## Supplementary Material for Slice Sampling on Hamiltonian Trajectories

## 1. Hamiltonian Dynamics for Non-Gaussian Distributions

Assume we have some prior  $\pi$  for f, where

$$-\log \pi(\mathbf{f}|\alpha) = \log Z - g(\mathbf{f}|\alpha), \tag{1}$$

for some hyperparameters  $\alpha$  and where Z is the normalizing constant. Then we can again set up a Hamiltonian,

$$H(\mathbf{f}, \mathbf{p}) = g(\mathbf{f}|\boldsymbol{\alpha}) + \frac{1}{2}\mathbf{p}^{T}M^{-1}\mathbf{p}.$$
 (2)

Hamilton's equations yield the system of differential equations

$$\nabla_t^2 \mathbf{f} = -M^{-1} \nabla_{\mathbf{f}} g(\mathbf{f} | \boldsymbol{\alpha}). \tag{3}$$

This can be solved exactly for some cases of  $g(\mathbf{f}|\alpha)$ ; the Gaussian example above is one such case.

When  $M = \operatorname{diag}(m_1, \ldots, m_d)$  and  $g(\mathbf{f}|\alpha) = \sum_{i=1}^d h_i(f_i)$ , the system is an uncoupled system and often has an analytic solution. In particular,  $f_i(t)$  is the solution to

$$\frac{1}{2} \left( \int \left[ c_1 - \frac{1}{m_i} h_i(f_i) \right]^{-1/2} df_i \right)^2 = (t + c_2)^2, \quad (4)$$

where  $c_2$  is determined by  $f_i(0)$  and

$$c_1 = \frac{1}{2} \left( \dot{f}_i(0) \right)^2 + \frac{1}{m_i} h_i(f_i(0)), \tag{5}$$

where  $\dot{f}_i(t)$  denotes the time-derivative of  $f_i$  at time t. For certain distributions, (4) has an analytic solution.

## **1.1.** Exp( $\lambda$ ) and Laplace( $\lambda$ )

Working with the general form of the solution can be difficult, especially when the sample space is constrained. The exponential distribution is one such example, but it induces a potential for which solutions are easy to obtain directly. Let the prior be such that the components of  $\mathbf{f}$  are mutually independent and  $f_i \sim \operatorname{Exp}(\lambda_i)$  so that

$$-\log \pi(f_i|\lambda_i) = -\log(\lambda_i) + \lambda_i \cdot f_i, \quad f_i > 0.$$
 (6)

Hamilton's equations are particularly simple in this case, which describes one-dimensional projectile motion, e.g. a

bouncing ball in a constant gravitational potential. The solution to Hamilton's equations is

$$f_t(t) = \frac{\lambda_i}{2m_i} t^2 + \dot{f}_i(0)t + f_i(0), \quad 0 \le t \le T_0.$$
 (7)

An example of such a trajectory is shown in Figure 1a.  $T_0>0$  is the time at which the particle has position coordinate equal to 0, at which point its momentum changes signs, i.e. it "bounces." For  $t>T_0$ , the particle repeatedly traces out the same trajectory. We find  $T_0$  as

$$T_0 = \frac{m_i}{\lambda_i} \dot{f}_i(0) + \sqrt{\frac{m_i^2}{\lambda_i^2} \dot{f}_i(0)^2 + \frac{2m_i}{\lambda_i} f_i(0)}.$$
 (8)

This yields the period of the trajectory,

$$T = 2\sqrt{\frac{m_i^2}{\lambda_i^2}\dot{f}_i(0)^2 + \frac{2m_i}{\lambda_i}f_i(0)}.$$
 (9)

Every time at which the particle reaches zero is then  $z_j = T_0 + (j-1)T$ . Hamilton's equations also yield the momentum,

$$p_i(t) = -\lambda_i t + m_i \dot{f}_i(0), \tag{10}$$

which we can use to find the momentum at the reflection point  $T_0$ , but before reflection,

$$p_i(T_0^-) = -\sqrt{m_i^2 \dot{f}_i(0)^2 + 2m_i \lambda_i f_i(0)}.$$
 (11)

After the first reflection, we have  $p_i(T_0^+) = -p_i(T_0^-)$ , and the dynamics proceed according to the equation

$$f_i(t) = -\frac{\lambda_i}{2m_i}(t - z_j)^2 + \frac{p_i(T_0^+)}{m_i}(t - z_j), \quad z_j \le t \le z_{j+1}.$$
(12)

We can use slightly different equations to describe the dynamics under a Laplace prior All that is required is some bookkeeping on the sign of the motion because the particle is not reflected at  $f_i = 0$ , but the sign on the potential switches. An example is shown in Figure 1b.

## **1.2.** Pareto $(x_m, \alpha)$ and GPD $(\mu, \sigma, \xi)$ via transformation

The Pareto and Generalized Pareto (denoted GPD) distributions are typically used to model processes with heavy

tails. The density of the Pareto distribution is

$$\pi(f_i|x_m,\alpha) = \alpha x_m^{\alpha} f_i^{-(\alpha+1)},$$
  
$$f_i \ge x_m, \quad x_m > 0, \quad \alpha > 0.$$
 (13)

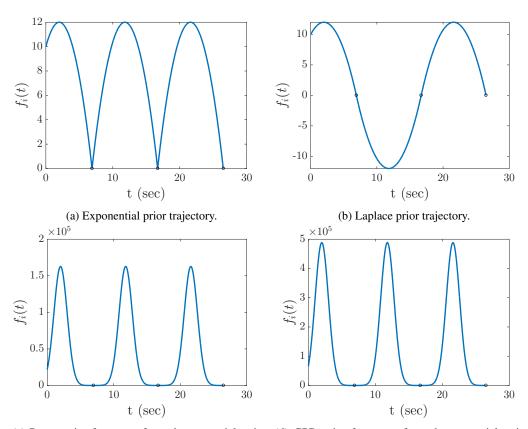
We can show that the random variable  $y_i := \log f_i - \log x_m$  is distributed as  $\operatorname{Exp}(\alpha)$ , for  $y_i > 0$ . Using this fact, we can generate analytic trajectories as in subsection 1.1 for  $y_i$  and slice sample from the resulting curves  $f_i(t) = x_m e^{y_i(t)}$ . An example is shown in Figure 1c.

The GPD has the density

$$\pi(f_i|\mu,\sigma,\xi) = \frac{1}{\sigma} \left( 1 + \xi \cdot \frac{f_i - \mu}{\sigma} \right)^{-(1+\xi^{-1})},$$

$$f_i \ge \mu, \quad \xi \ge 0. \tag{14}$$

(The GPD is also defined for  $\xi < 0$ , in which case  $\mu \le f_i \le \mu - \sigma/\xi$ . We focus on the  $\xi \ge 0$  case for now.) The random variable  $y_i := \log\left(1 + \xi \frac{f_i - \mu}{\sigma}\right)$  is distributed as  $\operatorname{Exp}(\xi^{-1})$  for  $y_i > 0$ , and we can again use the calculations from subsection 1.1 to slice sample from the curve  $f_i(t) = \mu + \frac{\sigma}{\xi}(e^{y_i(t)} - 1)$ . An example is shown in Figure 1d.



(c) Pareto prior from transformed exponential trajec- (d) GPD prior from transformed exponential trajectory.

Figure 1. Example trajectories under different priors.