Sturm-Liouville Theory

MSO-203B

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Overview

Sturm-Liouville Theory:-

• Introduction and Definitions.

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- Introduction and Definitions.
- Examples.

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Sturm-Liouville Theory:-

- Introduction and Definitions.
- Examples.
- Properties of Sturm-Liouville Problem.

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Introduction

Definition

Recall that a second order ordinary differential operator is of the form

$$L(y) = p_0 y'' + q_0 y' + r_0 y$$

for some p_0 , q_0 and r_0 are real and continuously differentiable function in I := [a, b].

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Remark

Note that this is true if $p'_0 = q_0$ and $q = r_0$

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Sturm-Liouville Problem

Sturm-Liouville Equation

We are interested in studying the BVP of the form

$$(py')' + qy + \lambda ry = 0$$
 (1)
 $B_1(y) = 0, B_2(y) = 0$

where p,q and r are real, continuous function such that p and r are strictly positive in I and λ is a real parameter.

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Eigenvalues and Eigenvectors

The values of λ for which the BVP (1) has a nontrivial solution are called the Eigenvalues of the operator L and the solutions are called eigenfunction.

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Boundary Conditions

Seperated Boundary Condition

If we have,

$$B_1(y) := \alpha_1 y(a) + \alpha_2 y'(a) = 0; \ \alpha_1^2 + \alpha_2^2 \neq 0$$

$$B_2(y) := \beta_1 y(b) + \beta_2 y'(b) = 0; \ \beta_1^2 + \beta_2^2 \neq 0$$
(2)

The BVP (1) with (2) is called a regular Sturm-Liouville Problem.

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The BVP (1) with (2) is called a regular Sturm-Liouville Problem.

Periodic Boundary Condition

If p(a) = p(b) and

$$B_1(y) := y(a) = y(b)$$
 (3)
 $B_2(y) := y'(a) = y'(b)$

Then the BVP (1) with (3) is called a periodic Sturm-Liouville Problem.

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Examples

Sturm-Liouville Problem

Consider the problem

$$y'' + \lambda y = 0 \text{ in } I = [0, \pi]$$
 (4)

Examples of Boundary Conditions

• y(0) = 0 and $y'(\pi) = 0$.

Then Equation (4) along with the above boundary condition is a Regular Sturm-Liouville Problem(RSLP).

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Examples

Sturm-Liouville Problem

Consider the problem

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Examples of Boundary Conditions

- y(0) = 0 and $y'(\pi) = 0$. Then Equation (4) along with the above boundary condition is a Regular Sturm-Liouville Problem(RSLP).
- ② $y(0) = y(\pi)$ and $y'(0) = y'(\pi)$. Then Equation (4) along with the above boundary condition is called a Periodic Sturm-Liouville Problem (PSLP).

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Example 1

RSL Problem

Consider the problem:

$$y'' + \lambda y = 0$$
, $y(0) = 0$, $y'(\pi) = 0$

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Example 1

RSL Problem

Consider the problem:

$$y'' + \lambda y = 0, \ y(0) = 0, \ y'(\pi) = 0$$

Solution for $\lambda < 0$

Let $\lambda < 0$ then $\lambda = -\mu^2$.

The general solution is given as

$$y(x) = Ae^{\mu x} + Be^{-\mu x} \tag{5}$$

y satisfies the boundary condition iff A = B = 0. So y(x) = 0 and hence there are no negative eigenvalues.

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Solution for $\lambda = 0$

y(x) = 0 is the only solution satisfying the equation

y'' = 0, y(0) = 0 $y'(\pi) = 0$. Hence 0 is not an eigenvalue.

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Solution for $\lambda = 0$

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y'' = 0, y(0) = 0 $y'(\pi) = 0$. Hence 0 is not an eigenvalue.

Solution for $\lambda > 0$

Let $\lambda > 0$ then $\lambda = \mu^2$ and the general solution is given by

 $y(x) = A\cos(\mu x) + B\sin\mu x.$

Now y satisfies the boundary condition iff A = 0 and $B \cos(\mu \pi) = 0$.

But, $B\cos(\mu\pi) = 0$ iff B = 0 or $\cos(\mu\pi) = 0$.

When A = B = 0 then y(x) = 0 is the only solution hence this condition doesn't give us eigenvalue.

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Solution for $\lambda > 0$

If $y \neq 0$ and $B \neq 0$ then $\cos(\mu \pi) = 0$ holds which implies $\mu = \frac{2n-1}{2}$ for $n \in \mathbb{Z}$. Thus the eigenvalues are given by $\phi_n(x) = B \sin(\frac{2n-1}{2}x)$ for $n \in \mathbb{Z}$.

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Remark

All the eigenvalues are real.

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Remark

- All the eigenvalues are real.
- Eigenfunctions corresponding to each eigenvalues form a vector space of dim 1.

Example 2

PSL Problem

Consider a problem $y'' + \lambda y = 0$, $y(0) - y(\pi) = 0$; $y'(0) - y'(\pi) = 0$

Solution for $\lambda < 0$

Let $\lambda < 0$ then $\lambda = -\mu^2$.

The general solution is given as

$$y(x) = Ae^{\mu x} + Be^{-\mu x} \tag{6}$$

y satisfies the boundary condition iff A = B = 0. So y(x) = 0 and hence there are no negative eigenvalues.

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Solution for $\lambda = 0$

If $\lambda = 0$ then the general solution is y(x) = A + Bx and since y satisfies the boundary conditions iff B=0. Thus A is arbitrary and hence 0 is an eigenvalue with eigenfunction being a non-zero constant.

Solution for $\lambda > 0$

If $\lambda > 0$ then $\lambda = \mu^2$ and the general solution is given by

$$y(x) = A\cos(\mu x) + B\sin(\mu x).$$

Since y satisfies the boundary condition if

$$A\sin(\mu\pi) + B(1 - \cos(\mu\pi)) = 0 \tag{7}$$

$$A(1-\cos(\mu\pi))-B\sin(\mu\pi)=0$$

Solution for $\lambda > 0$

Equation (7) has a nontrivial solution iff $\cos(\mu\pi) = 1$ which implies $\mu = \pm 2n$ with $n \in \mathbb{Z}$. Hence $\lambda = 4n^2$ for $n \in \mathbb{Z}$.

Thus positive eigenvalues are given by $\lambda_n = 4n^2$ for $n \in \mathbb{Z}$ and the eigenfunctions corresponding to λ_n are given by $\phi_n(x) = \cos(2nx)$ and $\chi_n(x) = \sin(2nx)$.

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Remark

- All eigenvalues are non-negative.
- Corresponding to the positive eigenvalues we have two linearly independent eigenfunctions.

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Theorem 1

The eigenvalues of a regular SLBVP

$$(py')' + qy + \lambda ry = 0$$

$$\alpha_1 y(a) + \alpha_2 y'(a) = 0; \ \alpha_1^2 + \alpha_2^2 \neq 0$$

$$\beta_1 y(b) + \beta_2 y'(b) = 0; \ \beta_1^2 + \beta_2^2 \neq 0$$
(8)

are real provided p and r are strictly positive in I.

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Proof of Theorem 1

Let $\lambda \in \mathbb{C}$ be a complex eigenvalue and u be the corresponding eigenfunction. Hence (λ, u) satisfies

$$(pu')' + qu = \lambda ru \tag{9}$$

$$\alpha_1 u(a) + \alpha_2 u'(a) = 0; \ \alpha_1^2 + \alpha_2^2 \neq 0$$
 (10)

$$\beta_1 u(b) + \beta_2 u'(b) = 0; \ \beta_1^2 + \beta_2^2 \neq 0$$
 (11)

Taking the complex conjugate we have,

$$(p\bar{u}')' + q\bar{u} = \bar{\lambda}r\bar{u} \tag{12}$$

$$\alpha_1 \bar{u}(a) + \alpha_2 \bar{u}'(a) = 0; \ \alpha_1^2 + \alpha_2^2 \neq 0$$
 (13)

$$\beta_1 \bar{u}(b) + \beta_2 \bar{u}'(b) = 0; \ \beta_1^2 + \beta_2^2 \neq 0$$
 (14)

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Proof of Theorem 1

Multiplying \bar{u} and u with (9) and (12) and subtracting we obtain by integrating on I

$$p(\bar{u}u'-u\bar{u}')|_a^b=(\bar{\lambda}-\lambda)\int_a^b ru\bar{u}$$

which implies $(\bar{\lambda} - \lambda) \int_a^b r |u|^2 = 0$ since $p(\bar{u}u' - u\bar{u}')|_a^b = 0$.

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Conclusion

Since r is strictly positive and u is a non-trivial solution we have $\bar{\lambda} = \lambda$.

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Theorem 2

Eigenfunctions corresponding to distinct eigenvalues are orthogonal w.r.t the weight function r(x) on I i.e, if u and v are eigenfunctions corresponding to distinct eigenvalues λ and μ respectively then

$$\int_{a}^{b} r(x)uv = 0$$

Proof of Theorem 2

Following the proof of the previous theorem for (λ, u) and (μ, v) we have,

$$\int_a^b [p(uv'-vu')]' + \int_a^b (\lambda-\mu)ruv = 0$$

Incorporating the boundary condition we get, $p(uv'-vu')|_a^b=0$ and hence

$$\int_{a}^{b} ruv = 0$$

since $\lambda \neq \mu$.

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Theorem 3

The eigenvalues of a regular SLBVP are simple i.e, An eigenfunction corresponding to an eigenvalue is unique upto a constant multiple.

Conclusion

- The dimension of the space of eigenfunctions corresponding to an eigenvalue is 1.
- This is not true if we have Periodic Boundary condition.

Proof of Theorem 3

Let u and v be two eigenvalues corresponding to λ . Following the proof of the previous theorems we can get

$$p(uv'-vu')|_a^b=c$$

where c is a constant.

The above expression implies that $p(x)\mathbb{W}(u,v)(x)=c$ for all $x\in I$.

Again from the boundary conditions we get, $\mathbb{W}(u, v)(a) = 0$.

This implies that W(u, v)(x) = 0 for all $x \in I$.

Hence, u and v are linearly independent.

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