmatrix} 1 & -1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} [0 \quad 1] \begin{bmatrix} -1 \\ 1 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} \end{aligned}$$

$$\begin{aligned} \langle 1 | (\sigma_y H) | 1 \rangle &= [0 \quad 1] \left(\begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \right) \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} [0 \quad 1] \begin{bmatrix} -i & i \\ i & i \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \frac{i}{\sqrt{2}} \end{aligned}$$

$$\begin{aligned} \langle 1 | (\sigma_z H) | 1 \rangle &= [0 \quad 1] \left(\begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \right) \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} [0 \quad 1] \begin{bmatrix} 1 & 1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} \end{aligned}$$

$$\begin{aligned} \langle 0 | (\sigma_x H) | 0 \rangle &= [1 \quad 0] \left(\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \right) \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} [1 \quad 0] \begin{bmatrix} 1 & -1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} \end{aligned}$$

$$\begin{aligned} \langle 0 | (\sigma_y H) | 0 \rangle &= [1 \quad 0] \left(\begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \right) \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} [1 \quad 0] \begin{bmatrix} -i & i \\ i & i \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= \frac{-i}{\sqrt{2}} \end{aligned}$$

$$\begin{aligned}
\langle 0 | (\sigma_z H) | 0 \rangle &= \begin{bmatrix} 1 & 0 \end{bmatrix} \left(\begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \right) \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\
&= \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} 1 & 1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\
&= \frac{1}{\sqrt{2}}
\end{aligned}$$