A QUICK CELESTIAL TOUR

Raju Mandal August 30, 2024

SPS Day, 2024



VACUUM EINSTEIN EQUATION



$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 0 \tag{1}$$

MAXIMALLY SYMMETRIC SOLUTIONS:

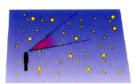
Anti-de Sitter(AdS)

 $\Lambda < 0$



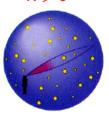
Negatively curved universe

Minkowski $\Lambda = 0$



Flat universe

de Sitter(dS) $\Lambda > 0$

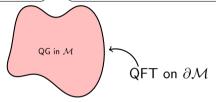


Positively curved



HOLOGRAPHY? INTRODUCTION

Quantum theory of gravity in bulk \equiv QFT on the bdy without dynamical gravity



Examples? $\Lambda < 0$

time anti-de Sitter space conformal boundary

Type IIB superstring theory in $AdS_5 \times S_5$

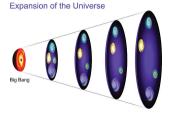
$$\equiv$$
 ($\mathcal{N}=$ 4) SYM in (3+1)D

AdS-CFT duality



Maldacena'97

- 1 But the universe that we live in is not AdS!
- **2** Our universe is "asymptotically de Sitter" ! $\Lambda > 0$



- **3** How general is the holographic principle?
- **4** Can we apply it for flat($\Lambda = 0$) and expanding($\Lambda > 0$) universe also?



6 Current topics of research in string theory



What is this tour about?

 \diamond Recent advancements towards $\Lambda = 0$ holography

Why asymptotically flat spacetime? Motivations

- Study of isolated systems(e.g a BH) in GR, far away from the source but dist. < cosmological scale.
- 2 Collider physics
- **3** Rich asymptotic structure
- Observable effects

S matrix

BMS'60 Lorentz ⋉ supertranslations

Memory effect

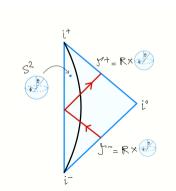
ASYMPTOTICALLY FLAT SPACETIME(AFS)

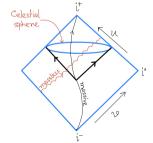
Defn:

$$g_{\mu\nu} = \eta_{\mu\nu} + \mathcal{O}(\frac{1}{r})$$
 as $r \to \infty$ (2)

Minkowski:

$$ds^{2} = -du^{2} - 2dudr + r^{2} \frac{4dzd\bar{z}}{(1+z\bar{z})^{2}}$$
 (3)





CELESTIAL HOLOGRAPHY

Quantum Gravity in (3+1)D AFS \equiv 2D Celestial CFTs on CS^2 at null infinity



image courtesy: Andrea Puhm

Pioneers:



Sabrina Pasterski



Shu-Heng Shao

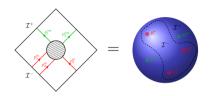


Andrew Strominger → ■ → → ■ → ○ ○ ○





DUALITY



Lorentz transformations

Global conformal transfns.

$$\begin{array}{ccc} \mathbf{x}^{\mu} \longrightarrow \mathbf{\Lambda}^{\mu}{}_{\nu}\mathbf{x}^{\nu} & z \longrightarrow \frac{a\mathbf{z}+b}{c\mathbf{z}+d}, \text{ c.c} \\ |\vec{p},\sigma\rangle = |\omega,\sigma,z,\bar{z}\rangle & \xrightarrow{\int_{0}^{\infty}d\omega\omega^{\Delta-1}...} |\Delta,\sigma,z,\bar{z}\rangle \\ & \text{momentum} & \text{Boost} \end{array}$$

states:

$$\langle out | S | in \rangle \xrightarrow{\int_0^\infty d\omega \omega^{\Delta-1} ...} \langle \mathcal{O}_1(z_1, \bar{z}_1) ... \mathcal{O}_N(z_N, \bar{z}_N) \rangle$$
 (4)

Scattering amplitude

Celestial amplitude conformal invariance manifest

symmetry:

translation invariance manifest

WHAT MAKES US HAPPY? FINDING A CFT MAYBE

- States $|\Delta, \sigma, z, \overline{z}\rangle$ transform as the primary states of a 2D CFT under Lorentz transformations.
- **2** Celestial amplitudes transform nicely under Lorentz transformations/ $SL(2,\mathbb{C})$ transformations.
- We have found a 2D CFT on the Celestial sphere !!
- **4** CFT data? Given $\{\Delta_i, \sigma_i\}$ and C_{ijk} , one can build higher point correlation functions from the OPE(Operator product expansion),

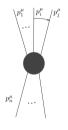
$$\mathcal{O}_{i}(z_{1},\bar{z}_{1})\mathcal{O}_{j}(z_{2},\bar{z}_{2}) = \sum_{k} c_{ij}^{k}(z_{12},\bar{z}_{12})\mathcal{O}_{k}(z_{2},\bar{z}_{2})$$
(5)

Then how do you find the OPE on celestial sphere?



CELESTIAL OPE FROM COLLINEAR FACTORIZATION

Celestial amplitudes get factorized upon taking collinear limit.





$$\mathcal{A}(1^{-a_1}, 2^{+a_2}, 3^{+a_3}, 4^{-a_4}) \xrightarrow[p_3 + p_4 = 0]{z_{34} \to 0} - \frac{f^{a_3 a_4 x}}{z_{34}} B(\Delta_3 - 1, \Delta_4 + 1) \mathcal{A}(1^{-a_1}, 2^{+a_2}, 4^{-x}) + \text{subleading in } z_{34} + \dots$$
(6)

OPE at leading order:

$$\mathcal{O}_{\Delta_{3},+}^{a_{3}}(z_{3},\bar{z}_{3})\mathcal{O}_{\Delta_{4},-}^{a_{3}}(z_{4},\bar{z}_{4}) \sim -\frac{f^{a_{3}a_{4}x}}{z_{34}}B(\Delta_{3}-1,\Delta_{4}+1)\mathcal{O}_{\Delta_{3}+\Delta_{4}-1,-}^{x}(z_{4},\bar{z}_{4}) \tag{7}$$

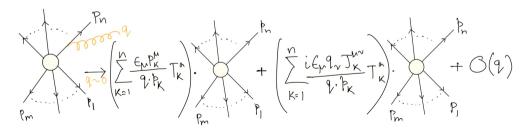
More generally,

$$\mathcal{O}_{\Delta_{1}}^{a,+}(z_{1},\bar{z}_{1})\mathcal{O}_{\Delta_{2}}^{b,+}(z_{2},\bar{z}_{2}) \sim -\frac{if_{c}^{ab}}{z_{12}}\sum_{n=0}^{\infty} B(\Delta_{1}+n-1,\Delta_{2}-1)\frac{\bar{z}_{12}^{n}}{n!}\bar{\partial}_{2}^{n}\mathcal{O}_{\Delta_{1}+\Delta_{2}-1}^{c,+}(z_{2},\bar{z}_{2}).$$
(8)

(Guevara, Himwich, Pate, Strominger '21)

SOFT FACTORIZATION THEOREM

Energetic soft theorems for gauge theory amplitudes



2 Infite tower of soft gluons:

$$R^{k,a}(z,\bar{z}) = \lim_{\Delta \to k} (\Delta - k) O_{\Delta}^{a,+}(z,\bar{z}), \qquad k = 1, 0, -1, ...$$
 (9)

 $\underbrace{ \text{gluon-gluon OPE (8)} }_{\text{mode expansions}} \underbrace{ \text{current algebra, S}}_{\text{(Strominger and collaborators '21)}}$

$$[S_{\alpha,m}^{p,a}, S_{\beta,n}^{q,b}] = -if^{abc} S_{\alpha+\beta,m+n}^{p+q-1,c}$$
(10)

 $\underbrace{ \text{graviton-graviton OPE} }_{\text{mode expansions}} \underbrace{ \text{current algebra, } w_{1+\infty} }$

$$[w_{\alpha,m}^p, w_{\beta,n}^q] = [m(q-1) - n(p-1)]w_{\alpha+\beta,m+n}^{p+q-2}$$
(11)

These are new infinite global symmetries of gauge theories and gravity in AFS.

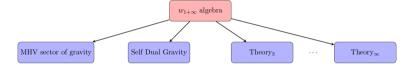


NISER's contributions to Celestial Holography?

- OPE = (singular part of the OPE)+ (what else ?)
- (Banerjee, Paul, Panda, Misra, RM'23)

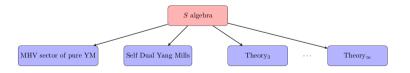
- **2** Find the "null states" of these algebras. But why?
- 1 They put constraints on the celestial amplitudes \longrightarrow Banerjee-Ghosh eqns. 202

4



(Banerjee, Kulkarni and Paul '23)

6



(Banerjee, Paul, Panda, Misra, RM'23)

6 Looking beyond MHV theories (e.g. NMHV gravitons),...

CELESTIAL HOLOGRAPHY GROUP IN NISER



Shamik Banerjee



Nishant Gupta



Mousumi Maitra



Sagnik Misra



Suman Guchait



KM 4 D > 4 D > 4 E > 4 E > E 9 Q C

Thank you for your attention !!

Keywords: celestial holography celestial cft asymptotic symmetries

asymptotically flat spacetime conformally soft theorems celestial ope

Flatspace holography gauge-gravity duality MHV gluons MHV gravitons

w1+00 algebra S algebra celestial ope BG equations