

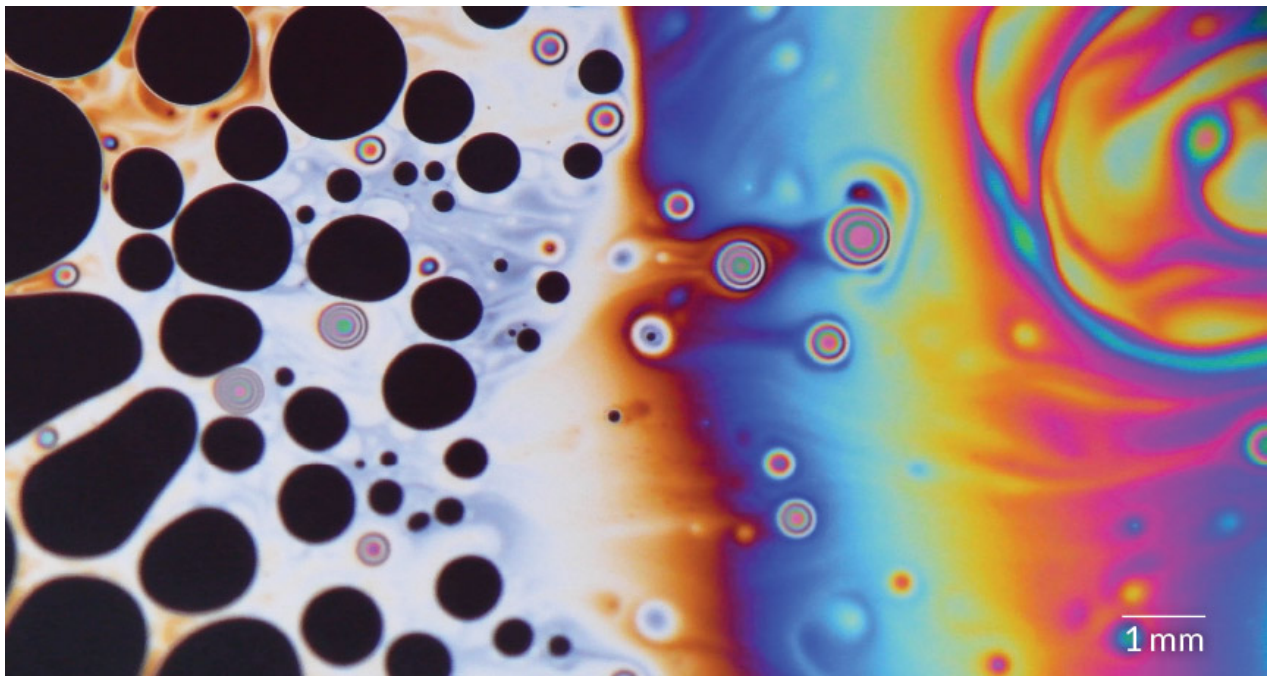
ME3 Literature Research Project

# Stability Characteristics and Death: Thin Films and Soap Bubbles

Pierre Tharreau

CID: 01506504

2nd December, 2020



*Experimentally observed pattern formations on a soap bubble film close to rupture [1]*

**Abstract**

# Project Objectives

- Define characteristics and fundamental interactions in thin films
  - Role of surface tension
  - Role and importance of surfactants
  - Classification (bubble, anti-bubble, thin vs. thick film)
  - Use and applications
- Discuss the interactions with the external environment
  - Stable bubbles (liquid structures, e.g. Marangoni effect, effect of EM fields)
  - Bursting of bubbles (ligaments formation, holes)
- Discuss how stability characteristics could be improved/changed

# Table of Contents

<b>1</b>	<b>Introduction</b>	<b>1</b>
<b>2</b>	<b>Stable curved films</b>	<b>2</b>
2.1	Surface tension . . . . .	2
2.2	Role of surfactant . . . . .	2
2.2.1	"Rigid" and "Mobile" film surfaces . . . . .	3
2.3	The Marangoni Effect . . . . .	4
2.4	Bubble Coalescence . . . . .	4
2.5	Behaviour in Electromagnetic fields . . . . .	4
2.6	Applications . . . . .	4
2.7	Discussion . . . . .	4
<b>3</b>	<b>Bubble Death and film rupture</b>	<b>5</b>
3.1	Film thinning mechanisms and rupture initialisation . . . . .	5
3.1.1	Gravity and capillary drainage . . . . .	5
3.1.2	Marginal Regeneration . . . . .	7
3.1.3	Newton Black films . . . . .	9
3.1.4	Rupture initialisation . . . . .	11
3.2	Hole Development . . . . .	13
3.2.1	Rim Velocity . . . . .	14
3.2.2	Appearance of instabilities . . . . .	15
3.3	The aftermath . . . . .	15
3.4	Applications . . . . .	15
3.5	Discussion . . . . .	15
<b>4</b>	<b>Improving stability</b>	<b>15</b>
<b>5</b>	<b>Conclusions</b>	<b>15</b>
	<b>References</b>	<b>16</b>

# 1 Introduction

A multitude of questions come to mind when observing the life of a soap bubble; how does it manage to maintain stability despite appearing so fragile? what effects explain its ephemeral nature? and what happens during rupture and after the bubble has burst? This report will attempt to summarize the mechanisms and structures which soap bubbles exhibit in their life time, while gauging their effect on the stability of these curious thin films.

Strictly, a thin film is a material layer bounded on either side a fluid or solid. Although thin films involving solids are well studied and play a significant role in a wide range of industries such as electronics and bio-systems [2], this report will focus on the properties and dynamics of their fluid-only counterparts. This will mainly include bubbles (or circular films, as cited in older works [3, 4]), but as many of their effects are shared and easier to observe with flat films and foams [5], those will also often be referenced. Indeed, foams are essentially composed of a multitude of small bubbles, each sharing film surfaces, and meeting at "Plateau" borders where fluid tends to accumulate [6].

Firstly, we will establish the different notions and key behaviours pertinent for stable films such as surface tension and the Marangoni effect, and their intrinsic link with surfactant molecules. These tend to adsorb to fluid boundaries [7] and form stabilising layer structures [8]. Building from this, we will essentially construct a timeline including different phenomena the bubble will experience during its lifetime: Starting with the drainage mechanisms caused by capillary and hydrostatic pressures [9], causing slow Poiseuille flow when the bubble's surfaces are considered "immobile" [10, 11, 12], to the much more common and complex Marangoni-induced marginal regeneration flows which further act to thin the film [13, 14, 11].

Pausing to look at a curious metastable equilibrium state, the "Newton" or "common" black film [15], which allows the film to exist at an extremely thin level, down to two surfactant atoms thick [16]. The nucleation and position of the hole that eventually forms on the surface will then be studied [17, 18], before its opening dynamics [19], rim destabilisation (causing interesting structures like ligaments, or cracks and folds to appear [20, 21]), and atomisation into a mist of thin aerosol particles [12] are finally discussed.

## 2 Stable curved films

### 2.1 Surface tension

Surface tension is a measurement for how strongly molecules of the same substance are attracted to each other at a fluid boundary, via intermolecular forces. It appears as a parameter or non-dimensional number in virtually every experiment, model and energetic analysis which attempts to describe the dynamics and behaviours of thin films [22]. Film and bubble persistence [12], critical thickness [23, 9] and hole rim velocity [19] are just a few characteristics which utilise it. It is the driving potential behind the Marangoni effect, studied in section 2.3, which as we will see can act either to stabilise or destabilise the film in certain situations.

### 2.2 Role of surfactant

Surfactants play a large role in the stability of thin films, we can all recognise this naturally as in every day life, we only really see them when soap is involved (soap contains surfactants). To understand their importance, we will describe their effect on surface tension, the structures they adopt to stabilise thin films, and how their role in thin film stability is evaluated.

Surfactants, or surface-active-agents, are molecules which lower the surface tension of fluid boundaries. They achieve this through their amphiphilic structure; their molecular composition admits a combination of hydrophobic and hydrophilic groups. This causes them to migrate, orient themselves and settle at liquid boundaries, displacing some of the liquid's molecules [7]. The molecules are said to have 'adsorbed' to the surface. Since intermolecular forces between water molecules are much stronger than forces between water and surfactant molecules, this will have the effect of reducing surface tension. This change is often qualified as "film" or "surface" pressure, which is equal to the difference in surface tension between the clean ("bare") liquid surface, and the surfactant laden surface [14] (de Gennes calls it "Langmuir" pressure in [24]).

Although the addition of surfactants can bring new thinning mechanisms such as marginal regeneration [13, 10], which act to weaken thin films, it has been seen time and time again [14, 17, 12, 9, 25], and is widely accepted that the presence of these amphiphiles will always increase film stability, when compared to their "bare" counterparts [18]. Indeed, as Lhuissier notes, "an astonishingly small amount of surface-active components" can increase a bubble's lifetime from practically 0 to times longer than 1s [9, 26]. On top of reducing surface tension, surfactant molecules will naturally arrange themselves and form stabilising structures. Karol

J. Mysels, a central player in the understanding of soap films in the 1960s, describes these as 'monolayers' and 'bilayers' [8]. In the case of a thin film surrounded by air for example, the surfactant molecules arrange themselves along the boundaries, forming two layers exhibiting a myriad of complex surface stresses, which have significant consequences for the film's dynamics and stability. In most cases, the film is not thin enough for the two boundaries to interact via molecular forces, in which case, studies will often describe them as a set of two surfactant monolayers. However in specific situations, the formation of extremely thin but stable films (5 to 10 nm [15]) is possible, in which case these interactions become significant and the film is said to admit a surfactant bilayer. These interesting films, called "common" or "Newton" black films, are further looked into in section 3.1.3.

The increased stability of surfactant laden thin films can be evaluated in many ways, such as critical thickness, lifetime and "fragility". Critical thickness and lifetime are intrinsically linked by drainage mechanisms that are described further in section 3.1. In an early paper, Manev E. and others observed a relationship between the minimum stable film thickness (before bursting) and surfactant concentration [23]. Experimentally, they found that an increased surfactant concentration will lead to a smaller critical thickness up to a certain threshold after which concentration has little to no effect. Although they do not mention it, this is recognised to match with the critical micelle concentration (cmc) of the surfactant used. This concentration had first been observed in 1926, and meticulously reviewed by P. Mukerjee and K. J. Mysels [27] to be the concentration at which adding surfactant has little to no effect on the surface tension of the fluid. At that point, the amphiphiles start to aggregate and group together to form 'micelles' on the surface.

In a recent 2016 paper, Champougny L. and others introduced bubble "fragility" as an interesting and unique quantification for stability [17]. They qualified it to be how sensitive bubbles are to external perturbations such as solid objects or dust. Carrying out an experiment in which bubbles of different soap concentration were formed and put in contact with a solid object, they found that films with dilute surfactant concentrations ( $< 0.8$  cmc) would burst on contact whereas for  $c = \text{cmc}$  and above, objects could penetrate without damage. Although this was the only paper I have seen to study this effect, I believe this result is sensible and further affirms the stabilising power surfactants have on thin films.

### **2.2.1 "Rigid" and "Mobile" film surfaces**

Surfactants with different solubility, surface viscosity and concentrations will tend to form two main classes of surfaces; "rigid" and "mobile" layers, as originally classified by Mysels in the

1960s [13]. For surfactants at lower concentrations (below the cmc), the adsorbed molecules are able to move easily along the surface, resulting in a "mobile" surface. These sometimes called "simple" mobile films [28] will exhibit complex turbulent motions, for instance seen in the colourful swirls appearing in common soap bubbles. On the other hand, at high concentrations, the molecules will form "rigid" and compact layers. These surfaces, deemed as "immobile" or static, will lead to slower drainage rates, and different stabilising effects [14]. Structures formed by insoluble surfactants at high concentration, and soluble surfactants at high concentrations are slightly different, but exhibit similar phenomena [14, 17]. Overall, the nature of surfactants used to stabilise a film has profound effects on the different drainage and rupture mechanisms it will experience, these include the possibility of marginal regeneration, and formation of a black film, which are discussed further in sections 3.1.2 and 3.1.3.

## **2.3 The Marangoni Effect**

## **2.4 Bubble Coalescence**

## **2.5 Behaviour in Electromagnetic fields**

## **2.6 Applications**

## **2.7 Discussion**



### **3 Bubble Death and film rupture**

Thin films will experience a multitude of effects from the moment they are created, that weaken and push them to become unstable and eventually burst. In the following sections, these effects will be addressed chronologically: starting with initial drainage mechanisms which weaken the films, the moment rupture nucleates, to the dynamics and interesting fluid structures that develop around the newly formed accelerating hole, and finally what becomes of the fluid once the films have burst. As they exhibit many more interesting drainage and rupture mechanisms, this section will focus predominantly on surface bubbles in contact with their founding bulk fluid, while making sure to note when the effects are also relevant to "free" bubbles and films, and foams. The thin film part of the surface bubble is referenced as the "cap" and the edge at which it makes contact with the water is referenced as the "meniscus", "periphery" or "base" throughout.

#### **3.1 Film thinning mechanisms and rupture initialisation**

All bubbles ranging from the ones formed with minimal surfactant concentrations (like tap water), which have lifetimes in the order of seconds [29, 9] to ones stabilised with high surfactant concentrations (well above the CMC) with lifetimes in the order of minutes or even hours, will exhibit film thickness reduction over time due to different drainage mechanisms. These effects cause fluid in the bubble's cap to be escape down into the bulk fluid on which the bubble is standing. These phenomena, dependant on the film's age, size and surfactant, will eventually lead the film thickness to reach some critical value and cause a hole to nucleate, sealing it's fate and beginning its annihilation.

##### **3.1.1 Gravity and capillary drainage**

Gravity and capillary drainage are thinning mechanisms with occur regardless of the presence of surfactants adsorbed to the film and in the bulk solution. They typically appear for younger bubbles, right at the moment of their creation, before any stabilising surface tension gradients can be established [9]. The gravity driven drainage can simply be understood as the hydrostatic pressure generated by the mass of fluid in the bubble's cap, which will tend to push fluid out through the bubble's periphery and into the bulk fluid. Capillary drainage, driven by the capillary pressure inside the cap, acts more as a suction force which pulls liquid out of the cap. This effect occurs due to the thin nature of the film, and is caused by the surface tension of the fluid.

As pointed out by Lhuissier and Villiermaux [9], the relative importance of these two effects depends on the surface bubble's size. Smaller bubbles tend to be more spherical and sit lower into the bulk fluid, whereas larger bubbles will tend to the half-sphere shape, with their thin film cap being mostly above the bulk fluid level. Recalling the weight of fluid in the cap is very small, and often neglected in models [22], capillary pressure will therefore dominate for most bubbles which follow  $R \leq 5a$  [9], where  $a$  corresponds to the capillary constant and  $R$  is bubble radius. This commonly used capillary length constant is a function of both gravity and surface tension. It can be noted here that this same capillary drainage mechanism also applies in foams [5], where instead of draining into the bulk liquid, fluid will drain from the films into what are called Plateau borders [6]. These are the edges at which the soap films, which constitute the foam, meet.

Regardless of whether hydrostatic or capillary forces are more significant, in both cases the drainage rates are relatively easy to compute, as they are simple pressure driven flows. A pressure gradient will drive liquid out of the cap, through the meniscus and into the bulk liquid (or directly into a Plateau border in the case of foams). In this regard, the Hagen-Poiseuille law or even very simple plug flow equations are used to determine the flowrate out of the cap, and therefore the thickness evolution of the film. Plug flow assumes a constant fluid velocity across the thickness of the cap, and was used by numerous authors, especially when negligible surfactant concentrations are considered [18, 9, 26, 17]. Poiseuille equations on the other hand were used in the case where surfactants have formed a pair of rigid and immobile monolayers [10, 11, 12], these are described in section 2.2.1.

The importance of surfactants in extending the lifetime of water surface bubbles specifically can now be clearly quantified. When no surfactants are present, the capillary and hydrostatic pressures lead to a very rapid drainage of water out of the cap. The only effect that can slow down this drainage is viscous stretching, where viscous forces oppose the thinning and deformation of the cap. However, this still leads to the drainage having characteristic times calculated using  $\mu/\rho g L$ , which yield values in the order of  $10^{-4}$  s [26, 9]. The "bare" water bubbles therefore appear to burst immediately. Noting how viscosity appears as a factor in the drainage time, G. Debregeas et al [18] carried out drainage and bursting experiments on very viscous polymer melt and glass bubbles. This allowed them to study bubbles with significantly longer lifetimes (in the order of minutes), that admit essentially no surfactants as these are less easily adsorbed from the environment. Examining the cap thickness evolution for differently sized bubbles, they explained and effectively corroborated their results with plug flow equations utilising the capillary length. Extending these equa-

tions down to bare water bubbles, they also found similarly small characteristic times found by Breward and Lhuissier [26, 9]. Overall this shows that for non-surfactant laden films, the thickness evolution can effectively be modeled by following the gravity and capillary drainage mechanisms.

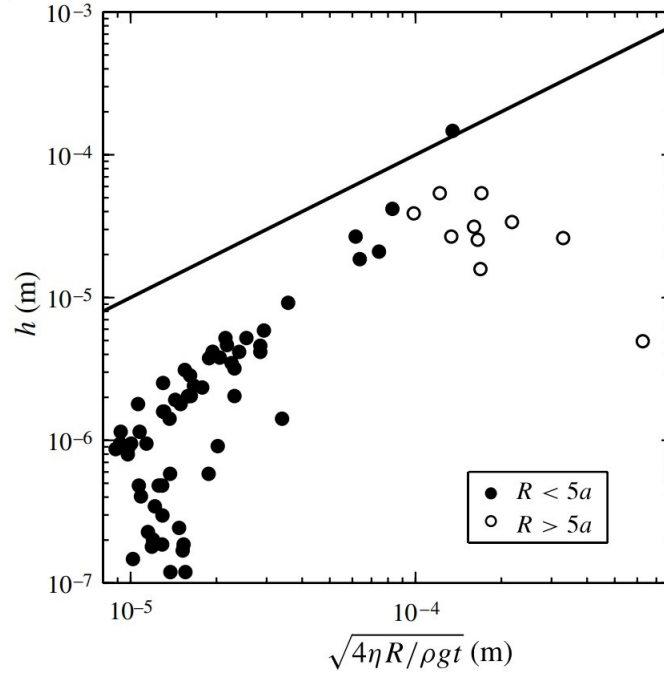


Figure 3.1: Comparison of the drainage law which only includes Poiseuille flow, with experiments [9]. The law used clearly underestimates the actual physical drainage rates.

However, as seen in figure 3.1, hydrostatic and capillary pressure driven drainage alone fail to model, and drastically underestimate, the thinning behaviour of films which do contain surfactants. Indeed, drainage seems to occur more rapidly than the expected Poiseuille flow. This discrepancy comes from an additional draining effect, marginal regeneration, which is studied in the next section (3.1.2).

### 3.1.2 Marginal Regeneration

Marginal regeneration is a film drainage mechanism caused by the presence of free moving particles, like soluble surfactants, adsorbed at the surface of a thin film. First observed and introduced by Mysels in 1959 [13], it results from the formation and movement of asymmetrical plumes of liquid along the base of films (and Plateau borders in foams [26]), as seen in figure 3.2. Observing the light interference fringes which appear when looking at a bubble from straight above, S. Bhamla et al were also able to clearly show these [14]. Essentially, these additional fluid motions cause an additional liquid flux out of the bubble's cap, increasing its drainage rate. At relatively low surfactant concentrations, below the CMC, marginal

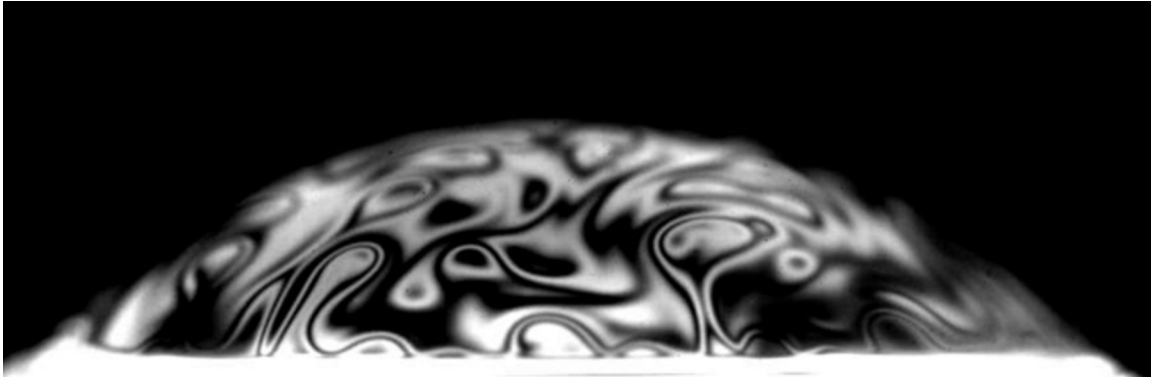


Figure 3.2: View of a Bubble (radius  $R = 5$  mm) under monochromatic lighting (sodium lamp at 589 nm) [9]. Fringes seen correspond to zones of constant thickness. Although they are usually studied in flat thin films, as it is easier to have a consistent light source pass through, here the convection plumes at the base of the film causing marginal regeneration can clearly be seen.

regeneration is seen to dominate over the slower Poiseuille flows studied in section 3.1.1 [9, 14].

Physical explanations for these flows, and how they actually act to thin the bubble caps vary slightly for different authors [10, 14, 9, 17, 11]. Originally, Mysels [13] believed these were caused by small thickness variations caused by thermal effects. However, more recent studies all agree on one point: the complex flows observed are driven by surface tension gradients. S. Bhamla et al [14] explain the surface flow by arguing that as a plume is created along the surface of the film, its head causes the film to dilate and "creates a new surface area", which decreases the surface concentration of surfactants, hence creating a zone of higher surface tension. This will result in more liquid being pulled from the bulk towards the depleted zone, essentially further fueling the motion. Although this explanation doesn't outwardly specify it, it is important to note that these plumes are zones of decreased thickness, which act to drain the film as thicker zones are drawn into the bulk fluid [9, 10], as seen in figure 3.3. R. Bruinsma et al [11] attribute the movement to buoyancy forces. They also underline that the flow from the film to the meniscus not only deposits fluid, but also surfactant molecules. Furthermore, fluid pulled from the bulk liquid up into the film has a lower surfactant concentration, because as mentioned before surfactants are drawn to fluid boundaries. The surface concentration of surfactant molecules is generally higher than the concentration in the bulk fluid [7].

The actual inception of these plumes, or "patches" as characterised by Champougny et al [30], can be attributed to the existence of a pinching region, seen in figure 3.3. This thin zone quickly appears along the base of the bubble cap by a process called 'marginal pinching' [31, 32], and its destabilisation causes the plumes to be formed. Simply put, the pressure

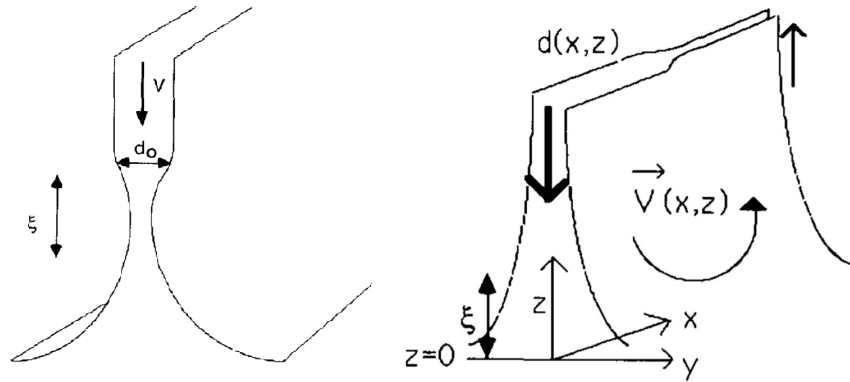


Figure 3.3: On the left: Diagram of the thinned zone caused by marginal pinching, with characteristic lengths used to setup the drainage laws. On the right: Marginal regeneration flows at the bottom of a film, thicker regions move down into the meniscus while thinner region (plumes) rise up into the film. [11]

inside of the meniscus (or Plateau border) is lower than the capillary pressure exerted by the thin film. This difference is what generates the flow out of the cap, but it also causes marginal pinching. According to Lhuissier and Villerraux [9], the relatively large surface tension gradient difference between the two sides of the thinned zone is what causes the Benard-Marangoni destabilisation triggering marginal regeneration.

Overall, models which take the marginal regeneration effect into account can be established, either by coupling it directly with capillary pressure [9], by considering mass balances for both surface, intralamellar (between the two surfactant monolayers) and bulk fluid [10], or by applying boundary conditions and assumptions to existing surface models (like the Boussinesq-Scriven model) [14]. These generally yield solutions which are much closer to experimental values.

### 3.1.3 Newton Black films

You might have encountered this curious state when observing floating soap bubbles. In the seconds before their death, they sometimes becomes practically invisible, no longer exhibiting colourful interference patterns. In certain cases before rupture, as the drainage mechanisms cause liquid to leave the film, there will be a limit at which thickness becomes very small (5 to 10nm), and the film stabilises. This resulting extremely thin metastable equilibrium is referenced as Newton (NBF) or common "Black" film (CBF), as A. Vrij and J. Th. G. Overbeek explain: "black, because interference between light reflected from the front side and that reflected from the back side leads to nearly complete extinction of the reflected light" [3]. The films aren't actually black per se, they just happen to reflect such little light that they appear invisible [28], or black in front of a black background. Here, this equilibrium

state is qualified as metastable as it is very sensitive to external fluctuations. These films are found to appear in higher surfactant concentration bubbles, above the CMC [14, 23], in patches located towards the apex of bubbles before rupture. This is understandable, as the film is usually thinnest at those locations. In a recent research article attempting to develop a model for describing these patterns observed at the surface of surfactant-laden bubbles, L. Shen et al [1] describe these regions as "black holes". They characterise these to be in a close packing configuration, where the maximum amount of surfactant molecules have adsorbed to the surface. This can be recognised as the critical micelle concentration described earlier, and agrees with the previous works which noticed the formation of black films at similarly high concentrations [14, 23].

This unexpected stability is often said to come from a balance of the capillary suction with intermolecular interactions (they become non-negligible at this thickness) [26, 30, 33, 3, 4]. The intermolecular forces at work are: the electrostatic repulsion and van der Waals attractions of the molecules from the two boundaries of the surfactant bilayer, and the "hard-core" repulsion of the molecules themselves when the film is at its thinnest [16]. Note that the overall sign of the intermolecular interactions should be repulsive for black films to form. If they are not, they will accelerate the thinning and cause rupture, desabilising the film instead of stabilising it [4, 28, 18]. It should be restated here that black films can be categorised in two forms: the thicker "common" black film (10 nm), and the thinnest possible "Newton" black film, which has a typical thickness of only two surfactant molecules (roughly 4-5 nm) [16, 15]. Interestingly, V. Casteletto et al [16] noticed and studied an apparent hysteresis in the CBF-NBF transition. Indeed, by varying the internal pressure of a soap film created on the end of a small hole (2 mm), and measuring its thickness via interferometry, they were able to find a range of pressures at which either of CBF and NBF could be observed. This previously undocumented non-linear nucleation effect could be linked back to the spontaneous thickness changes which were observed and analysed by A. Vrij and J. Th. G. Overbeek 50 years prior [3].

Most experiments described in this report neglect, or do not take into account the effect of ambient humidity and evaporation in their models [23, 29]. Only a few mention in passing its "importance" or controlling it, but without citing its effects [28]. Intuitively, it could be perceived that at such small thicknesses, this external flux could potentially have an impact. In a recent 2018 paper, L. Champougny et al [30] studied this effect by measuring the maximum length reached by a flat film pulled at constant velocity until rupture, for different ambient humidity. Surprisingly, they found a positive correlation between humidity and critical film

length. However by closer observing the thinning dynamics, they did not find evaporation to affect these, this lead them to believe evaporation only starts to become significant very close to rupture.

#### **3.1.4 Rupture initialisation**

Now that we have examined the thinning effect of drainage mechanisms, the following questions arise: what triggers the film to actually begin rupturing, why don't all films thin down to black film thickness, and why can't they remain so indefinitely? Overall, the theories outlined below all attempt to find a value for the "critical thickness" of the film at the onset of bursting. This somewhat simple value, when linked with the thickness evolution equations derived from the drainage mechanisms, could be used to predict the film's lifetime.

Black films are believed to rupture unpredictably due to the lack of control over the environment: thermal shocks, vibrations, dust, or as recently observed, evaporation [30]. Recently, Lhuissier and Villermaux [9] considered a statistical approach by admitting a frequency of thermal fluctuation events from the environment, where an activation energy would cause rupture (this is the nucleation energy for the creation of a hole at the surface). Although this was able to successfully estimate the lifetime for black films, it yielded non-physical values for thicker films. For these non-black films, the energy of activation for nucleating a hole is considered too large for these phenomena to have an effect [4]. Instead, thermal fluctuations in the environment were believed to cause surface "corrugations" (variations in thickness) to appear [3, 4], which grow over time, and when a critical wavelength is reached, to cause spontaneous rupture. Vrij developed a popular model for this wavelength, based on surface tension and the "Free energy of interaction", being the long-range intermolecular forces acting in the surfactant bilayer. Although this model was able to somewhat accurately predict critical thicknesses for certain experiments at the time, it incorrectly models the effect of bubble radius and surface tension for more recent experiments [28]. As Vrij admits, his model did not include the complex movement of surfactant molecules at the surface, and their role in surface elasticity. The latter was later proven to be an important factor in describing critical thickness by D. Langevin and A.A. Sonin in 1994 [28].

More recent advancements made in understanding the marginal regeneration effect for example [14], have led alternative explanations to be explored as a substitute for thermal effects. These include its role in catalysing turbulent flows along the surface of bubbles which, following an energy scale analysis, could show rupture is a possible consequence (albeit unconvincingly [9]). Alternatively, Lhuissier and Villermaux also observed that nucleation of

holes in tap water bubbles occurred predominantly at "convection cells" [9], these are the rising plumes described in section 3.1.2. This seems rather evident, as these are known to be zones of reduced thickness. They developed a model which revolves around probabilities for a thin enough plume to rupture, which turned out to be a function of the capillary length scale. Although this model does fit well with experiments [12], it still relies on a probability factor they call the "puncture efficiency". This value has no current physical relationship and is experimentally derived. This makes their solution appear more like probability fitting rather than being rooted in physically explained mechanisms. In parallel, more recent studies have considered models based on molecular dynamics simulations [34], or further investigations into the thermal effects [35]. However, although experiments over the years have provided a thorough quantitative understanding of the phenomena, none have been able to provide a solely qualitative model to explain the critical thickness and onset of rupture in thin films.

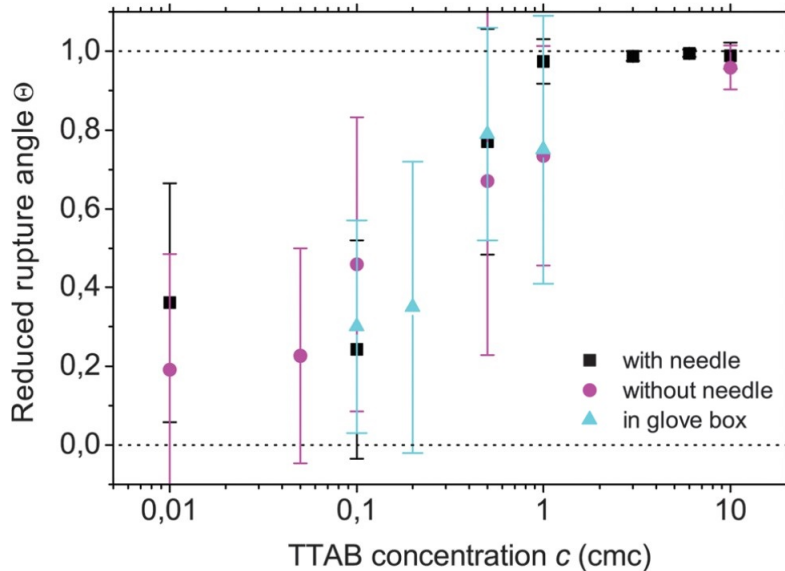


Figure 3.4: Reduced rupture angle  $\Theta$  (0 corresponds to a rupture nucleated at the meniscus, tends to 1 as the point approaches the apex of the bubble) as a function of TTAB surfactant concentration, normalized with respect with its cmc. Each symbol corresponds to a different experimental condition, and each point shows the average and standard deviation for 4 to 20 bubbles [17].

Despite this, an interesting remark can be made for the position of the nucleated hole along the bubble surface. Overall, it seems that when marginal regeneration is the driving drainage mechanism, bubbles will tend to burst from a point close to their periphery [14, 9]. On the other hand, "bare" bubbles, and bubbles with immobile surfactant surfaces (section 2.2.1) tend to burst at their apex. Champougny et al [17] carried out experiments to study this effect, which yielded figure 3.4. This data clearly shows with relative certainty that high surfactant concentration bubbles do indeed burst at their apex. This can be explained by



the monolayers becoming "saturated", resulting in immobile surfaces which only exhibit the symmetric gravity and capillary drainage mechanisms [14]. This symmetry implies the hole can only appear at the apex. Values for lower concentrations however come with large uncertainties, and are reminiscent of the stochastic nucleation behaviour caused by marginal regeneration. In this aspect, a probabilistic model is required to predict the position of the nucleation hole. In short, the rupture behaviours for mobile surfactant films have yet to be fully defined, and it is difficult to do so due to the stochastic nature of the problem.

### 3.2 Hole Development

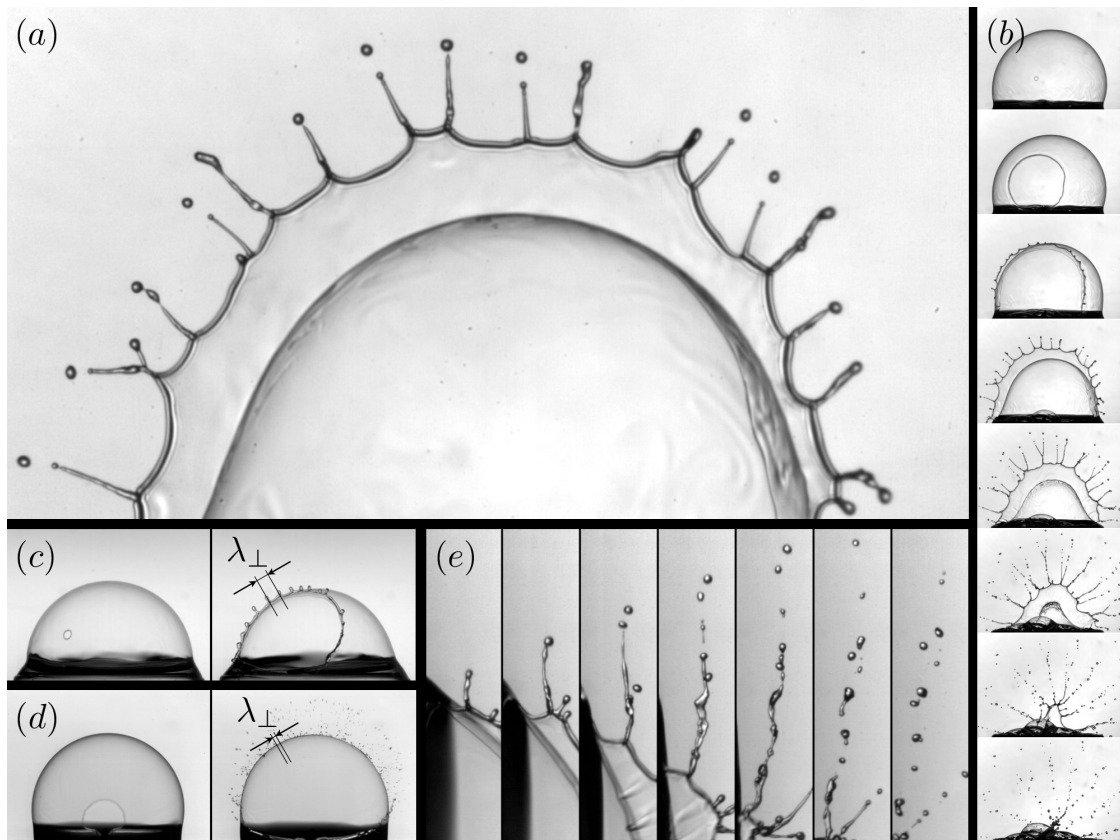


Figure 3.5: Fluid structures emerging from the opening rim of the nucleated hole in a soap bubble [21].

Now that a hole has nucleated somewhere along the bubble's surface, it will start to grow and widen, accumulating fluid in its rim until instabilities grow and cause it to fragment into a multitude of fluid structures and droplets. These steps can be admired in figure 3.5. Although this behaviour might not seem to be directly linked to stability (as the bubble is already doomed), the dynamics of film rupture are very interesting and have significant impacts for their applications. Indeed, bubbles play a key role in the production of natural and man-made aerosols [9, 12]. On top of this, if the dynamics of hole opening can be well defined, this could give more insight into the process causing its nucleation.

### 3.2.1 Rim Velocity

$$v = \sqrt{\frac{2\sigma}{h\rho}} \quad (3.1)$$

The rim velocity pertains to the velocity of the initially circular opening hole now present in the film. Primitively derived by Rayleigh using a simple energy balance [36], it was first improved by Taylor and Culick independently to include the inertia of the liquid accumulating at the "rolling edge" [37, 19]. This known as the Taylor-Culick velocity, available in equation 3.1, where  $\sigma$  is surface tension,  $h$  is the film thickness at rupture (assumed uniform), and  $\rho$  is the liquid density. Although, as we will see, this model is known to overestimate rim velocities over time, it is a good approximation for the immediate velocity of the opening hole (for relatively low viscosity fluids, like water). More recent papers still utilize it often, for example to verify the critical film thickness [12, 18].

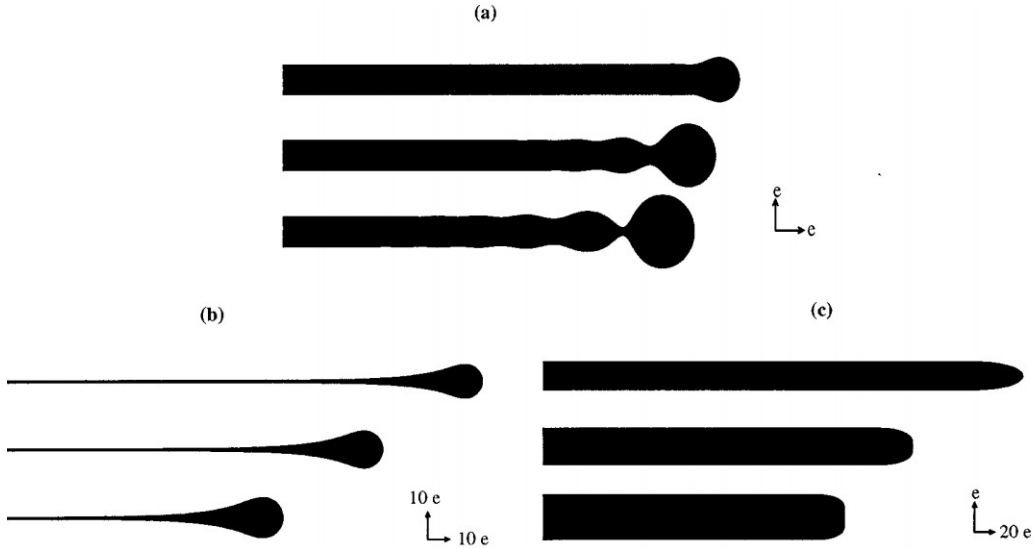


Figure 3.6: Film retraction profiles for (a)  $Oh = 10^{-6}$  (b)  $Oh = 10^2$  (c)  $Oh = 10^8$ . Arrow lengths are given in terms of film thickness  $e$ . Here, the absence of fluid accumulation for high  $Oh$ , and appearance of capillary wave disturbances ahead of the edge for low  $Oh$  can clearly be seen [38].

A known fact about the Taylor-Culick velocity is that it neglects the viscosity, since the Reynolds number is generally very high [38]. However, for high viscosity fluid films, or "sheets", this is not the case, and retraction behaviours deviate from the usual. First looked into by Debrégeas and de Gennes in 1998 [18], these have been studied extensively since [38, 39]. Interestingly, for bare high viscosity films, fluid doesn't accumulate in the rim and the radius of the hole actually grows exponentially [18]. It is useful to introduce a non-dimensional number which appears commonly in thin film problems: the Ohnesorge number  $Oh$ , which expresses the relative effects of viscosity and surface tension. This number al-

lows us to determine which retraction regime a bubble might have: at low  $Oh < 0.1$ , "capillary wave disturbances" can be observed ahead of the rim which accumulates liquid, whereas at high  $Oh > 10$ , viscosity dominates, and no rim is formed [39, 38]. These effects are shown schematically in figure 3.6.

### **3.2.2 Appearance of instabilities**

## **3.3 The aftermath**

## **3.4 Applications**

## **3.5 Discussion**

Marginal regeneration was only specifically observed in surface bubble and foams, and examples of it occurring in free floating spherical bubbles couldn't be found. However, knowing that fluid generally builds up at the base of concave films [1], we could expect a similar drainage flow to appear at that point. Instead of forming a convecting flow with its bulk fluid, or a plateau border, it could form this mass flux with this built-up fluid mass.

Overall, drainage mechanisms and film rupture are a priori linked, as thinner films are more likely to rupture. Nevertheless, the onset of bursting, although quantitatively well described is not yet fully understood. Further research is required to truly understand the physical phenomena causing thicker films to rupture (thicker than black films), and develop satisfactory models which can be applied generally.

# **4 Improving stability**

# **5 Conclusions**

## References

- [1] Shen L, Denner F, Morgan N, van Wachem B, Dini D. Transient structures in rupturing thin films: Marangoni-induced symmetry-breaking pattern formation in viscous fluids. *Science Advances*. 2020;6(28). Available from: <https://advances.sciencemag.org/content/6/28/eabb0597>.
- [2] Benelmekki M, Erbe A. Chapter 1 - Nanostructured thin films—background, preparation and relation to the technological revolution of the 21st century. In: Benelmekki M, Erbe A, editors. *Nanostructured Thin Films*. vol. 14 of *Frontiers of Nanoscience*. Elsevier; 2019. p. 1 – 34. Available from: <http://www.sciencedirect.com/science/article/pii/B9780081025727000015>.
- [3] Vrij A, Overbeek JTG. Rupture of thin liquid films due to spontaneous fluctuations in thickness. *Journal of the American Chemical Society*. 1968;90(12):3074–3078. Available from: <https://doi.org/10.1021/ja01014a015>.
- [4] Vrij A. Possible mechanism for the spontaneous rupture of thin, free liquid films. *Discuss Faraday Soc*. 1966;42:23–33. Available from: <http://dx.doi.org/10.1039/DF9664200023>.
- [5] Braun RJ, Snow SA, Naire S. Models for gravitationally-driven free-film drainage. *Journal of Engineering Mathematics*. 2002 Aug;43(2):281–314. Available from: <https://doi.org/10.1023/A:1020320404339>.
- [6] Almgren FJ, Taylor JE. The Geometry of Soap Films and Soap Bubbles. *Scientific American*. 1976;235(1):82–93. Available from: <http://www.jstor.org/stable/24950395>.
- [7] Adamson A, Gast A. *Physical chemistry of surfaces*. New York: Wiley; 1997.
- [8] Mysels KJ. Soap films and some fundamentals of thin-liquid films. *Journal of the American Oil Chemists' Society*. 1968;45(3):139–140. Available from: <https://aocs.onlinelibrary.wiley.com/doi/abs/10.1007/BF02915337>.
- [9] Lhuissier H, Villerraux E. Bursting bubble aerosols. *Journal of Fluid Mechanics*. 2011;696:5–44.
- [10] Nierstrasz VA, Frens G. Marginal Regeneration and the Marangoni Effect. *Journal of Colloid and Interface Science*. 1999;215(1):28 – 35. Available from: <http://www.sciencedirect.com/science/article/pii/S0021979799962264>.
- [11] Bruinsma R. Theory of hydrodynamic convection in soap films. *Physica A: Statistical Mechanics and its Applications*. 1995;216(1):59 – 76. Available from: <http://www.sciencedirect.com/science/article/pii/037843719500023Z>.
- [12] Modini RL, Russell LM, Deane GB, Stokes MD. Effect of soluble surfactant on bubble persistence and bubble-produced aerosol particles. *Journal of Geophysical Research: Atmospheres*. 2013 feb;118(3):1388–1400. Available from: <https://doi.org/10.1002/jgrd.50186>.
- [13] Mysels KJ. *Soap films : studies of their thinning and a bibliography*. Lond. : Pergamon Press; 1959.
- [14] Bhamla MS, Chai C, Álvarez Valenzuela MA, Tajuelo J, Fuller GG. Interfacial mechanisms for stability of surfactant-laden films. *PLOS ONE*. 2017 05;12(5):1–14. Available from: <https://doi.org/10.1371/journal.pone.0175753>.

- [15] Jang SS, Goddard WA. Structures and Properties of Newton Black Films Characterized Using Molecular Dynamics Simulations. *The Journal of Physical Chemistry B*. 2006;110(15):7992–8001. PMID: 16610899. Available from: <https://doi.org/10.1021/jp056685c>.
- [16] Casteletto V, Cantat I, Sarker D, Bausch R, Bonn D, Meunier J. Stability of Soap Films: Hysteresis and Nucleation of Black Films. *Phys Rev Lett*. 2003 Jan;90:048302. Available from: <https://link.aps.org/doi/10.1103/PhysRevLett.90.048302>.
- [17] Champougny L, Roché M, Drenckhan W, Rio E. Life and death of not so “bare” bubbles. *Soft Matter*. 2016;12:5276–5284. Available from: <http://dx.doi.org/10.1039/C6SM00178E>.
- [18] Debrégeas G, de Gennes PG, Brochard-Wyart F. The Life and Death of “Bare” Viscous Bubbles. *Science*. 1998;279(5357):1704–1707. Available from: <https://science.sciencemag.org/content/279/5357/1704>.
- [19] Culick FEC. Comments on a Ruptured Soap Film. *Journal of Applied Physics*. 1960;31(6):1128–1129. Available from: <https://doi.org/10.1063/1.1735765>.
- [20] Bico J. Cracks in bursting soap films. *Journal of Fluid Mechanics*. 2015;778:1–4.
- [21] Lhuissier H, Villermaux E. Bursting bubbles. *Physics of Fluids*. 2009;21(9):091111. Available from: <https://doi.org/10.1063/1.3200933>.
- [22] Ida MP, Miksis MJ. The Dynamics of Thin Films I: General Theory. *SIAM Journal on Applied Mathematics*. 1998;58(2):456–473. Available from: <http://www.jstor.org/stable/118341>.
- [23] Manev E, Scheludko A, Exerowa D. Effect of surfactant concentration on the critical thicknesses of liquid films. *Colloid and Polymer Science*. 1974 jul;252(7-8):586–593. Available from: <https://doi.org/10.1007%2Fb01558157>.
- [24] de Gennes PG. “Young” Soap Films. *Langmuir*. 2001;17(8):2416–2419. Available from: <https://doi.org/10.1021/la001538l>.
- [25] Pfeiffer P, Zeng Q, Tan BH, Ohi CD. Merging of soap bubbles and why surfactant matters. *Applied Physics Letters*. 2020 mar;116(10):103702. Available from: <https://doi.org/10.1063%2F1.5135764>.
- [26] Breward CJW, Howell PD. The drainage of a foam lamella. *Journal of Fluid Mechanics*. 2002;458:379–406.
- [27] Mukerjee P, Mysels KJ. Critical micelle concentrations of aqueous surfactant systems. National Standard reference data system; 1971.
- [28] Langevin D, Sonin AA. Thinning of soap films. *Advances in Colloid and Interface Science*. 1994;51:1 – 27. Available from: <http://www.sciencedirect.com/science/article/pii/0001868694800332>.
- [29] Zheng Q, Klemas V, Hsu a. Laboratory Measurement of Water Surface Bubble Life Time. *Journal of Geophysical Research*. 1983 01;88:701–706.
- [30] Champougny L, Miguët J, Henaff R, Restagno F, Boulogne F, Rio E. Influence of Evaporation on Soap Film Rupture. *Langmuir*. 2018;34(10):3221–3227. PMID: 29498527. Available from: <https://doi.org/10.1021/acs.langmuir.7b04235>.

- [31] Joye JL, Hirasaki GJ, Miller CA. Asymmetric Drainage in Foam Films. *Langmuir*. 1994;10(9):3174–3179. Available from: <https://doi.org/10.1021/la00021a046>.
- [32] Aradian A, Raphaël E, Gennes PGd. “Marginal pinching” in soap films. *Europhysics Letters (EPL)*. 2001 Sep;55(6):834–840. Available from: <http://dx.doi.org/10.1209/epl/i2001-00356-y>.
- [33] Cantat I. *Les mousses structure et dynamique*. Paris: Belin; 2010.
- [34] Nguyen TD, Fuentes-Cabrera M, Fowlkes JD, Rack PD. Coexistence of spinodal instability and thermal nucleation in thin-film rupture: Insights from molecular levels. *Phys Rev E*. 2014 Mar;89:032403. Available from: <https://link.aps.org/doi/10.1103/PhysRevE.89.032403>.
- [35] Shah MS, van Steijn V, Kleijn CR, Kreutzer MT. Thermal fluctuations in capillary thinning of thin liquid films. *Journal of Fluid Mechanics*. 2019;876:1090–1107.
- [36] Rayleigh JWS. *Scientific Papers by Lord Rayleigh, Vol. III, 1887-1892 (1902)*. Dover Publications; 1964. Available from: <https://books.google.co.uk/books?id=qE0mygEACAAJ>.
- [37] Taylor GI. The dynamics of thin sheets of fluid. III. Disintegration of fluid sheets. *Proceedings of the Royal Society of London Series A Mathematical and Physical Sciences*. 1959;253(1274):313–321. Available from: <https://royalsocietypublishing.org/doi/abs/10.1098/rspa.1959.0196>.
- [38] Brenner MP, Gueyffier D. On the bursting of viscous films. *Physics of Fluids*. 1999;11(3):737–739. Available from: <https://doi.org/10.1063/1.869942>.
- [39] SAVVA N, BUSH JWM. Viscous sheet retraction. *Journal of Fluid Mechanics*. 2009;626:211–240.