

(2)

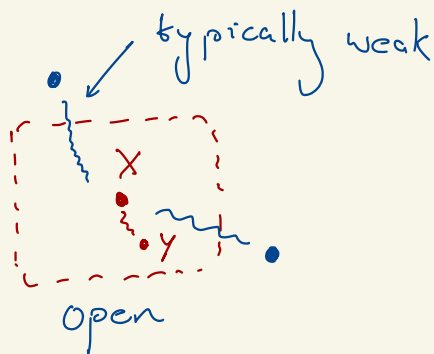
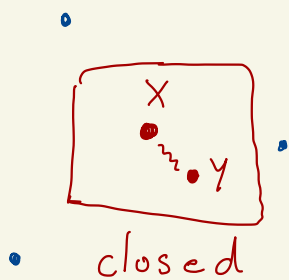
Postulate 2: The evolution of a closed \*

quantum system is described by a unitary transformation. Specifically, the state  $|\psi\rangle$  of the system at time  $t_1$  is related to the state  $|\psi'\rangle$  at  $t_2$  by a unitary operator  $U$  which depends only on  $t_1$  &  $t_2$ .

$$|\psi(t_2)\rangle = U(t_1, t_2) |\psi(t_1)\rangle$$

$t_1 \rightarrow t_2$

\* NB: The assumption of a closed quantum system is a simplification. In general quantum systems will interact with their environment.



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We can write down a refined version of postulate 2...

Postulate 2\*: The time evolution of the state of a closed quantum system is described by the Schrödinger equation...

$$i\hbar \frac{d|\psi(t)\rangle}{dt} = H|\psi(t)\rangle \quad \text{(SE)}$$

where

- $\hbar$  is Planck's constant (normally set  $\hbar=1$ )

- $H$  is a Hermitian operator called the system Hamiltonian.

↳ If we know the Hamiltonian & initial state  $|\psi(0)\rangle$  we can solve (SE) to find  $|\psi(t)\rangle$  for all  $t$ .

We can now obtain postulate 2 from postulate 2\*... (4)

↳ for simplicity we assume  $H$  is time independent...

↳ In this case we can solve (SE) to obtain

$\hbar=1$

$$|\psi(t_2)\rangle = e^{-iH(t_2-t_1)} |\psi(t_1)\rangle$$

$$= U(t_1, t_2) |\psi(t_1)\rangle$$

$$\text{Crux: } U(t_1, t_2) = e^{-iH(t_2-t_1)}$$

Matrix exponential!

## Exercises

(1) Prove  $U(t_1, t_2)$  is unitary.

(2) Write down an expression for  $U(t_1, t_2)$  when  $H = H(t)$  is time dependent.

# Matrix exponential

⑤

$$\left[ e^X := \sum_{k=0}^{\infty} \frac{1}{k!} X^k \right] \text{Definition}$$

↳ Properties

- $e^0 = \mathbb{1}$
- $e^{X^T} = (e^X)^T$
- if  $Y$  is invertible  $e^{YXY^{-1}} = Y e^X Y^{-1}$
- $[X, Y] = 0 \Rightarrow e^{X+Y} = e^X e^Y$

DANGER!  $[X, Y] \neq 0 \xrightarrow[\text{be}]{\text{can}} e^{X+Y} \neq e^X e^Y$

↓  
commutator!

↳ •  $[X, Y] = XY - YX$

•  $\{X, Y\} = XY + YX \rightarrow \text{anti-commutator}$

# Terminology

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↳ if  $\underbrace{[X, Y] = 0}$  we say  $X$  &  $Y$  commute

↳  $XY = YX \Rightarrow$  order of application does n't matter.

→ if  $\{X, Y\} = 0$  we say  $X$  &  $Y$  anti-commute.

Before moving on let's check our soln to (SE)

↳ Claim:  $|\psi(t)\rangle = e^{-iHt} |\psi(0)\rangle$  is soln to (SE)

Proof: 
$$i \frac{d|\psi(t)\rangle}{dt} = i \frac{d}{dt} \left( e^{-iHt} |\psi(0)\rangle \right)$$

$$\boxed{\hbar = 1}$$

$$= i(-iH) \underbrace{\left[ e^{-iHt} |\psi(0)\rangle \right]}_{|\psi(t)\rangle}$$

$$= H|\psi(t)\rangle$$

□

Note:

⑦

- ① For all Hermitian operators  $H$  &  $\forall t \geq 0$  the operator

$$U(t) = e^{-iHt}$$

is unitary.

- ② For all unitary operators  $U(t)$  there exists some Hermitian operator  $H$  such that

$$U(t) = e^{-iH(t)}$$

\*

Crux: We can think of all unitary evolutions as arising from evolution under some Hamiltonian  $H$

↳ Designing quantum gates means engineering system Hamiltonians!

\* Think about group  $SU(n)$  with lie-algebra  $su(n)$

**NB** → Figuring out the Hamiltonian of a system, or designing a system with a specific Hamiltonian, is not easy!

→ Requires intuition + experiments!

## Examples

- magnet
- molecule
- laser
- superconductor
- photosynthetic complex

Guess  
+  
experiment

Developing methods to learn Hamiltonians from system measurements is an active field!!

Recall that the Hamiltonian  $H$  is  
a Hermitian operator...

⑨

∴ we can write

normalized eigenvector

$$H = \sum_i E_i |E_i\rangle \langle E_i|$$

real eigenvalues

**NB**  $H|E_i\rangle = E_i |E_i\rangle$

We call  $|E_i\rangle$  a stationary state :  
∴ set  $|\psi(0)\rangle = |E_i\rangle \dots$

$$\begin{aligned} |\psi(t)\rangle &= e^{-iHt} |E_i\rangle \\ &= \sum_{k=0}^{\infty} \frac{1}{k!} (-i)^k H^k |E_i\rangle \\ &= \sum_k \frac{1}{k!} (-i)^k E_i^k |E_i\rangle \\ &= \underline{e^{-iE_i t}} |E_i\rangle \end{aligned}$$

∴ Global phase - cannot be observed!



# Observables

①

↳ In quantum mechanics observables are described by Hermitian matrices acting on the state space of the system.

⇒ Given a Hermitian operator  $O$  we can always write

$$\begin{aligned} O &= \sum_m E_m \underbrace{|E_m\rangle\langle E_m|}_{\substack{\text{projector onto} \\ m\text{'th eigenspace}}} \\ &= \sum_m E_m \underbrace{P_m}_{\substack{\text{projector onto} \\ m\text{'th eigenspace}}} \end{aligned}$$

where  $\{P_m\}$  is a valid projective measurement.

// By "measure  $O$ " we mean "do the projective measurement  $\{P_m\}$ ".

↳ get eigenvalue  $E_m$  with prob  $\|P_m|\psi\rangle\|^2$

// We often set  $E_m = m \Rightarrow O = \sum_m m P_m$

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We will often care about the average value when measuring  $\hat{O}$ .

↳ call it expectation value of  $\hat{O}$ .

$$\hookrightarrow E(\hat{O}) = \sum_m m p(m)$$

$$= \sum_m m \langle \psi | P_m | \psi \rangle$$

$$= \langle \psi | \sum_m m P_m | \psi \rangle$$

$$= \langle \psi | \hat{O} | \psi \rangle$$

$$\left. \begin{aligned} \hat{O} &= \sum_m m P_m \\ &= \sum_m m |E_m\rangle \langle E_m| \end{aligned} \right\}$$

$$\boxed{E(\hat{O}) = \langle \psi | \hat{O} | \psi \rangle \equiv \langle \hat{O} \rangle}$$

We also care about the standard deviation...

$$[\Delta(\hat{O})]^2 = \langle (\hat{O} - \langle \hat{O} \rangle)^2 \rangle$$

$$= \langle \hat{O}^2 \rangle - \langle \hat{O} \rangle^2$$

Exercise: Suppose we have an observable (13)  
 $O$  &  $| \psi \rangle$  is an eigensate of  $O$  with eigenvalue  $m$ . What is  $\langle O \rangle$  &  $\Delta(O)$  when measuring  $| \psi \rangle$ ?

Example:  $Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$$= 1 \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} - 1 \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$= 1 P_+ + (-1) P_-$$

↳ "Measure  $Z$ " means do the projective measurement  $\{P_+, P_-\}$

Note:  $P_+ |1\rangle = 0$                        $P_+ |0\rangle = |0\rangle$

$P_- |1\rangle = |1\rangle$                        $P_- |0\rangle = 0$

if  $| \psi \rangle = | + \rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$

$$\langle Z \rangle = \langle \psi | Z | \psi \rangle$$

$$= \frac{1}{\sqrt{2}} (\langle 0 | + \langle 1 |) (|0\rangle + |1\rangle) \frac{1}{\sqrt{2}}$$

$$= \frac{1}{2} \langle 0 | 0 \rangle - \frac{1}{2} \langle 1 | 1 \rangle$$

$$= \underline{0}$$

Let's go back to Hamiltonians...

(14)

↳ Hamiltonians are Hermitian

↳  $\therefore$  They are observables!

$$H = \sum_i E_i \underbrace{|E_i\rangle\langle E_i|}_{\substack{\text{stationary state} \\ \text{or} \\ \text{energy eigenstate}}}$$

When we measure  $H$  we get  $E_i$  with  $p(E_i) = \| |E_i\rangle\langle E_i| \|^2$

↳ The smallest result we can get is

$$E_g = \min_i (E_i)$$

↳ ground state energy!

↳ at 0K the system will be in  $|E_g\rangle$

$\therefore$  Physically very relevant!!