

Condensed Matter Physics

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References

Condensed Matter Physics – Michael P. Marder

Density-Functional Theory of Atoms and Molecules

– Robert G. Parr and Weitao Yang

Introduction To Solid State Physics – Charles Kittel



The Free Fermi Gas and Single Electron Model

Condensed Matter Physics – Michael P. Marder

Chapter 6



The Hamiltonian

Much of condensed matter physics lies within a Hamiltonian that one easily can write down in a single line. It is

$$\hat{\mathcal{H}} = \sum_l \frac{\hat{p}_l^2}{2M_l} + \frac{1}{2} \sum_{l \neq l'} \frac{q_l q_{l'}}{|\hat{R}_l - \hat{R}_{l'}|}.$$



The single-electron model

$$\sum_{l=1}^N \left(\frac{-\hbar^2 \nabla_l^2}{2m} + U(\vec{r}_l) \right) \Psi(\vec{r}_1 \dots \vec{r}_N) = \mathcal{E} \Psi(\vec{r}_1 \dots \vec{r}_N)$$

$$\left(\frac{-\hbar^2 \nabla^2}{2m} + U(\vec{r}) \right) \psi_l(\vec{r}) = \mathcal{E}_l \psi_l(\vec{r})$$



The free Fermi gas

$$\frac{-\hbar^2}{2m} \sum_{l=1}^N \nabla_l^2 \Psi(\vec{r}_1 \dots \vec{r}_N) = \mathcal{E} \Psi(\vec{r}_1 \dots \vec{r}_N)$$

To simplify further we impose periodic boundary conditions

$$\Psi(x_1 + L, y_1, z_1 \dots, z_N) = \Psi(x_1, y_1, z_1 \dots z_N)$$

$$\Psi(x_1, y_1 + L, z_1 \dots, z_N) = \Psi(x_1, y_1, z_1 \dots z_N)$$

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One Free Fermion

$$\psi_{\vec{k}} = \frac{1}{\sqrt{\mathcal{V}}} e^{i\vec{k} \cdot \vec{r}}$$

$$L^3 = \mathcal{V}$$

$$\vec{k} = \frac{2\pi}{L} (l_x, l_y, l_z)$$

l_x, l_y , and l_z are integers ranging from $-\infty$ to ∞

$$\epsilon_{\vec{k}}^0 = \frac{\hbar^2 k^2}{2m}$$



Many Free Fermions

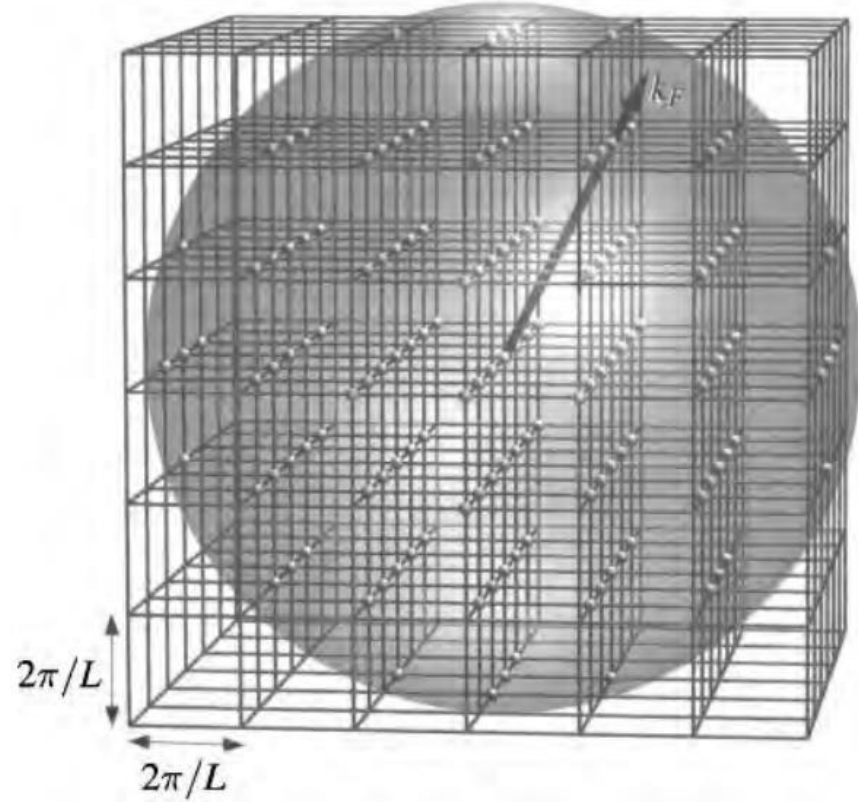
The ground state of electrons obeying free Fermi gas assumption is constructed from products of the one-electron wave functions. The Pauli exclusion principle forbids any given state from being occupied more than once, and therefore any given state indexed by \vec{k} is able to host no more than two electrons, one for each value of spin.

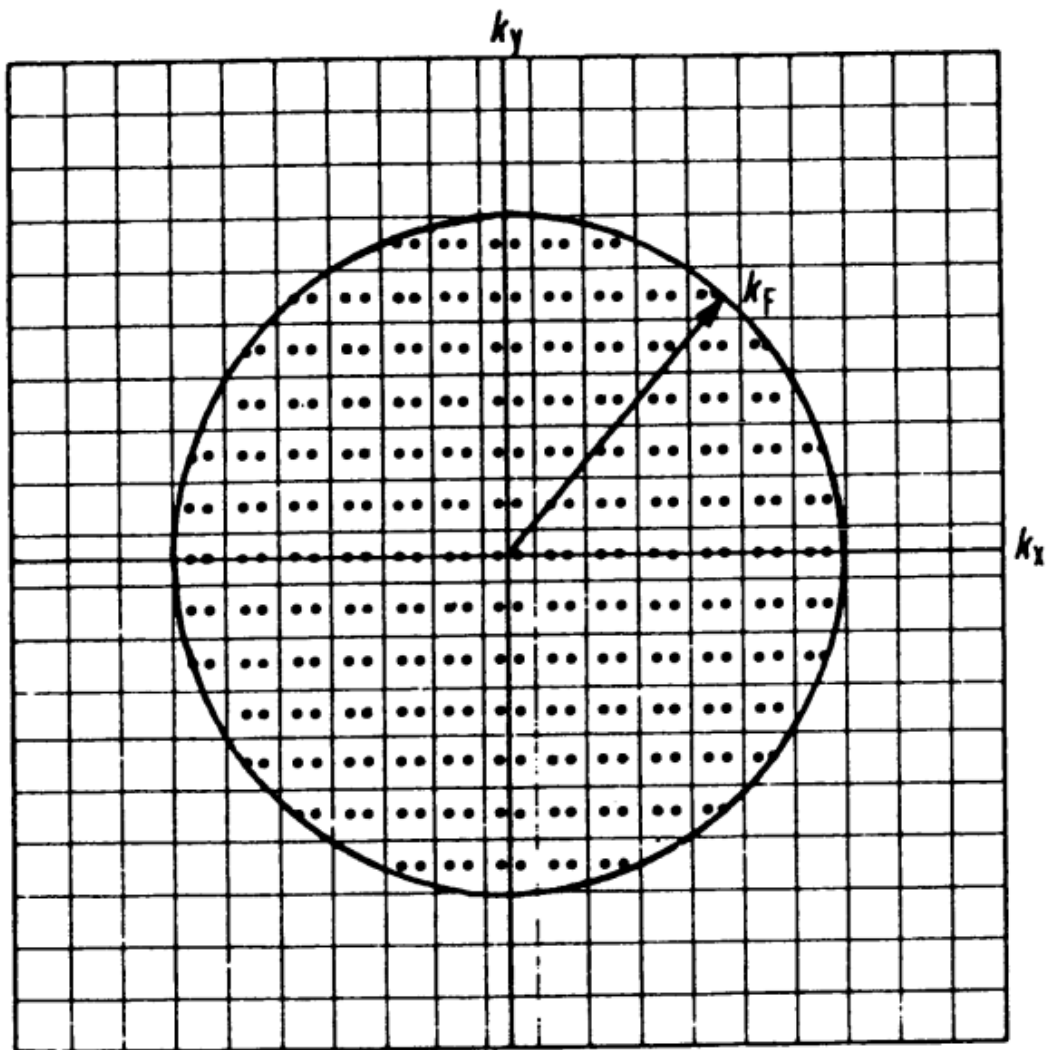


Many Free Fermions

$$\vec{k} = \frac{2\pi}{L} (l_x, l_y, l_z)$$

l_x, l_y , and l_z are integers ranging from $-\infty$ to ∞





Density of States

$$D_{\vec{k}} = 2 \frac{1}{(2\pi)^3}$$

For each wave vector Pauli's exclusion principle allows two electrons, one with spin up and the other with spin down.

$$\int [d\vec{k}] \equiv \frac{2}{V} \sum_{\vec{k}} = \int d\vec{k} D_{\vec{k}} = \frac{2}{(2\pi)^3} \int d\vec{k}$$



Energy Density of States

$$D(\mathcal{E}) = \int [d\vec{k}] \delta(\mathcal{E} - \mathcal{E}_{\vec{k}})$$

The units of densities of states are able to change without much warning. Often they are expressed in units of 1/[eV atom], which means they are related to the function defined by above equation by a factor of density n .

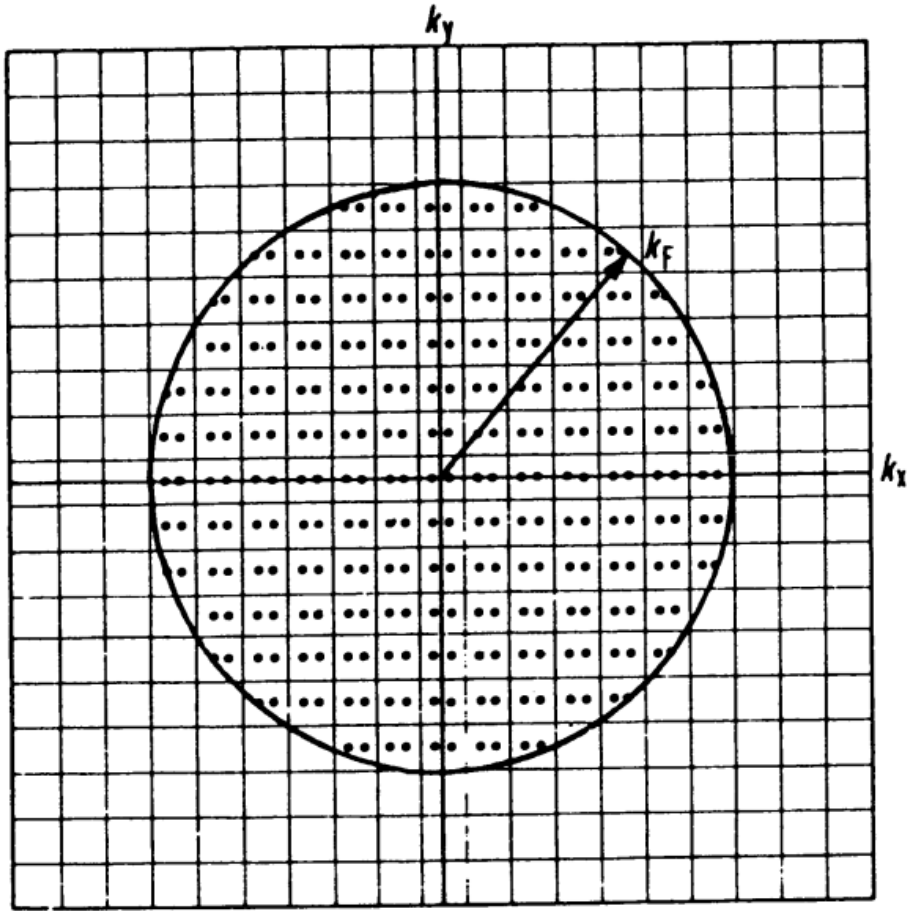


Energy Density of States

$$\begin{aligned} D(\mathcal{E}) &= \int [d\vec{k}] \delta(\mathcal{E} - \mathcal{E}_{\vec{k}}^0) \\ &= 4\pi \frac{2}{(2\pi)^3} \int_0^\infty dk k^2 \delta(\mathcal{E} - \mathcal{E}_{\vec{k}}^0) \\ &= \frac{1}{\pi^2} \int_0^\infty \frac{d\mathcal{E}^0}{|d\mathcal{E}^0/dk|} \frac{2m\mathcal{E}^0}{\hbar^2} \delta(\mathcal{E} - \mathcal{E}^0) \\ &= \frac{m}{\hbar^3 \pi^2} \sqrt{2m\mathcal{E}} \\ &= 6.812 \cdot 10^{21} \sqrt{\mathcal{E}/\text{eV}} \text{ eV}^{-1} \text{ cm}^{-3}. \end{aligned}$$

For the free
Fermi gas





$$4\pi \int_0^{k_F} g(k) k^2 dk = \frac{N}{V_g} = n$$

$$k_F = (3\pi^2 n)^{1/3}, \quad E_F = \frac{\hbar^2}{2m} (3\pi^2 n)^{2/3}$$

$$k_F = (3\pi^2 n)^{1/3} = 3.09 [n \cdot \text{\AA}^3]^{1/3} \text{\AA}^{-1}$$

$$\mathcal{E}_F = \frac{\hbar^2 k_F^2}{2m} = 36.46 [n \cdot \text{\AA}^3]^{2/3} \text{eV}$$

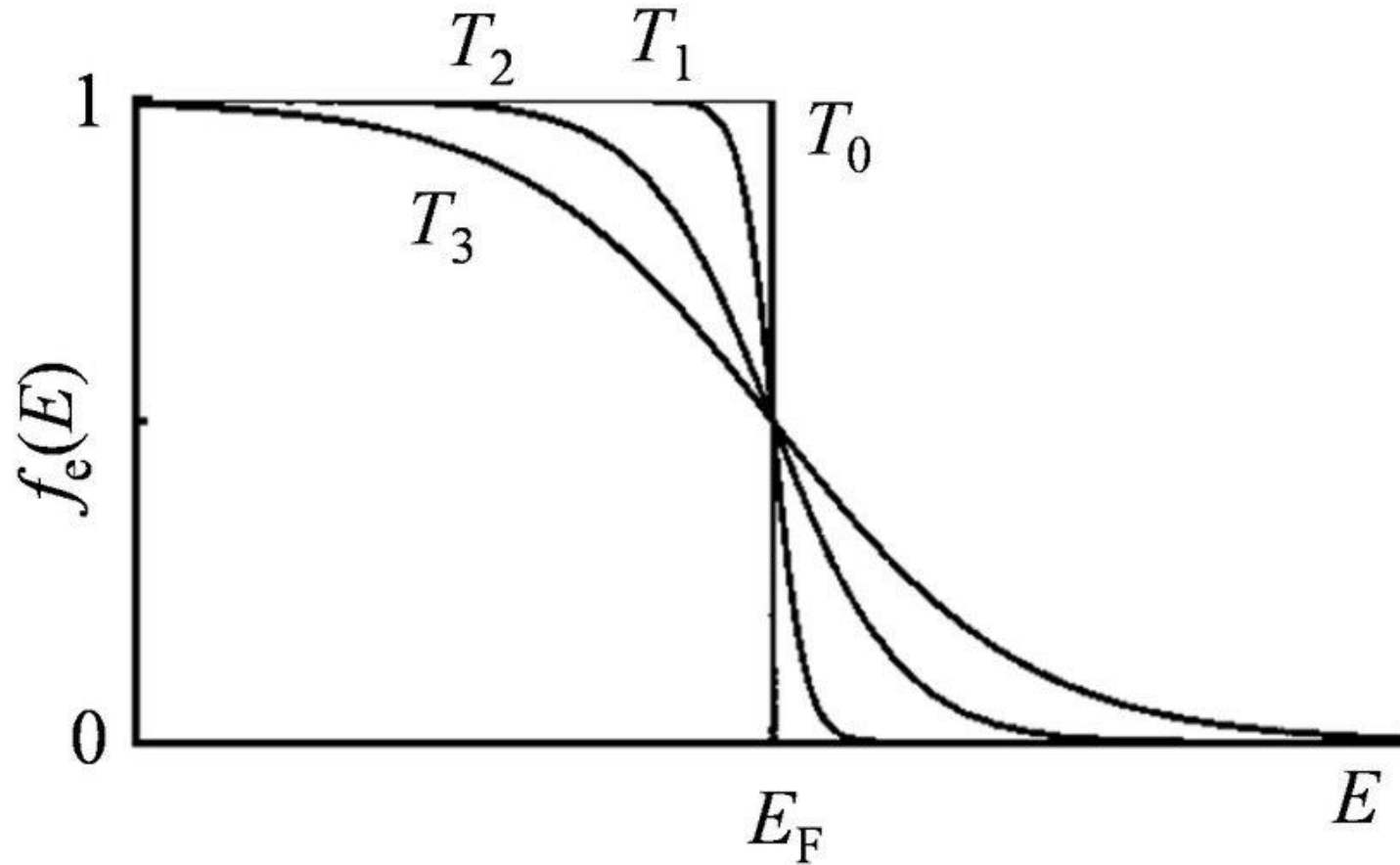
For the free
Fermi gas

$$v_F = \hbar k_F / m = 3.58 [n \cdot \text{\AA}^3]^{1/3} \cdot 10^8 \text{ cm s}^{-1}$$

$$D(\mathcal{E}_F) = \frac{3}{2} \frac{n}{\mathcal{E}_F} = 4.11 \cdot 10^{-2} [n \cdot \text{\AA}^3] \text{ eV}^{-1} \text{\AA}^{-3}$$



Fermi Distribution Function



Fermi Distribution Function

$$f(\mathcal{E}) = \frac{1}{e^{\beta(\mathcal{E}-\mu)} + 1}$$

