



2023

# MHV Gluon Scattering in the Massive Scalar Background and Celestial OPE



Shamik Banerjee<sup>1,2</sup>

<sup>1</sup>HBNI, Mumbai

Raju Mandal<sup>1,2</sup>

<sup>2</sup>NISER, Bhubaneswar

Akavoor Manu<sup>1,3</sup>

<sup>3</sup>IOPB, Bhubaneswar

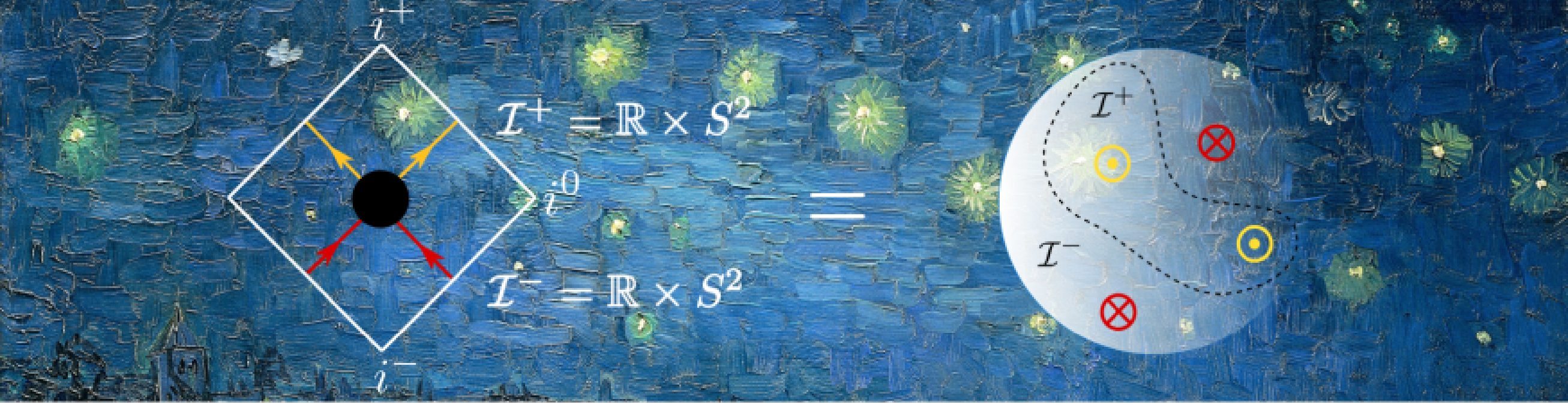
Partha Paul<sup>4</sup>

<sup>4</sup>IISC, Bangalore

## Introduction to Celestial Holography

### Conjecture :

Quantum theories of Gravity in (3+1)-D AFS  $\iff$  2D CCFTs on the Celestial sphere at null infinity.



[Picture Courtesy: Andrea Puhm]

## Asymptotic Symmetries of Gravity and Gauge Theories in AFS

### Gravity

✍ '60s : Poincaré  $\implies$  BMS = Lorentz  $\times$  Supertranslations (Bondi,Burg,Metzner and Sachs)

✍ '10s : BMS  $\longrightarrow$  Virasoro  $\times$  Supertranslations Extended BMS  
(using asymptotic symmetry analysis and soft theorems.)

Lorentz group  $\longrightarrow$  Virasoro

Kapec,Mitra,Strominger,...

### Another extension :

Lorentz group  $\longrightarrow$   $SL(2, \mathbb{C})$  current algebra Banerjee,Ghosh,Paul '20

Guevara-Himwich-Pate-Strominger '21

wedge subalgebra of  $w_{1+\infty}$  algebra

### Gauge Theories

✍ Asymptotic symmetry algebra is HSA  $\times$  Poincaré.

Guevara-Himwich-Pate-Strominger '21

## Mellin Transform and the Celestial Amplitude

Scattering amplitudes in boost eigenbasis.

$$\tilde{\mathcal{A}}_n(\{z_i, \bar{z}_i, h_i, \bar{h}_i, a_i\}) = \langle \prod_{i=1}^n \mathcal{O}_{h_i, \bar{h}_i}^{a_i}(z_i, \bar{z}_i) \rangle = \prod_{i=1}^n \int_0^\infty d\omega_i \omega_i^{\Delta_i-1} \mathcal{S}_n(\{\omega_i, z_i, \bar{z}_i, \sigma_i, a_i\}) \quad (1)$$

Example : [S. Banerjee and S. Ghosh ; arXiv:2011.00017]

$$\tilde{\mathcal{A}}(1^{-a_1}, 2^{+a_2}, 4^{-x}) \sim f^{a_1 a_2 x} \frac{z_{14}^3}{z_{12}^2 z_{24}} \underbrace{\delta(\bar{z}_{14}) \delta(\bar{z}_{24})}_{\text{delta distributions ??}} \prod_{i=1}^3 \Theta(\epsilon_i \sigma_{i,1}) \quad (2)$$

Under  $SL(2, \mathbb{C})/\mathbb{Z}_2$  or Lorentz transformations,

$$\tilde{\mathcal{A}}_n(\{z_i, \bar{z}_i, h_i, \bar{h}_i, a_i\}) = \prod_{i=1}^n \frac{1}{(cz_i + d)^{2h_i}} \frac{1}{(\bar{c}\bar{z}_i + \bar{d})^{2\bar{h}_i}} \tilde{\mathcal{A}}_n\left(\left\{\frac{az_i + b}{cz_i + d}, \frac{\bar{a}\bar{z}_i + \bar{b}}{\bar{c}\bar{z}_i + \bar{d}}, h_i, \bar{h}_i, a_i\right\}\right). \quad (3)$$

## Momentum Space Soft Theorems

### Gauge theories at tree level

$$\mathcal{M}_{n+1}^a(p_1, \dots, p_n, q) \xrightarrow{\omega \rightarrow 0} (S_{(0)}^a + S_{(1)}^a) \mathcal{M}_n(p_1, \dots, p_n) + \mathcal{O}(q) \quad (4)$$

$\mathcal{O}(\frac{1}{\omega}) \quad \mathcal{O}(1)$

no further factorisation  
( $q \sim \omega$ ).

where,

$$S_{(0)}^a \sim \sum_{k=1}^n \frac{\epsilon_{\mu} p_k^{\mu}}{q \cdot p_k} T_k^a \quad S_{(1)}^a = \sum_{k=1}^n i \frac{\epsilon_{\mu} q_{\nu} J_k^{\mu\nu}}{q \cdot p_k} T_k^a \quad (\text{universal}) \quad (5)$$

and  $J_k^{\mu\nu}$  is the total angular momentum of the k'th particle.

## Conformal Soft Theorems in Mellin Basis

CONFORMAL SOFT THEOREMS  $\equiv$  WARD IDENTITIES FOR ASYMPTOTIC SYMMETRIES

Leading Conformal Soft Theorem  $\implies$  level-zero Kac-Moody Algebra

$$\left\langle R^{1,a}(z) \prod_{i=1}^n \mathcal{O}_{h_i, \bar{h}_i}^{a_i}(z_i, \bar{z}_i) \right\rangle = - \sum_{k=1}^n \frac{T_k^a}{z - z_k} \left\langle \prod_{i=1}^n \mathcal{O}_{h_i, \bar{h}_i}^{a_i}(z_i, \bar{z}_i) \right\rangle \quad (6)$$

where leading conformal soft gluon operator is defined as,

$$R^{1,a}(z) = \lim_{\Delta \rightarrow 1} (\Delta - 1) \mathcal{O}_{\Delta,+}^a(z, \bar{z}) \quad (7)$$

Sub-leading Conformal Soft Theorem  $\implies$  Current Algebra

$$\left\langle R^{0,a}(z, \bar{z}) \prod_{i=1}^n \mathcal{O}_{h_i, \bar{h}_i}^{a_i}(z_i, \bar{z}_i) \right\rangle = - \sum_{i=1}^n \frac{\epsilon_k}{z - z_k} (-2\bar{h}_k + 1 + (\bar{z} - \bar{z}_k) \bar{\partial}_k) T_k^a P_k^{-1} \left\langle \prod_{i=1}^n \mathcal{O}_{h_i, \bar{h}_i}^{a_i}(z_i, \bar{z}_i) \right\rangle \quad (8)$$

where sub-leading conformal soft gluon operator is defined as,

$$R^{0,a}(z, \bar{z}) = \lim_{\Delta \rightarrow 0} \Delta \mathcal{O}_{\Delta,+}^a(z, \bar{z}) \quad (9)$$

Where,

$$R^{0,a}(z, \bar{z}) = \sum_{n=-\frac{1}{2}}^{\frac{1}{2}} \frac{R_n^{0,a}(z)}{\bar{z}^{n-\frac{1}{2}}} \quad (10)$$

## YM Theory Coupled with Massive Complex Scalar Background

- ✍ Tree level celestial amplitudes in trivial background have kinematic singularities due to scale and translation invariances.
- ✍ The YM theory chirally coupled with massive complex scalar background breaks translation as well as scale invariance.
- ✍ This theory indeed removes those singularities by replacing the delta distributions with some smooth functions.
- ✍ Leading conformal soft theorem and leading OPE structure were shown to remain unmodified. [Casali,Melton,Strominger'22]
- ✍ We have shown that the sub-leading soft gluon theorem and the OPE structure of two outgoing positive helicity gluons at  $\mathcal{O}(1)$  also remain same.
- ✍ As a consequence of this, the scattering amplitudes of this theory also satisfy the same BG equations.
- ✍ Calculation done in (2, 2) signature. Celestial torus

## Tree Level 5-point Amplitude

In momentum basis

$$\mathcal{A}(\phi, 1^{-a_1}, 2^{-a_2}, 3^{+a_3}, 4^{+a_4}, 5^{+a_5}) = (-i) \times \left\{ f^{a_1 a_2 x_1} f^{x_1 a_3 x_2} f^{x_2 a_4 a_5} A(\phi, 1^-, 2^-, 3^+, 4^+, 5^+) + f^{a_1 a_2 x_1} f^{x_1 a_4 x_2} f^{x_2 a_3 a_5} A(\phi, 1^-, 2^-, 4^+, 3^+, 5^+) \right. \\ \left. + f^{a_1 a_4 x_1} f^{x_1 a_3 x_2} f^{x_2 a_2 a_5} A(\phi, 1^-, 4^+, 3^+, 2^-, 5^+) + f^{a_1 a_3 x_1} f^{x_1 a_2 x_2} f^{x_2 a_4 a_5} A(\phi, 1^-, 3^+, 2^-, 4^+, 5^+) \right. \\ \left. + f^{a_1 a_4 x_1} f^{x_1 a_2 x_2} f^{x_2 a_3 a_5} A(\phi, 1^-, 4^+, 2^-, 3^+, 5^+) + f^{a_1 a_3 x_1} f^{x_1 a_4 x_2} f^{x_2 a_2 a_5} A(\phi, 1^-, 3^+, 4^+, 2^-, 5^+) \right\} \quad (11)$$

In boost eigenbasis

$$\tilde{\mathcal{A}}^\phi(1^{-a_1}, 2^{-a_2}, 3^{+a_3}, 4^{+a_4}, 5^{+a_5}) = \langle \mathcal{O}_{\Delta_1,-}^{a_1}(z_1, \bar{z}_1) \mathcal{O}_{\Delta_2,-}^{a_2}(z_2, \bar{z}_2) \mathcal{O}_{\Delta_3,+}^{a_3}(z_3, \bar{z}_3) \mathcal{O}_{\Delta_4,+}^{a_4}(z_4, \bar{z}_4) \mathcal{O}_{\Delta_5,+}^{a_5}(z_5, \bar{z}_5) \rangle_\phi \\ = (-i) \left\{ f^{a_1 a_2 x_1} f^{x_1 a_3 x_2} f^{x_2 a_4 a_5} \frac{z_{12}^4}{z_{12} z_{23} z_{34} z_{45} z_{51}} + f^{a_1 a_2 x_1} f^{x_1 a_4 x_2} f^{x_2 a_3 a_5} \frac{z_{12}^4}{z_{12} z_{24} z_{43} z_{35} z_{51}} + f^{a_1 a_4 x_1} f^{x_1 a_3 x_2} f^{x_2 a_2 a_5} \frac{z_{12}^4}{z_{14} z_{43} z_{32} z_{25} z_{51}} \right. \\ \left. + f^{a_1 a_3 x_1} f^{x_1 a_2 x_2} f^{x_2 a_4 a_5} \frac{z_{12}^4}{z_{13} z_{32} z_{24} z_{45} z_{51}} + f^{a_1 a_4 x_1} f^{x_1 a_2 x_2} f^{x_2 a_3 a_5} \frac{z_{12}^4}{z_{14} z_{42} z_{23} z_{35} z_{51}} + f^{a_1 a_3 x_1} f^{x_1 a_4 x_2} f^{x_2 a_2 a_5} \frac{z_{12}^4}{z_{13} z_{34} z_{42} z_{25} z_{51}} \right\} \times \\ \frac{N_5}{2(2\pi)^4} \times \Gamma(\Delta_1 + 1) \Gamma(\Delta_2 + 1) \Gamma(\Delta_3 - 1) \Gamma(\Delta_4 - 1) \Gamma(\Delta_5 - 1) \times f(\beta_5) \times \\ \int d^3 \hat{x} (-q_1 \cdot \hat{x})^{-\Delta_1-1} (-q_2 \cdot \hat{x})^{-\Delta_2-1} (-q_3 \cdot \hat{x})^{-\Delta_3+1} (-q_4 \cdot \hat{x})^{-\Delta_4+1} (-q_5 \cdot \hat{x})^{-\Delta_5+1} \quad (12)$$

## Sub-leading Conformal Soft Theorem and the Structure of OPE

Check for sub-leading conformal soft theorem

$$\lim_{\Delta_5 \rightarrow 0} \Delta_5 \langle \mathcal{O}_{\Delta_1,-}^{a_1}(z_1, \bar{z}_1) \mathcal{O}_{\Delta_2,-}^{a_2}(z_2, \bar{z}_2) \mathcal{O}_{\Delta_3,+}^{a_3}(z_3, \bar{z}_3) \mathcal{O}_{\Delta_4,+}^{a_4}(z_4, \bar{z}_4) \mathcal{O}_{\Delta_5,+}^{a_5}(z_5, \bar{z}_5) \rangle_\phi \\ = - \sum_{k=1}^4 \frac{\epsilon_k}{z_5 - \bar{z}_k} \left( -2\bar{h}_k + 1 + (\bar{z}_5 - \bar{z}_k) \bar{\partial}_k \right) T_k^{a_5} P_k^{-1} \langle \mathcal{O}_{\Delta_1,-}^{a_1}(z_1, \bar{z}_1) \mathcal{O}_{\Delta_2,-}^{a_2}(z_2, \bar{z}_2) \mathcal{O}_{\Delta_3,+}^{a_3}(z_3, \bar{z}_3) \mathcal{O}_{\Delta_4,+}^{a_4}(z_4, \bar{z}_4) \rangle_\phi \quad (13)$$

Extraction of OPE at  $\mathcal{O}(1)$  from collinear limit for two positive-helicity outgoing gluons

Collinear limit  $z_4 \rightarrow z_5, \bar{z}_4 \rightarrow \bar{z}_5$

OPE factorization at sub-leading order,  $\mathcal{O}(1)$

$$\tilde{\mathcal{A}}^\phi(1_{\Delta_1}^{-a_1}, 2_{\Delta_2}^{-a_2}, 3_{\Delta_3}^{+a_3}, 4_{\Delta_4}^{+a_4}, 5_{\Delta_5}^{+a_5}) \Big|_{\mathcal{O}(1)} = \frac{1}{2} \times B(\Delta_4 - 1, \Delta_5 - 1) \left[ - \frac{(\Delta_4 - 1)}{(\Delta_4 + \Delta_5 - 2)} (i f^{x a_4 a_5} \mathcal{L}_{-1}(5)) \tilde{\mathcal{A}}^\phi(1_{\Delta_1}^{-a_1}, 2_{\Delta_2}^{-a_2}, 3_{\Delta_3}^{+a_3}, 5_{\Delta_4+\Delta_5-1}^{+a_4+a_5}) \right. \\ \left. + \frac{(\Delta_5 - 1)}{(\Delta_4 + \Delta_5 - 2)} \mathcal{J}_{-1}^{a_4}(5) \tilde{\mathcal{A}}^\phi(1_{\Delta_1}^{-a_1}, 2_{\Delta_2}^{-a_2}, 3_{\Delta_3}^{+a_3}, 5_{\Delta_4+\Delta_5-1}^{+a_4+a_5}) \right. \\ \left. + \frac{(\Delta_4 - 1)}{(\Delta_4 + \Delta_5 - 2)} \mathcal{J}_{-1}^{a_5}(5) \tilde{\mathcal{A}}^\phi(1_{\Delta_1}^{-a_1}, 2_{\Delta_2}^{-a_2}, 3_{\Delta_3}^{+a_3}, 5_{\Delta_4+\Delta_5-1}^{+a_4+a_5}) \right] \quad (14)$$

$\mathcal{O}(1)$  term in the OPE of two gluon primaries

$$\mathcal{O}_{\Delta_4,+}^{a_4+1}(z_4, \bar{z}_4) \mathcal{O}_{\Delta_5,+}^{a_5+1}(z_5, \bar{z}_5) \Big|_{\mathcal{O}(1)} \sim \frac{1}{2} \times B(\Delta_4 - 1, \Delta_5 - 1) \left[ - \frac{(\Delta_4 - 1)}{(\Delta_4 + \Delta_5 - 2)} i f^{x a_4 a_5} L_{-1} + \right. \\ \left. \left( \frac{(\Delta_5 - 1)}{(\Delta_4 + \Delta_5 - 2)} \delta^{a_4 y} \delta^{a_5 x} + \frac{(\Delta_4 - 1)}{(\Delta_4 + \Delta_5 - 2)} \delta^{a_5 y} \delta^{a_4 x} \right) R_{-1,0}^{1,y} \right] \mathcal{O}_{\Delta_4+\Delta_5-1}^{x, \epsilon_5}(z_5, \bar{z}_5). \quad (15)$$

## BG Equations and 3-point Amplitude

The general form of the color-ordered  $SL(2, \mathbb{C})$  covariant 3-point amplitude is,

$$\tilde{\mathcal{M}}_3(1_{\Delta_1}^-, 2_{\Delta_2}^-, 3_{\Delta_3}^+) = C(\Delta_1, \Delta_2, \Delta_3) z_{12}^{h_3-h_1-h_2} z_{13}^{h_2-h_1-h_3} z_{23}^{h_1-h_2-h_3} z_{12}^{h_3-h_1-h_2} z_{13}^{h_2-h_1-h_3} z_{23}^{h_1-h_2-h_3} \quad (16)$$

BG equations

$$\left( \partial_3 - \frac{\Delta_3}{z_{13}} - \frac{1}{z_{23}} \right) \tilde{\mathcal{M}}_3(1_{\Delta_1}^-, 2_{\Delta_2}^-, 3_{\Delta_3}^+) + \epsilon_1 \epsilon_3 \frac{\Delta_1 - \sigma_1 - 1 + \bar{z}_{13} \bar{\partial}_1}{z_{13}} \tilde{\mathcal{M}}_3(1_{\Delta_1-1}^-, 2_{\Delta_2}^-, 3_{\Delta_3+1}^+) = 0 \\ \left( \partial_3 - \frac{\Delta_3}{z_{23}} - \frac{1}{z_{13}} \right) \tilde{\mathcal{M}}_3(1_{\Delta_1}^-, 2_{\Delta_2}^-, 3_{\Delta_3}^+) + \epsilon_2 \epsilon_3 \frac{\Delta_2 - \sigma_2 - 1 + \bar{z}_{23} \bar{\partial}_2}{z_{23}} \tilde{\mathcal{M}}_3(1_{\Delta_1}^-, 2_{\Delta_2-1}^-, 3_{\Delta_3+1}^+) = 0 \quad (17)$$

$\implies$  Recursion relation for the coefficients  $C(\Delta_1, \Delta_2, \Delta_3)$ .

solution :

$$C(\Delta_1, \Delta_2, \Delta_3) = \mathcal{N}_3 \Gamma\left(\frac{\Delta_1 + \Delta_2 - \Delta_3 + 3}{2}\right) \Gamma\left(\frac{\Delta_1 - \Delta_2 + \Delta_3 - 3}{2}\right) \Gamma\left(-\frac{\Delta_1 + \Delta_2 + \Delta_3 - 1}{2}\right) f(\beta) \quad (18)$$

where,  $\beta = \sum_{i=1}^3 \Delta_i$  and  $\mathcal{N}_3 = \prod_{j=1}^3 (-i \epsilon_j)^{\Delta_j - \sigma_j}$ .

## References

1. S. Banerjee, R. Mandal, A. Manu and P. Paul, "MHV Gluon Scattering in the Massive Scalar Background and Celestial OPE," [arXiv:2302.10245 [hep-th]]

Email : raju.mandal@niser.ac.in

