# EE 102 Week 2, Lecture 2 (Fall 2025)

**Instructor: Ayush Pandey** 

Date: September 10, 2025

## 1 Goals

- The timeless trio of signals: the complex exponential, the unit step, and impulse.
- The unit step function as a switch and an accumulator.
- The unit impulse function as an exciter and a sampler.
- The complex exponential signal as a sinusoid and a phasor.
- Applications of the timeless trio in real-world signal processing.

## 2 The unit step function

When we defined signals informally, we discussed how any mathematical function from your calculus class could be a signal as long as it represents something physical. Then it will not come as a surprise that for physical system applications, we would usually be interested only in positive values of time and we would want our signal to take the value of zero for all t < 0. Since we are extending the general concept of mathematical functions (which are defined for all t), it is important to have a mathematical way to write signals that are zero for t < 0. The unit step function does exactly that! Formally, we define the unit step function as

**Definition 1.** The unit step is a function defined as

$$u(t) = \begin{cases} 1, & t \ge 0 \\ 0, & t < 0 \end{cases}.$$

## 2.1 The unit step as a switch

As you can see in Figure 1, it is a discontinuous function that "steps" from 0 to 1 at t = 0. You may find the unit step function with different names: like the Heaviside function, or the ultrasensitive switch function.

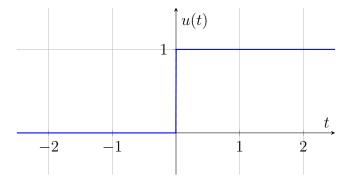


Figure 1: The unit step function u(t).

So, any function that is zero for t < 0 can be written as the product of the unit step function and another function that defines the behavior of the signal for  $t \ge 0$ . For example, if we have a signal that is zero for t < 0 and equals f(t) for  $t \ge 0$ , we can write it as x(t) = f(t)u(t).

### 2.2 Example: A sinusoidal audio wave that starts at zero time

If  $f(t) = A\sin(\omega t + \phi)$ , then we can write the signal as

$$x(t) = A\sin(\omega t + \phi)u(t).$$

This signal is zero for t < 0 and equals a sinusoidal wave for  $t \ge 0$ . The unit step function effectively "switches on" the sinusoidal wave at t = 0. But did we *really* need the step function here? You can argue that we could have defined all signals using two cases,

$$x(t) = \begin{cases} f(t), & t \ge 0\\ 0, & t < 0 \end{cases}$$

but this type of definition will quickly get cumbersome. So, yet again, we are introducing a mathematical object to make our lives easier, at least in the long run (at the moment, it may seem that we are making our life harder by learning another new function). For the specific sinusoidal example, the alternative way to define it is using a piecewise function that would need two different cases:

$$x(t) = \begin{cases} 0, & t < 0 \\ A\sin(\omega t + \phi), & t \ge 0 \end{cases}.$$

We would prefer  $x(t) = A\sin(\omega t + \phi)u(t)$  over the piecewise definition so that we can work with just a single expression.

#### 2.3 Unit step as a general switch

Beyond the switching behavior of unit step, we can also use it to define arbitrary "pulse" signals and also other piecewise continuous signals. For example, we can define a rectangular pulse of width  $\tau$  as

$$p_{\tau}(t) = u(t) - u(t - \tau).$$

This pulse is 1 for  $0 \le t < \tau$  and zero otherwise (can you prove this without relying on sketching?). It is also possible to write the equation of the pulse signal by using a time reversal:

$$p_{\tau}(t) = u(t) + u(\tau - t) - 1$$

The sketch in both cases is the same: a pulse that starts at t = 0 and ends at  $t = \tau$ . More generally, we can write a pulse that starts at  $t_1$  and ends at  $t_2$  as

$$p_{t_1,t_2}(t) = u(t-t_1) - u(t-t_2).$$

See Figure 2 for a sketch of the rectangular pulse that starts at  $t_1$  and ends at  $t_2$ .

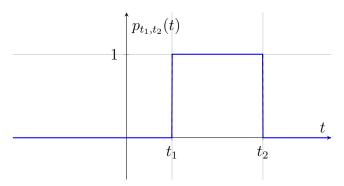


Figure 2: A rectangular pulse  $p_{t_1,t_2}(t)$  that is 1 for  $t \in [t_1,t_2)$  and 0 otherwise.

The pulse function provides us with more control over when the signal "turns on" and "turns off" — in Figure 2, it turns on at  $t = t_1$  and turns off at  $t = t_2$ . We would expect such knobs of any modern switching system!

We can also define a triangular pulse of width  $2\tau$  that is zero for all t < 0 and  $t > 2\tau$ , and peaks at  $t = \tau$ . This triangular pulse would increase linearly from 0 to 1 in the interval  $[0,\tau]$  and then decrease linearly from 1 to 0 in the interval  $[\tau, 2\tau]$ . By writing the equations of the two linear segments and then combining with unit step to decide when would each linear segment be "active", we can write the equation of the triangular pulse. Let's build it step-by-step. We write the two parts that "activate" the two linear segments using unit step functions:

$$w_1(t) = \underbrace{u(t) - u(t - \tau)}_{\text{active on } [0,\tau)}, \qquad w_2(t) = \underbrace{u(t - \tau) - u(t - 2\tau)}_{\text{active on } [\tau,2\tau)}.$$

now, we define the two linear segments:

$$r(t) = \underbrace{\frac{t}{\tau}}_{\text{linear rise}}, \qquad f(t) = \underbrace{2 - \frac{t}{\tau}}_{\text{linear fall}}.$$

Then, the triangular pulse is obtained by gating each linear segment with its corresponding window and then adding the two gated segments:

$$tr_{\tau}(t) = \underbrace{r(t)}_{\text{rise}} \underbrace{w_1(t)}_{\text{gate }[0,\tau)} + \underbrace{f(t)}_{\text{fall}} \underbrace{w_2(t)}_{\text{gate }[\tau,2\tau)}.$$

$$tr_{\tau}(t) = \frac{t}{\tau} \left[ u(t) - u(t - \tau) \right] + \left( 2 - \frac{t}{\tau} \right) \left[ u(t - \tau) - u(t - 2\tau) \right].$$

See Figure 3 for a sketch of the triangular pulse.

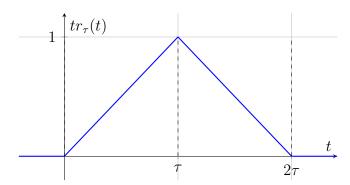


Figure 3: Triangular pulse of width  $2\tau$ : zero outside  $[0, 2\tau]$ , peak 1 at  $t = \tau$ .

The triangular pulse provides a smoother transition between the off and the on states compared to the rectangular pulse and is constructed using the unit step function as well. So, the unit step function is not just a simple switch it is also a building block for constructing more complex signals. The next "pulse" type signal that you should try is a trapezoidal pulse!

## 2.4 The unit step as an accumulator

The unit step function can also be viewed as an accumulator. If we integrate the unit step function, we obtain the ramp function (Problem 2.1 in HW #2):

$$r(t) = \int_{-\infty}^{t} u(\tau)d\tau = \begin{cases} t, & t \ge 0\\ 0, & t < 0 \end{cases}.$$

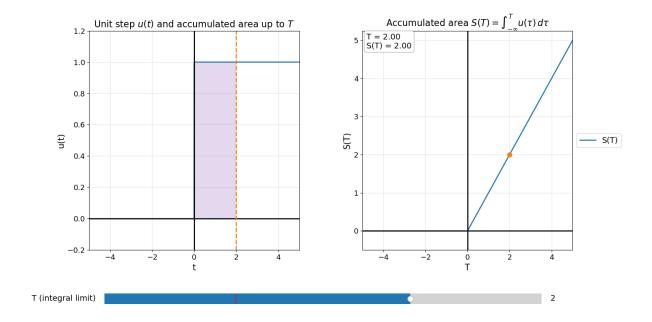


Figure 4: The unit step function u(t) (left) and its accumulated area, the ramp function r(t) (right). The slider below sets the upper limit of integration T.

This ramp function increases linearly for  $t \geq 0$  and is zero for t < 0. It effectively accumulates the area under the unit step function. See Figure 4 for a sketch of the unit step function and its accumulated area (the ramp function) on the right side. A virtual manipulative is available for you to explore this concept interactively on the course Github page<sup>1</sup>.

## 3 The unit impulse function

The unit impulse function is our second of the "timeless trio" of signals — one of the three most important signals in signal processing. In other fields, such as physics, it is also known as the Dirac delta function. In discrete mathematics, it is known as the Kronecker delta function. So, the same mysterious function has many different names! The reason is clear — it is a very useful mathematical object, while not even being a function in the traditional sense! We will define it formally later, but for now, let's understand it informally. You only need to remember two properties about the impulse function:

• It is zero at all times except at t=0, that is  $\delta(t)=0$  for all  $t\neq 0$ .

<sup>&</sup>lt;sup>1</sup>Here is the link for virtual manipulative on Github for unit step as an accumulator

• It has an area of 1 under its curve, that is,  $\int_{-\infty}^{\infty} \delta(t)dt = 1$ .

How is that possible? If a signal is zero everywhere except at one point, then there must be something unique happening at that point. To build intuition for this function, consider the rectangular pulse we defined earlier, with a slight modification. Let's define a rectangular pulse of width  $\epsilon$  and height  $\frac{1}{\epsilon}$  around the origin:

$$p_{\epsilon}(t) = \frac{1}{\epsilon} \left[ u \left( t + \frac{\epsilon}{2} \right) - u \left( t - \frac{\epsilon}{2} \right) \right].$$

A sketch of this pulse is shown in Figure 5.

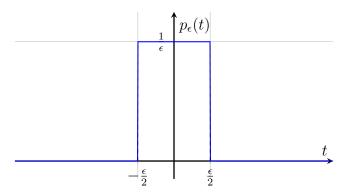


Figure 5: Rectangular pulse  $p_{\epsilon}(t) = \frac{1}{\epsilon} \left[ u \left( t + \frac{\epsilon}{2} \right) - u \left( t - \frac{\epsilon}{2} \right) \right]$  of width  $\epsilon$  centered at the origin.

As you can see in the figure, the area of the rectangle is equal to 1 for all values of  $\epsilon$ . But this does not satisfy the two properties of the impulse function listed above as there are points  $t \neq 0$  where the pulse is non-zero. So, we need to modify this pulse further. As we make  $\epsilon$  smaller, visualize how the rectangle becomes taller and narrower, while still maintaining an area of 1. You can interactively explore this concept using the virtual manipulative available on the course Github page<sup>2</sup>.

## 3.1 Impulse as the limit of a rectangular pulse

Formally, we can write the unit impulse function as the limit of the rectangular pulse as  $\epsilon$  approaches zero:

$$\delta(t) = \lim_{\epsilon \to 0} p_{\epsilon}(t) = \lim_{\epsilon \to 0} \frac{1}{\epsilon} \left[ u \left( t + \frac{\epsilon}{2} \right) - u \left( t - \frac{\epsilon}{2} \right) \right].$$

Despite the definition above, it is not possible to write a closed-form expression for the impulse function because it is not defined at t=0 and is zero everywhere else. The only

<sup>&</sup>lt;sup>2</sup>Here is the link for virtual manipulative on Github for impulse approximation using a pulse. Additionally, you can approximate an impulse using a Gaussian function, see here.

quantifiable property that we know so far is that the area under the curve of a delta function is equal to 1. To visually describe an impulse function, we draw an arrow pointing upwards at the point at which the impuse is located. That is, if we have  $\delta(t)$ , we draw an arrow at t=0; if we have  $\delta(t-t_0)$ , we draw an arrow at  $t=t_0$ . The height of the arrow is not important, but we label it with a number to indicate the area under the impulse. For example, if we have  $A\delta(t-t_0)$ , we draw an arrow at  $t=t_0$  and label it with A to indicate that the area under the impulse is equal to A.

In practice, we can never generate a true impulse function, we can only get infinitesimally close to it as we keep making the width of the rectangular pulse infinitesimally close to zero and the height of the pulse close to infinity. Even though this may sound needlessly confusing, it provides us with a very powerful mathematical tool. One example is discussed next.

#### 3.2 Impulse as a time-sampler

The impulse function can sample any test function at a specific point in time. Note that if we write  $f(t)\delta(t)$ , the product will be zero for all  $t \neq 0$  because  $\delta(t)$  is zero for all  $t \neq 0$ . The only point where the product is non-zero is at t = 0. So we can write  $f(t)\delta(t) = f(0)\delta(t)$ . This is also an impulse function located at t = 0 with an area of f(0). Now if we integrate this product over all time, we get

$$\int_{-\infty}^{\infty} f(t)\delta(t)dt = f(0)\int_{-\infty}^{\infty} \delta(t)dt = f(0) \cdot 1 = f(0).$$

This property is known as the sifting property of the impulse function. By integrating any test function multiplied by an impulse function, we can extract the value of the test function at the location of the impulse  $\rightarrow$  we have sampled that function! More generally, if we have an impulse located at  $t = t_0$ , we can write this impulse as  $\delta(t - t_0)$ . Then, by the same reasoning as above, we can show that

$$f(t_0) = \int_{-\infty}^{\infty} f(t)\delta(t - t_0) dt.$$

To prove the above, write the integral as

$$\int_{-\infty}^{\infty} f(t)\delta(t-t_0) dt = f(t_0) \int_{-\infty}^{\infty} \delta(t-t_0) dt = f(t_0) \cdot 1 = f(t_0).$$

We obtained the latter equality by observing that  $\delta(t-t_0)$  is zero at every point other than  $t=t_0$ . Since  $f(t_0)$  is independent of t, we can take it outside the integral. The remaining integral is equal to 1 because the area under the impulse function is equal to 1. This property makes the impulse function a powerful tool in signal processing and system analysis.

In formal mathematical analysis, the above is not seen as a property but is instead used to define the impulse function.

**Definition 2.** The unit impulse function  $\delta(t)$  is defined as the function for which the area under curve of its product with any test function f(t) that is continuous at t = 0, is equal to the value of the test function at the time instant at which the impulse is located  $(t = 0 \text{ for } \delta(t))$ . That is,

$$\int_{-\infty}^{\infty} f(t)\delta(t)dt = f(0). \tag{1}$$

### 3.3 Relationship between unit step and unit impulse

If you look back at the unit step function, you will notice that it is discontinuous at t = 0. So, can we define the differentiation of unit step function with time, du/dt? Generally, the answer will be no since the derivative of a function is not defined at points where the function is discontinuous. Let's try an alternate approach. Consider the following integral:

$$\int_{-\infty}^{\infty} f(t) \frac{du(t)}{dt} dt.$$

We can evaluate this integral using integration by parts to write

$$\int_{-\infty}^{\infty} f(t) \frac{du(t)}{dt} dt = f(t)u(t) \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} u(t) \frac{df(t)}{dt} dt$$
$$= f(\infty) - \int_{0}^{\infty} \frac{df(t)}{dt} dt$$
$$= f(\infty) - [f(t)]_{0}^{\infty}$$
$$= f(\infty) - f(\infty) + f(0).$$

So, we derived that

$$\int_{-\infty}^{\infty} f(t) \frac{du(t)}{dt} dt = f(0), \tag{2}$$

which is the same as the definition of the impulse function discussed earlier — the area under the product of any test function and the impulse function is equal to the value of the test function at the location of the impulse. So, we can conclude by comparing equations (1) and (2) that

$$\frac{du(t)}{dt} = \delta(t).$$

Using a similar approach, you can also show that (HW #2)

$$u(t) = \int_{-\infty}^{t} \delta(\tau) d\tau.$$

# 4 Complex exponential signals

#### Continuous time

$$x(t) = A e^{j(\omega_0 t + \phi)} = A \cos(\omega_0 t + \phi) + j A \sin(\omega_0 t + \phi).$$

Real and imaginary parts are orthogonal sinusoids. Fundamental period  $T_0 = \frac{2\pi}{\omega_0}$ .

#### Discrete time

$$x[n] = A e^{j(\Omega_0 n + \phi)}.$$

This is periodic iff  $\frac{\Omega_0}{2\pi} = \frac{M}{N}$  with integers M,N coprime. Then the fundamental period is  $N_0 = N$ . Otherwise, it is *aperiodic* on  $\mathbb{Z}$ .

### Geometric phasor

The complex exponential traces a circle of radius A in the complex plane at angular speed  $\omega_0$  (continuous) or advances by a fixed angle  $\Omega_0$  per sample (discrete). The real part is the projection on the horizontal axis and the imaginary part is the vertical projection.