

BUILDING A SIMULATOR

① FORWARD DYNAMICS : given $\ddot{\theta}$ find \ddot{q}

$$(1) \begin{bmatrix} M & -J_c^T \\ J_c & 0 \end{bmatrix} \begin{bmatrix} \ddot{\theta} \\ \ddot{q} \end{bmatrix} = \begin{bmatrix} S^T \ddot{\theta} - R \\ -J \ddot{q} \end{bmatrix}$$

(2) \rightarrow RIGID CONTACT ASSUMPTION

We have also contact forces f as unknown

let's find a way to eliminate them ...

② solve (1) for \ddot{q} and insert it into (2)

$$M \ddot{q} + h = S^T \ddot{\theta} + J_c^T f$$

$$\ddot{q} = M^{-1} (-R + S^T \ddot{\theta} + J_c^T f)$$

$$J_c M^{-1} (-R + S^T \ddot{\theta} + J_c^T f) + J \ddot{q} = 0$$

③ solve for contact forces f

$$J_c M^{-1} J_c^T f + J_c N^{-1} (-h + S^T \dot{z}) + \ddot{J} \dot{q} = 0$$

$$f = (J_c M^{-1} J_c^T)^{-1} [-\ddot{J} \dot{q} + J_c N^{-1} (R - S^T \dot{z})]$$

Λ_c : joint inertia reflected at the contact

④ plug f back into dynamics (1)

$$M \ddot{q} + R = S^T \dot{z} + J^T \Lambda [-\ddot{J} \dot{q} + J M^{-1} (R - S^T \dot{z})]$$

$$(3) M \ddot{q} + R = \underline{S^T \dot{z}} + J^T (J M^{-1} J^T)^{-1} \ddot{J} \dot{q} + J^T (J M^{-1} J^T)^{-1} \underline{J M^{-1} R}$$

$$J^T = A = \begin{bmatrix} \end{bmatrix} \quad A^\# = \begin{bmatrix} \end{bmatrix}$$

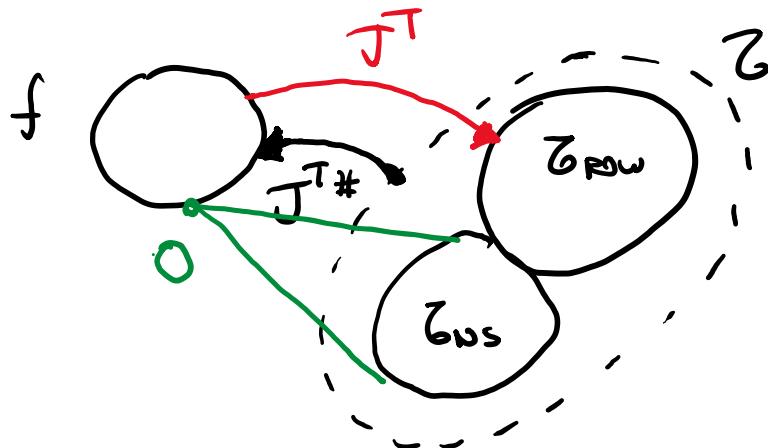
$$A^\# = (A^T A)^{-1} A^T$$

$$A_w^\# = (A^T W A)^{-1} A^T W \rightarrow J_w^{T\#} = (J W J^T)^{-1} J W$$

$$\text{cf. } A A_w^\# = I_m \quad \text{(ok)}$$

$J^T (J M^{-1} J^T)^{-1} J M^{-1} S^T \dot{z}$
 $J^T \#$ weighted with M^{-1}
(DYNAMICALLY CONSISTENT PSEUDOINVERSE) $J_M^{T\#}$

For forces?



$$z_{\text{ROW}} = J^T J^{T\#} z$$

$$z_{\text{NULL}} = (I - J^T J^{T\#}) z$$

$$N = I - J^T J^{T\#}$$

NULL SPACE
PROJECTOR

Torques That do not produce any contact force but only joint accelerations

Note N requires J to be full rank

The weighted version N_w is not symmetric
($N_w \neq N_w^T$)

⑤ From (3) collect $N_{\bar{R}}$

$$M\ddot{q} + R \underbrace{[I - J^T J_{\bar{R}}^{T\#}]}_{N_{\bar{R}}} + J^T \Lambda (J\dot{q}) = \underbrace{[I - J^T J_{\bar{R}}^{T\#}]}_{N_{\bar{R}}} S^T z$$

$$M\ddot{q} + N_{\bar{R}} R + J^T \Lambda (J\dot{q}) = N_{\bar{R}} S^T z$$

CONTACT
CONSISTENT
JOINT SPACE
DYNAMICS (1)

6) Replacing $-J\ddot{q} = \ddot{S}\ddot{q}$ [HUTTER PHD THESIS]
2.3

$$M\ddot{\ddot{q}} + N_{\bar{H}} h + J^T (J M^{-1} J^T)^{-1} J \ddot{\ddot{q}} = N_{\bar{H}} S^T \zeta$$

$$\underbrace{[I - J^T (J M^{-1} J^T)^{-1} J M^{-1}]}_{N_{\bar{M}}} M\ddot{\ddot{q}} + N_{\bar{H}} h = N_{\bar{H}} S^T \zeta$$

$$N_{\bar{H}} (M\ddot{\ddot{q}} + h) = N_{\bar{H}} S^T \zeta \quad (2)$$

↳ projection of dynamics through N

⊕ The contact force disappeared from equation

$$N_{\bar{H}} \underbrace{J^T f}_{\text{now space}} = 0$$

Note: $N\bar{\mu}$ is a linear operator that maps the dynamics onto a constrained manifold such that EOM is independent of constraint forces \Rightarrow any operator s.t. $PJ^T = 0$ would work

7) Integrate

$$\ddot{q} = (N \ M)^{-1} N^T (S^T z - h) \quad \text{we cannot do that!}$$

↳ is not invertible, most of times is singular

- if disappears but we also lose information on accelerations that we cannot recover because we project them onto a lower dimensional space with (2), and the result won't be correct.

for now we use:

$$\ddot{\ddot{q}} = M^{-1} [N(S^T \dot{z} - R) - J^T \Lambda J \dot{q}]$$

- with this we do not lose informations because we multiply $\ddot{\ddot{q}}$ for M That is positive definite

WITH forward Euler integration:

$$\dot{\dot{q}}_{k+1} = \dot{\dot{q}}_k + dT \ddot{\ddot{q}}_k$$

$$q_{k+1} = q_k + dT \dot{q}_{k+1} + \frac{dT^2}{2} \ddot{\ddot{q}}_{k+1}$$

NOTE : $\ddot{q} \approx 0$ for non redundant manipulators
because They are completely constrained

NOTE : \dot{q}^+ ? constraint consistent joint
velocity?

⊖ ... we need to consider impact dynamics

LIFT OFF: $F \downarrow 0$, $\ddot{x}_c > 0$

TOUCH DOWN: $x_c = \text{const}$, $\dot{x}_c, \ddot{x}_c = 0$, $F \uparrow 0$

↳ collision at landing:

- foot point is immediately brought to rest
- instantaneous change in velocity To avoid penetration
- To compute it we need to integrate the dynamics of collision (inelastic impact)

CONTACT DYNAMICS

① interpretation of Lagrangian dynamics on single contact

$$\int_{t^-}^{t^+} [M \ddot{q} + R - J^T F] dt = M(\dot{q}^+ - \dot{q}^-) - J^T F_c = 0 \quad (1)$$

\downarrow impulse

- all non impulsive forces drop out
- after collision the foot is at rest

$$\dot{x}_f^+ = \emptyset = J \dot{q}^+ \quad (\text{inelastic impact}) \quad (2)$$

② solving for \dot{q}^+ and inserting into (2)

$$\dot{q}^+ = \dot{q}^- + M^{-1} J^T F_c$$

$$\dot{x}_f^+ = J(M^{-1} J^T F_c + \dot{q}^-) = 0$$

③ we can compute the impulse:

$$F_c = - (J M^{-1} J^T)^{-1} J \dot{q}^- \quad F_c \geq 0$$

④ we compute post-impact joint velocity

$$\dot{\bar{q}}^+ = M^{-1} J^T \left[-(J M^{-1} J^T)^{-1} J \dot{q}^- \right] + \dot{q}^-$$

$$= \left[I - M^{-1} J^T (J M^{-1} J^T)^{-1} J \right] \dot{q}^-$$

$$I - J_{\bar{q}}^{\#} J = N_{\text{rel}}$$

$$\boxed{\dot{\bar{q}}_c^+ = N_{\text{vel}} \dot{q}^-}$$

ONLY AT IMPACT!

Note:

This Null space projector is acting at the velocity level and is different from the previous one that acts at Torque level

ENERGY LOST IN COLLISIONS

⑥

$$E_{\text{Loss}} = E_{\text{kin}} = \frac{1}{2} \Delta \dot{q}^T M \Delta \dot{q}$$

↓ Task space

$$\Delta \dot{x}_c = 0 - \dot{x}_c = J \Delta \dot{q}$$

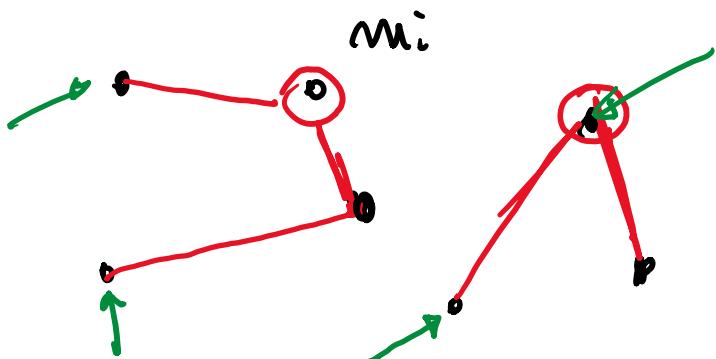
$$\Delta \ddot{q} = J^{-1} \Delta \dot{x}_c$$

$$E_{\text{Loss}} = E_{\text{kin}} = \frac{1}{2} \Delta \dot{x}_c^T J^{-T} M J^{-1} \Delta \dot{x}_c$$

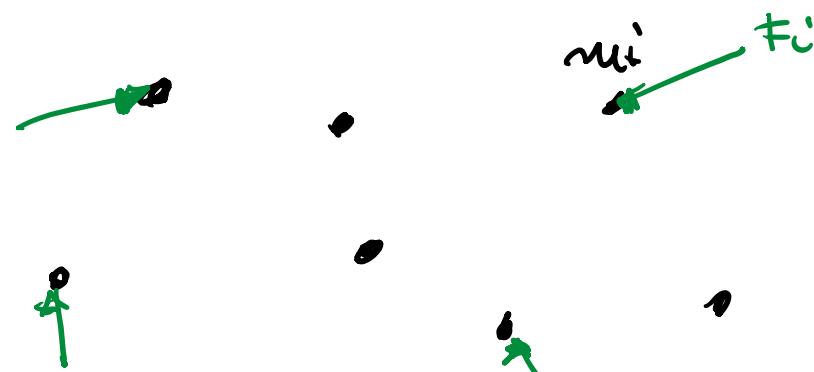
$$= \frac{1}{2} \Delta \dot{x}_c^T \Lambda_c \Delta \dot{x}_c$$

↑ end-effector inertia

GAUSS PRINCIPLE OF LEAST CONSTRAINT



collection of particles with
some constraints



unconstrained equivalent

\Rightarrow The accelerations deviate from the unconstrained ones in a LS sense:

$$\ddot{q}_c = \min_{\ddot{q}} \frac{1}{2} [\ddot{q} - \ddot{q}_{uc}]^T M(q) [\ddot{q} - \ddot{q}_{uc}]$$

$$\text{s.t. } J\ddot{q} + \dot{J}\dot{q} = 0 \quad \text{with } \dot{q}_{uc} = M^{-1}(S^T z - R)$$

$$\rightarrow \frac{1}{2} \ddot{q}^T M \ddot{q} - \underbrace{\frac{1}{2} \ddot{q}^T M q_{uc} - \frac{1}{2} \ddot{q}_{uc}^T M \ddot{q}}_{- \ddot{q}_{uc}^T M \ddot{q}} + \dot{q}_{uc}^T \dot{q}_{uc}$$

SOLVE THE QP

$$\begin{aligned}\ddot{q}_c = \min_{\dot{x}} \quad & \frac{1}{2} \dot{x}^T G \dot{x} + g \cdot \dot{x} \\ \text{s.t.} \quad & A \dot{x} = b\end{aligned}$$

WITH :

$$x = \ddot{q}$$

$$G = M$$

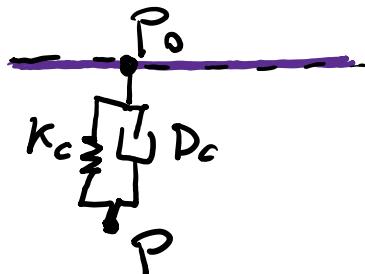
$$g = -\ddot{q}_{uc}^T M$$

$$A = J$$

$$b = -J \ddot{q}$$

SOFT CONTACT MODEL

Many simulation environments use them
The contact force is given by:



$$f = K_c(P_0 - P) - D_c \dot{P}$$

$$M\ddot{\bar{q}} + R = S^T \bar{z} + J^T f$$

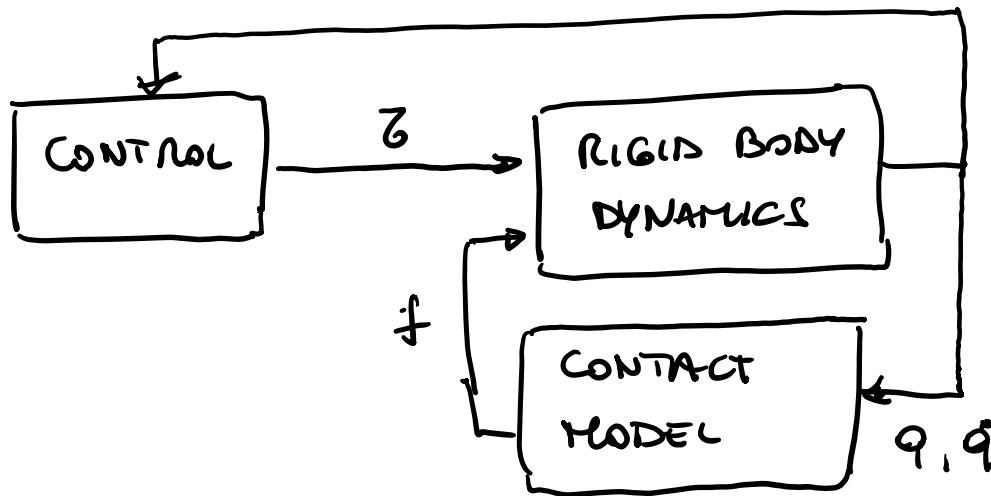
- different directions:

$$f_z = K (P_{0z} - P_z) - D \dot{P}_z$$

$$f_x = K (P_{0x} - P_x) - D \dot{P}_x$$

- generic law extension (e.g. nonlinear spring)

$$f = \text{Law}(P, \dot{P})$$



COMPLIANT VS RIGID MODELS

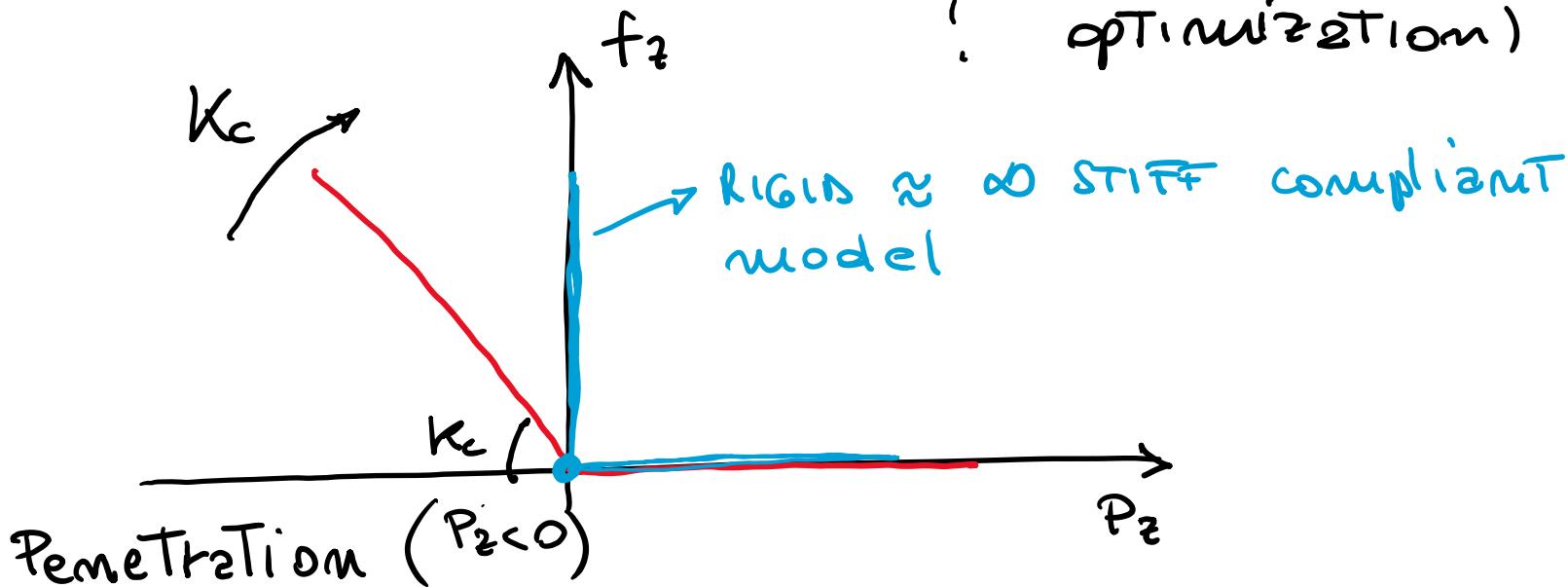
COMPLIANT

- + easy to compute
- stiff springs \Rightarrow ODE numerically stiff \Rightarrow small integration steps

SLOW SIMULATIONS

RIGID

- + straight forward to integrate analyze
- + preferred in stability analysis
- pathological to differentiate (e.g. in optimization)



References:

- M. Hutter, Design and control of legged robots with compliant actuation, PhD Thesis, 2013.
- F. Aghili, A Unified Approach for Inverse and Direct Dynamics of Constrained Multibody Systems Based on Linear Projection Operator: Applications to Control and Simulation, 2005.
- A. Del Prete, PhD Thesis, 2013.