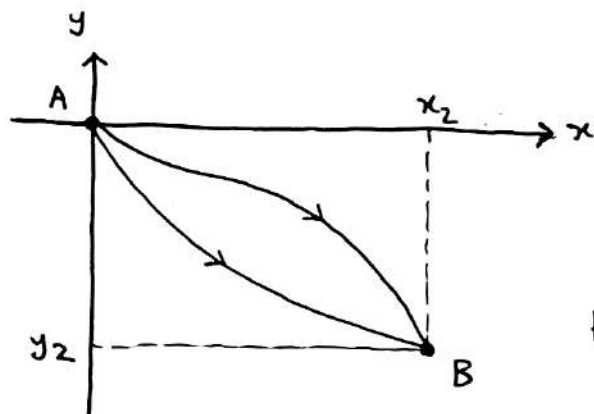


Section 0 Motivation

Example 0.1: Brachistochrone problem (Johann Bernoulli, 1696)

Particle slides on wire under influence of gravity, between 2 fixed points, A and B. What shape should the wire be for the shortest travel time, starting from rest?



Travel time

$$T = \int dt = \int_A^B \frac{dl}{v(x,y)}$$

E conserved:  $T + V = \text{const.}$

$$\frac{1}{2}mv^2 + mgy = mgy_1 = 0 \quad \text{at A: } y_1 = 0$$

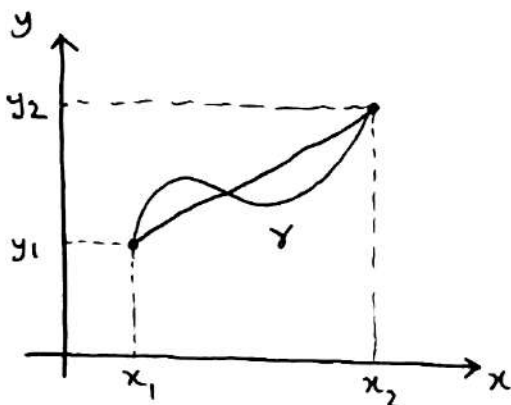
$$\text{so } v = \sqrt{2g} \sqrt{-y}$$

$$\text{So minimise } T(y) = \frac{1}{\sqrt{2g}} \int_0^{x_2} \frac{\sqrt{1+(y')^2}}{\sqrt{-y}} dx \quad (\text{line elt. } dl)$$

subject to  $y(0) = 0, \quad y(x_2) = y_2.$

Example 0.2: Geodesics (shortest path  $\gamma$  between 2 points on a surface  $\Sigma$ , if they exist)

Take  $\Sigma = \mathbb{R}^2$  (a plane, Pythagorean thm holds)



Distance along  $\gamma$ :

$$D(y) = \int_A^B dl = \int_{x_1}^{x_2} \sqrt{1+(y')^2} dx$$

Minimise  $D$  by varying path.

In general we want to minimise / maximise some function

$$F(y) = \int_{x_1}^{x_2} f(x, y(x), y'(x)) dx \quad (0.1)$$

among all functions with  $y(x_1) = y_1, y(x_2) = y_2$

Note  $y = y(x)$ .

(0.1) is a functional (a function on the space of functions).

Functions: numbers  $\rightarrow$  numbers

Functionals: functions  $\rightarrow$  numbers (example: area under graph)

Calculus of variations: finding extrema of functionals on spaces of functions.

Notation  $C(\mathbb{R})$  space of continuous functions on  $\mathbb{R}$

$C^k(\mathbb{R})$  space of <sup>functions on  $\mathbb{R}$  with</sup> continuous  $k^{\text{th}}$ -derivatives

$C^k_{(\alpha, \beta)}(\mathbb{R})$  space of continuous  $k^{\text{th}}$ -derivatives s.t.  $f(\alpha) = f(\beta)$   
<sup>functions on  $\wedge$   
 $\mathbb{R}$  with</sup>

Need to specify the function space beforehand.

Example 0.3 Fermat's principle

Light between two points travels along paths which require the least time.

Example 0.4 Principle of least (stationary) action

$$S(\gamma) = \int_{t_1}^{t_2} (T - V) dt \quad \text{for motion of particle}$$

(e.g.  $m\ddot{x} = -\nabla V$  so Newton's eqns should follow)

Variational Principles - Lecture 2Section 1 Calculus for functions on  $\mathbb{R}^n$ 

$$f \in C^2(\mathbb{R}^n), \quad f: \mathbb{R}^n \rightarrow \mathbb{R} \quad (\text{cts 2nd derivatives})$$

The point  $\underline{a} \in \mathbb{R}^n$  is stationary if  $\nabla f(\underline{a}) = \underline{0}$

Expand near  $\underline{x} = \underline{a}$ :

$$f(\underline{x}) = f(\underline{a}) + \underbrace{(\underline{x} - \underline{a}) \cdot \nabla f|_{\underline{a}}}_{\substack{0 \text{ as } \underline{a} \text{ stationary} \\ \nabla f|_{\underline{a}} = \underline{0}}} + \frac{1}{2} (x_i - a_i)(x_j - a_j) \partial^2_{ij} f|_{\underline{a}} + O(|\underline{x} - \underline{a}|^2)$$

The Hessian matrix is given by  $H_{ij} = \partial_i \partial_j f = H_{ji}$

Shift origin to set  $\underline{a} = \underline{0}$ . Diagonalise  $H(\underline{0})$  by an orthogonal transformation.

$$H' = R^T H(\underline{0}) R = \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_n \end{pmatrix}$$

$$f(\underline{x}') - f(\underline{0}) = \frac{1}{2} \sum_i \lambda_i (x'_i)^2 + O(|\underline{x}'|^2)$$

(note 1st order term is 0 as  $\underline{0}$  is stationary)

(i) If all  $\lambda_i > 0$  then  $f(\underline{x}') > f(\underline{0})$  in all directions  
local minimum

(ii) all  $\lambda_i < 0$  : local maximum

(iii) some  $\lambda_i > 0$ , some  $\lambda_i < 0$  : saddle point

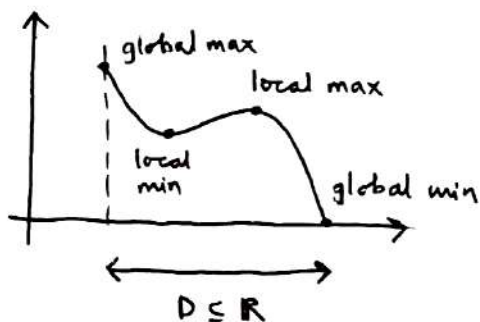
(iv) some  $\lambda_i = 0$  : need higher order terms.

Special case  $n=2$  (may be easier than finding evals)

$$\det(H) = \lambda_1 \lambda_2 \quad \text{tr}(H) = \lambda_1 + \lambda_2$$

- $\det H > 0$  and  $\text{tr} H > 0$  : local minimum
- $\det H > 0$  and  $\text{tr} H < 0$  : local maximum
- $\det H < 0$  : saddle point
- $\det H = 0$  : higher order derivatives

Remarks If we have  $f: D \rightarrow \mathbb{R}$  where  $D \subseteq \mathbb{R}^n$  then may have extremum on boundary of domain



Here we cannot find the global max/min via derivatives.

Remark 2 If  $f$  is harmonic on  $\mathbb{R}^2$  :  $f_{xx} + f_{yy} = 0$   
with  $f$  defined on  $D \subseteq \mathbb{R}^2$

Then anywhere in  $D$ ,  $\text{tr} H = 0$  so critical points must be saddle points and the min/max is on the boundary (see reason above).

Example 1.1  $f(x, y) = x^3 + y^3 - 3xy$

$$\nabla f = (3x^2 - 3y, 3y^2 - 3x) = (0, 0) \text{ at critical pt}$$

$$x^2 - y = 0, \quad y^2 - x = 0 \Rightarrow y^4 = y$$

So critical points are  $(0, 0)$  and  $(1, 1)$ .

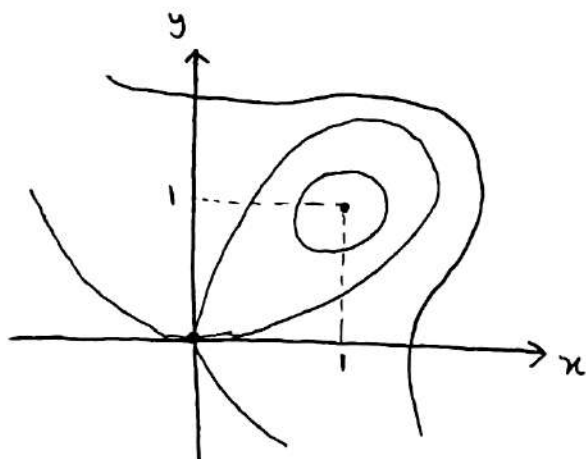
$$H = \begin{pmatrix} 6x & -3 \\ -3 & 6y \end{pmatrix}$$

At  $(0, 0)$   $\det H = -9 < 0$   $\text{tr } H = 0$

so have a saddle point, where  $f = 0$ .

At  $(1, 1)$   $\det H = 27 > 0$   $\text{tr } H = 12 > 0$

so have a local minimum, where  $f = -1$ .

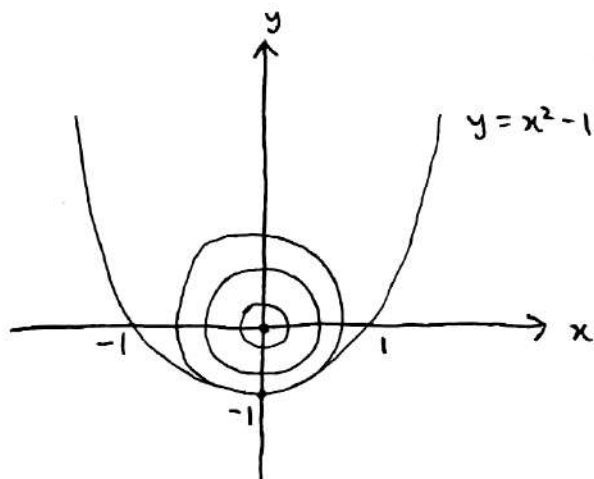


Near  $(0, 0)$   $f \approx -3xy$   
as  $x^3, y^3$  small  
so decreases on  $y = x$   
and increases on  $y = -x$   
(no global min/max)

(only a rough sketch)

### Section 1.1 Constraints and Lagrange multipliers

Example 1.2 Find the circle centred at  $(0, 0)$  with smallest radius, which intersects the parabola  $y = x^2 - 1$ .



Direct method: solve the constraints

Minimise  $x^2 + y^2$  subject to  $y = x^2 - 1$

so minimise

$$x^2 + (x^2 - 1)^2 = x^4 - x^2 + 1$$

$$\Rightarrow 4x^3 - 2x = 0$$

2 solns:  $x = \pm \frac{1}{\sqrt{2}}, y = -\frac{1}{2} \Rightarrow$  radius is  $\frac{\sqrt{3}}{2}$

$x = 0, y = -1 \Rightarrow$  radius is 1

so  $\frac{\sqrt{3}}{2}$  is smallest.

But what if we can't solve analytically?

### Lagrange Multipliers

Define new function  $h(x, y, \lambda) = f(x, y) - \lambda g(x, y)$

where  $f$  is the function we want to min/max

and  $g(x, y) = 0$  is the constraint.  $\lambda$  is Lagrange multiplier.

Here  $h = x^2 + y^2 - \lambda(y - x^2 + 1)$

Then extremise over 3 variables with no constraint.

$$\frac{\partial h}{\partial x} = 2x + 2\lambda x = 0 \quad \frac{\partial h}{\partial y} = 2y - \lambda = 0 \quad \frac{\partial h}{\partial \lambda} = y - x^2 + 1 = 0$$

(constraint)

Combining first 2:  $2x + 4xy = 0 \Rightarrow x = 0$  or  $y = -\frac{1}{2}$

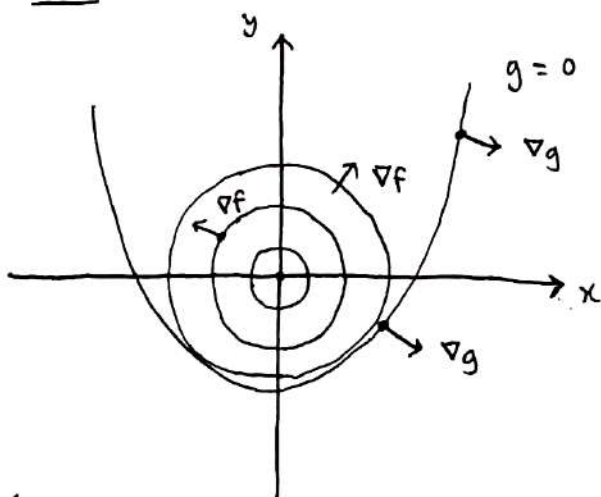
so using  $\frac{\partial h}{\partial \lambda}$  have either  $x = 0, y = -1$   
or  $x = \pm \frac{1}{\sqrt{2}}, y = -\frac{1}{2}$  as before.

$(0, -1) \rightarrow f = 1 \quad (\lambda = -2)$

$(\pm \frac{1}{\sqrt{2}}, -\frac{1}{2}) \rightarrow f = \frac{3}{4} \quad (\lambda = -1)$

Finding the critical points without solving analytically.

Why does this work?



$\nabla g$  perpendicular to  
"surface"  $g = 0$ .



Representing circles by  $f = \text{constant}$ ,

have also  $\nabla f$  perpendicular to  $f = \text{constant}$ .

At the extremum,  $\nabla f$  and  $\nabla g$  are parallel

$$\text{i.e. } \nabla f = \lambda \nabla g \quad \text{i.e. } \underline{\nabla(f - \lambda g) = \underline{0}}$$

so we consider the extrema of

$$\underline{h = f - \lambda g} \quad \text{to find the solution(s).}$$

For multiple constraints:

Extremise  $f: \mathbb{R}^n \rightarrow \mathbb{R}$  subject to  $g_\alpha(\underline{x}) = 0$   
( $g_\alpha: \mathbb{R}^n \rightarrow \mathbb{R}, \alpha = 1, \dots, k$ )

$$\text{Then define } \underline{h(x_1, \dots, x_n, \lambda_1, \dots, \lambda_k) = f - \sum_{\alpha=1}^k \lambda_\alpha g_\alpha}$$

a function of  $n+k$  vars, with  $k$  Lagrange multipliers.

So we work with

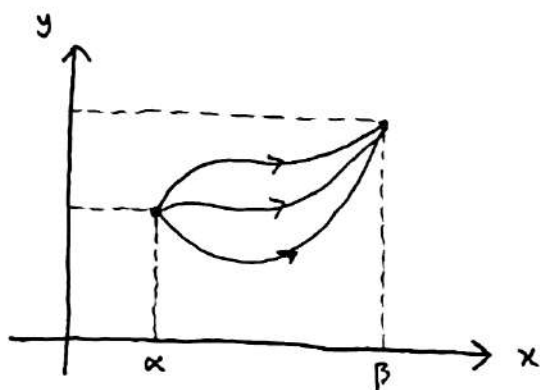
$$\frac{\partial h}{\partial x_i} = 0, \quad \frac{\partial h}{\partial \lambda_\alpha} = 0, \quad \text{eliminate } \lambda_\alpha \text{ and solve for } \underline{x}.$$

The method works also if the constraints cannot be eliminated algebraically.

Section 2 Euler-Lagrange equations

Extremise functional 0.1

$$F(y) = \int_{\alpha}^{\beta} f(x, y, y') dx \quad (2.1)$$



$f$  is given,  $\alpha, \beta$  are fixed  
Functional depends on  $y$ .

Small perturbation  $y \rightarrow y + \varepsilon \eta(x)$  in (2.1)

Compute  $F(y + \varepsilon \eta(x))$  with  $\eta(\alpha) = \eta(\beta) = 0$

(so perturbed function goes through  
same fixed pts)

Lemma If  $g: [\alpha, \beta] \rightarrow \mathbb{R}$  is continuous on  $[\alpha, \beta]$  and  
 $\int_{\alpha}^{\beta} g(x) \eta(x) dx = 0$  for all continuous  $\eta(x)$  on  $[\alpha, \beta]$   
s.t.  $\eta(\alpha) = \eta(\beta) = 0$ , then  $g(x) = 0 \quad \forall x \in [\alpha, \beta]$ .

Proof Suppose  $\exists \bar{x} \in (\alpha, \beta)$  s.t.  $g(\bar{x}) \neq 0$ . WLOG  
suppose that  $g(\bar{x}) > 0$ . Then  $\exists$  interval  $[x_1, x_2] \subseteq (\alpha, \beta)$   
s.t.  $g(x) > c$  on  $[x_1, x_2]$  for some  $c > 0$ .

$$\text{Set } \eta(x) = \begin{cases} (x-x_1)(x_2-x) & x \in [x_1, x_2] \\ 0 & \text{otherwise} \end{cases} \quad (2.2)$$

Then  $\eta(x)$  is continuous and we see  $\int_{x_1}^{x_2} g(x) \eta(x) dx > 0$   
so  $\int_{\alpha}^{\beta} g(x) \eta(x) dx > 0$  (as  $\eta$  is 0 elsewhere). ✗

(This was on 1A Analysis I ES4).



Remark  $\eta$  is a bump function (in 2.2)

A general form for  $C^k$  bump functions ( $x \in [x_1, x_2]$   
bump interval)

$$\text{is } \eta = \begin{cases} ((x-x_1)(x_2-x))^{k+1} & x \in [x_1, x_2] \\ 0 & \text{otherwise} \end{cases}$$

Now back to 2.1. Expand in  $\epsilon$ :

$$F(y + \epsilon \eta) = \int_{\alpha}^{\beta} f(x, y + \epsilon \eta, y' + \epsilon \eta') dx$$

$$= F(y) + \epsilon \int_{\alpha}^{\beta} \left( \frac{\partial f}{\partial y} \eta + \frac{\partial f}{\partial y'} \eta' \right) dx + o(\epsilon^2)$$

(Note: we get this by expanding the  $y + \epsilon \eta$  and  $y' + \epsilon \eta'$  as these are the only dependencies of  $f$  that depend on  $\epsilon$ .)

Can check this works and the  $F(y)$  comes from the 1st order term)

At extremum, have  $F(y + \epsilon \eta) = F(y) + o(\epsilon^2)$

$$\text{i.e. } \frac{dF}{d\epsilon} \Big|_{\epsilon=0} = 0 \quad (\text{1st derivative term must be 0 at extremum}).$$

Integrating by parts on the  $\epsilon$ -term, we want

$$\begin{aligned} 0 &= \int_{\alpha}^{\beta} \left( \frac{df}{dy} \eta - \frac{d}{dx} \left( \frac{\partial f}{\partial y'} \right) \eta \right) dx + \left[ \frac{\partial f}{\partial y'} \eta \right]_{\alpha}^{\beta} \\ &= \int_{\alpha}^{\beta} \left( \frac{\partial f}{\partial y} - \frac{d}{dx} \left( \frac{\partial f}{\partial y'} \right) \right) \eta dx \end{aligned}$$

$(\eta(\alpha) = \eta(\beta) = 0)$

apply lemma with this  $\underline{0}$ , as  $\eta(\alpha) = \eta(\beta) = 0$ .

$$\text{So } \boxed{\frac{d}{dx} \left( \frac{\partial f}{\partial y'} \right) = \frac{\partial f}{\partial y}}$$

(2.3)

(necessary condition for extremum)

Euler-Lagrange equation

### Remarks

- (2.3) is a 2nd order ODE for  $y(x)$  with boundary conditions  $y(\alpha) = y_1$ ,  $y(\beta) = y_2$ .
- $\frac{d}{dx} \left( \frac{\partial f}{\partial y'} \right) - \frac{\partial f}{\partial y}$  is called a functional derivative  
denoted  $\frac{\delta F(y)}{\delta y(x)}$
- Sometimes  $\delta y = \varepsilon \eta(x)$  is written, called a "small variation" and they write  $F(y + \delta y) = F(y) + \delta F(y)$  where  
$$\delta F = \int_{\alpha}^{\beta} \left( \frac{\delta F(y)}{\delta y(x)} \delta y(x) \right) dx.$$
- Other boundary conditions are possible e.g.  $\left. \frac{\partial f}{\partial y'} \right|_{\alpha, \beta} = 0$
- We consider  $x, y, y'$  to be independent vars when taking partial derivatives.  
Total derivative: 
$$\frac{dh}{dx} = \frac{\partial h}{\partial x} + \frac{\partial h}{\partial y} \frac{dy}{dx} + \frac{\partial h}{\partial y'} \frac{d^2 y}{dx^2}$$

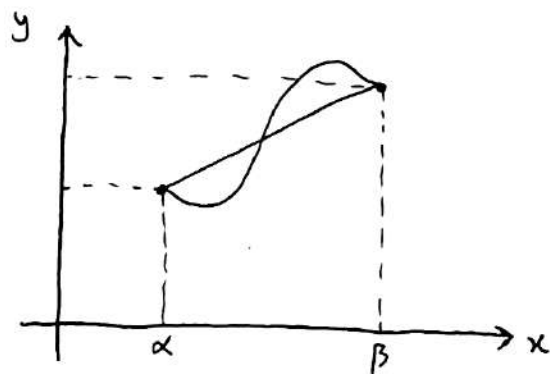
### Section 2.1 First integrals of the Euler - Lagrange equation

In some cases 2.3 (E-L eqn) can be integrated once to get a 1st order ODE - "first integral"

Cases:

- (a)  $f$  does not explicitly depend on  $y$ , so  $\frac{\partial f}{\partial y} = 0$   
Then we get  $\frac{d}{dx} \left( \frac{\partial f}{\partial y'} \right) = 0$  i.e.  $\frac{\partial f}{\partial y'} = c$  (constant)
- 
- (2.4)

Example 2.1 Geodesics on Euclidean plane



$$F(y) = \int_{\alpha}^{\beta} \sqrt{dx^2 + dy^2}$$

↓  
arc length functional

$$\text{so } F(y) = \int_{\alpha}^{\beta} \underbrace{\sqrt{1 + (y')^2}}_{f(y')} dx$$

$f(y')$  doesn't depend on  $y$  explicitly

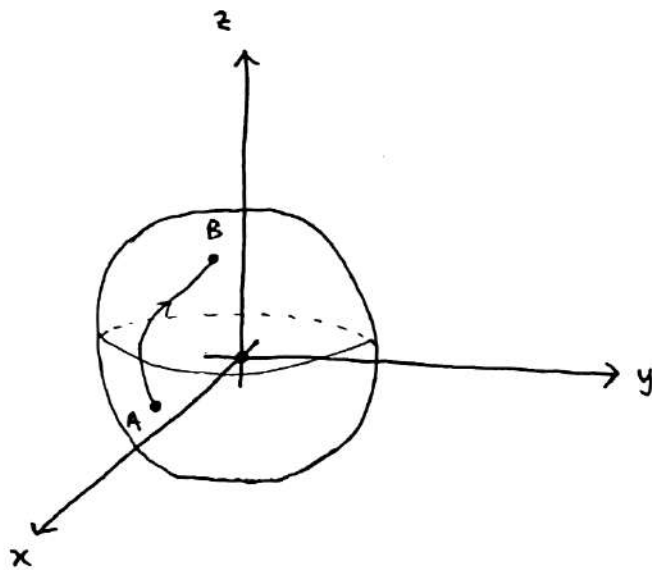
So use 2.4:  $\frac{\partial f}{\partial y} = 0$  so  $\frac{y'}{\sqrt{1 + (y')^2}} = \text{constant}$

so  $(y')^2 = c + c(y')^2$  for some  $c$

$$\Rightarrow (y')^2 = \frac{c}{1-c} \Rightarrow y' = \pm \sqrt{\frac{c}{1-c}}$$

So  $y'$  must be constant so have  $y' = m$  for some  $m$   
 $\Rightarrow \underline{y = mx + k}$  a straight line.

Example 2.2 Geodesics on a sphere  $S^2 \subseteq \mathbb{R}^3$



Parametrise sphere:

$$x = \sin \theta \sin \phi$$

$$y = \sin \theta \cos \phi$$

$$z = \cos \theta$$

$$0 \leq \theta \leq \pi, \quad 0 \leq \phi \leq 2\pi$$

take as unit sphere

$$ds^2 = dx^2 + dy^2 + dz^2 = d\theta^2 + \sin^2 \theta d\phi^2$$

Then we parametrise as  $\phi = \phi(\theta)$  and can write

$$ds = \sqrt{1 + \sin^2 \theta (\phi')^2} d\theta$$

$$\text{Then } F(\phi) = \int_{\theta_1 = \alpha}^{\theta_2 = \beta} \sqrt{1 + \sin^2 \theta (\phi')^2} d\theta$$

Integrand  $f = f(\theta, \phi')$  doesn't depend on  $\phi$  itself

$$\frac{\partial f}{\partial \phi} = 0 \quad \Rightarrow \quad \frac{\partial f}{\partial \phi'} = K \quad (\text{constant})$$

Evaluating  $\frac{\partial f}{\partial \phi'}$  we have

$$\frac{\phi' \sin^2 \theta}{\sqrt{1 + \sin^2 \theta (\phi')^2}} = K$$

We square and solve for  $(\phi')^2$ :

$$(\phi')^2 = \frac{K^2}{\sin^2 \theta (\sin^2 \theta - K^2)}$$

$$\text{So } \phi = \pm \int \frac{K}{\sin \theta \sqrt{\sin^2 \theta - K^2}} d\theta$$

(2 solns, each going one way around sphere)

Make substitution  $u = \cot \theta$

$$\text{Then get } \pm \frac{\sqrt{1-K^2}}{K} \cos(\phi - \phi_0) = \cot \theta$$

a great circle. (Geodesics are segments of great circles)

(b) Consider for general  $f(x, y, y')$ :

$$\begin{aligned} \frac{d}{dx} \left( f - y' \frac{\partial f}{\partial y'} \right) &= \frac{\partial f}{\partial x} + y' \frac{\partial f}{\partial y} + \cancel{y'' \frac{\partial f}{\partial y'}} \\ &\quad - \cancel{y'' \frac{\partial f}{\partial y'}} - y' \frac{d}{dx} \left( \frac{\partial f}{\partial y'} \right) \\ &= y' \left( \cancel{\frac{\partial f}{\partial y}} - \frac{d}{dx} \frac{\partial f}{\partial y'} \right) + \frac{\partial f}{\partial x} = \frac{\partial f}{\partial x} \end{aligned}$$

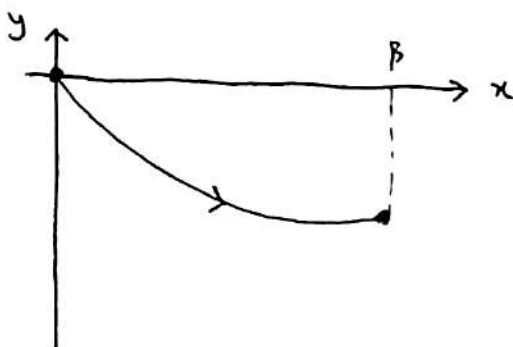
(by E-L)

So if  $f$  does not explicitly depend on  $x$ , then  $\frac{\partial f}{\partial x} = 0$   
 so we have  $\frac{d}{dx} \left( f - y' \frac{\partial f}{\partial y'} \right) = 0$

$$\text{so } \underline{f - y' \frac{\partial f}{\partial y'} = c} \quad (2.5)$$

and we can reduce the order to get a 1st order ODE.

Example 2.3 Brachistochrone problem





Recall from section 0 that the functional is

$$F(y) = \frac{1}{\sqrt{2g}} \int_0^B \frac{\sqrt{1+(y')^2}}{\sqrt{-y}} dx$$

depends on  $y, y'$  but not  $x$ : use 2.5.

From  $f - y' \frac{\partial f}{\partial y'} = c$  we get

$$\frac{\sqrt{1+(y')^2}}{\sqrt{-y}} - y' \frac{y'}{\sqrt{1+(y')^2} \sqrt{-y}} = c$$

Rearranging gives  $\frac{1}{\sqrt{1+(y')^2}} = c \sqrt{-y}$

$$\Rightarrow y' = \pm \frac{\sqrt{1+c^2 y^2}}{c \sqrt{-y}} \Rightarrow x = \pm c \int \frac{\sqrt{-y}}{\sqrt{1+c^2 y^2}} dy$$

Set  $y = -\frac{1}{c^2} \sin^2 \frac{\theta}{2}$ ,  $dy = -\frac{1}{c^2} \sin \frac{\theta}{2} \cos \frac{\theta}{2} d\theta$

Then have 
$$x = \pm c \int (-1) \frac{1}{c^2} \frac{\sin^2 \frac{\theta}{2} \cos \frac{\theta}{2}}{\sqrt{1 - \sin^2 \frac{\theta}{2}}} d\theta$$

$$= \mp \frac{1}{2c^2} \int (1 - \cos \theta) d\theta$$

$$= \mp \frac{1}{2c^2} (\theta - \sin \theta) \quad (\text{curve goes through } (0,0) \text{ so constant of integration is } 0)$$

Then we get 
$$x = \frac{\theta - \sin \theta}{2c^2}, \quad y = -\frac{1}{c^2} \sin^2 \frac{\theta}{2}$$

---

The parametrisation of a cycloid (curve traced out by a point on a circle as it rolls without slipping)

## Section 2.2 Fermat's principle

Light/sound travels along the path between two points that takes the least time.

Ray: path  $y = y(x)$  speed  $c(x, y)$

$$F(y) = \int \frac{dl}{c} = \int_{\alpha}^{\beta} \frac{\sqrt{1+(y')^2}}{c(x, y)} dx$$

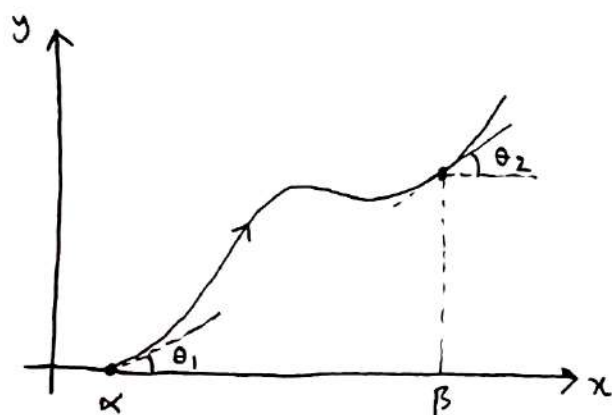
Assume  $c = c(x)$  only. Then  $\frac{\partial f}{\partial y} = 0 \Rightarrow \frac{\partial f}{\partial y'} = \text{constant}$

$$\frac{y'}{\sqrt{1+(y')^2} c(x)} = \text{constant}$$

$\tan \theta = y'$  so making  
this sub, we get

$$\frac{\sin \theta_1}{c(x_1)} = \frac{\sin \theta}{c(x)} \quad (2.6)$$

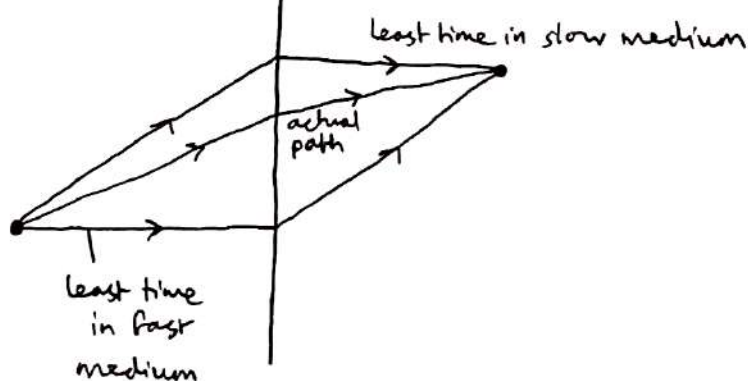
Snell's law in optics



Fast ( $c_F$  constant)

Slow ( $c_S$  constant)

$$c_S < c_F$$



### Section 3 Extensions of the Euler-Lagrange equations

#### 3.1 Euler-Lagrange equations with constraints

Extremise  $F(y) = \int_{\alpha}^{\beta} f(x, y, y') dx$  subject to

$$G(y) = \int_{\alpha}^{\beta} g(x, y, y') dx = k \quad (\text{constant}).$$

Use Lagrange multipliers: extremise

$$\underline{\Phi(y; \lambda) = F(y) - \lambda G(y)}$$

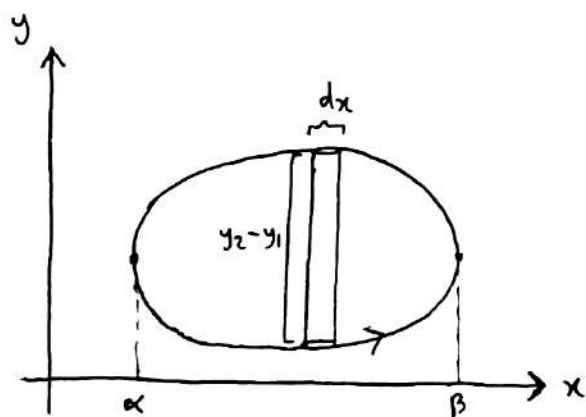
Replace  $f$  in E-L eqn  
by  $f - \lambda g$ :

$$\frac{d}{dx} \left( \frac{\partial}{\partial y'} (f - \lambda g) \right) = \frac{\partial}{\partial y} (f - \lambda g) = 0 \quad (3.1)$$

#### Example 3.1 Dido problem / Isoperimetric problem

What simple, closed plane curve of fixed length  $L$  maximises the enclosed area?

WLOG we assume convexity.



Find area functional:

$x$  monotonically increases  
from  $\alpha$  to  $\beta$ , decreases  
from  $\beta$  to  $\alpha$ . (not words)

Given  $x$ , there exists  $(y_1, y_2)$   
on the curve with  $y_1(x) = y_2(x)$   
and  $y_2 > y_1$

$$\text{so } dA = [y(x)]_{x_1}^{x_2} dx$$

Area functional

$$A(y) = \int_{\alpha}^{\beta} (y_2(x) - y_1(x)) dx = \oint_C y(x) \frac{dx}{dy_1}$$

Constraint

$$L(y) = \oint dl = \oint_C \sqrt{1+(y')^2} dx = L$$

Let  $\underline{h = y - \lambda \sqrt{1+(y')^2}}$  (Lagrange multiplier function)

Using  $\frac{d}{dx} \left( \frac{\partial h}{\partial y'} \right) = \frac{\partial h}{\partial y}$

note  $\frac{\partial h}{\partial x} = 0$  so can use

first integral of E-L.  
(see L4)

So  $h - y' \frac{\partial h}{\partial y'} = k$  for some constant  $k$

Writing this out explicitly, we get

$$k = y - \frac{\lambda}{\sqrt{1+(y')^2}} \Rightarrow \underline{(y')^2 = \frac{\lambda^2}{(y-k)^2} - 1}$$

Solving this gives  $(x-x_0)^2 + (y-y_0)^2 = \lambda^2$

And circumference  $2\pi\lambda = L \Rightarrow \lambda = \frac{L}{2\pi}$

This gives a circle of radius  $\frac{L}{2\pi}$ .

Example 3.2 The Sturm-Liouville problem

Let function  $\rho(x) > 0$  for  $x \in [\alpha, \beta]$ . Define also  $\sigma = \sigma(x)$ .

Then define

$$F(y) = \int_{\alpha}^{\beta} \left( \rho \cdot (y')^2 + \sigma \cdot (y^2) \right) dx$$

and extremise  $F$  subject to

$$G(y) = \int_{\alpha}^{\beta} y^2 dx = 1.$$

$G(y) = 1$   
rewrite 1 using  
an integrand:

$$\mathcal{E}(y; \lambda) = F(y) - \lambda (G(y) - 1)$$

$$\int_{\alpha}^{\beta} \frac{1}{\beta - \alpha} dx = 1$$

Write  $h = p \cdot (y')^2 + \sigma \cdot (y^2) - \lambda \left( y^2 - \frac{1}{\beta - \alpha} \right)$

$$\frac{\partial h}{\partial y'} = 2p y', \quad \frac{\partial h}{\partial y} = 2\sigma y - 2\lambda y$$

Applying E-L, rearranging gives

$$-\frac{d}{dx}(p y') + \sigma y = \lambda y \quad (3.2)$$

$\mathcal{L}(y)$

$\mathcal{L}$  is called the Sturm-Liouville operator

Then  $\mathcal{L}(y) = \lambda y$  is an eigenvalue problem  
(note: similar form to TISE)

Note: if  $\sigma > 0$  then  $F(y) > 0$ .

Claim The smallest ~~positive~~ eigenvalue is equal to the positive minimum.

Proof Take (3.2)  $\times y$  and IBP from  $\alpha$  to  $\beta$ :

$$F(y) - \underbrace{\left[ y \cdot y' p \right]_{\alpha}^{\beta}}_0 = \lambda \underbrace{G(y)}_1 = \lambda$$

(boundary term is 0 -  
fixed end problem)

So lowest eigenvalue = minimum of  $\frac{F(y)}{G(y)}$ .

□



### 3.2 Several dependent variables

$$\underline{y}(x) = (y_1(x), y_2(x), \dots, y_n(x)) \quad \text{Extremise}$$

$$F(\underline{y}) = \int_{\alpha}^{\beta} f(x, y_1, \dots, y_n, y_1', \dots, y_n') \, dx$$

Perturbation  $y_i \rightarrow y_i(x) + \varepsilon \eta_i(x)$

with  $i = 1, \dots, n$  and  $\eta_i(\alpha) = \eta_i(\beta) = 0$ .

Using same derivation as that of E-L eqn, we get

$$F(\underline{y} + \varepsilon \underline{\eta}) - F(\underline{y}) = \int_{\alpha}^{\beta} \sum_{i=1}^n \eta_i \left( \frac{d}{dx} \left( \frac{\partial f}{\partial y_i'} \right) - \frac{\partial f}{\partial y_i} \right) dx \\ + \text{boundary terms} + O(\varepsilon^2)$$

Recall the "fundamental lemma" from L3: by setting all the  $\eta_i$ 's but one to zero in turn, we get again

$$\underline{\frac{d}{dx} \left( \frac{\partial f}{\partial y_i'} \right) = \frac{\partial f}{\partial y_i}} \quad (3.3)$$

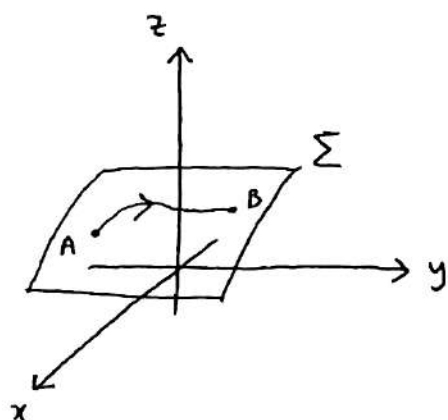
a system of  $n$  2<sup>nd</sup> order ODEs.

#### First integrals of 3.3

\* if  $\frac{\partial f}{\partial y_j} = 0$  for some  $1 \leq j \leq n$  then  $\frac{\partial f}{\partial y_j'} = \text{constant}$

\* if  $\frac{\partial f}{\partial x} = 0$  then  $f - \sum_i y_i' \frac{\partial f}{\partial y_i'} = \text{constant}$

Example Geodesics on surfaces



$\Sigma \subseteq \mathbb{R}^3$  is a surface given by  $g(x, y, z) = 0$ .

Find shortest path on surface between 2 points, if one exists.

Take  $t$  to be a parameter on the curve:

$$A = \underline{x}(0), \quad B = \underline{x}(1) \quad \underline{x} = (x, y, z) = \underline{x}(t)$$

$$\Phi(\underline{x}, \lambda) = \int_0^1 \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} - \lambda(t) g(x, y, z) dt$$

(Note  $\lambda = \lambda(t)$  as we need the curve to lie on the surface i.e.  $g(x, y, z) = 0$  everywhere - if don't have  $\lambda = \lambda(t)$  then we are just saying  $g(x, y, z)$  integrates to 0 over curve, not necessarily  $g(x, y, z) = 0$ )

Write integrand  $\sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} - \lambda(t) g = \underline{h(x, y, z, \dot{x}, \dot{y}, \dot{z}, \lambda)}$

Use E-L with  $h$

• Variation wrt  $\lambda$  :  $\frac{d}{dt} \left( \frac{\partial h}{\partial \dot{\lambda}} \right) = \frac{\partial h}{\partial \lambda} \Rightarrow \frac{g(x, y, z) = 0}{\forall t}$   
 0 ( $h$  does not depend on  $\dot{\lambda}$ )

Variation wrt  $x_i$  ( $x, y, z$ ) :

$$\frac{d}{dt} \left( \frac{\dot{x}_i}{\sqrt{\dot{x}_1^2 + \dot{x}_2^2 + \dot{x}_3^2}} \right) + \lambda \frac{\partial g}{\partial x_i} = 0 \quad i=1,2,3$$

a system of ODEs to be solved.

Alternatively, solve the constraint  $g=0$ , as we did in Ex 2.2 (geodesics on sphere).

### 3.3 Several independent variables

In general with  $\Phi: \mathbb{R}^n \rightarrow \mathbb{R}^m$  : for  $n > 1$ , E-L eqns become PDEs.

Consider case where  $n=3$ ,  $m=1$  so  $\Phi: \mathbb{R}^3 \rightarrow \mathbb{R}$

$$F(\phi) = \iiint_D f(x, y, z, \phi, \phi_x, \phi_y, \phi_z) dx dy dz$$

(note  $\phi_x := \frac{\partial \phi}{\partial x}$ )

This is a volume integral over a domain  $D \subseteq \mathbb{R}^3$

Assume there is an extremum, consider perturbation

$$\phi \rightarrow \phi(x, y, z) + \varepsilon \eta(x, y, z) \quad \text{where } \underline{\eta = 0 \text{ on } \partial D}$$

$$\begin{aligned} F(\phi + \varepsilon \eta) - F(\phi) &= \varepsilon \int_D \left( \eta \frac{\partial f}{\partial \phi} + \eta_x \frac{\partial f}{\partial \phi_x} + \eta_y \frac{\partial f}{\partial \phi_y} + \eta_z \frac{\partial f}{\partial \phi_z} \right) dx dy dz \\ &\quad + o(\varepsilon^2) \\ &= \varepsilon \int_D \left( \eta \frac{\partial f}{\partial \phi} + \underbrace{\nabla \cdot \left( \eta \left( \frac{\partial f}{\partial \phi_x}, \frac{\partial f}{\partial \phi_y}, \frac{\partial f}{\partial \phi_z} \right) \right)}_{\substack{\text{use div. thm: } \eta = 0 \text{ on } \partial D \\ \text{so this becomes } 0}} - \eta \nabla \cdot \left( \frac{\partial f}{\partial \phi_x}, \frac{\partial f}{\partial \phi_y}, \frac{\partial f}{\partial \phi_z} \right) \right) dx dy dz \\ &\quad + o(\varepsilon^2) \end{aligned}$$

Then we get

$$F(\phi + \varepsilon \eta) - F(\phi) = \varepsilon \int \eta \left( \frac{\partial f}{\partial \phi} - \nabla \cdot \left( \frac{\partial f}{\partial \phi_x}, \frac{\partial f}{\partial \phi_y}, \frac{\partial f}{\partial \phi_z} \right) \right) dx dy dz + o(\varepsilon^2)$$

Then apply lemma:

$$\frac{\partial f}{\partial \phi} - \sum_{i=1}^3 \frac{\partial}{\partial x_i} \left( \frac{\partial f}{\partial (\partial_i \phi)} \right) = 0 \quad (3.4)$$

In general we have this for  $n$  rather than 3 specifically:

$$\frac{\partial f}{\partial \phi} - \sum_{i=1}^n \frac{\partial}{\partial x_i} \left( \frac{\partial f}{\partial (\partial_i \phi)} \right) = 0 \quad (3.4.1)$$

Example Extremise "potential energy"  $n=2$

$$F(\phi) = \iint_{D \subseteq \mathbb{R}^2} \frac{1}{2} (\phi_x^2 + \phi_y^2) dx dy \quad f = \text{integrand}$$

$$\frac{\partial f}{\partial \phi} = 0, \quad \frac{\partial f}{\partial \phi_x} = \phi_x, \quad \frac{\partial f}{\partial \phi_y} = \phi_y \quad \text{so (3.4.1) gives}$$

$$\frac{\partial}{\partial x} \phi_x + \frac{\partial}{\partial y} \phi_y = 0 \quad \text{so} \quad \underline{\phi_{xx} + \phi_{yy} = 0}$$

Laplace equation

Example Minimal surfaces

Minimise area of  $\Sigma \subseteq \mathbb{R}^3$  subject to boundary conditions of 2 curves given as the boundary of the surface.



Take  $\Sigma = \{ \underline{x} \in \mathbb{R}^3 \text{ s.t. } g(x, y, z) = 0 \}$

Assume (note: can use implicit function theorem) that we solved  $g=0$  to give  $z = \phi(x, y)$

$$ds^2 = dx^2 + dy^2 + dz^2, \quad dz = \phi_x dx + \phi_y dy$$

$$\text{so } ds^2 = (1 + \phi_x^2) dx^2 + (1 + \phi_y^2) dy^2 + 2\phi_x \phi_y dx dy$$

This is called the first fundamental form or Riemannian metric

$$ds^2 = \sum_{i,j=1}^2 \bar{g}_{ij}(x, y) dx_i dx_j \quad \bar{g} = \begin{pmatrix} 1 + \phi_x^2 & \phi_x \phi_y \\ \phi_x \phi_y & 1 + \phi_y^2 \end{pmatrix}$$

$$\text{Area element } \sqrt{\det \bar{g}} dx dy$$

Then we get the area functional

$$A(\phi) = \int_D \underbrace{\sqrt{1 + \phi_x^2 + \phi_y^2}}_h dx dy \quad \text{apply E-L to } h \quad (3.4.1)$$

$$\frac{\partial h}{\partial \phi_x} = \frac{\phi_x}{\sqrt{1 + \phi_x^2 + \phi_y^2}} \quad \frac{\partial h}{\partial \phi_y} = \frac{\phi_y}{\sqrt{1 + \phi_x^2 + \phi_y^2}}$$

Since  $\frac{\partial h}{\partial \phi} = 0$ , applying 3.4.1 gives

$$\partial_x \left( \frac{\partial h}{\partial \phi_x} \right) + \partial_y \left( \frac{\partial h}{\partial \phi_y} \right) = 0 \text{ and expanding derivatives}$$

$$\text{gives } \underline{(1 + \phi_y^2) \phi_{xx} + (1 + \phi_x^2) \phi_{yy} - 2\phi_x \phi_y \phi_{xy} = 0} \quad (3.5)$$

the minimal surface equation.

Assume circular symmetry  $z = \phi(r)$ ,  $r = \sqrt{x^2 + y^2}$

$\Rightarrow$  3.5 becomes an ODE

$$\underline{r z'' + z' + (z')^3 = 0}$$

(check by chain rule - find derivatives of  $\phi$ ).