# Astro 507; Problem Set 2

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# 1. Heat Capacities

The definition of heat capacity gives that

$$C \equiv \frac{\mathrm{d}Q}{\mathrm{d}T} \tag{1}$$

 $C_P$  and  $C_V$  are the same as C but with the pressure, P, and volume, V, each held constant respectively. Applying the first law of thermodynamics, this gives

$$C_V = \left(\frac{\mathrm{d}U}{\mathrm{d}T}\right)_{V,N} \qquad C_P = \left(\frac{\mathrm{d}U}{\mathrm{d}T}\right)_P + P\left(\frac{\mathrm{d}V}{\mathrm{d}T}\right)_P$$
 (2)

Now let's apply the equipartition theorem and therefore write that

$$U = \frac{1}{2} N f k_B T \tag{3}$$

where f is the number of degrees of freedom. Now let's do some derivatives and apply the ideal gas law  $(V = Nk_BT/P)$ 

$$C_V = \frac{\mathrm{d}}{\mathrm{d}T} \left( \frac{1}{2} N f k_B T \right) \tag{4}$$

$$=Nk_B \cdot \frac{f}{2} \tag{5}$$

$$C_P = \frac{\mathrm{d}}{\mathrm{d}T} \left( \frac{1}{2} N f k_B T \right) + P \frac{\mathrm{d}}{\mathrm{d}T} \left( \frac{N k_B T}{P} \right) \tag{6}$$

$$=\frac{1}{2}Nfk_B+Nk_B\tag{7}$$

$$=Nk_B\left(1+\frac{f}{2}\right) \tag{8}$$

Therefore, the difference between the two is

$$C_P - C_V = Nk_B$$
 (10)

(9)

And dividing the two gives the adiabatic index

$$\boxed{\frac{C_P}{C_V} = 1 + \frac{2}{f} = \gamma} \tag{11}$$

We are given the following information:

$$V = (40 \,\mathrm{kpc})^3, \qquad n = 0.01 \,\mathrm{cm}^{-3}, \qquad T = 10^6 \,\mathrm{K}$$
 (12)

This means that the total number of particles in this volume is approximately

$$N = nV = 1.88 \times 10^{67} \tag{13}$$

I think that this means that the fastest particle would have a probability of 1/N. We of course know that the Maxwell-Boltzmann distribution is

$$f(v) = \left(\frac{m}{2\pi k_B T}\right)^{3/2} 4\pi v^2 \exp\left(-\frac{mv^2}{2k_B T}\right)$$
 (14)

In order to find the maximum velocity we can set the following integral equal to 1/N since we need to find the velocity above which the probability of finding a particle is 1/N (such that we expect to find a single particle).

$$\frac{1}{N} = \int_{v_{\text{max}}}^{\infty} f(v) \, \mathrm{d}v \tag{15}$$

$$= \int_{v_{\text{max}}}^{\infty} dv \left(\frac{m}{2\pi k_B T}\right)^{3/2} 4\pi v^2 \exp\left(-\frac{mv^2}{2k_B T}\right)$$
 (16)

$$= \sqrt{\frac{2}{\pi}} \int_{v_{\text{max}}}^{\infty} dv \left(\frac{m}{2k_B T}\right)^{3/2} v^2 \exp\left(-\frac{mv^2}{2k_B T}\right)$$
 (17)

$$= \sqrt{\frac{2}{\pi}} \int_{v_{\text{max}}}^{\infty} dv \, A^{3/2} v^2 \exp\left(-Av^2\right) \tag{18}$$

$$\approx \sqrt{\frac{2}{\pi}} \cdot \frac{\sqrt{A}}{2} v_{\text{max}} \exp\left(-Av_{\text{max}}^2\right)$$
 (apply hint)

$$\frac{1}{N} \approx \sqrt{\frac{A}{2\pi}} v_{\text{max}} \exp\left(-A v_{\text{max}}^2\right) \tag{19}$$

At this point I wasn't entirely sure how to solve this so I plugged the whole lot into mathematica (using the values for N and T above plus assuming  $m = m_e$ ). After also asserting that  $v_{\text{max}}$  is large, this gives

$$v_{\text{max}} = 6.9 \times 10^7 \,\text{m}\,\text{s}^{-1}$$
 (20)

Given that  $0.99c \approx 3 \times 10^8$ ,  $v_{\text{max}}$  is approximately 23% of the speed of cosmic rays. This implies to me that cosmic rays **cannot** be thermal in origin.

# 3. Atmospheric Escape

#### 3a. Compute number density

We know that the mean free path is defined as

$$\lambda_{\rm esc} \equiv \frac{1}{n_{\rm esc}\sigma},\tag{21}$$

thus we can simply set this equal to the scale height and solve for the number density.

$$\lambda_{\rm esc} = H \tag{22}$$

$$\frac{1}{n_{\rm esc}\sigma} = \frac{k_B T}{mg}$$

$$\boxed{n_{\rm esc} = \frac{mg}{\sigma k_B T}}$$
(23)

$$n_{\rm esc} = \frac{mg}{\sigma k_B T} \tag{24}$$

### 3b. Particle flux expression

Integration time! Basically just applying the double angle rule and doing some integrals.

$$\Phi(v) dv = \frac{f(v)}{4\pi} dv \int_0^{2\pi} d\phi \int_0^{\pi/2} v \cos\theta \sin\theta d\theta$$
 (25)

$$= \frac{vf(v)}{4\pi} dv \int_0^{2\pi} d\phi \frac{1}{2} \int_0^{\pi/2} \sin(2\theta) d\theta$$
 (26)

$$= -\frac{vf(v)}{4\pi} dv \int_0^{2\pi} d\phi - \frac{1}{4} [\cos(2\theta)]_0^{\pi/2}$$
 (27)

$$= -\frac{vf(v)}{4\pi} dv \int_{0}^{2\pi} d\phi - \frac{1}{4} \cdot (-1 - 1)$$
 (28)

$$= \frac{vf(v)}{8\pi} \,\mathrm{d}v \int_0^{2\pi} \mathrm{d}\phi \tag{29}$$

$$\Phi(v) dv = \frac{vf(v)}{4} dv$$
(30)

#### 3c. Total particle flux

Right, let's see if I can remember how to do integrals! Let's start by performing a substitution.

$$\Phi = \int_{V_{\text{cor}}}^{\infty} \Phi(v) \, \mathrm{d}v \tag{31}$$

$$= \frac{1}{4} \int_{v}^{\infty} v f(v) \, \mathrm{d}v \tag{32}$$

$$= \sqrt{\frac{1}{8\pi}} \int_{\text{vesc}}^{\infty} \left(\frac{m}{2k_B T}\right)^{3/2} v^3 \exp\left(-\frac{mv^2}{2k_B T}\right) dv \tag{33}$$

$$= \sqrt{\frac{1}{8\pi}} \int_{V_{\text{opt}}}^{\infty} A^{3/2} v^3 \exp(-Av^2) \, dv$$
 (34)

$$u = v^2, \qquad \frac{\mathrm{d}u}{\mathrm{d}v} = 2v \tag{35}$$

$$\Phi = \sqrt{\frac{1}{8\pi}} \int_{v_{\text{esc}}}^{\infty} A^{3/2} v^3 \exp\left(-Av^2\right) \frac{\mathrm{d}u}{2v}$$
(36)

$$= \sqrt{\frac{1}{8\pi}} \int_{v_{\text{enc}}^2}^{\infty} \frac{1}{2} A^{3/2} u \exp(-Au) du$$
 (37)

$$= \sqrt{\frac{A^3}{32\pi}} \int_{v_{\text{esc}}^2}^{\infty} u \exp(-Au) \, du \tag{38}$$

Okay let's pause for a second. That covers the substitution part, now we need to do integration by parts. Since I'm already using u let's call the two variables in integration by parts a and b such that

$$\int a \, \mathrm{d}b = ab - \int b \, \mathrm{d}a \tag{39}$$

In our case, we have that

$$a = u, \qquad \mathrm{d}b = \exp(-Au)$$
 (40)

and thus we can also write that

$$da = 1, b = -\frac{1}{A}\exp(-Au) (41)$$

So now putting that all together, we can write that the integral in full is

$$= \left[ u \cdot -\frac{1}{A} \exp(-Au) \right]_{\text{v}_{\text{esc}}}^{\infty} - \int_{\text{v}_{\text{esc}}}^{\infty} -\frac{1}{A} \exp(-Au) \, du$$
 (42)

$$= \left[ -\frac{u}{A} \exp(-Au) \right]_{\mathbf{v}_{\text{asc}}^2}^{\infty} + \frac{1}{A} \int_{\mathbf{v}^2}^{\infty} \exp(-Au) \, \mathrm{d}u$$
 (43)

$$= \left(0 + \frac{v_{\rm esc}^2}{A} \exp(-Av_{\rm esc}^2)\right) - \left[\frac{\exp(-Au)}{A^2}\right]_{v_{\rm esc}^2}^{\infty} \tag{44}$$

$$= \frac{v_{\rm esc}^2}{A} \exp\left(-Av_{\rm esc}^2\right) + \frac{\exp(-Av_{\rm esc}^2)}{A^2} \tag{45}$$

Okay phew, we made it. Now let's go and plug that back in!

$$\Phi = \sqrt{\frac{A^3}{32\pi}} \left( \frac{v_{\rm esc}^2}{A} \exp\left(-Av_{\rm esc}^2\right) + \frac{\exp(-Av_{\rm esc}^2)}{A^2} \right) \tag{46}$$

$$= \frac{(1 + Av_{\rm esc}^2)e^{-Av_{\rm esc}^2}}{\sqrt{32\pi A}} \tag{47}$$

Now let's put this into the format that you asked for. My constant A is related to  $v_s$  as  $A = v_s^{-2}$  and thus we can write that  $Av_{\rm esc}^2 = \lambda_{\rm esc}$ . Additionally, we want this in terms of the total number of particles and so we multiply by the number density.

$$\Phi = \frac{n_{\rm esc}v_s(1+\lambda_{\rm esc})e^{-\lambda_{\rm esc}}}{\sqrt{32\pi}}$$
(48)

Let's just quickly write down the constants that we've been given:

$$m \approx 2 \text{ amu} = 3.32 \times 10^{-27} \text{ kg}, \quad T = 1000 K, \quad \sigma = \pi \text{Å}^2 = \pi \times 10^{-20} \text{ m}^2$$
 (49)

where the mass has come from them being hydrogen molecules. Now let's evaluate  $\lambda_{\rm esc}$ ,  $n_{\rm esc}$ ,  $v_s$  so that we can find  $\phi$ .

$$v_{\rm esc} = 11200 \,\mathrm{m \, s^{-1}}$$
 (50)

$$n_{\rm esc} = \frac{m_{\rm H2} \cdot g}{\sigma k_B T} = \frac{3.35 \times 10^{-27} \,\mathrm{kg} \cdot 10 \,\mathrm{m} \,\mathrm{s}^{-2}}{\pi \times 10^{-20} \,\mathrm{m}^2 \cdot 1.38 \times 10^{-23} \,\mathrm{m}^2 \,\mathrm{kg} \,\mathrm{s}^{-2} \,\mathrm{K}^{-1} \cdot 1000 \,\mathrm{K}} = 7.66 \times 10^{13} \,\mathrm{m}^{-3} \qquad (51)$$

$$v_s = \sqrt{\frac{2k_BT}{m_{\rm H2}}} = \sqrt{\frac{2 \cdot 1.38 \times 10^{-23} \,\mathrm{m}^2 \,\mathrm{kg} \,\mathrm{s}^{-2} \,\mathrm{K}^{-1} \cdot 1000 \,\mathrm{K}}{3.35 \times 10^{-27} \,\mathrm{kg}}} = 2883 \,\mathrm{m} \,\mathrm{s}^{-1}$$
 (52)

$$\lambda_{\rm esc} = \left(\frac{11200}{2883}\right)^2 = 15.1\tag{53}$$

Cool, now it's time to just plug those numbers in to find about  $\phi$ .

$$\phi_{\rm H2} = 9.91 \times 10^{10} \,\mathrm{m}^{-2} \mathrm{s}^{-1} \tag{54}$$

Finally, we need to multiply this by the area through which the particles move. Since we are given that  $H \ll R$ , we can say the surface at this height is essentially the same as the surface area of the Earth.

$$\mathcal{R}_{H2} = \phi \cdot 4\pi R_{\oplus}^2 \tag{55}$$

$$= 9.91 \times 10^{10} \,\mathrm{m}^{-2} \mathrm{s}^{-1} \cdot 4\pi (6.4 \times 10^6 \,\mathrm{m})^2 \tag{56}$$

$$\mathcal{R}_{H2} = 5.1 \times 10^{25} \,\text{molecules s}^{-1} \tag{57}$$

Therefore, if we were to evolve forwards for 1 Gyr, the total number of molecules that would be lost are

$$N_{\rm H2,lost} = 1.6 \times 10^{42}$$
 (58)

and so this is comparable to the current hydrogen content of the Earth's atmosphere, approximately a factor of 6 lower.

# 3d. Earth Oxygen Loss

This calculation will be exactly the same except we now change the mass from  $m \approx 2$  amu to  $m \approx 32$  amu. This drastically changes things due to the exponential terms and so the rate of oxygen loss is an extremely low

$$\mathcal{R}_{O2} = 1.6 \times 10^{-71} \,\text{molecules s}^{-1}$$
 (59)

Meaning that over a billion years we wouldn't expect the loss of *any* oxygen. I therefore conclude that over time the abundance of atmospheric oxygen relative to hydrogen will increase (assuming no other sources or sinks).