

Kronecker delta: 4 rules you need to know

It is impossible to imagine theoretical physics without the Kronecker delta. You will encounter this relatively simple, yet powerful tensor practically in all fields of theoretical physics. For example, it is used to

- Write long expressions more compactly.
- Simplify complicated expressions.

In combination with the *Levi-Civita tensor*, the two tensors are very powerful! That's why it's worth understanding how the Kronecker delta works.

1.1 Definition and Examples

Kronecker delta δ_{ij} - is a small greek letter delta, which yields either 1 or 0, depending on which values its two indices i and j take on. The maximal value of an index corresponds to the considered dimension, so in three-dimensional space i and j run from 1 to 3.

Kronecker delta is equal to 1, if i and j are equal. And Kronecker delta is 0, if i and j are not equal.

Definition

$$\delta_{ij} = \begin{cases} 1 & \text{if } i = j \\ 0 & \text{if } i \neq j \end{cases} \quad (1)$$

Examples

- $\delta_{11} = 1$ - because both indices are the same.
- $\delta_{23} = 0$ - because both indices are different.
- $a \delta_{33} = a \cdot 1 = a$
- $\delta_{23} \delta_{22} = 1 \cdot 0 = 0$

1.2 Einstein's Summation convention

In order to represent an expression like this

$$\sum_{j=1}^3 \delta_{ij} \delta_{jk} = \delta_{i1} \delta_{1k} + \delta_{i2} \delta_{2k} + \delta_{i3} \delta_{3k} \quad (2)$$

or an expression like this

$$\sum_{i=1}^3 a_i b_i = a_1 b_1 + a_2 b_2 + a_3 b_3 \quad (3)$$

compactly, we agree on the following Summation convention rule:

Summation convention

We omit the sum sign, but keep in mind that if *two equal indices* appear in an expression, then we *sum over that index*.

Example

In the following scalar product we sum over i :

$$\sum_{i=1}^3 a_i b_i = a_1 b_1 + a_2 b_2 + a_3 b_3 \quad (4)$$

We omit the sum sign and keep in mind that we sum over i :

$$\sum_{i=1}^3 a_i b_i \rightarrow a_i b_i \quad (5)$$

Thus, using Einstein's summation convention, the scalar product is:

$$a_i b_i = a_1 b_1 + a_2 b_2 + a_3 b_3 \quad (6)$$

Another advantage of the sum convention (in addition to compactness) is formal *commutativity*. For example, you may write down the expression $\varepsilon_{ijk} \hat{\mathbf{e}}_i s_j \delta_{km}$ in a different order as you wish, for example like this: $\varepsilon_{ijk} \delta_{km} s_j \hat{\mathbf{e}}_i$. This might help you see what can be shortened or summarized further.

But be careful! There are exceptions, for example with the differential operator ∂_j , which acts on a successor. You can't move something before the derivative that is supposed to be differentiated. So you should be careful with operators in index notation.

1.3 4 Rules for Kronecker delta

Rule 1

Indices ij may be interchanged:

$$\delta_{ij} = \delta_{ji} \quad (7)$$

Why is that? According to the definition, if the indices i and j are equal, then δ_{ij} is equal to

1. But then, also δ_{ji} is equal to 1. And if the indices are unequal, you have: δ_{ij} is equal to zero. And δ_{ji} is equal to zero. So as you can see: Kronecker delta is *symmetric*!

Rule 2

If the product of two or more Kronecker deltas contains a summation index j , then the product can be shortened, such that the summation index j disappears:

$$\delta_{ij} \delta_{jk} = \delta_{ik} \quad (8)$$

Why it is so? Let's consider for example the case where the indices i and j are equal: $i = j$ and the indices j and k are not equal: $j \neq k$. Then it follows that i and k must also be unequal: $i \neq k$. So δ_{jk} is zero and therefore the whole term on the left hand side is zero: $\delta_{ij} \delta_{jk} = 0$. δ_{ik} on the left-hand side is also zero, because i and k are different. The equation is fulfilled. You can proof all other possible cases in the same way.

So instead of writing two deltas you can just write δ_{ik} . We say: The summation index j is *contracted*.

Example

Consider $\delta_{km} \delta_{mn}$. The summation index here is m , so you can eliminate it by contracting it. You get δ_{kn} .

Example

Consider $\delta_{ij} \delta_{kj} \delta_{in}$. Here you have two summation indices i and j . So in principle you can eliminate both of them.

First possible way of contraction: From the first rule you know, that Kronecker delta is symmetric, so you can swap k and j in δ_{kj} and then contract the index j first. You get: $\delta_{ik} \delta_{in}$. And then you contract the index i . The simplified result is δ_{kn} .

Second possible way of contraction: First reorder the product to $\delta_{kj} \delta_{ij} \delta_{in}$. Contract the summation index i first. You get: $\delta_{kj} \delta_{jn}$. Contract the second summation index j . You get: δ_{kn} .

Remember that the contraction order is not important here. *In both cases* you get the same result: δ_{kn} . So which way of simplification you take doesn't matter!

Rule 3

If the index in a_j also occurs in Kronecker delta δ_{jk} , then the Kronecker delta disappears and the factor a_j gets the other index k :

$$a_j \delta_{jk} = a_k \quad (9)$$

Why is that? This rule is basically another case of index contraction. This rule tells you that you can also contract summation indices that don't have to be carried by a Kronecker delta.

Let's make another example. Consider $\Gamma_{jmk} \delta_{nk}$. The summation index is k , so you can eliminate it. The result is Γ_{jmn} .

Rule 4

If j runs from 1 to n , then:

$$\delta_{jj} = n \quad (10)$$

Why is that? According to the summation convention (2), the summation is carried out over j here. So δ_{jj} is equal to δ_{11} plus δ_{22} plus δ_{33} and so on up to n . And each Kronecker delta yields 1, because the index values are equal. So $1 + 1 + 1$ and so on, results in n :

$$\begin{aligned} \delta_{jj} &= \delta_{11} + \delta_{22} + \dots + \delta_{nn} \\ &= 1 + 1 + \dots + 1 \\ &= n \end{aligned} \quad (11)$$

1.4 The 3 most common mistakes you should avoid making

If you use summation convention and the above rules, you must also pay attention to the *correct notation*: Summation is done if an index occurs exactly *twice* on *one* side of the equation.

Mistake 1

$$v_i = b_i c_i \quad (12)$$

Why? Because on the right hand side you have double index i , so you sum over i . But on the left hand side there is also the summation index i - and this of course makes no sense... To correct this expression, you can rename the summation index to j .

Mistake 2

$$\delta_{jk} v_k = \delta_{mk} r_k \quad (13)$$

Because the summation index k appears on *both* sides of the equation. To correct this, rename one of the summation indices, for example to j .

Mistake 3

$$\delta_{i1} \delta_{1k} = \delta_{ik} \quad (14)$$

Why? Because here you try to sum over an ordinary *number*, as if this number would be a summation index. Number '1' occurs twice, but it is *not a variable index* over which you can sum. Therefore you can not use the contraction rule on ordinary numbers.

1.5 Scalar product with Kronecker delta

Consider a three-dimensional vector with the components x , y and z :

$$\mathbf{v} = \begin{bmatrix} x \\ y \\ z \end{bmatrix} \quad (15)$$

You can represent this vector \mathbf{v} in an orthonormal basis as follows:

$$\mathbf{v} = x \hat{\mathbf{e}}_x + y \hat{\mathbf{e}}_y + z \hat{\mathbf{e}}_z \quad (16)$$

here $\hat{\mathbf{e}}_x$, $\hat{\mathbf{e}}_y$ and $\hat{\mathbf{e}}_z$ are three basis vectors which are orthogonal to each other and normalized. In this case they span an orthogonal three-dimensional coordinate system.

The Kronecker delta needs vectors written in **index notation**. Here we do not denote the vector components with different letters x, y, z , but we choose one letter (here the letter v) and then number the vector components consecutively. The vector components are then called v_1 , v_2 and v_3 and the basis expansion looks like this:

$$\begin{aligned} \mathbf{v} &= (v_1, v_2, v_3) \\ &= v_1 \hat{\mathbf{e}}_1 + v_2 \hat{\mathbf{e}}_2 + v_3 \hat{\mathbf{e}}_3 \end{aligned} \quad (17)$$

One of the advantages of index notation is that this way you will never run out of letters for the vector components. Just imagine a fifty-dimensional vector. There aren't even that many letters to give each component $v_1, v_2 \dots v_{50}$ of the vector a unique letter! It gets worse when you want to write out this fifty-dimensional vector as in (??)...

Another advantage of the index notation is that by numbering the vector components in this

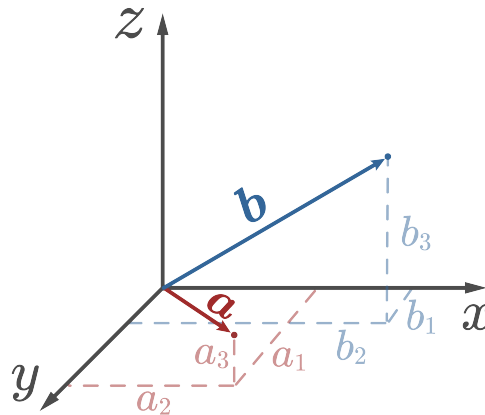


Abbildung 1: Two example vectors \mathbf{a} and \mathbf{b} and their components.

way, you can use the sum sign to represent (??) more compactly. It becomes even more compact if we omit the big sum sign according to the summation convention. Look how compact the vector \mathbf{v} can be represented in a basis:

Vector in index notation

$$\mathbf{v} = v_j \hat{\mathbf{e}}_j \quad (18)$$

here, as you know, we sum over index j . Whether you call the index j , i or k or any other letter is of course up to you.

Now that you know how a vector is represented in index notation, we can analogously write the *scalar product* $\mathbf{a} \cdot \mathbf{b}$ of two vectors $\mathbf{a} = (a_1, a_2, a_3)$ and $\mathbf{b} = (b_1, b_2, b_3)$ in index notation. For this we use the index representation of a vector (18):

$$\mathbf{a} \cdot \mathbf{b} = a_i \hat{\mathbf{e}}_i \cdot b_j \hat{\mathbf{e}}_j \quad (19)$$

In index notation, you may sort the factors in (19) as you like. This is the advantage of index notation, where the commutative law applies. Let's take advantage of that and put parentheses around the basis vectors to emphasize their importance in introducing the Kronecker delta:

$$\mathbf{a} \cdot \mathbf{b} = a_i b_j (\hat{\mathbf{e}}_i \cdot \hat{\mathbf{e}}_j) \quad (20)$$

The basis vectors $\hat{\mathbf{e}}_i$ and $\hat{\mathbf{e}}_j$ are **orthonormal** (i.e. orthogonal and normalized). Recall what the property of being orthonormal means for two vectors. Their scalar product $\hat{\mathbf{e}}_i \cdot \hat{\mathbf{e}}_j$ yields:

- $\hat{\mathbf{e}}_i \cdot \hat{\mathbf{e}}_j = 1$, if $i = j$.

- $\hat{\mathbf{e}}_i \cdot \hat{\mathbf{e}}_j = 0$, if $i \neq j$.

Doesn't that look familiar to you? The scalar product of two orthonormal vectors behaves *exactly like Kronecker delta*! Therefore replace the scalar product of two basis vectors with a Kronecker delta:

Scalar product of two orthonormal vectors

$$\hat{\mathbf{e}}_i \cdot \hat{\mathbf{e}}_j = \delta_{ij} \quad (21)$$

Thus we can write the scalar product (20) using Kronecker delta:

Scalar product with Kronecker delta

$$\mathbf{a} \cdot \mathbf{b} = a_i b_j \delta_{ij} \quad (22)$$

If you remember the third rule you can contract the index j if you want:

$$\mathbf{a} \cdot \mathbf{b} = a_i b_i \quad (23)$$

And you get exactly the definition of the scalar product, where the vector components are summed component-wise:

$$\mathbf{a} \cdot \mathbf{b} = a_1 b_1 + a_2 b_2 + a_3 b_3 \quad (24)$$

Check: Write out the sum

We can write out the double summation over i and j in (22) for practice. In other words, we have to go through all possible combinations of the indices i and j :

$$\begin{aligned} a_i b_j \delta_{ij} = & a_1 b_1 \delta_{11} + a_1 b_2 \delta_{12} + a_1 b_3 \delta_{13} \\ & + a_2 b_1 \delta_{21} + a_2 b_2 \delta_{22} + a_2 b_3 \delta_{23} \\ & + a_3 b_1 \delta_{31} + a_3 b_2 \delta_{32} + a_3 b_3 \delta_{33} \end{aligned} \quad (25)$$

As you can see - because of the definition of Kronecker delta - only 3 components of 9 in total are not zero, where $i = j$. So you may omit all summands with unequal indices. Together with the third rule you have:

$$a_i b_i = a_1 b_1 \delta_{11} + a_2 b_2 \delta_{22} + a_3 b_3 \delta_{33} \quad (26)$$

Using the definition of Kronecker delta, $\delta_{11} = \delta_{22} = \delta_{33} = 1$, you get the scalar product you are familiar with:

$$a_i b_i = a_1 b_1 + a_2 b_2 + a_3 b_3 \quad (27)$$