

116217 - Solid State Physics

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Abstract

1 Introduction

Phase are distinguished by symmetries, e.g., liquid and solid have difference in transnational symmetry. Those phase transitions happens together with spontaneous symmetry breakage. For example, magnet "chooses" poles when transits to magnetic phase.

1.1 Heat capacity

Most of matereals have $C = 3k_B$, except diamond in room temperature and pressure. In low temperatures though, materials have $C = \alpha T + \gamma T^3$. Thus, most of materials get to saturation in room temperature, while diamond doesn't. The linear factor appears only in conductors, thus should come from electrons.

Equipartition theorem

$$\left\langle x_m \frac{\partial H}{\partial x_n} \right\rangle = \delta_{mn} k_B T$$

in particular, for quadratic degree of freedom we get $\frac{1}{2} k_B T$.

Thus, assuming each atom is harmonic oscillator, we have

$$\langle E \rangle = 3k_B T$$

and

$$c_V = 3k_B$$

Einstein model Now assume the oscillators are quantum:

$$E_n = \hbar\omega \left(n + \frac{1}{2} \right)$$

The partition function

$$Z = \sum_n e^{-\beta E_n} = \sum_n e^{-\beta \hbar\omega \left(n + \frac{1}{2} \right)} = e^{-\frac{\beta \hbar\omega}{2}} \frac{1}{1 - e^{-\beta \hbar\omega}} = \frac{e^{-\frac{\beta \hbar\omega}{2}}}{e^{-\frac{\beta \hbar\omega}{2}} \left[e^{\frac{\beta \hbar\omega}{2}} - e^{-\frac{\beta \hbar\omega}{2}} \right]} = \frac{1}{2 \sinh\left(\frac{\beta \hbar\omega}{2}\right)}$$

$$\langle E \rangle = -\frac{\partial}{\partial \beta} \ln Z_N = -\frac{1}{Z_N} \frac{\partial}{\partial \beta} Z = N \cdot 2 \sinh\left(\frac{\beta \hbar\omega}{2}\right) \cdot \frac{\hbar\omega}{2} \frac{\cosh\left(\frac{\beta \hbar\omega}{2}\right)}{\sinh^2\left(\frac{\beta \hbar\omega}{2}\right)} = \frac{N \hbar\omega}{2} \frac{1}{\tanh\left(\frac{\beta \hbar\omega}{2}\right)}$$

$$\langle E \rangle = N \hbar\omega \left(n_B(\beta \hbar\omega) + \frac{1}{2} \right)$$

$$n_B(x) = \frac{1}{e^x - 1}$$

Now

$$c_V = \frac{\partial \langle E \rangle}{\partial T} = 3N k_B (\beta \hbar\omega)^2 \frac{e^{\beta \hbar\omega}}{(e^{\beta \hbar\omega} - 1)^2}$$

$$\lim_{T \rightarrow \infty} c_V = 3k_B N$$

$$\lim_{T \rightarrow 0} c_V = 0$$

which is expected, however, we got exponential dependence on T and not T^3 .

1.2 Debye model

Debye proposed to model atoms in solid as sound waves. Plank already quantized EM-waves, the only differences is number of polarization options and speed of wave.

We'll often use periodic bound conditions, since this preserves the symmetry of the system.

Now, for wave e^{ikx} we get $e^{ikx} = e^{ik(x+L)}$, thus we get a condition on wavenumber

$$k = \frac{2\pi}{L}n$$

We also want to switch to integral on k .

The connection between frequency and wavenumber is $\omega(\vec{k}) = v \cdot |\vec{k}|$. The energy per wavenumber

$$\begin{aligned}\langle E \rangle &= 3 \sum_{\vec{k}} \hbar \omega(\vec{k}) \left(n_B(\beta \hbar \omega(\vec{k})) + \frac{1}{2} \right) \\ \langle E \rangle &= 3 \frac{L^3}{(2\pi)^3} \int d^3k \hbar \omega(\vec{k}) \left(n_B(\beta \hbar \omega(\vec{k})) + \frac{1}{2} \right) = 3 \frac{4\pi L^3}{(2\pi)^3} \int \frac{\omega^2 d\omega}{v^3} \hbar \omega \left(n_B(\beta \hbar \omega) + \frac{1}{2} \right)\end{aligned}$$

Now,

$$g(\omega) = L^3 \cdot \frac{12\pi\omega^2}{(2\pi)^3 v^3} = N \frac{12\pi\omega^2}{(2\pi)^3 v^3 \rho}$$

is density of states.

Define $\omega_d^3 = 6\pi^2 \rho v^3$, acquiring $g(\omega) = N \frac{9\omega^2}{\omega_d^3}$

$$\langle E \rangle = \frac{9N\hbar}{\omega_d^3} \int \omega^3 d\omega \frac{1}{e^{\beta\hbar\omega} - 1} + \text{const}$$

Substituting $\beta\hbar\omega = x$:

$$\begin{aligned}\langle E \rangle &= \frac{9N\hbar}{\omega_d^3 (\beta\hbar)^4} \underbrace{\int dx x^3 \frac{1}{e^x - 1}}_{\frac{\pi^4}{15}} \\ \frac{\partial \langle E \rangle}{\partial T} &= \frac{12\pi^4}{5} \frac{nk_B (k_B T)^3}{(\hbar\omega_d)^3}\end{aligned}$$

The reason we get divergence in high temperatures is that we assume any frequency is possible, however, the wavelength can't be less than distance between atoms.

The total number of possible frequencies equals to number of degrees of freedom:

$$\begin{aligned}\int_0^{\omega_{cutoff}} d\omega g(\omega) &= 3N \\ \int_0^{\omega_{cutoff}} d\omega N \frac{9\omega^2}{\omega_d^3} &= N \cdot \frac{3\omega^3}{\omega_d^3} = 3N\end{aligned}$$

i.e., $\omega_{cutoff} = \omega_d$:

$$\langle E \rangle = \int_0^{\omega_d} d\omega g(\omega) \frac{\hbar\omega}{e^{\beta\hbar\omega} - 1}$$

In low temperatures our calculation is yet good, in high temperatures, using Taylor expansion of $e^{\beta\hbar\omega}$

$$\langle E \rangle = \int_0^{\omega_d} d\omega g(\omega) \frac{\hbar\omega}{\beta\hbar\omega} = \int_0^{\omega_d} d\omega g(\omega) k_B T = 3N k_B T$$

We define some different quantities based on Debye frequency $T_d = \frac{\hbar\omega_d}{k_B}$ $k_d = \frac{\omega_d}{v}$

1.3 Phonons

Suppose we have a one-dimensional chain of atoms of mass m in position

$$x_n^{eq} = na$$

with oscillations

$$u = x_n - x_n^{eq}$$

If u is small it can be approximated with quadratic potential, i.e., a chain of springs.

$$V_{tot} = \sum_j V(x_{j+1} - x_j) = \sum_j \frac{K}{2} (x_{j+1} - x_j - a)^2 = \sum_j \frac{K}{2} (u_{j+1} - u_j)^2$$

$$F_n = -\frac{\partial V_{tot}}{\partial u_n} = K(u_{n+1} - u_n) + K(u_{n-1} - u_n)$$

We get system of equations

$$m\ddot{u}_n = F_n = K(u_{n+1} + u_{n-1} - 2u_n)$$

Or, in matrix form

$$m \begin{pmatrix} \ddot{u}_1 \\ \ddot{u}_2 \\ \vdots \\ \ddot{u}_n \end{pmatrix} = K \begin{pmatrix} -2 & 1 & 0 & \dots & 0 \\ 1 & -2 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & -2 \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \\ \vdots \\ u_n \end{pmatrix}$$

Guessing a solution

$$u_n = Ae^{i\omega t} e^{-ikan}$$

We get

$$-m\omega^2 u_n = KAe^{i\omega t} (e^{-ika(n+1)} + e^{-ika(n-1)} - 2e^{-ikan}) = Ku_n (e^{-ika} + e^{ika} - 2)$$

$$-m\omega^2 = 2K(\cos(ak) - 1)$$

i.e.,

$$m\omega^2 = 4K \sin^2\left(\frac{ka}{2}\right)$$

or,

$$\omega(k) = 2\sqrt{\frac{K}{m}} \left| \sin\left(\frac{ka}{2}\right) \right|$$

The dispersion has minimum and maximum. We can interval of allowed frequencies "band" and its size is called "bandwidth". The dispersion is periodic with period of $\frac{2\pi}{a}$, since the system is symmetric to translations of a .

Solutions of the equations Since translation operator commutes with Hamiltonian, the eigenfunctions of Hamiltonian got to be eigenfunctions of the translation operator, which are

$$u_N = Ae^{i\omega t - ikna}$$

1.4 Reciprocal lattice

Reciprocal space Space of wavenumbers. We found equivalent points in inverse space. We call all the points equivalent to $k = 0$ reciprocal lattice, i.e., $k = \frac{2\pi}{a}p$ for $p \in \mathbb{Z}$.

First Brillouin zone is a unit cell centered in $k = 0$.

Finite system If the system is finite with size L ,

$$e^{ikna} = e^{ik\left(n + \frac{L}{a}\right)a}$$

Thus

$$k = \frac{2\pi}{L} \cdot p$$

i.e., the number of modes

$$N_{modes} = \frac{\frac{2\pi}{a}}{\frac{2\pi}{L}} = \frac{L}{a} = N_{atoms}$$

Properties of dispersion

Sound waves

$$\omega(k) \approx \sqrt{\frac{K}{m}} |k| a = v |k|$$

Short wavelengths

$$k \approx \frac{\pi}{a}$$

Velocity

$$v_{group} = \frac{d\omega}{dk}$$
$$v_{phase} = \frac{\omega}{k}$$

1.5 Quantized phonons

We quantize our oscillator, with energy

$$E_N = \hbar\omega \left(n + \frac{1}{2} \right)$$

Crystal momentum We define crystal momentum, momentum modulo $\hbar \frac{2\pi}{a}$. Since the lattice can give or take momentum in portions $\hbar \frac{2\pi}{a}$.

1.6 Two-atom chain of atoms

Suppose we have two types of atoms, x and y , with locations

$$x_n = na + \frac{a}{10}$$

$$y_n = na + \frac{8a}{10}$$

Suppose $m_1 = m_2$ but $k_1 \neq k_2$, (k_1 connects x_n with y_n and k_2 connects y_n with x_{n+1}).

$$\delta \ddot{x}_n = k_2(\delta y_n - \delta x_n) + k_1(\delta y_{n-1} - \delta x_n)$$

$$\delta \ddot{x}_n = k_1(\delta x_{n+1} - \delta y_n) + k_2(\delta x_n - \delta y_n)$$