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An Infinite Family of S Invariant Theories on the Celestial Sphere

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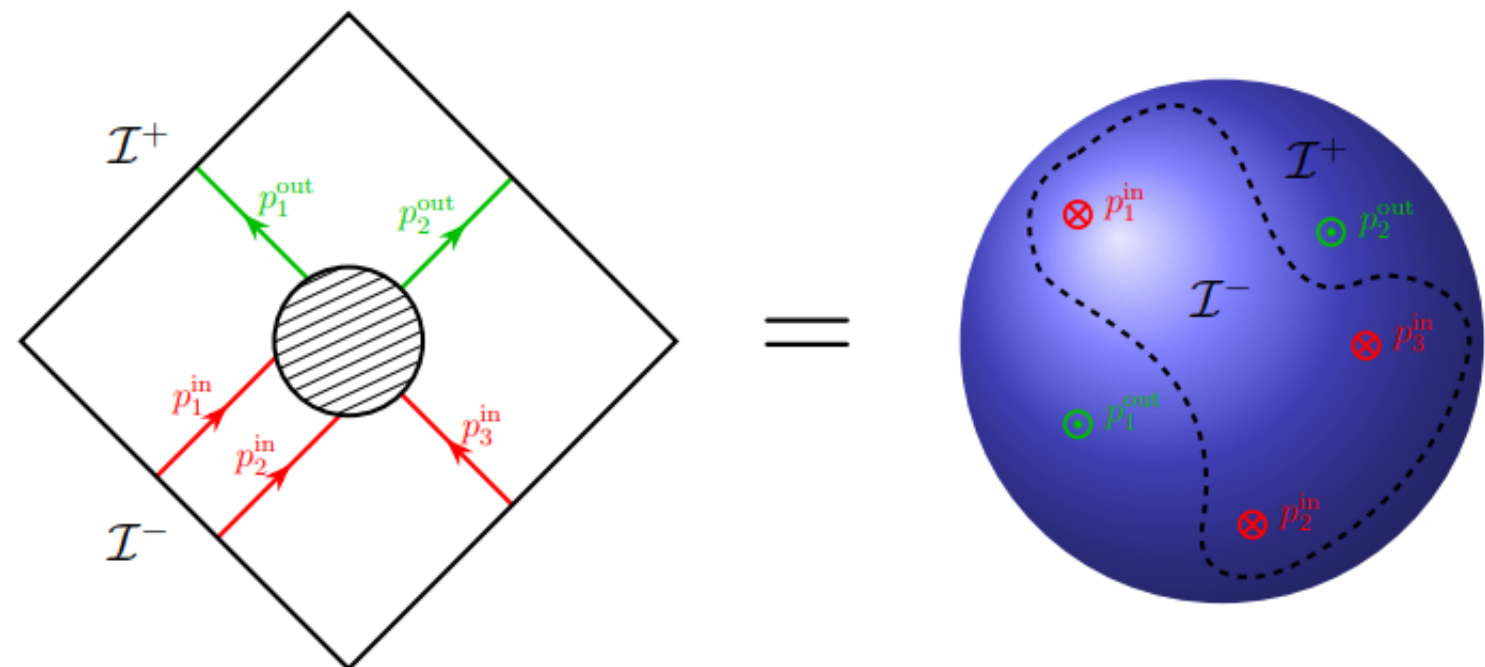
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Celestial CFT : A Putative Dual for QG in AFS

A quantum theory of gravity in (3+1)-D AFS \iff A 2D CCFT on the Celestial sphere at null infinity.



$$\langle out | S | in \rangle = \langle \mathcal{O}_{\Delta_1}^{\pm}(z_1, \bar{z}_1) \dots \mathcal{O}_{\Delta_n}^{\pm}(z_n, \bar{z}_n) \rangle \quad (1)$$

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Holographic Symmetry Algebras for Gravity and Gauge Theories in AFS

- Conformally soft theorems \longrightarrow Infinite dimensional non-trivial symmetries in Gauge theories and Gravity in 4D AFS.
- Gravity \longrightarrow wedge subalgebra of $w_{1+\infty}$ algebra. (Guevara, Himwich, Pate and Strominger '21)
- Gauge Theories \longrightarrow S algebra

Goal and the Motivation

- Goal** \Rightarrow Our goal is to classify the theories which are invariant under **S algebra** and also to find the KZ-type null states of these theories. 2311.16796 (Banerjee, RM, Misra, Panda and Paul '23)
- Banerjee, Kulkarni and Paul '23 \Rightarrow computed G^+G^+ OPE and KZ-type null states of any $w_{1+\infty}$ invariant theories 2301.13225; 2311.06485 \Rightarrow existence of a discrete infinite family of $w_{1+\infty}$ invariant theories on the celestial sphere.

Gluons and the S Algebra

The **S algebra** is obtained from the singular part of the OPE i.e.

Guevara, Himwich, Pate and Strominger '21

$$\mathcal{O}_{\Delta_1}^{a,+}(z_1, \bar{z}_1) \mathcal{O}_{\Delta_2}^{b,+}(z_2, \bar{z}_2) \sim -\frac{i f^{ab}}{z_{12}} \sum_{n=0}^{\infty} B(\Delta_1 + n - 1, \Delta_2 - 1) \frac{\bar{z}_{12}^n}{n!} \bar{\partial}_2^n \mathcal{O}_{\Delta_1 + \Delta_2 - 1}^{c,+}(z_2, \bar{z}_2). \quad (2)$$

Soft gluons :

$$R^{k,a}(z, \bar{z}) = \lim_{\Delta \rightarrow k} (\Delta - k) \mathcal{O}_{\Delta}^{a,+}(z, \bar{z}), \quad k = 1, 0, -1, \dots \quad (3)$$

Holomorphic soft gluon currents :

$$R^{k,a}(z, \bar{z}) = \sum_{n=\frac{k-1}{2}}^{\frac{1-k}{2}} \frac{R_n^{k,a}(z)}{\bar{z}^{n+\frac{k-1}{2}}} \quad (4)$$

Modes of the Holomorphic currents :

$$R_n^{k,a}(z) = \sum_{\alpha \in \mathbb{Z} - \frac{k+1}{2}} \frac{R_{\alpha,n}^{k,a}}{z^{\alpha+\frac{k+1}{2}}} \quad (5)$$

Algebra :

$$[R_{\alpha,m}^{k,a}, R_{\beta,n}^{l,b}] = -i f^{ab} \frac{(\frac{1-k}{2} - m + \frac{1-l}{2} - n)! (\frac{1-k}{2} + m + \frac{1-l}{2} + n)!}{(\frac{1-k}{2} - m)! (\frac{1-l}{2} - n)! (\frac{1-k}{2} + m)! (\frac{1-l}{2} + n)!} R_{\alpha+\beta, m+n}^{k+l-1, c} \quad (6)$$

Redefinition :

$$S_{\alpha,m}^{q,a} = (q - m - 1)! (q + m - 1)! R_{\alpha,m}^{3-2q,a} \quad (7)$$

S Algebra :

$$[S_{\alpha,m}^{p,a}, S_{\beta,n}^{q,b}] = -i f^{abc} S_{\alpha+\beta, m+n}^{p+q-1, c} \quad (8)$$

General Structure of the Gluon-Gluon OPE

$$\begin{aligned} & \mathcal{O}_{\Delta_1}^{a,+}(z_1, \bar{z}_1) \mathcal{O}_{\Delta_2}^{b,+}(z_2, \bar{z}_2) \\ &= -\frac{i f^{ab}}{z_{12}} \sum_{n=0}^{\infty} B(\Delta_1 + n - 1, \Delta_2 - 1) \frac{\bar{z}_{12}^n}{n!} \bar{\partial}_2^n \mathcal{O}_{\Delta_1 + \Delta_2 - 1}^{c,+}(z_2, \bar{z}_2) + \sum_{p,q=0}^{\infty} \sum_{k=1}^{\tilde{n}_{p,q}} z_{12}^p \bar{z}_{12}^q \mathcal{C}_{p,q}^k(\Delta_1, \Delta_2) \tilde{\mathcal{O}}_{k,p,q}^{ab}(z_2, \bar{z}_2). \end{aligned} \quad (9)$$

Task is to determine:

- The OPE coefficients $\mathcal{C}_{p,q}^k$ and
- the S-algebra descendants $\tilde{\mathcal{O}}_{k,p,q}^{ab}$ of a positive helicity soft gluon.

Strategy

- We consider S invariant theories for all of which S-algebra is universal \Rightarrow Existence of a Master OPE .
- We know that tree-level MHV sector of the pure YM theory is an example of such S invariant theories.
- This Master OPE inserted in a MHV gluon scattering amplitude \Rightarrow known MHV OPE.
- Master OPE = MHV-sector OPE + R**
- R should vanish inside MHV scattering amplitude \Rightarrow R is a lin. combination of MHV null states.
- R consists only non-singular terms.

OPEs in Terms of MHV Null States

Using the above arguments we can rewrite (9) as,

$$\begin{aligned} \mathcal{O}_{\Delta_1}^{a,+}(z_1, \bar{z}_1) \mathcal{O}_{\Delta_2}^{b,+}(z_2, \bar{z}_2) |_{\text{Any Theory}} &= \mathcal{O}_{\Delta_1}^{a,+}(z_1, \bar{z}_1) \mathcal{O}_{\Delta_2}^{b,+}(z_2, \bar{z}_2) |_{\text{MHV}} \\ &+ \sum_{p,q=0}^{\infty} z_{12}^p \bar{z}_{12}^q \sum_{i=1}^{\tilde{n}_{p,q}} \tilde{\mathcal{C}}_{p,q}^k(\Delta_1, \Delta_2) M_{k,p,q}^{a,b}(\Delta_1, \Delta_2, z_2, \bar{z}_2). \end{aligned} \quad (10)$$

MHV Null States at $\mathcal{O}(1)$

The general null state at $\mathcal{O}(1)$ in the MHV-sector is given by

$$\Psi_j^{ab}(\Delta) = R_{\frac{j-1}{2}, \frac{j+1}{2}}^{-j,a} \mathcal{O}_{\Delta+j,+}^b - \frac{(-1)^j j}{\Gamma(j+2)} \frac{\Gamma(\Delta+j-1)}{\Gamma(\Delta-2)} R_{-1,0}^{1,a} \mathcal{O}_{\Delta-1,+}^b - \frac{(-1)^j}{\Gamma(j+1)} \frac{\Gamma(\Delta+j-1)}{\Gamma(\Delta-1)} R_{-1/2, 1/2}^{0,a} \mathcal{O}_{\Delta,+}^b \quad (11)$$

where $j = 1, 2, 3, \dots$

Let's consider the following basis :

$$M_k^{ab}(\Delta) = \sum_{i=1}^k \frac{1}{\Gamma(k-i+1)} \frac{\Gamma(\Delta+k-1)}{\Gamma(\Delta+i-1)} \Psi_i^{ab}(\Delta). \quad (12)$$

Action of the S algebra on the MHV Null States

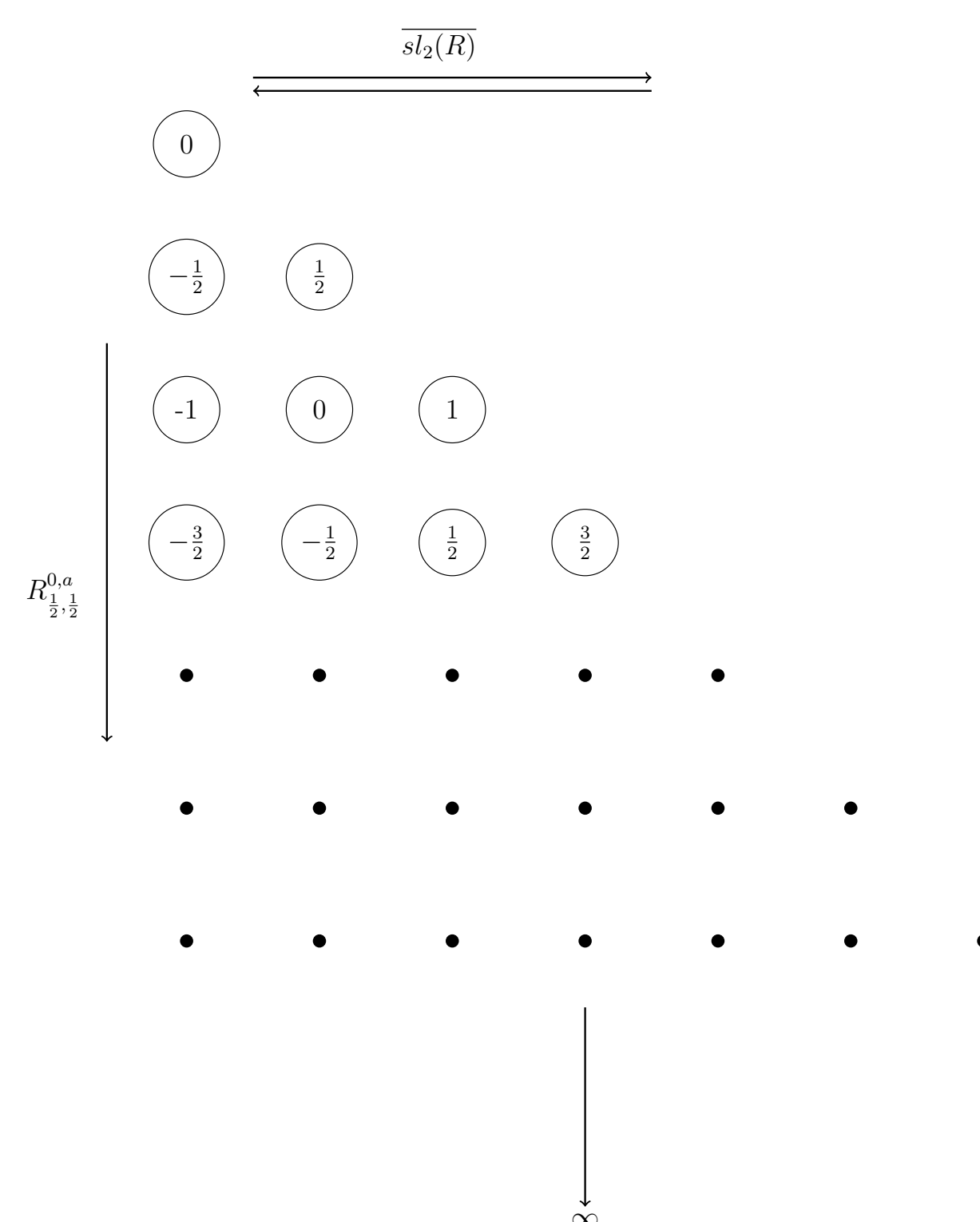
Action of the Leading Soft Gluon modes:

$$R_{0,0}^{1,a} M_k^{bc}(\Delta) = -i f^{abd} M_k^{dc}(\Delta) - i f^{acd} M_k^{bd}(\Delta) \quad (13)$$

$$R_{n,0}^{1,a} M_k^{bc}(\Delta) = 0, n > 0 \quad (14)$$

Action of the Subleading Soft Gluon mode :

$$[R_{1/2, 1/2}^{0,a}, M_k^{bc}(\Delta)] = -i f^{abd} (k+2) M_{k+1}^{dc}(\Delta-1) + (\Delta+k-2) \left\{ i f^{acd} M_k^{bd}(\Delta-1) + i f^{abd} M_k^{dc}(\Delta-1) \right\}. \quad (15)$$



Building up S invariant OPEs

Observation :

\Rightarrow Consider the following set of null states,

$$M_k^{bc}(\Delta), k = 1, 2, 3, \dots, n. \quad (16)$$

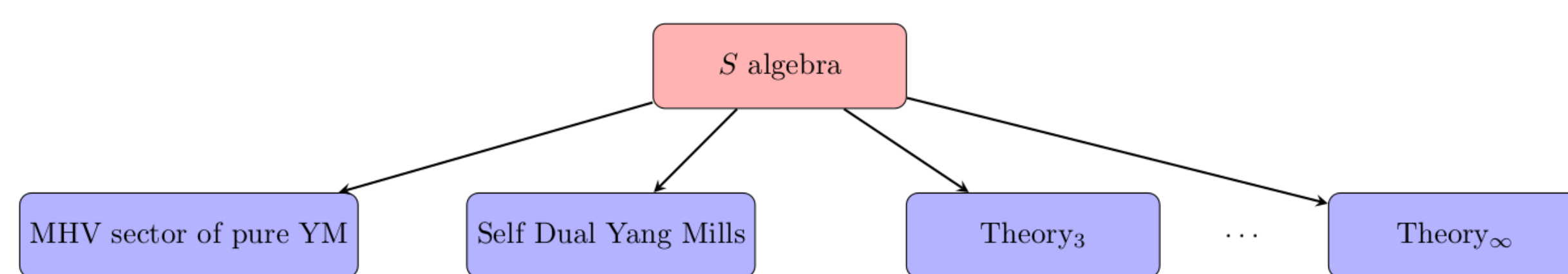
\Rightarrow Action of $R_{1/2, 1/2}^{0,a}$ on the null states is closed if we set

$$M_{k+1}^{ab}(\Delta) = 0, k \geq n \geq 0. \quad (17)$$

Inference :

\Rightarrow We can get an S invariant OPE if we consider the finite set of null states (16).

$$\mathcal{O}_{\Delta_1,+}^a(z, \bar{z}) \mathcal{O}_{\Delta_2,+}^b(0, 0) |_{\mathcal{O}(1)} = \mathcal{O}_{\Delta_1,+}^a(z, \bar{z}) \mathcal{O}_{\Delta_2,+}^b(0, 0) \Big|_{\mathcal{O}(1)}^{MHV} + \sum_{k=1}^n B(\Delta_1 + k, \Delta_2 - 1) M_k^{ab}(\Delta_1 + \Delta_2) \quad (18)$$



Knizhnik-Zamolodchikov Type Null States

$$K^a(\Delta) = \xi^a(\Delta) - i \sum_{k=1}^n M_k^a(\Delta+1), \quad (19)$$

where

$$\xi^a(\Delta) = C_A L_{-1} \mathcal{O}_{\Delta}^{a,+} - (\Delta+1) R_{-1,0}^{1,b} R_{0,0}^{1,b} \mathcal{O}_{\Delta}^{a,+} - R_{-1/2, 1/2}^{0,b} R_{0,0}^{1,b} \mathcal{O}_{\Delta+1}^{a,+} \quad (20)$$

and

$$M_k^a(\Delta) = f^{abc} M_k^{bc}. \quad (21)$$

References

S. Banerjee, R. Mandal, S. Misra, S. Panda and P. Paul, "All OPEs invariant under the infinite symmetry algebra for gluons on the celestial sphere," [PhysRevD.110.026020]

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