

Basic MHD Equations

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This is the lecture notes for Magnetohydrodynamics (MHD) equations, including the physical discussions and equations derivations.

1 Equations and Derivations

1.1 Electromagnetic Equations

Maxwell's Equations

$$\nabla \times \mathbf{B} = \mu \mathbf{j} + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} \quad (1)$$

$$\nabla \cdot \mathbf{B} = 0 \quad (2)$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad (3)$$

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}, \quad (4)$$

where \mathbf{E} is the *electric field*, \mathbf{B} is the *magnetic field*, \mathbf{j} is the *current density*, ρ is the *charge density*, μ is the *magnetic permeability*, ϵ_0 is the *permittivity of free space*, and c is the *speed of light* in vacuum ($\sim 2.998 \times 10^8$ m/s). The four Maxwell equations represent different physical phenomena, the first equation suggests that either currents or time-varying electric fields can generate magnetic fields, the second equation indicates that there are no magnetic monopoles and a magnetic flux tube has a constant strength along its length, and the third and fourth equations imply that either time-varying magnetic fields or electric charges can give rise to electric fields.

Here we also introduce the concept of *Debye Length*, which is a measure of a charge carrier's net electrostatic effect in a solution and how far its electrostatic effect persists. The potential function for a single electron is

$$\Phi(r) = \frac{e}{4\pi\epsilon_0 r}.$$

Under the influence of a positive charge q , the charge distribution satisfies Boltzmann's distribution:

$$n_e = n_{e0} \exp\left(-\frac{e\Phi(r)}{k_B T_e}\right) \approx n_{e0} \left(1 + \frac{e\Phi(r)}{k_B T_e}\right)$$

Assume $n_e = n_{e0} + \delta n_e$, we have:

$$\delta n_e = n_{e0} \frac{e\Phi(r)}{k_B T_e}.$$

Poisson's equation claims:

$$\nabla^2 \Phi(r) = -\frac{\rho}{\epsilon_0} = \frac{e^2 n_{e0}}{k_B T_e \epsilon_0} \Phi(r).$$

Combining the boundary conditions $\Phi(r \rightarrow \infty) = 0$ and $\Phi(r \rightarrow 0) = \frac{q}{4\pi\epsilon_0 r}$, we have the Debye potential function Φ :

$$\Phi(r) = \frac{q}{4\pi\epsilon_0 r} \exp\left(-\frac{r}{\lambda_D}\right), \quad (5)$$

where

$$\lambda_D = \sqrt{\frac{k_B T_e \epsilon_0}{e^2 n_{e0}}} \quad (6)$$

is the *Debye Length*.

Generalized Ohm's Law

Ohm's Law asserts that the current density is proportional to the total electric field (in a frame of reference moving with the frame), it can be written as:

$$\mathbf{j} = \sigma(\mathbf{E} + \mathbf{v} \times \mathbf{B}), \quad (7)$$

where σ is the electrical conductivity.

However, a generalization of Ohm's Law may be more appropriate in some regions of the space. We start from equations of motion for electrons and ions (assumed to be protons):

$$\frac{\partial}{\partial t}(n_e m_e \mathbf{u}_e) + \nabla \cdot (n_e m_e \mathbf{u}_e \mathbf{u}_e) = -\nabla \cdot \mathbf{P}_e - en_e(\mathbf{E} + \mathbf{u}_e \times \mathbf{B}) + \mathbf{R}_e,$$

$$\frac{\partial}{\partial t}(n_i m_i \mathbf{u}_i) + \nabla \cdot (n_i m_i \mathbf{u}_i \mathbf{u}_i) = -\nabla \cdot \mathbf{P}_i + en_i(\mathbf{E} + \mathbf{u}_i \times \mathbf{B}) + \mathbf{R}_i.$$

Multiply the ion equation by m_e/m_i , then subtract the electron equation from the ion equation, we have:

$$\begin{aligned} m_e \frac{\partial}{\partial t}(n_i \mathbf{u}_i - n_e \mathbf{u}_e) + m_e \nabla \cdot (n_i \mathbf{u}_i \mathbf{u}_i - n_e \mathbf{u}_e \mathbf{u}_e) = \\ \nabla \cdot (\mathbf{P}_e - \frac{m_e}{m_i} \mathbf{P}_i) + e[(n_e + \frac{m_e}{m_i} n_i) \mathbf{E} + (n_e \mathbf{u}_e + \frac{m_e}{m_i} n_i \mathbf{u}_i) \times \mathbf{B}] - (\mathbf{R}_e - \frac{m_e}{m_i} \mathbf{R}_i), \end{aligned}$$

where $\mathbf{R}_i = -\mathbf{R}_e$. Note that in MHD regime we have $n_e \approx n_i \approx n$, and we neglect all the $O(m_e/m_i)$ terms, we have:

$$\frac{m_e}{e} \frac{\partial \mathbf{j}}{\partial t} + m_e \nabla \cdot [n(\mathbf{u}_i \mathbf{u}_i - \mathbf{u}_e \mathbf{u}_e)] = \nabla \cdot \mathbf{P}_e + ne(\mathbf{E} + \mathbf{u}_e \times \mathbf{B}) - \mathbf{R}_e.$$

Replace u_e with $u_e = u_i - \frac{\mathbf{j}}{ne} \approx \mathbf{u} - \frac{\mathbf{j}}{ne}$, where u is the plasma velocity, we have:

$$\mathbf{E} + \mathbf{v} \times \mathbf{B} = \frac{m_e}{ne^2} \frac{\partial \mathbf{j}}{\partial t} - \frac{1}{ne} \nabla \cdot \mathbf{P}_e + \frac{\mathbf{j} \times \mathbf{B}}{ne} + \frac{1}{ne} \mathbf{R}_e + \frac{m_e}{ne^2} \nabla \cdot (\mathbf{j} \mathbf{u} + \mathbf{u} \mathbf{j} - \frac{\mathbf{j} \mathbf{j}}{ne}).$$

\mathbf{R}_e represents the rate of change of the electron momentum due to the collision with ion, which can be written as:

$$\mathbf{R}_e = nm_e \mu_c (\mathbf{u}_i - \mathbf{u}_e) = \frac{\mu_c m_e}{e} \mathbf{j},$$

where μ_c is the collision frequency. By defining the resistivity $\eta = \frac{\mu_c m_e}{ne^2}$, we can write the collision term as

$$\frac{\mathbf{R}_e}{ne} = \eta \mathbf{j}.$$

Finally, we can write the *generalized Ohm's Law* as:

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = \eta \mathbf{j} + \frac{1}{ne} \mathbf{j} \times \mathbf{B} - \frac{1}{ne} \nabla \cdot \mathbf{P}_e + \frac{m_e}{ne^2} [\frac{\partial \mathbf{j}}{\partial t} + \nabla \cdot (\mathbf{u} \mathbf{j} + \mathbf{j} \mathbf{u} - \frac{\mathbf{j} \mathbf{j}}{ne})]. \quad (8)$$

Induction Equation

Combining equations (1), (3), and (7), we can easily have:

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times (-\mathbf{v} \times \mathbf{B} + \mathbf{j}/\sigma) = \nabla \times (\mathbf{v} \times \mathbf{B}) - \nabla \times (\eta \nabla \times \mathbf{B}),$$

where $\eta = 1/(\mu\sigma)$ is the *magnetic diffusivity*. Identify that $\nabla \times (\nabla \times \mathbf{B}) = \nabla(\nabla \cdot \mathbf{B}) - \nabla^2 \mathbf{B}$, we find the *induction equation*:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B}. \quad (9)$$

1.2 Plasma Equations

Mass Continuity

The *equation of mass conservation* usually comes in two forms (equivalent to each other):

$$\frac{d\rho}{dt} + \rho \nabla \cdot \mathbf{v} = 0 \text{ (a)} \quad \text{or} \quad \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 \text{ (b)}, \quad (10)$$

where ρ is the *mass density* and \mathbf{v} is the *velocity* of the fluid.

Equation of Motion

Under the condition of electrical neutrality, the *equation of motion* can be written as:

$$\rho \frac{d\mathbf{v}}{dt} = -\nabla p + \mathbf{j} \times \mathbf{B} + \mathbf{F}, \quad (11)$$

where p is the *pressure*, $\mathbf{j} \times \mathbf{B}$ is the *Lorentz force*, and \mathbf{F} represents the *external force*. Usually, the external force consists of two parts: the *gravity* (\mathbf{F}_g) and *viscosity* (\mathbf{F}_v). The gravity force is:

$$\mathbf{F}_g = -\rho g(r) \hat{r},$$

where $g(r) = M(r)G/r^2$ is the *local gravitational acceleration*. The viscous force is

$$\mathbf{F}_v = \rho \nu [\nabla^2 + \frac{1}{3} \nabla (\nabla \cdot \mathbf{v})],$$

where ν is the *kinematic viscosity*. We can simplify the force to $\mathbf{F}_v = \rho \nu \nabla^2 \mathbf{v}$ when the flow is incompressible.

Perfect Gas Law

The gas pressure is determined by an equation of state, which is taken for simplicity as the *perfect gas law*:

$$p = \frac{\tilde{R}}{\tilde{\mu}} \rho T = \frac{k_B}{m} \rho T = n k_B T, \quad (12)$$

where \tilde{R} is the *specific gas constant*, $\tilde{\mu}$ is the *mean molecular weight*, and T is the *temperature*, n is the *total number* of particles per volume.

1.3 Energy Equations

Different Forms of the Heat Equation

The *heat equation* is the last fundamental equation, which can be written as:

$$\rho T \frac{ds}{dt} = -\mathcal{L}, \quad (13)$$

where \mathcal{L} is the *energy loss function*, representing the net effect of all the sinks and sources of energy in the system. s is the *entropy per unit mass*. The equation states that the rate of increase of heat for a unit volume as it moves in space is due to the net effect of the energy sinks and sources.