

MATB42 Week 2 Notes

1. HEAT Equation:

- The PDE $u_t = k u_{xx}$ is called the **heat eqn / diffusion eqn** where k is a constant called the rate of diffusion.

- The heat eqn describes the diffusion of heat through a 1-D medium, such as a metal rod. The unknown function $u(x, t)$ is the temperature at point x at time t .

- We aim to solve $u(x, t)$ in the finite domain $0 < x < l$ where l is the length of the domain.

- Initial Condition:

$u(x, 0) = \phi(x)$ where $\phi(x)$ describes the initial heat distribution at time $t=0$.

- Boundary Condition:

$$u_x(0, t) = 0$$

$$u_x(l, t) = 0$$

The boundary conditions indicate that no heat is leaving or entering the domain.

To solve the heat eqn, we'll use Separation of Variables.

Assume $u(x, t) = X(x) \cdot T(t)$

Then, the PDE $u_t = k u_{xx}$ becomes
 $X \cdot T' = k \cdot X'' \cdot T$

Recall: By convention, we move all the constants and terms with t to the LHS and all the terms with x to the RHS.

Thus, we get $\frac{T'}{k \cdot T} = \frac{X''}{X} = F$

The separation constant. Is always negative.

$$\frac{T'}{k \cdot T} = \frac{X''}{X} = -\lambda, \text{ where } \lambda > 0$$

$$\frac{T'}{k \cdot T} = -\lambda$$

$$T' = -\lambda k T$$

Recall: $(\ln(f(x)))' = \frac{f'(x)}{f(x)}$

$$\frac{T'}{T} = -\lambda k$$

Hence, $\frac{T'}{T} = (\ln(T))'$

$$(\ln(T))' = -\lambda k$$

$$\ln(T) = \int -\lambda k$$

$$= -\lambda k t + C$$

$$T = e^{-\lambda k t + C}$$

$$T = A e^{-\lambda k t}$$

To get rid of the derivative, we integrate both sides with respect to t .

Now, we have $\ln(T) = -\lambda k t$

To get rid of \ln , raise both sides by e .

$$e^{\ln f(x)} = f(x)$$

$$\frac{X''}{X} = -\lambda$$

$$X'' = -\lambda X$$

$$X'' + \lambda X = 0$$

$$r^2 + \lambda = 0$$

$$r^2 = -\lambda$$

$$r = \pm \sqrt{\lambda}i$$

$$\text{Take } r = \sqrt{\lambda}i$$

$$e^{rx} = e^{(\sqrt{\lambda}x)i}$$

$$= \cos(\sqrt{\lambda}x) + i\sin(\sqrt{\lambda}x)$$

$$X = C \cos(\sqrt{\lambda}x) + D \sin(\sqrt{\lambda}x)$$

$$u(x, t) = \underbrace{X(x) \cdot T(t)}_{\text{We found both terms}}$$

Now, we plug in the boundary equations.
Recall that the boundary conditions are

$$1. u_x(0, t) = 0$$

$$2. u_x(l, t) = 0$$

This means that we need to differentiate u with respect to (w.r.t) x first.

$$u_x = \frac{\partial u}{\partial x}$$

$$= \frac{\partial}{\partial x} (X(x) \cdot T(t))$$

$$= T(t) \cdot \frac{\partial}{\partial x} (X(x)) \quad \leftarrow T \text{ is treated as a constant.}$$

$$= T(t) \underbrace{(-C \sin(\sqrt{\lambda}x) + D \cos(\sqrt{\lambda}x))(\sqrt{\lambda})}_{x'}$$

$$u_x(0,t)=0 \rightarrow X'(0) \cdot T(t)=0 \rightarrow X'(0)=0$$

Recall: We ignore the case when $T(t)=0$ bc that gives us the trivial soln.

$$X'(0) = (-C \sin(0) + D \cos(0))(\sqrt{\lambda})$$

$$X'(0) = 0$$

$$D\sqrt{\lambda} = 0$$

Since we assumed that $\lambda > 0$, this means $D=0$.

Now, we'll plug in the other boundary condition.
 $u_x(l,t)=0 \rightarrow X'(l) \cdot T(t)=0 \rightarrow X'(l)=0$

$$X'(l) = -C \sin(\sqrt{\lambda} l) \sqrt{\lambda}$$

$$X'(l) = 0$$

$$-C \sin(\sqrt{\lambda} l) \sqrt{\lambda} = 0$$

$$C \sin(\sqrt{\lambda} l) = 0$$

Either $C=0$ or $\sin(\sqrt{\lambda} l)=0$.

We ignore the case when $C=0$. (Trivial soln)

$$\sin(\sqrt{\lambda} l) = 0 \rightarrow \sqrt{\lambda} l = n\pi, n \geq 0$$

$$\rightarrow \sqrt{\lambda} = \frac{n\pi}{l}$$

For each value of n , we have a soln, denoted as u_n .

$$u_n(x,t) = A_n \cos\left(\frac{n\pi x}{l}\right) e^{-\left(\frac{n\pi}{l}\right)^2 kt}$$

Recall:

$$X = C \cos(\sqrt{\lambda} x) + \underbrace{D \sin(\sqrt{\lambda} x)}_{D=0}$$

$$= C \cos\left(\frac{n\pi x}{l}\right) \leftarrow \sqrt{\lambda} = \frac{n\pi}{l}$$

$$T = A e^{-\lambda kt} \leftarrow \sqrt{\lambda} = \frac{n\pi}{l} \rightarrow \lambda = \left(\frac{n\pi}{l}\right)^2$$

$$= A e^{-\left(\frac{n\pi}{l}\right)^2 kt}$$

Note: Because $\cos(0) = 1 \neq 0$, it does not give the trivial soln and we can't ignore the case when $n=0$.

Since each value of n gives a soln to the PDE, the general soln is a linear comb of all the Unis.

$$u(x,t) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right) e^{-\left(\frac{n\pi}{L}\right)^2 kt}$$

↑
Notice we start at $n=0$.

We now apply the initial condition $u(x,0) = \phi(x)$.

$$u(x,t) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right) e^{-\left(\frac{n\pi}{L}\right)^2 kt}$$

$$u(x,0) = \phi(x) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right)$$

This is called the **Fourier (cosine) Series of $\phi(x)$** . The reason we got cosine instead of sine is bc our boundary conditions involve the derivative of u .

Recall the Orthogonal Relations for Fourier Series:

$$\int_0^l \cos\left(\frac{n\pi x}{l}\right) \cos\left(\frac{m\pi x}{l}\right) dx = \begin{cases} 0, & \text{if } m \neq n \\ \frac{l}{2}, & \text{if } m = n \neq 0 \end{cases}$$

Since we have $m = n \neq 0$, we have to distinguish when $m = 0$ and when $m \neq 0$.

1. When $m = 0$:

$$\phi(x) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{l}\right) \leftarrow \text{From before}$$

Multiply both sides by $\cos\left(\frac{m\pi x}{l}\right)$ and integrate from 0 to l .

$$\int_0^l \phi(x) \cos\left(\frac{m\pi x}{l}\right) dx = \sum_{n=0}^{\infty} A_n \underbrace{\int_0^l \cos\left(\frac{n\pi x}{l}\right) \cos\left(\frac{m\pi x}{l}\right) dx}_{\substack{\text{This integral equals 0} \\ \text{for when } n \neq 0, \text{ and} \\ \text{is equal to } l \text{ when} \\ n = 0}}$$

$$= A_n \int_0^l 1 + 0 + 0 + \dots + 0$$

$$= A_n l$$

$$A_n = \frac{1}{l} \int_0^l \phi(x) \underbrace{\cos\left(\frac{m\pi x}{l}\right)}_{\substack{\text{Equals 1 since } \cos(0) = 1}} dx, \text{ when } m = 0$$

$$= \frac{1}{l} \int_0^l \phi(x) dx$$

2. When $m \neq 0$:

$$\phi(x) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{l}\right) \leftarrow \text{From Before}$$

Multiply both sides by $\cos\left(\frac{m\pi x}{l}\right)$ and integrate from 0 to l .

$$\int_0^l \phi(x) \cos\left(\frac{m\pi x}{l}\right) dx = \sum_{n=0}^{\infty} A_n \underbrace{\int_0^l \cos\left(\frac{n\pi x}{l}\right) \cos\left(\frac{m\pi x}{l}\right) dx}_{\substack{\text{The integral equals 0} \\ \text{when } n \neq m \text{ and} \\ \frac{l}{2} \text{ when } n=m}}$$

$$= 0 + \dots + 0 + A_n \frac{l}{2} + 0 + \dots + 0$$

$$A_n = \frac{2}{l} \int_0^l \phi(x) \cos\left(\frac{m\pi x}{l}\right) dx$$

— Earlier, we said that the separation constant is always negative, and we denoted it as $-\lambda$, where $\lambda \geq 0$. We'll see why here.

If $\lambda = 0$:

$$T' = 0 \quad \text{and} \quad X'' = 0$$

$$T(t) = A \quad X(x) = Cx + D$$

$$u_x(0, t) = 0 \rightarrow X'(0) \cdot T(t) = 0 \rightarrow X'(0) = 0$$

$$X'(0) = C$$

$$\therefore C = 0$$

$$u_x(l, t) = 0 \rightarrow X'(l) \cdot T(t) = 0 \rightarrow X'(l) = 0$$

$$X'(l) = C$$

$$\therefore C = 0$$

Both boundary conditions show that $C = 0$.

Furthermore, we see no restrictions on A or D . This means that, when $\lambda=0$, we get a constant $A \cdot D$. This is the exact constant we get when $n=0$.

I.e.

$$U_0^{(x,t)} = A_0 e^0 \cdot D_0 \cos(0) \\ = A_0 \cdot D_0$$

If $\lambda < 0$:

Let $\lambda = -\mu$, where $\mu > 0$.

Now, we have: $T' - \mu k T = 0$ and $X'' - \mu X = 0$.

From $T' - \mu k T = 0$, we get $T(t) = A e^{\mu k t}$.

From $X'' - \mu X = 0$, we get $X(x) = C e^{\sqrt{\mu} x} + D e^{-\sqrt{\mu} x}$.

We now apply the boundary eqns.

First, we apply $U_x(0, t) = 0$.

$$U_x = X' \cdot T \\ = \underbrace{(C e^{\sqrt{\mu} x} - D e^{-\sqrt{\mu} x})}_{X'} (\sqrt{\mu}) \cdot T$$

$$U_x(0, t) = 0 \rightarrow X'(0) \cdot T(t) = 0 \rightarrow X'(0) = 0$$

$$X'(0) = (C - D) \sqrt{\mu}$$

$$(C - D) \sqrt{\mu} = 0$$

$$C - D = 0 \quad (\text{we know } \mu \neq 0)$$

$$C = D$$

Now, we use the other boundary condition.

$$U_x(l, t) = 0 \rightarrow X'(l) \cdot T(t) = 0 \rightarrow X'(l) = 0$$

$$X'(l) = (C e^{\sqrt{\mu} l} - D e^{-\sqrt{\mu} l}) \sqrt{\mu} = 0$$

$$\text{Using } C = D, \text{ we get } C(e^{\sqrt{\mu} l} - e^{-\sqrt{\mu} l}) = 0$$

This means that either $C=0$ or $e^{\sqrt{\lambda}l} - e^{-\sqrt{\lambda}l} = 0$.

$C=0$ gives us the trivial soln, so we ignore it.

$$e^{\sqrt{\lambda}l} - e^{-\sqrt{\lambda}l} = 0$$

$$e^{\sqrt{\lambda}l} = e^{-\sqrt{\lambda}l}$$

$$\sqrt{\lambda}l = -\sqrt{\lambda}l \quad (\text{we took the ln of both sides.})$$

This means that $\lambda=0$, which is a contradiction.

Hence, there are no non-trivial solns for $\lambda < 0$.

\therefore λ is never negative but may be equal to 0.

2. Important Formulas:

$$- u(x,t) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{l}\right) e^{-\left(\frac{n\pi}{l}\right)^2 kt}$$

$$- \text{If } m=0, A_n = \frac{1}{l} \int_0^l \phi(x) dx$$

$$- \text{If } m \neq 0, \text{ and } n=m, A_n = \frac{2}{l} \int_0^l \phi(x) \cos\left(\frac{n\pi x}{l}\right) dx$$

$$- \int_0^l \cos\left(\frac{n\pi x}{l}\right) \cos\left(\frac{m\pi x}{l}\right) dx = \begin{cases} 0, & \text{if } m \neq n \\ \frac{l}{2}, & \text{if } m=n \neq 0 \end{cases}$$

3. Examples

E.g. 1 Find the soln to the heat eqn on $0 < x < l$ with $U_x(0, t) = 0$, $U_x(l, t) = 0$ and $\phi(x) = \cos\left(\frac{2\pi x}{l}\right)$.

Soln:

$$u(x, t) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{l}\right) e^{-\left(\frac{n\pi}{l}\right)^2 kt}$$

Because n could equal m , which equals 0, we need to split this into 2 cases.

Case 1 ($n=m=0$):

$$\begin{aligned} A_n &= \frac{1}{l} \int_0^l \phi(x) dx \\ &= \frac{1}{l} \int_0^l \cos\left(\frac{2\pi x}{l}\right) dx \\ &= \frac{1}{l} \left(\frac{l}{2\pi}\right) \left[\sin\left(\frac{2\pi x}{l}\right) \Big|_0^l \right] \\ &= \frac{1}{2\pi} \left(\underbrace{\sin(2\pi)}_0 - \underbrace{\sin(0)}_0 \right) \\ &= 0 \end{aligned}$$

Case 2 ($n=m \neq 0$):

$$\begin{aligned} A_n &= \frac{2}{l} \int_0^l \phi(x) \cos\left(\frac{m\pi x}{l}\right) \\ &= \frac{2}{l} \int_0^l \cos\left(\frac{2\pi x}{l}\right) \cos\left(\frac{m\pi x}{l}\right) \\ &= 0 \text{ if } n \neq 2, \text{ and } 1, \text{ if } n=2. \end{aligned}$$

$$\therefore A_n = \begin{cases} 0, & \text{if } n=m=0 \\ 0, & \text{if } n \neq 2 \\ 1, & \text{if } n=2 \end{cases}$$

E.g. 2 Find the soln to the heat eqn on $0 < x < l$ with $u_x(0, t) = 0$, $u_x(l, t) = 0$ and $\phi(x) = 1$

Soln:

$$u(x, t) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{l}\right) e^{-\left(\frac{n\pi}{l}\right)^2 kt}$$

If $n=0$: ($n=m=0$)

$$\begin{aligned} A_n &= \frac{1}{l} \int_0^l \phi(x) dx \\ &= \frac{1}{l} \int_0^l 1 dx \\ &= \frac{1}{l} (l-0) \\ &= 1 \end{aligned}$$

If $n \neq 0$: ($n=m \neq 0$)

$$\begin{aligned} A_n &= \frac{2}{l} \int_0^l \phi(x) \cos\left(\frac{m\pi x}{l}\right) dx \\ &= \frac{2}{l} \int_0^l \cos\left(\frac{m\pi x}{l}\right) dx \\ &= \frac{2}{m\pi} \left(\sin(m\pi) - \sin(0) \right) \\ &= 0 \end{aligned}$$

$$\therefore A_n = \begin{cases} 1, & \text{if } n=0 \\ 0, & \text{if } n \neq 0 \end{cases}$$

E.g. 3 Find the soln to the heat eqn on $0 < x < l$ with $u(0,t) = 0$, $u(l,t) = 0$ and $u(x,0) = \phi(x)$.

Soln:

Notice that the boundary conditions do not differentiate u w.r.t x .

I.e. we have $u(0,t) = 0$ instead of $u_x(0,t) = 0$.

Assume $u(x,t) = X(x) \cdot T(t)$ \leftarrow Separation of var

The PDE $u_t = k u_{xx}$ now becomes

$$X \cdot T' = k \cdot X'' \cdot T$$

$$\frac{T'}{kT} = \frac{X''}{X} = -\lambda$$

$$T' = -\lambda k T$$

$$X'' = \lambda X$$

$$X'' - \lambda X = 0$$

$$X = C \cos(\sqrt{\lambda} x) + D \sin(\sqrt{\lambda} x)$$

$$T = A e^{-\lambda k t}$$

Now, let's plug in the boundary conditions.

$$u(0,t) = 0 \rightarrow X(0) \cdot T(t) = 0 \rightarrow X(0) = 0$$

$$X(0) = C$$

$$X(0) = 0$$

$$\therefore C = 0$$

$$u(l,t) = 0 \rightarrow X(l) \cdot T(t) = 0 \rightarrow X(l) = 0$$

$$X(l) = D \sin(\sqrt{\lambda} l)$$

$$X(l) = 0$$

$$D \sin(\sqrt{\lambda} l) = 0$$

$$\sqrt{\lambda} l = n\pi, \quad n > 0$$

$$\sqrt{\lambda} = \frac{n\pi}{l}$$

For each n , we have $U_n(x,t) = A_n \sin\left(\frac{n\pi x}{\ell}\right) e^{-\left(\frac{n\pi}{\ell}\right)^2 kt}$
 To get the general soln, we need to sum up all the U_n 's.

$$u(x,t) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{\ell}\right) e^{-\left(\frac{n\pi}{\ell}\right)^2 kt}$$

To solve for A_n , we'll plug in the initial condition.

$$u(x,0) = \phi(x)$$

$$\phi(x) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{\ell}\right)$$

Now, we multiply both sides by $\sin\left(\frac{m\pi x}{\ell}\right)$ and integrate from 0 to ℓ .

$$\int_0^{\ell} \phi(x) \sin\left(\frac{m\pi x}{\ell}\right) dx = \sum_{n=1}^{\infty} A_n \sin\left(\frac{m\pi x}{\ell}\right) \sin\left(\frac{n\pi x}{\ell}\right) dx$$

$$= A_n \cdot \frac{\ell}{2}$$

$$A_n = \frac{2}{\ell} \int_0^{\ell} \phi(x) \sin\left(\frac{n\pi x}{\ell}\right) dx$$

E.g. 4 Find the soln to the heat eqn on $0 < x < l$ with $u(0, t) = 0$, $u_x(l, t) = 0$ and $u(x, 0) = \phi(x)$.

Soln:

Note: This is sometimes called a **mixed boundary eqn**.

Assume $u(x, t) = X(x) \cdot T(t)$

The PDE $u_t = k u_{xx}$ is now: $T'X = k \cdot T \cdot X''$.

$$\frac{T'}{k \cdot T} = \frac{X''}{X} = -\lambda$$

$$T = A e^{-\lambda k t}$$

$$X = C \cos(\sqrt{\lambda} x) + D \sin(\sqrt{\lambda} x)$$

Now, let's plug in the boundary conditions.

First, we'll use $u(0, t) = 0$.

$$u(0, t) = 0 \rightarrow X(0) \cdot T(t) = 0 \rightarrow X(0) = 0$$

$$X(0) = C$$

$$X(0) = 0$$

$$\therefore C = 0$$

Next, we'll use $u_x(l, t) = 0$.

$$u_x = X'(x) \cdot T(t)$$

$$= \underbrace{D \cos(\sqrt{\lambda} x) (\sqrt{\lambda})}_{X'} \cdot T(t)$$

$$u_x(l, t) = 0 \rightarrow X'(l) \cdot T(t) = 0 \rightarrow X'(l) = 0$$

$$X'(l) = \sqrt{\lambda} D \cos(\sqrt{\lambda} l)$$

$$X'(l) = 0$$

$$\sqrt{\lambda} D \cos(\sqrt{\lambda} l) = 0$$

Since we assume that $\lambda > 0$, and $D \neq 0$, we get $\cos(\sqrt{\lambda} l) = 0$

$$\sqrt{\lambda} l = \frac{(2n+1)\pi}{2}, \quad n \geq 0$$

$$\sqrt{\lambda} = \frac{(2n+1)\pi}{2l}$$

Each value of n gets us a soln, denoted as $U_n(x, t)$.

$$U_n(x, t) = A_n \sin\left(\frac{(2n+1)\pi x}{2l}\right) e^{-\left(\frac{(2n+1)\pi}{2l}\right)^2 kt}$$

To get the general soln of the PDE, we need a linear combination of each of the U_n 's.

$$u(x, t) = \sum_{n=0}^{\infty} A_n \sin\left(\frac{(2n+1)\pi x}{2l}\right) e^{-\left(\frac{(2n+1)\pi}{2l}\right)^2 kt}$$

To solve for A_n , we'll use the initial condition.
 $u(x, 0) = \phi(x)$

$$u(x, 0) = \sum_{n=0}^{\infty} A_n \sin\left(\frac{(2n+1)\pi x}{2l}\right)$$

$$\phi(x) = \sum_{n=0}^{\infty} A_n \sin\left(\frac{(2n+1)\pi x}{2l}\right)$$

Now, we'll multiply both sides by

$$\sin\left(\frac{(2m+1)\pi x}{2l}\right) \text{ and integrate from } 0 \text{ to } l.$$

$$\int_0^l \phi(x) \sin\left(\frac{(2m+1)\pi x}{2l}\right) = \sum_{n=0}^{\infty} A_n \sin\left(\frac{(2n+1)\pi x}{2l}\right) \sin\left(\frac{(2m+1)\pi x}{2l}\right)$$

$$= A_n \cdot \frac{l}{2}$$

$$A_n = \frac{2}{l} \int_0^l \phi(x) \sin\left(\frac{(2m+1)\pi x}{2l}\right)$$

Note: Since we have a sine function, the above also includes the case when $n=m=0$.