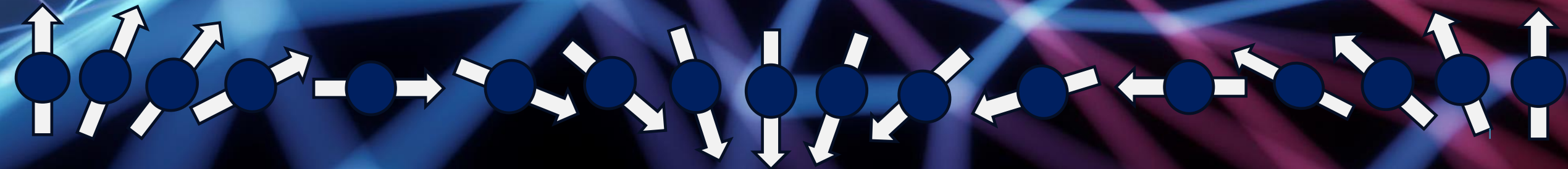


GENERATING SPIN SPIRALS USING RYDBERG ATOMS

Reet Mhaske

Supervisor: Prof Matthias Weidemuller



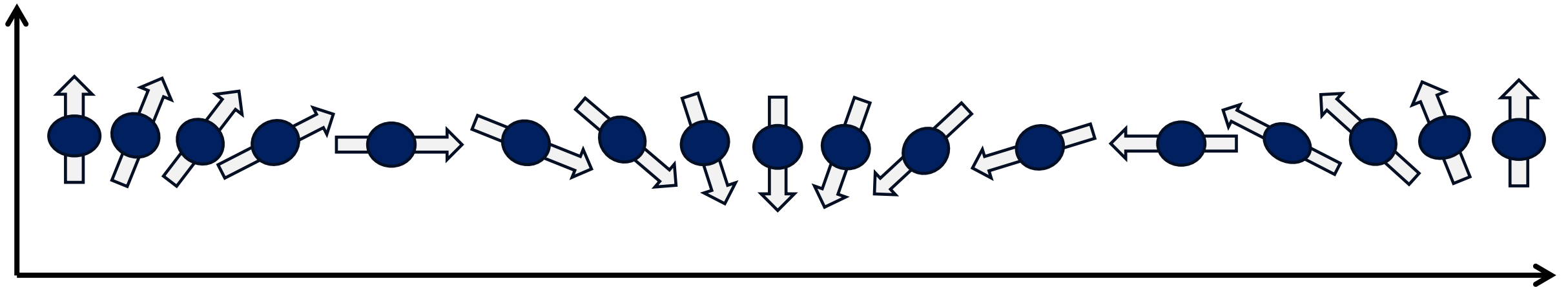
OVERVIEW

- What are spiral states?
- Why Spin Spiral?
- How to create spin spirals?
- Mapping to a spin Hamiltonian

1. WHAT ARE SPIN SPIRAL STATES?

$$|\Psi_{spiral}\rangle = \prod_i \sin\left(\frac{2\pi}{\lambda} x_i\right) \hat{x}_i + \cos\left(\frac{2\pi}{\lambda} x_i\right) \hat{y}_i$$

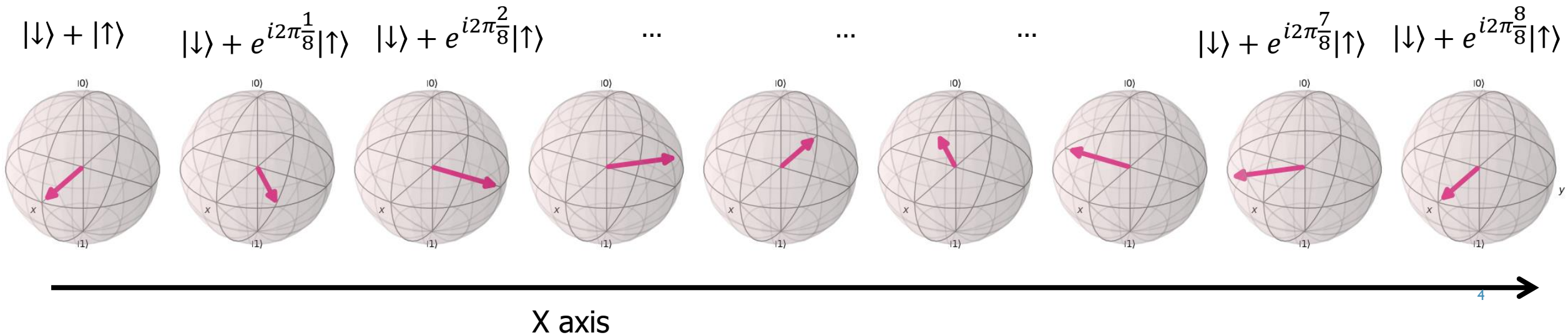
Y axis



X axis

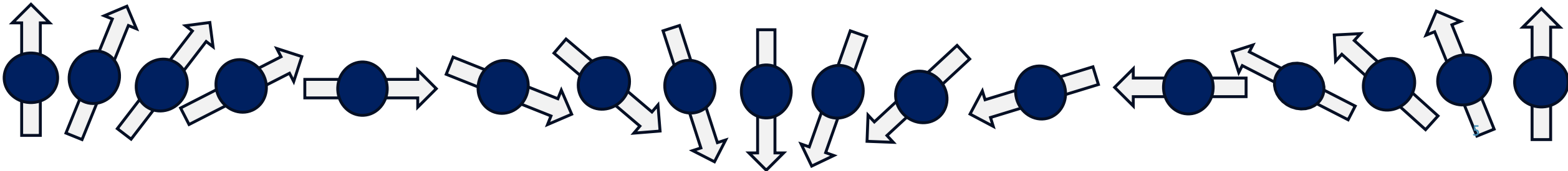
1. WHAT ARE SPIN SPIRAL STATES?

$$|\Psi_{\text{spiral}}\rangle = \prod_i |\downarrow\rangle_i + e^{\frac{i2\pi}{\lambda}x_i} |\uparrow\rangle_i$$

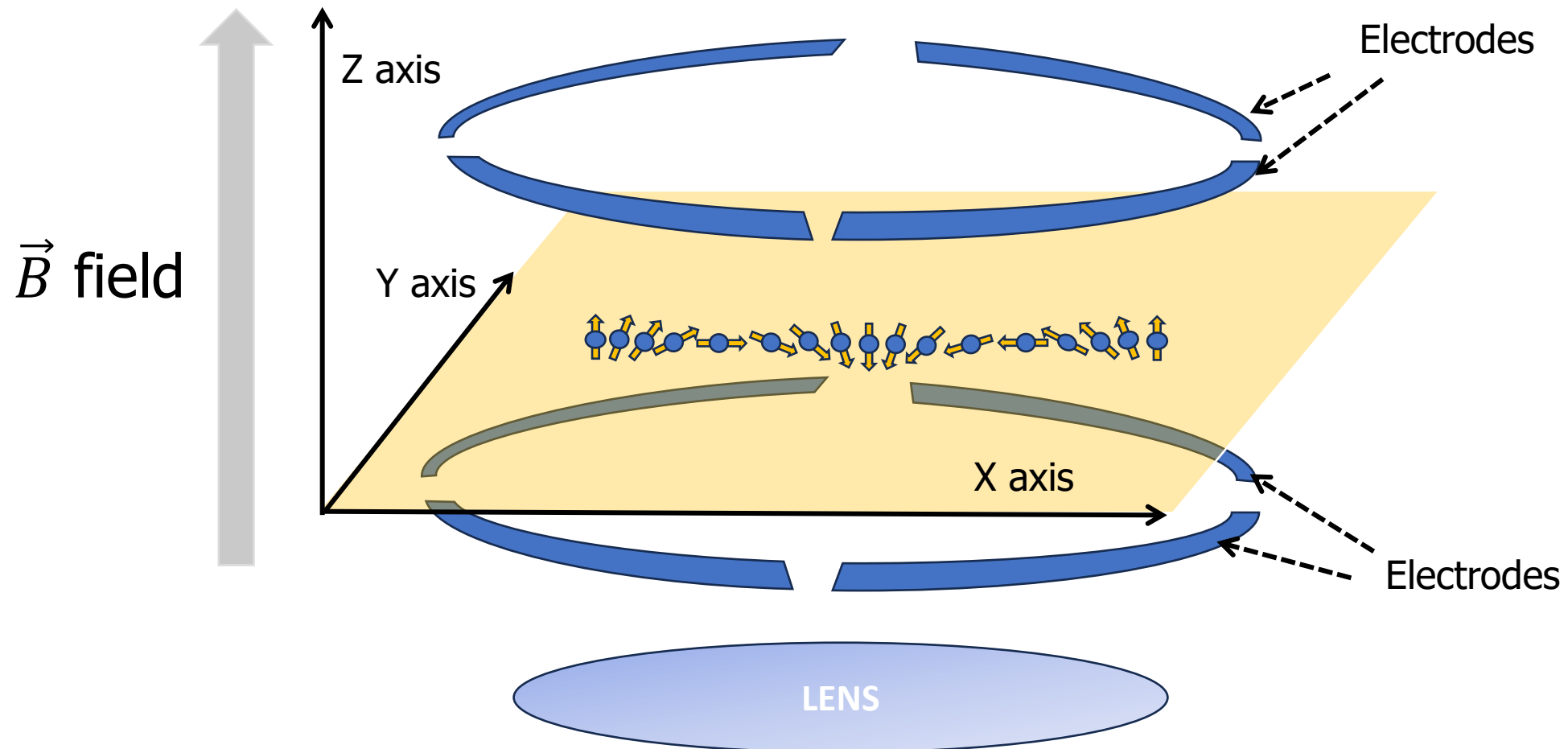


2. WHY SPIN SPIRALS ?

- Introduces a new length scale (λ) to the system
- Can study spin-spin correlation
- Understand spin dynamics better
- How far-from equilibrium systems relax



"SEEING" SPIRALS USING FLUORESCENCE IMAGING



3. HOW TO CREATE SPIN SPIRALS?

- The Protocol
- Choosing the two energy levels
- Setting the electrode voltages

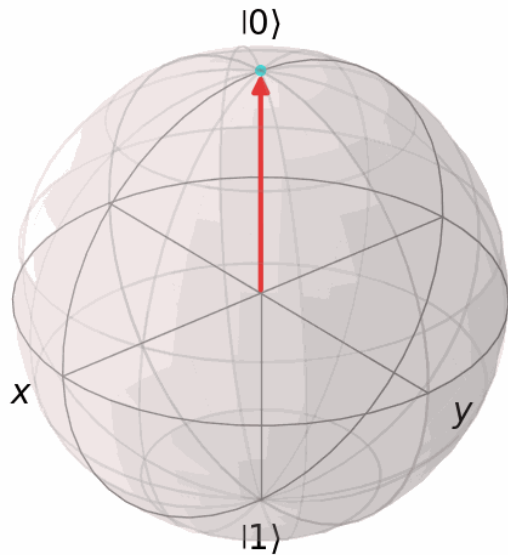


3. HOW TO CREATE SPIN SPIRALS

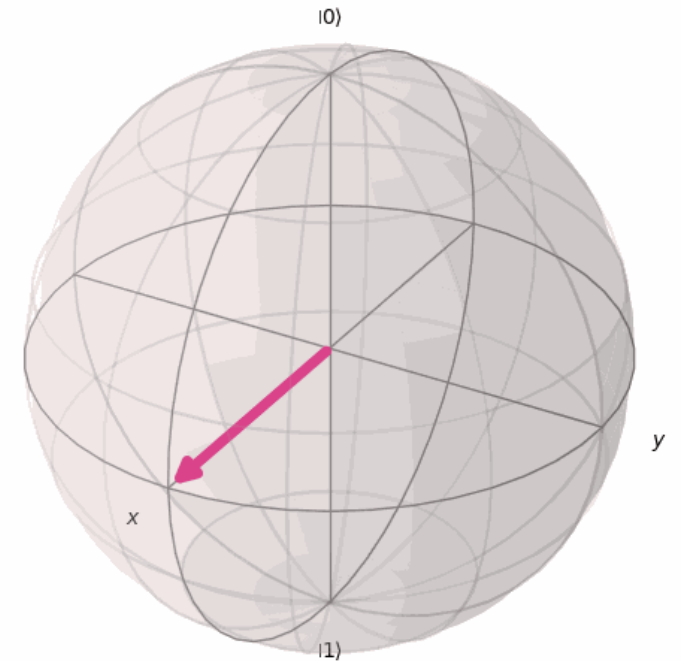
A. THE PROTOCOL

3A. THE PROTOCOL

Consider a single qubit



under $H = |\downarrow\rangle\langle\downarrow| + \hbar\omega |\uparrow\rangle\langle\uparrow|$

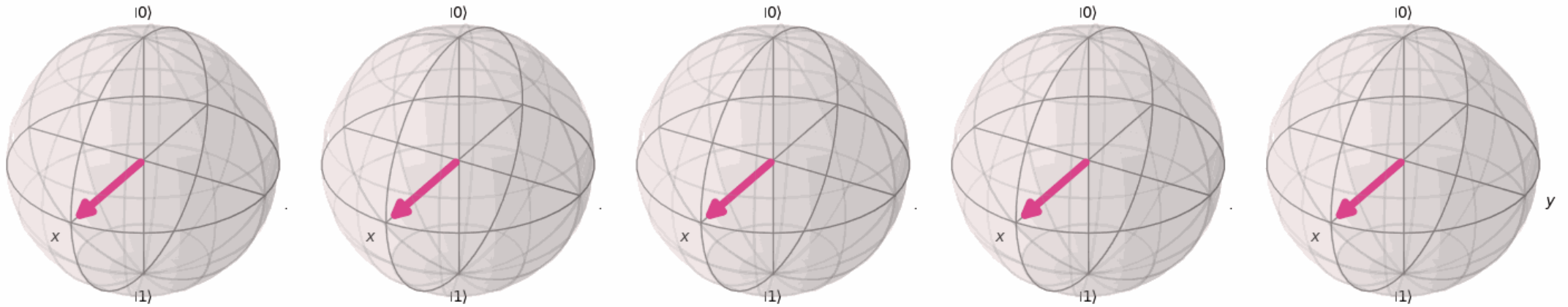


A single qubit initialized in $|\uparrow\rangle + e^{i\theta}|\downarrow\rangle$

Evolves as $|\uparrow\rangle + e^{i(\theta+\omega t)}|\downarrow\rangle$

3A. THE PROTOCOL

For multiple qubits, create position dependent differences in ω



Energy gaps:

ω_0

$\omega_0 + \delta\omega$

$\omega_0 + 2\delta\omega$

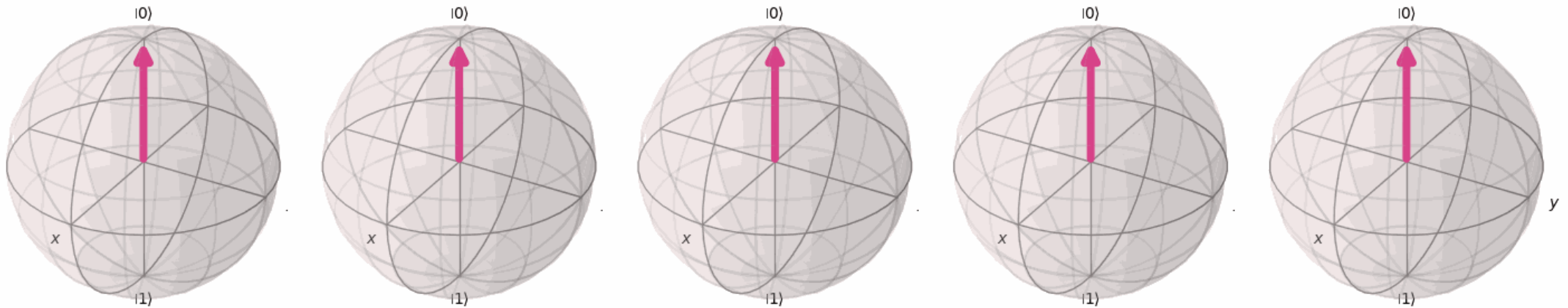
$\omega_0 + 3\delta\omega$

$\omega_0 + 4\delta\omega$

X axis

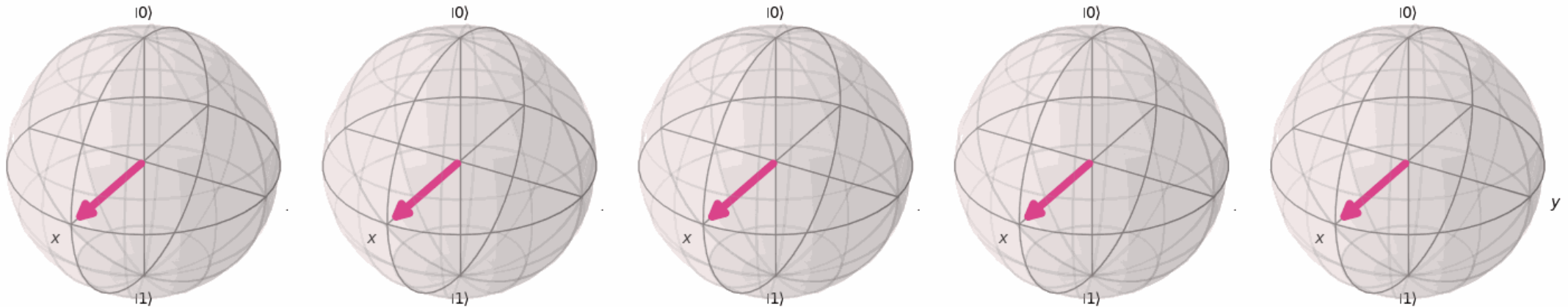
SO THE PROTOCOL TO CREATE SPIRAL STATES WOULD BE...

1. Initialize all spins in x-y plane



SO THE PROTOCOL TO CREATE SPIRAL STATES WOULD BE...

2. Create the position dependent frequency difference



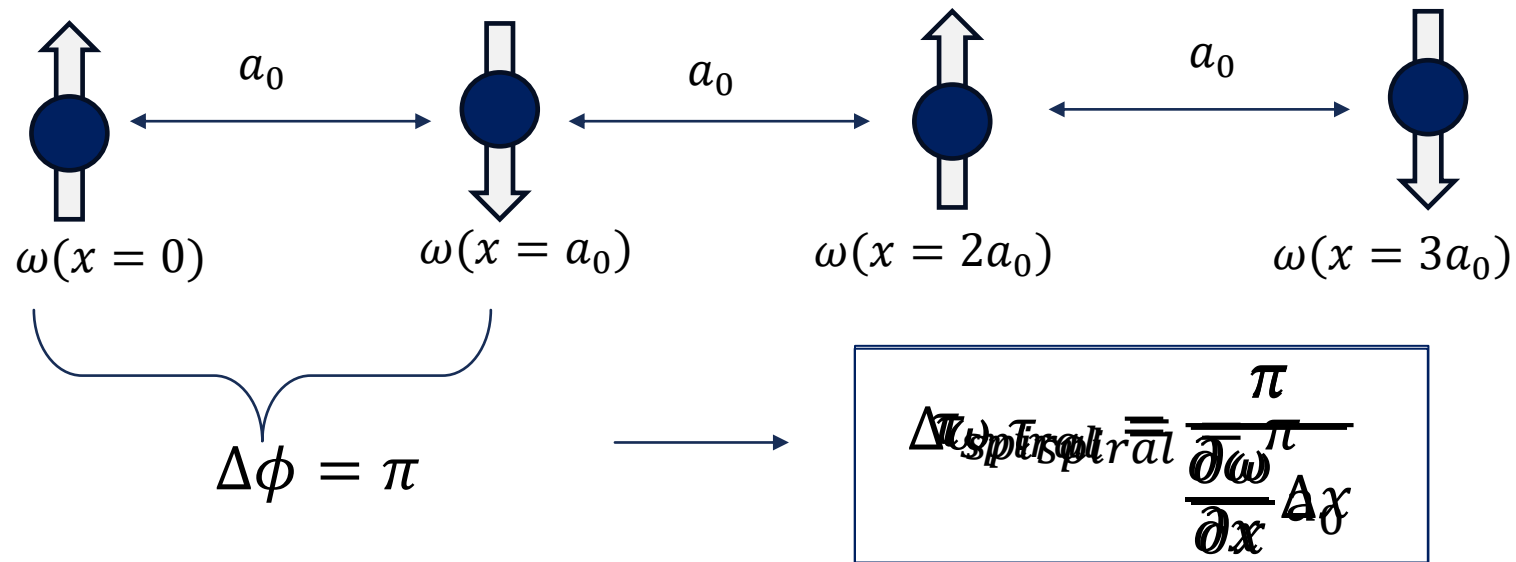
3. Remove the position dependent frequency difference, when desired state is achieved (after τ_{spiral})

SPIRALIZATION TIME

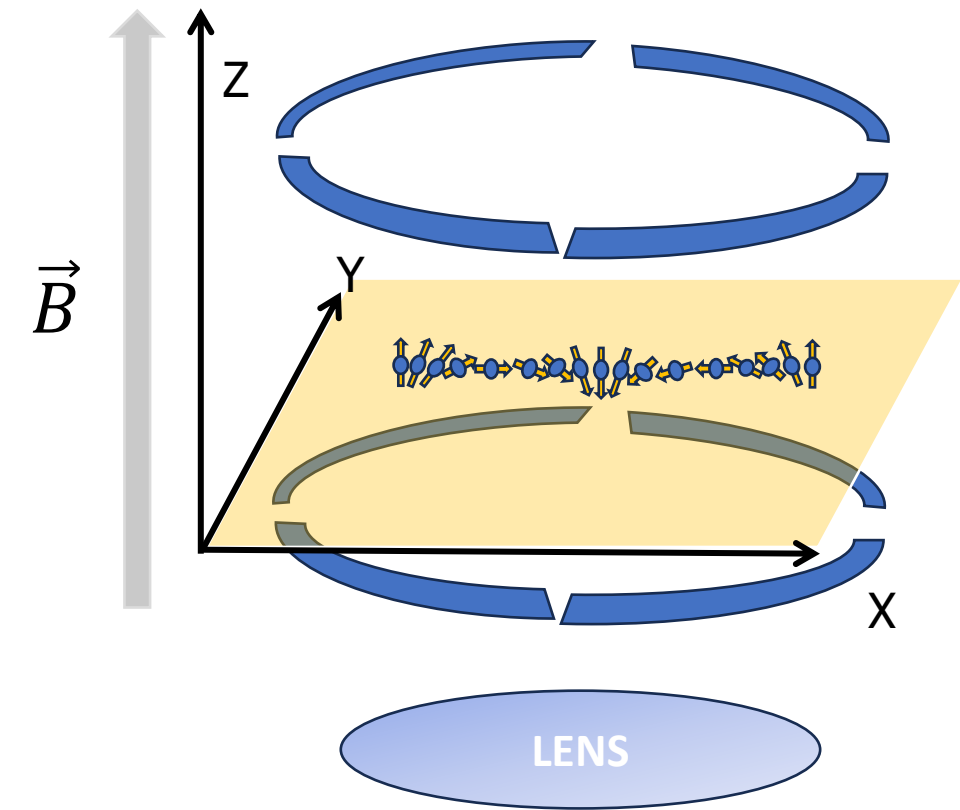
Time to reach desired state (τ_{spin})

Maximum τ_{spin} needed:

For Neel state ($\lambda = 2a_0$):



CREATING A FREQUENCY VARIATION WITH POSITION $\left(\frac{\delta\omega}{\delta x}\right)$



Rydberg Hamiltonians in external fields

$$H(\vec{E}, \vec{B}) = H_o - \hat{d} \cdot \vec{E} - \hat{\mu} \cdot \vec{B} + \frac{1}{8m_e} |\hat{d} \times \vec{B}|^2$$

Electric dipole operator

Magnetic dipole operator

Better experimental control on \vec{E} than on \vec{B}

CREATING A FREQUENCY VARIATION WITH POSITION $\left(\frac{\delta\omega}{\delta x}\right)$

For a spiral state:

$$\begin{aligned}\phi(x) &\propto x \\ |\Psi_{spiral}\rangle &= \prod_i |\downarrow\rangle_i + e^{\frac{i2\pi}{\lambda}x_i} |\uparrow\rangle_i \\ \phi(x) &= \omega(x)\tau_{spiral}\end{aligned}$$

$$\omega(x) = \frac{\partial\omega(x)}{\partial x} x$$

$$\Delta\omega = \frac{\partial\omega(x)}{\partial x} \Delta x = \frac{\partial\omega(\vec{E})}{\partial|\vec{E}_{plane}|} \frac{\partial|\vec{E}_{plane}|}{\partial x} \Delta x$$

ACHIEVING LINEAR VARIATION OF ω

$$\Delta\omega = \frac{\partial\omega(\vec{E})}{\partial|\vec{E}_{plane}|} \frac{\partial|\vec{E}_{plane}|}{\partial x} \Delta x$$

Choosing two energy levels:

- $\frac{\partial\omega}{\partial|\vec{E}_{plane}|} \sim \text{constant}$

Adjusting electrodes voltages:

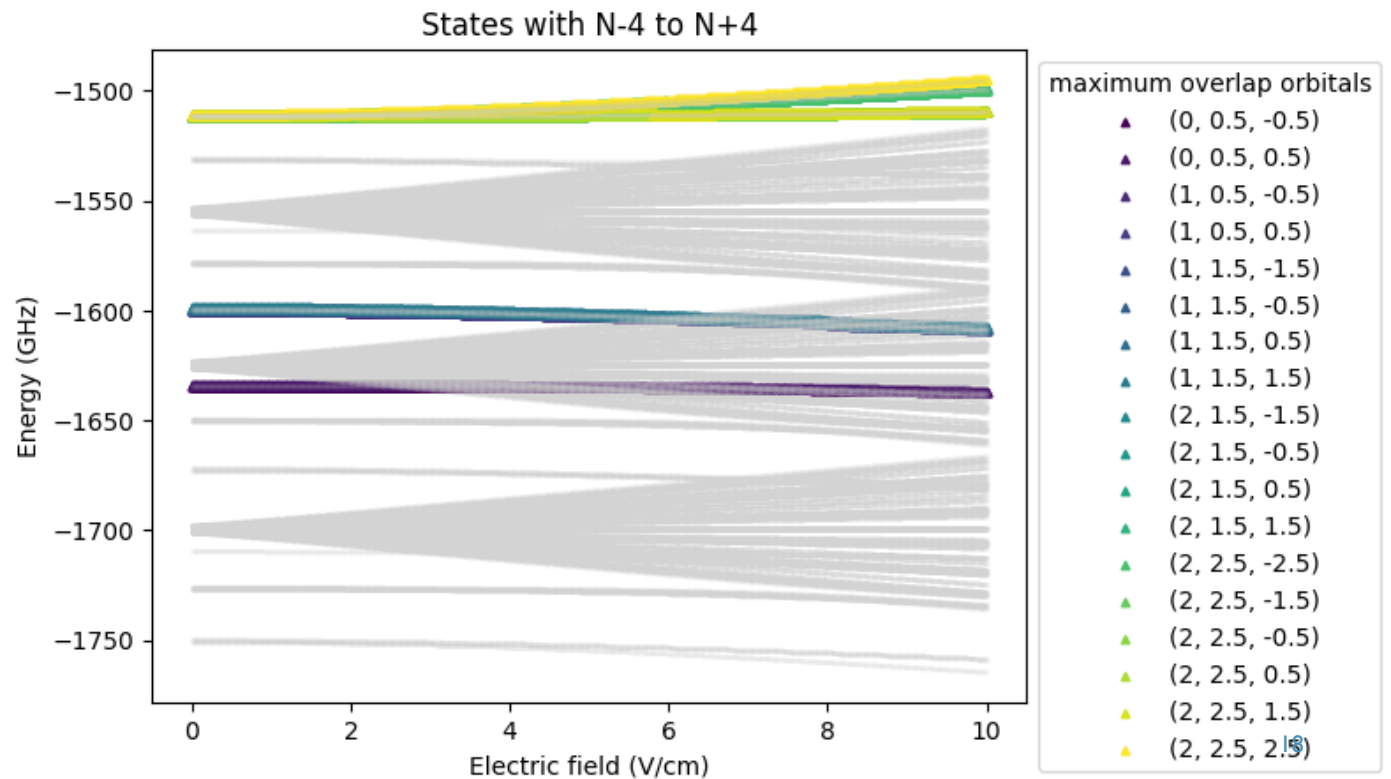
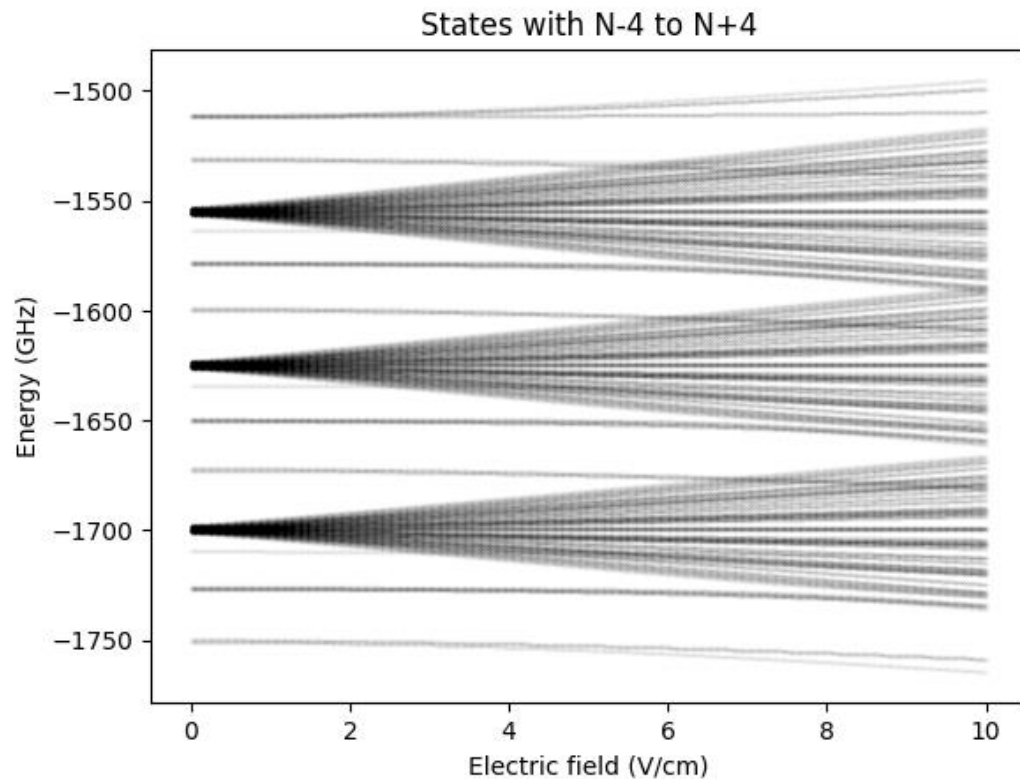
- $\frac{\partial|\vec{E}_{plane}|}{\partial x} \sim \text{constant}$



3.B CHOOSING THE TWO LEVELS

RYDBERG ATOMS IN ELECTRIC FIELD $\left(\frac{\delta\omega}{\delta|\vec{E}|}\right)$

At $\vec{B} = 160 \text{ G } \hat{z}$ variation of energy levels with \vec{E}



CRITERIA FOR CHOOSING ENERGY LEVELS

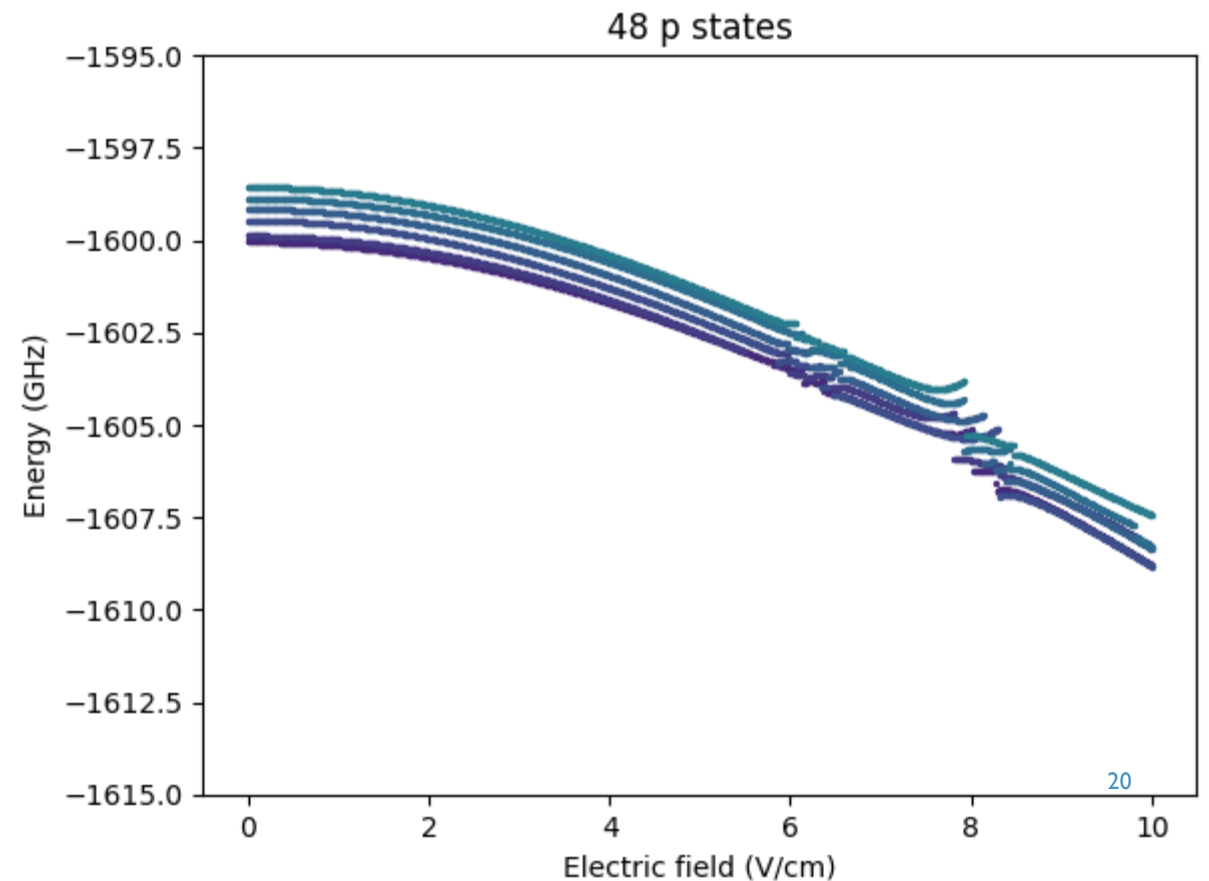
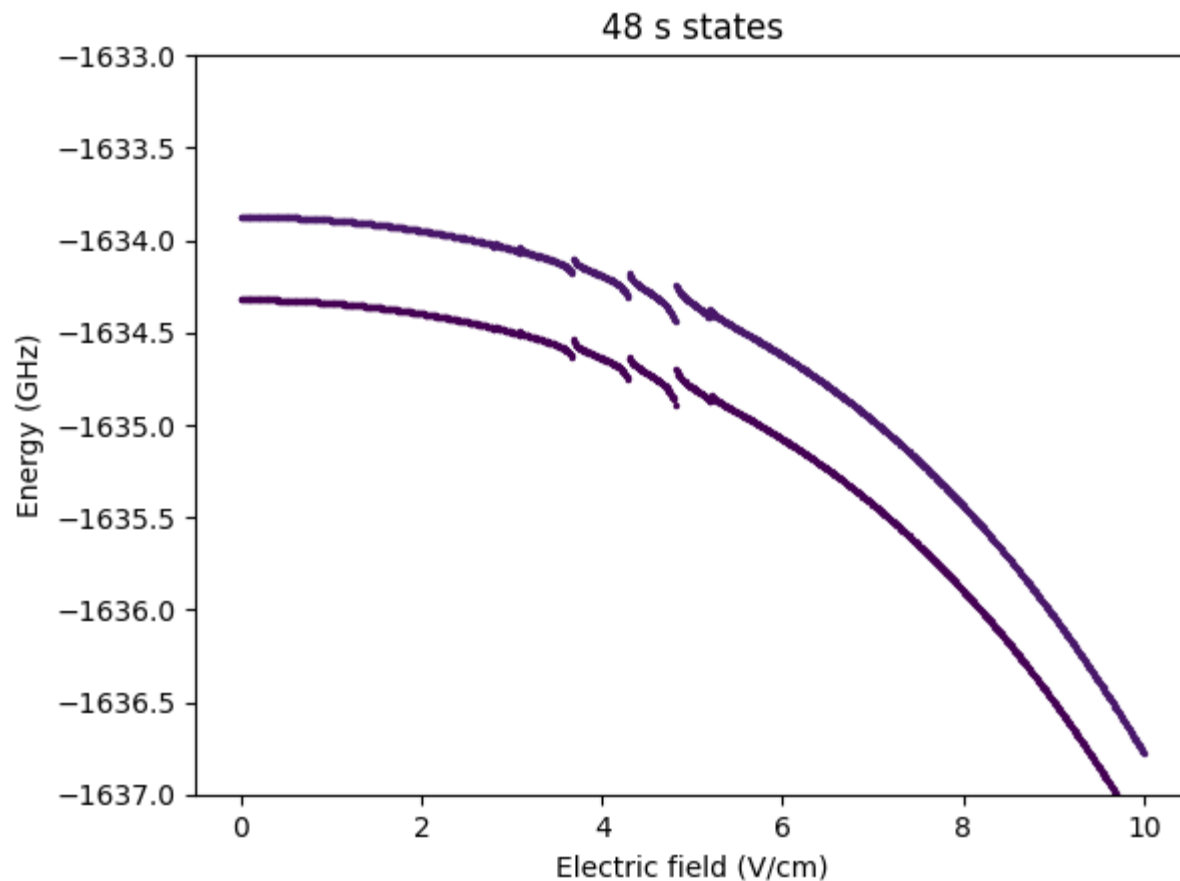
- $\frac{\partial \omega}{\partial |\vec{E}_{plane}|} \sim \text{constant for considerable range}$
- $\tau_{spiral} \ll T_{interaction}$
- $\frac{\partial \omega}{\partial |\vec{E}_{plane}|}$ large enough to get small τ_{spiral}
- Large electric dipole moment

Reference dipole moment : 1.38 GHz (cm/V)

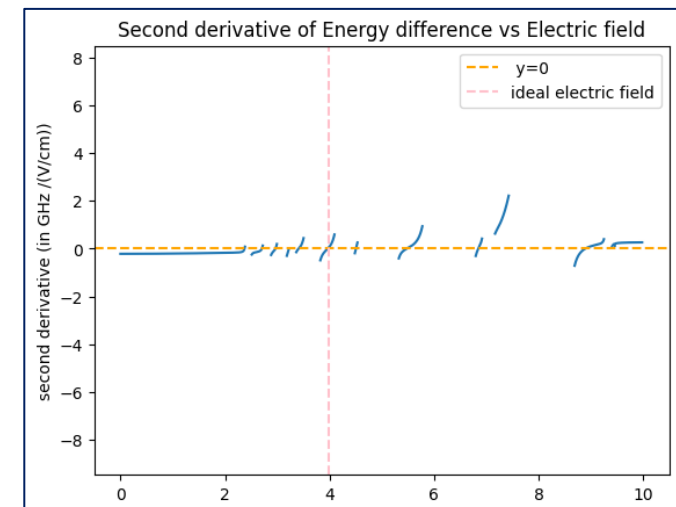
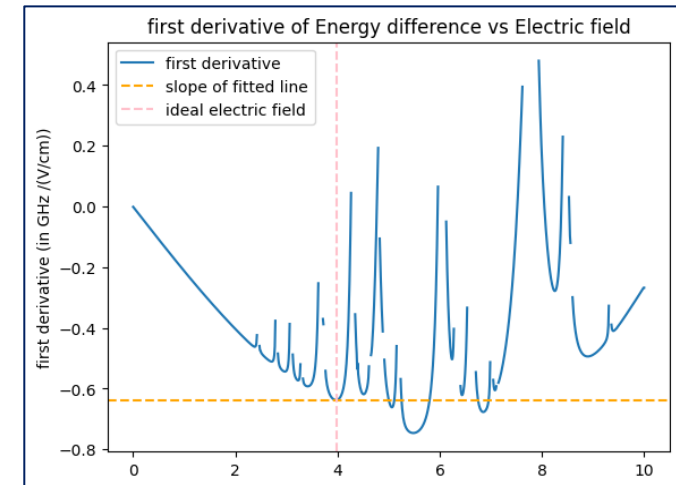
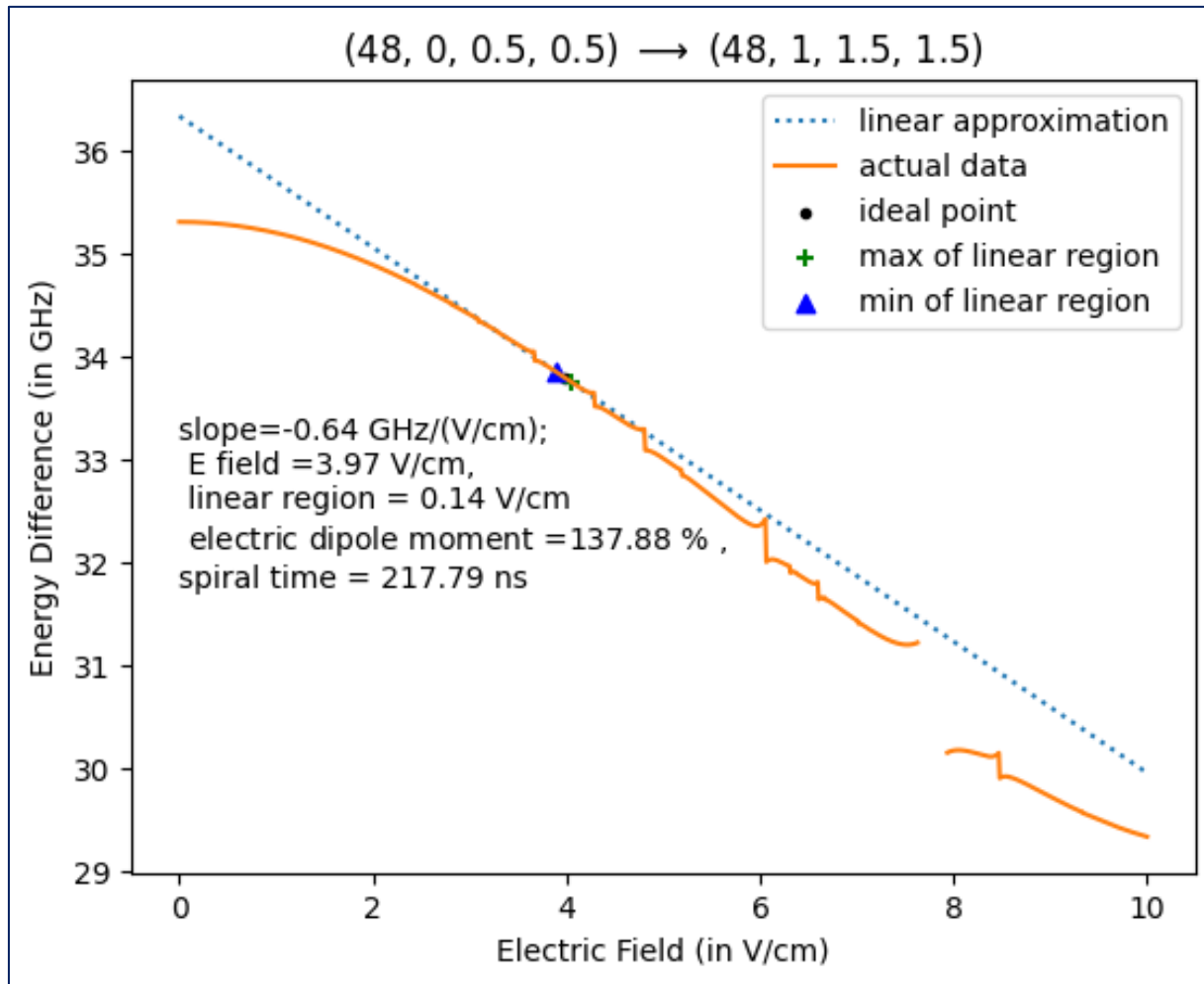
[for $|48\ s\ 0.5\ 0.5\rangle \rightarrow |48\ p\ 1.5\ 0.5\rangle$]

A CLOSER LOOK AT STARK MAPS

At $\vec{B} = 160 \text{ G } \hat{z}$ variation of energy levels with \vec{E}



CHOOSING THE TWO LEVEL SYSTEM



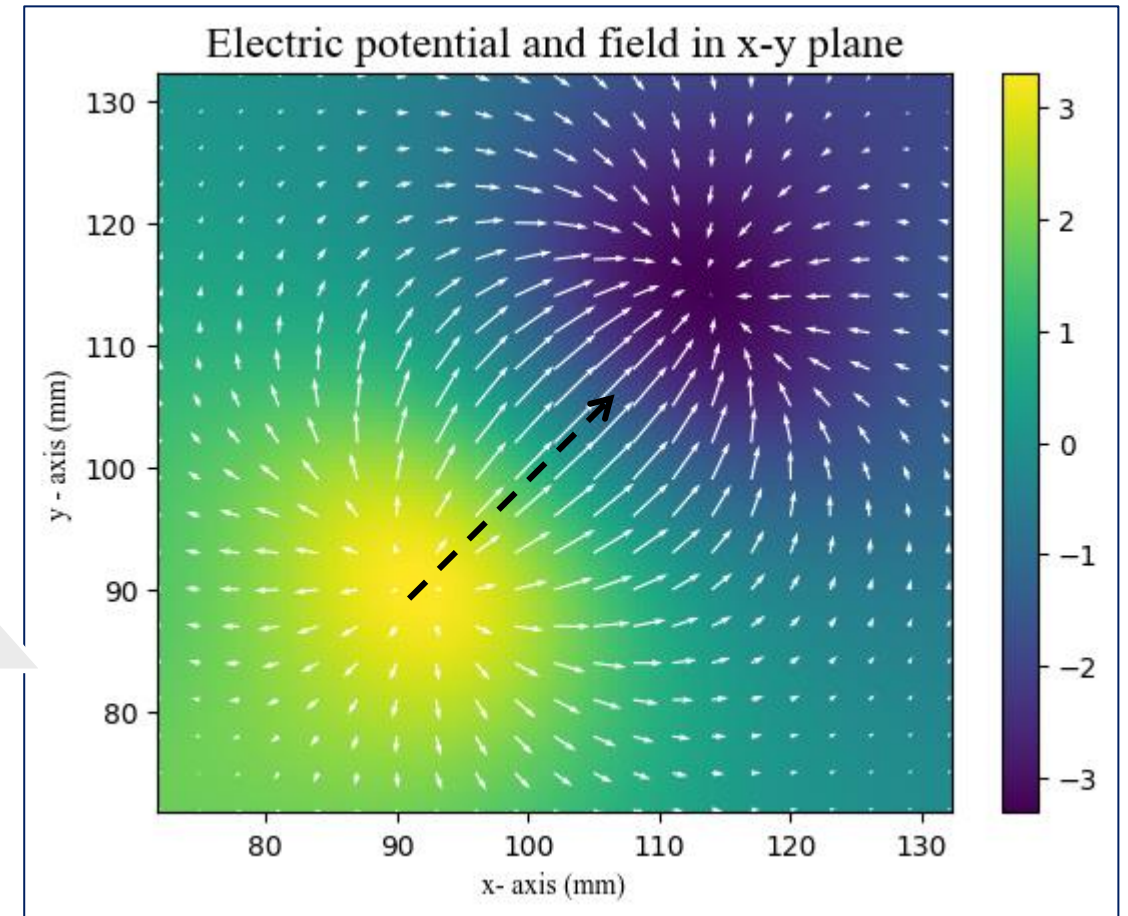
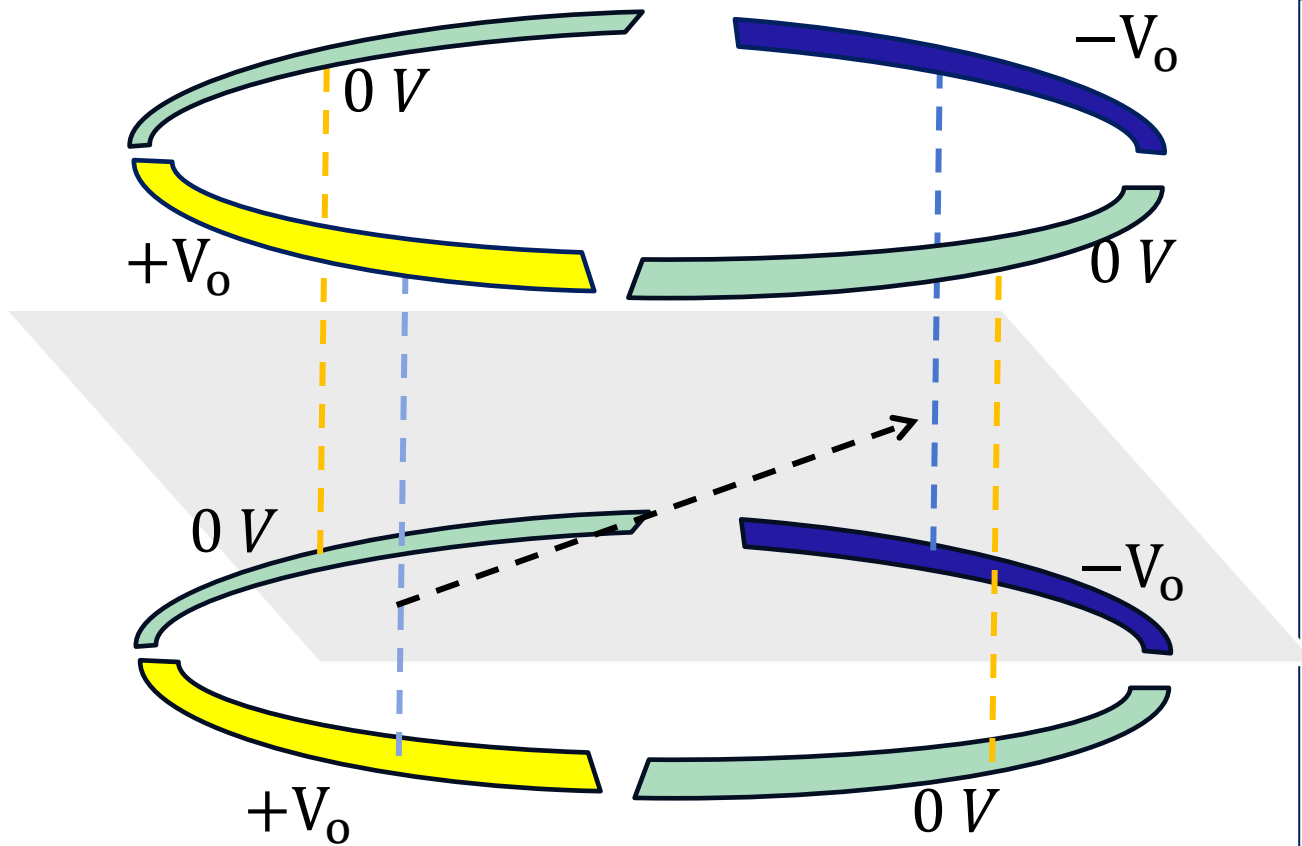
3C. SETTING ELECTRODE VOLTAGES

- A. Generating offset electric field
- B. Generating $\frac{\partial |\vec{E}_{plane}|}{\partial r}$
- C. Combining these two

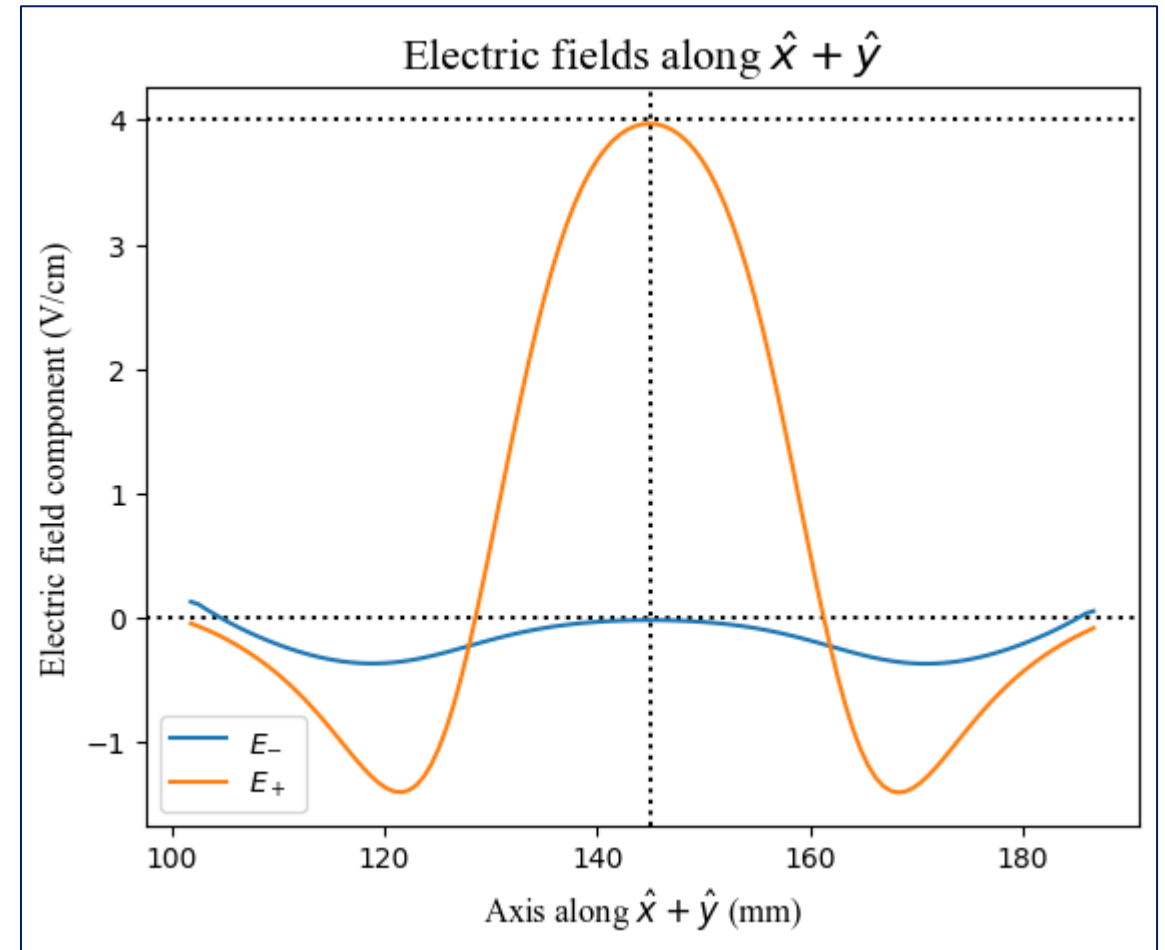
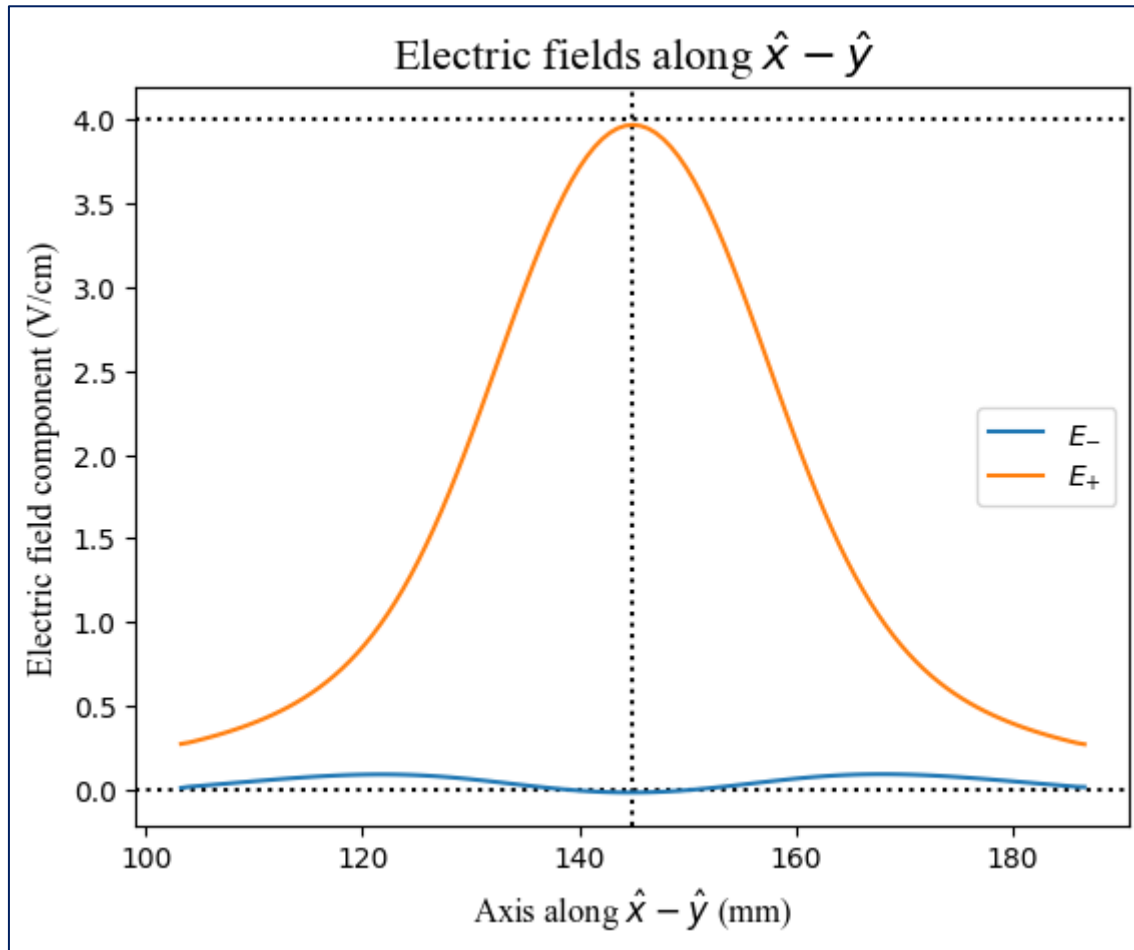
DESIRED ELECTRIC FIELDS

- $\frac{\partial |\vec{E}_{plane}|}{\partial r} = \text{constant}$
- $\frac{\partial |\vec{E}_{plane}|}{\partial r_{\perp}} \sim 0$
- $\frac{\partial |\vec{E}_{plane}|}{\partial z} \sim 0$
- Electrode potentials < 40V (to be able to ramp faster)

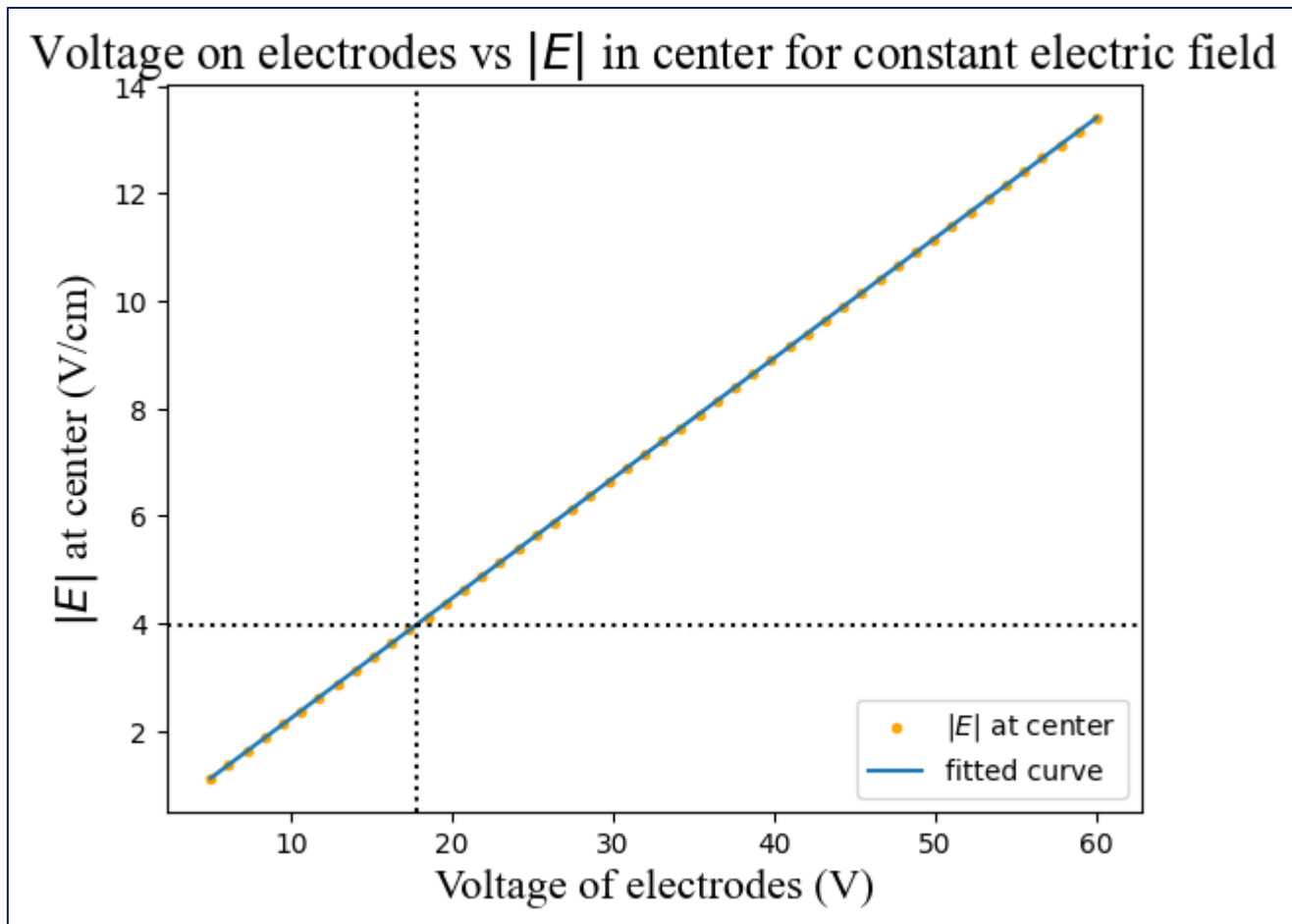
A. GENERATING CONSTANT $|\vec{E}_{plane}|$



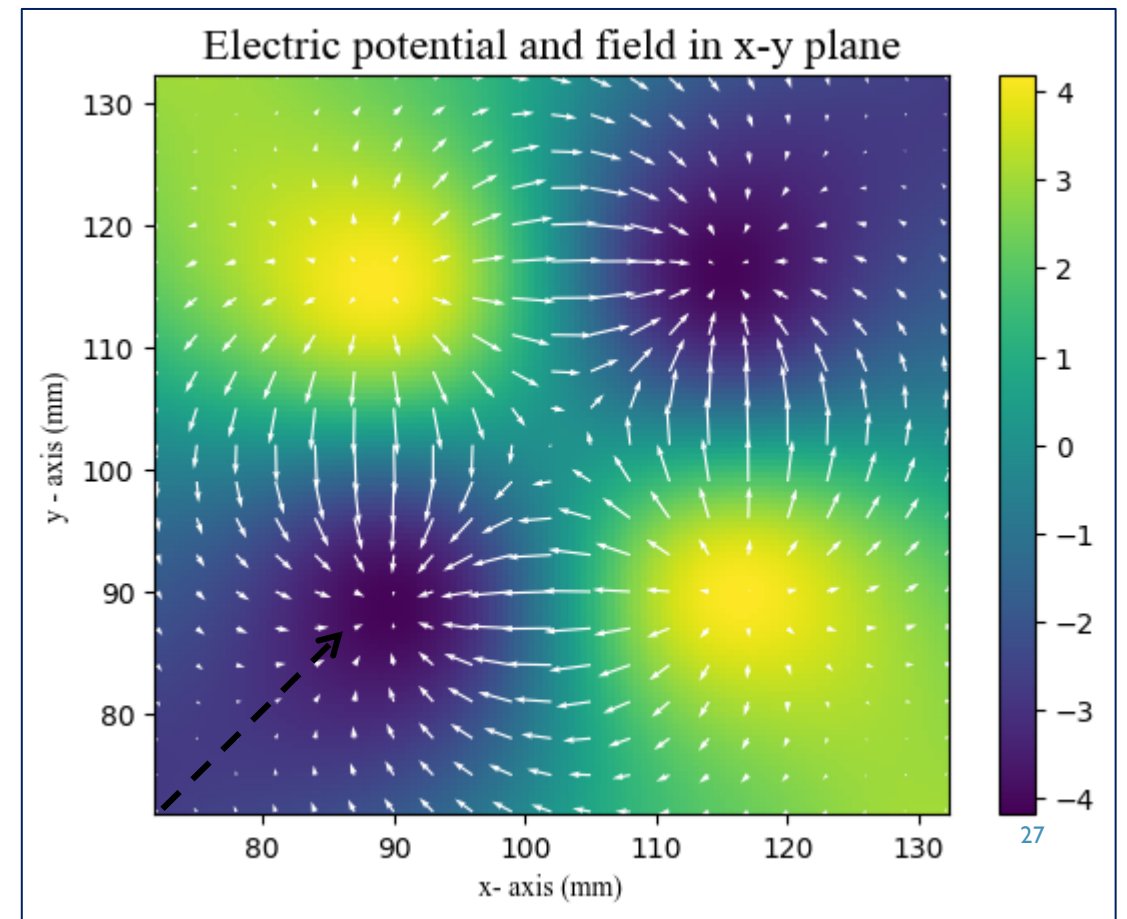
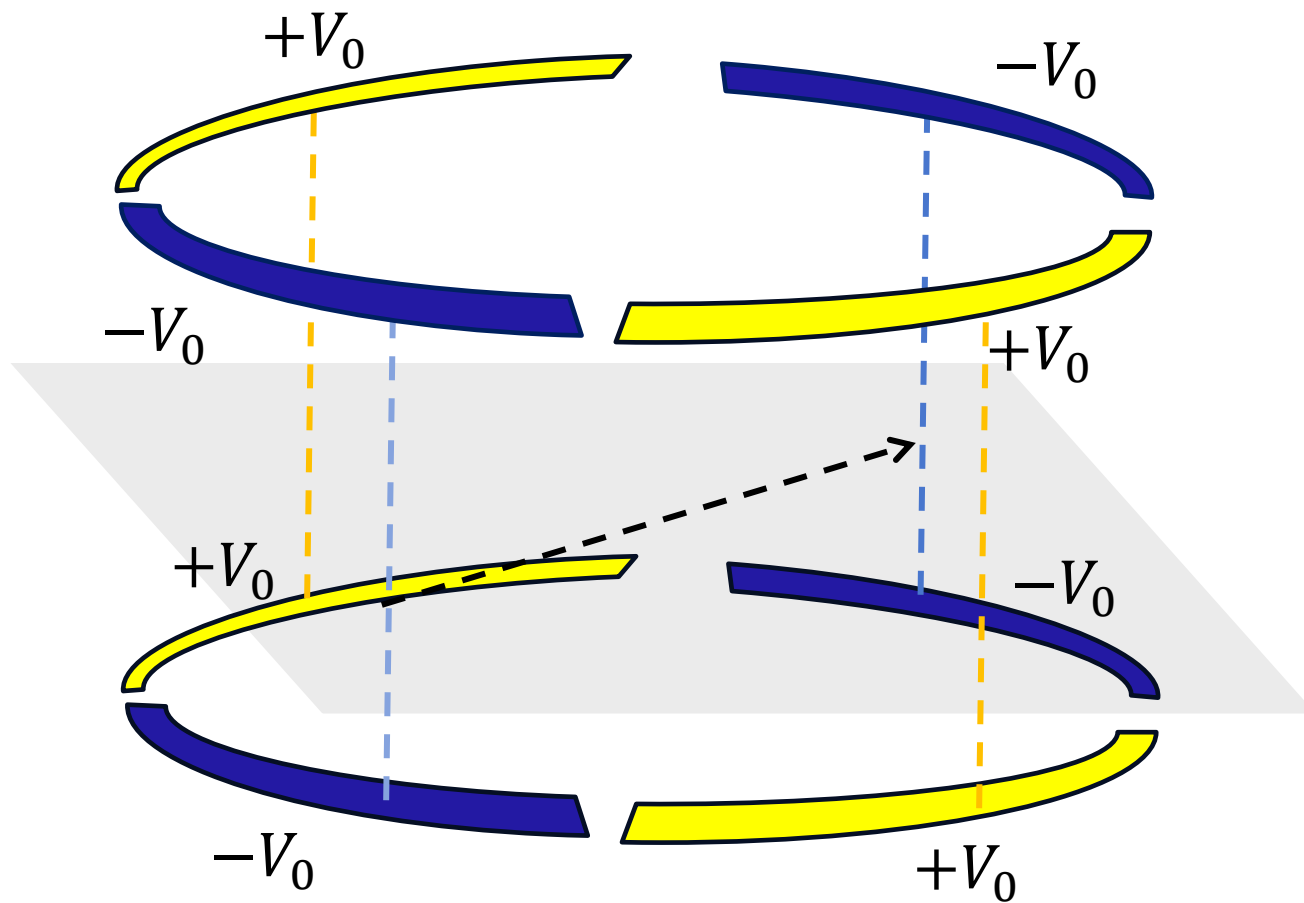
THE OFFSET FIELDS GENERATED



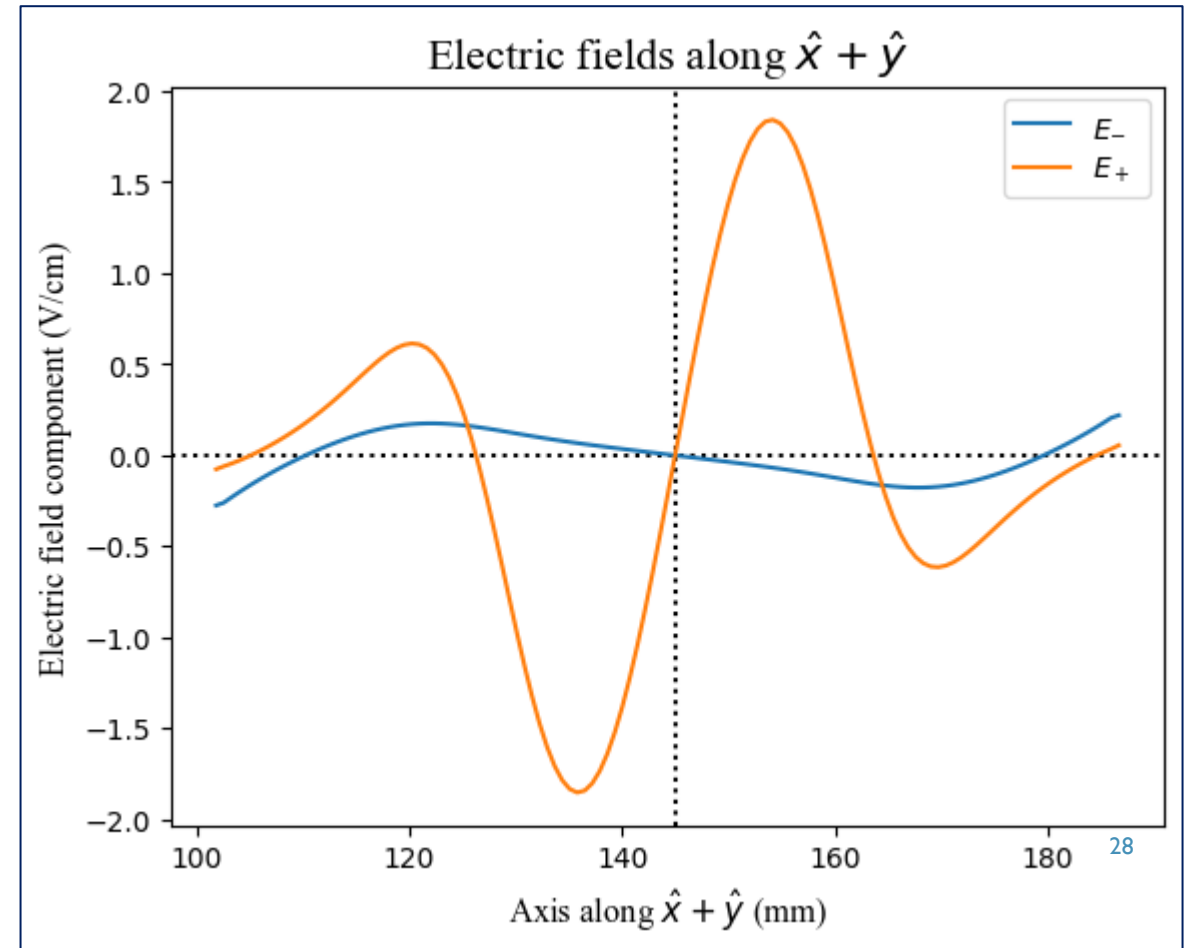
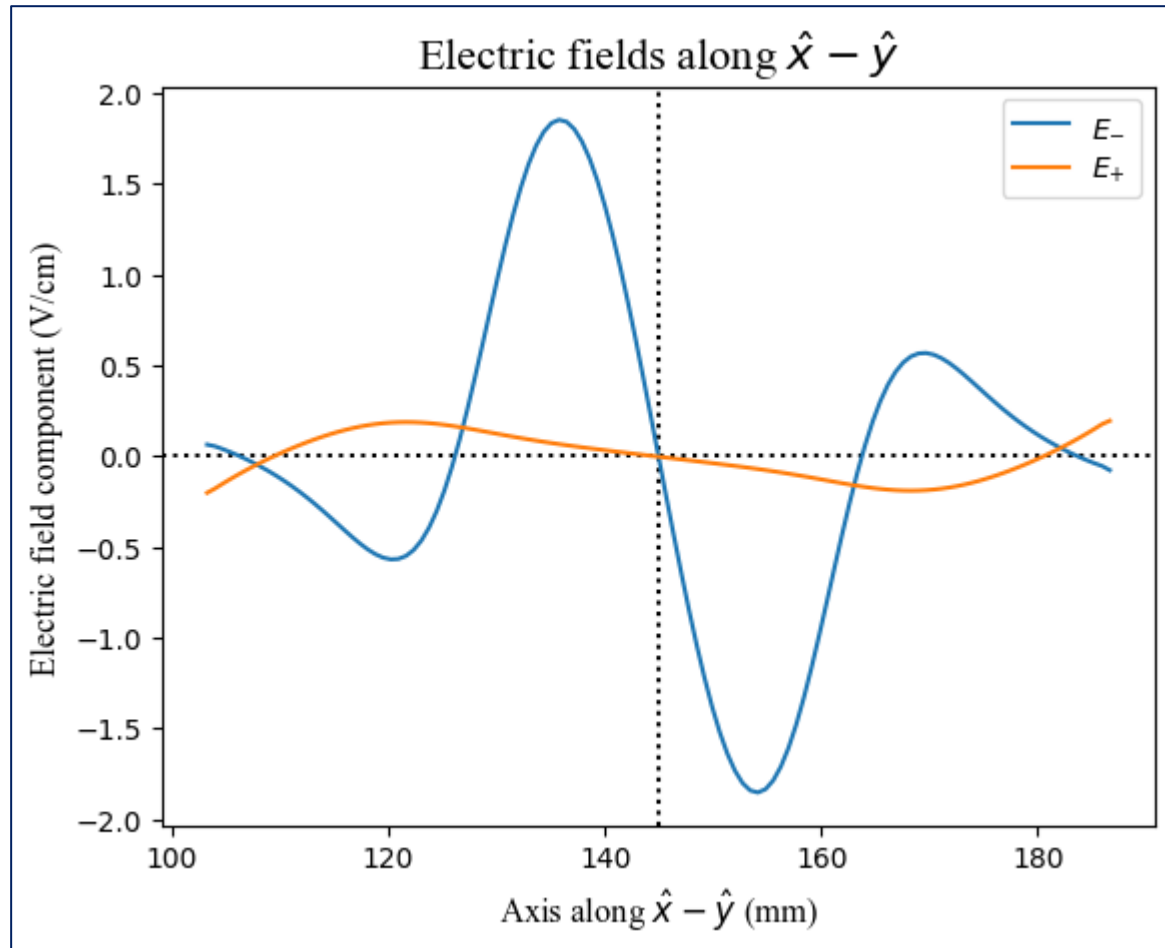
FINALIZING V_0



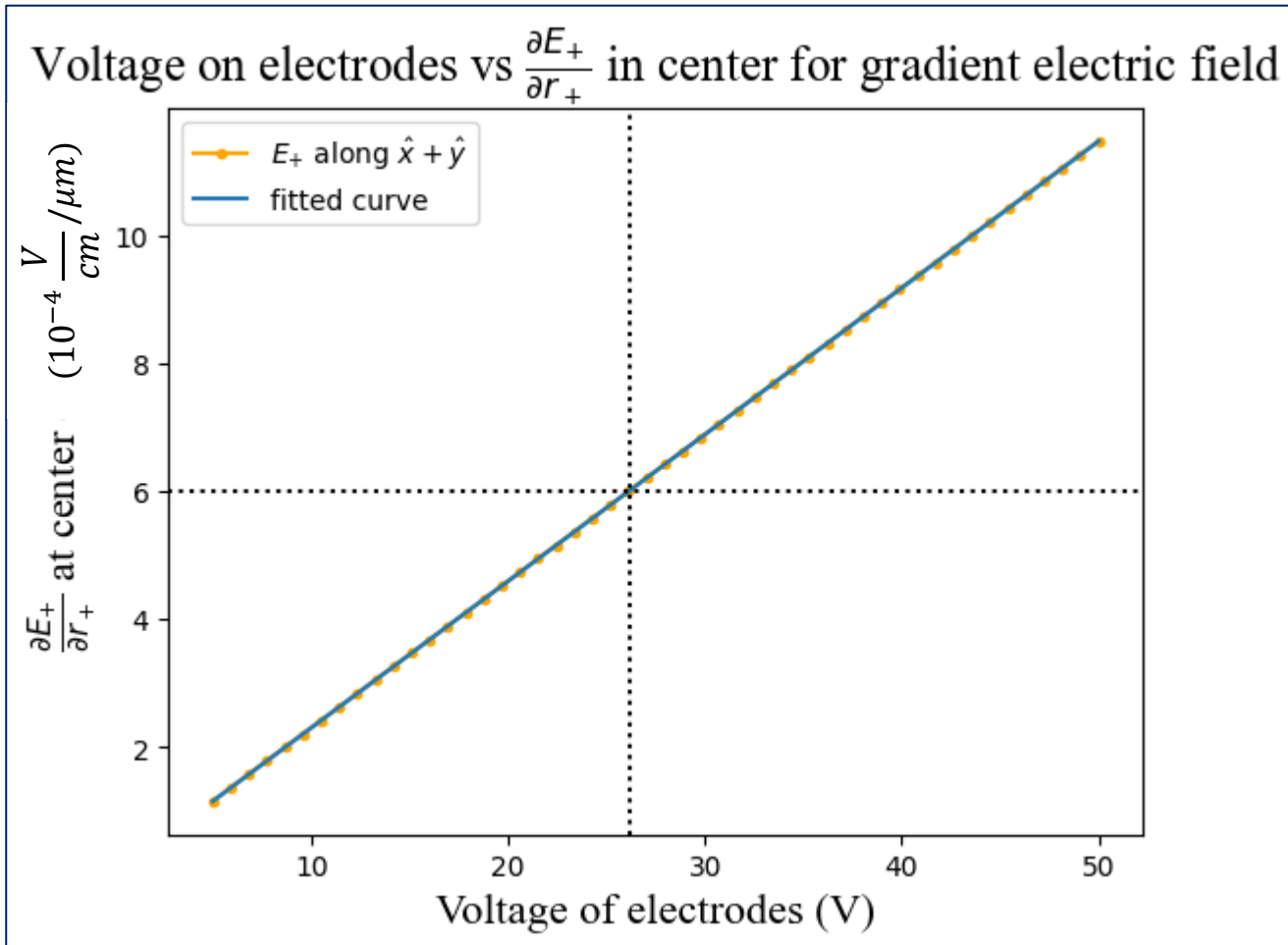
B. GENERATING $\frac{\partial |\vec{E}_{plane}|}{\partial r}$



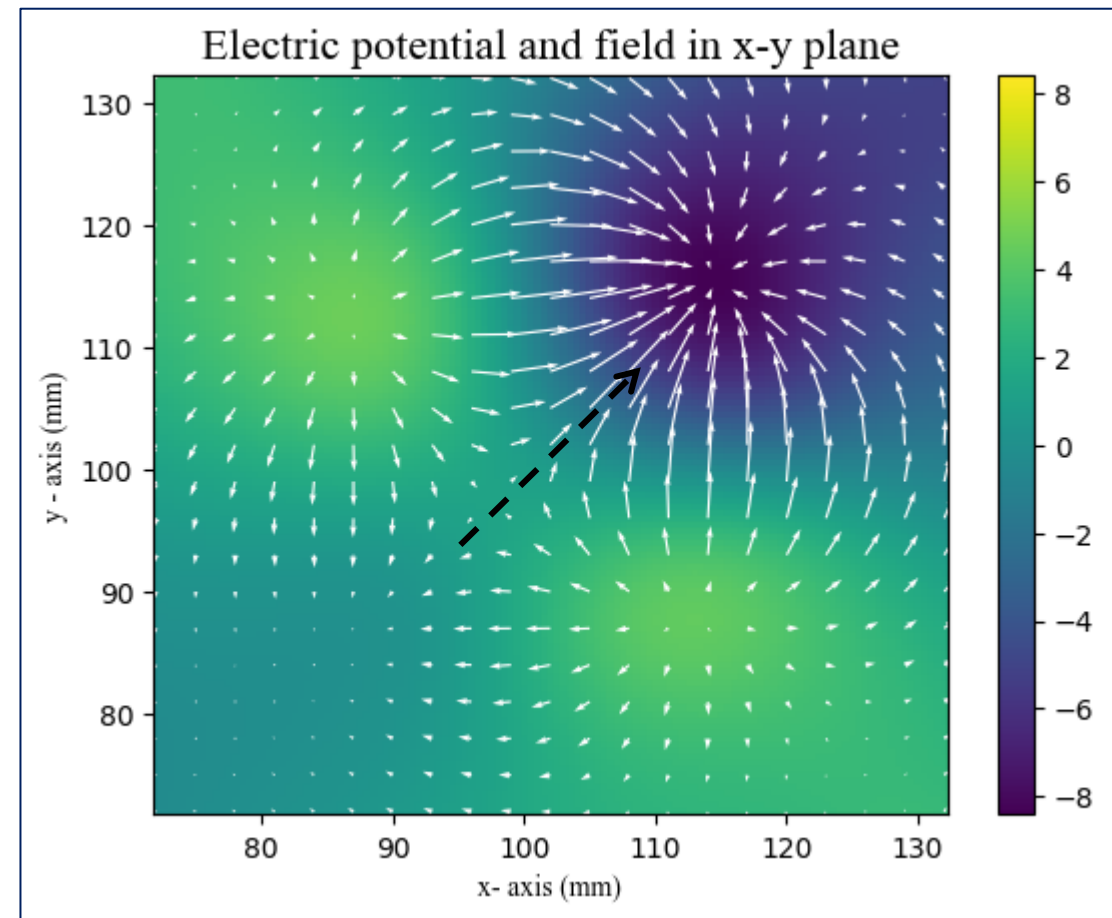
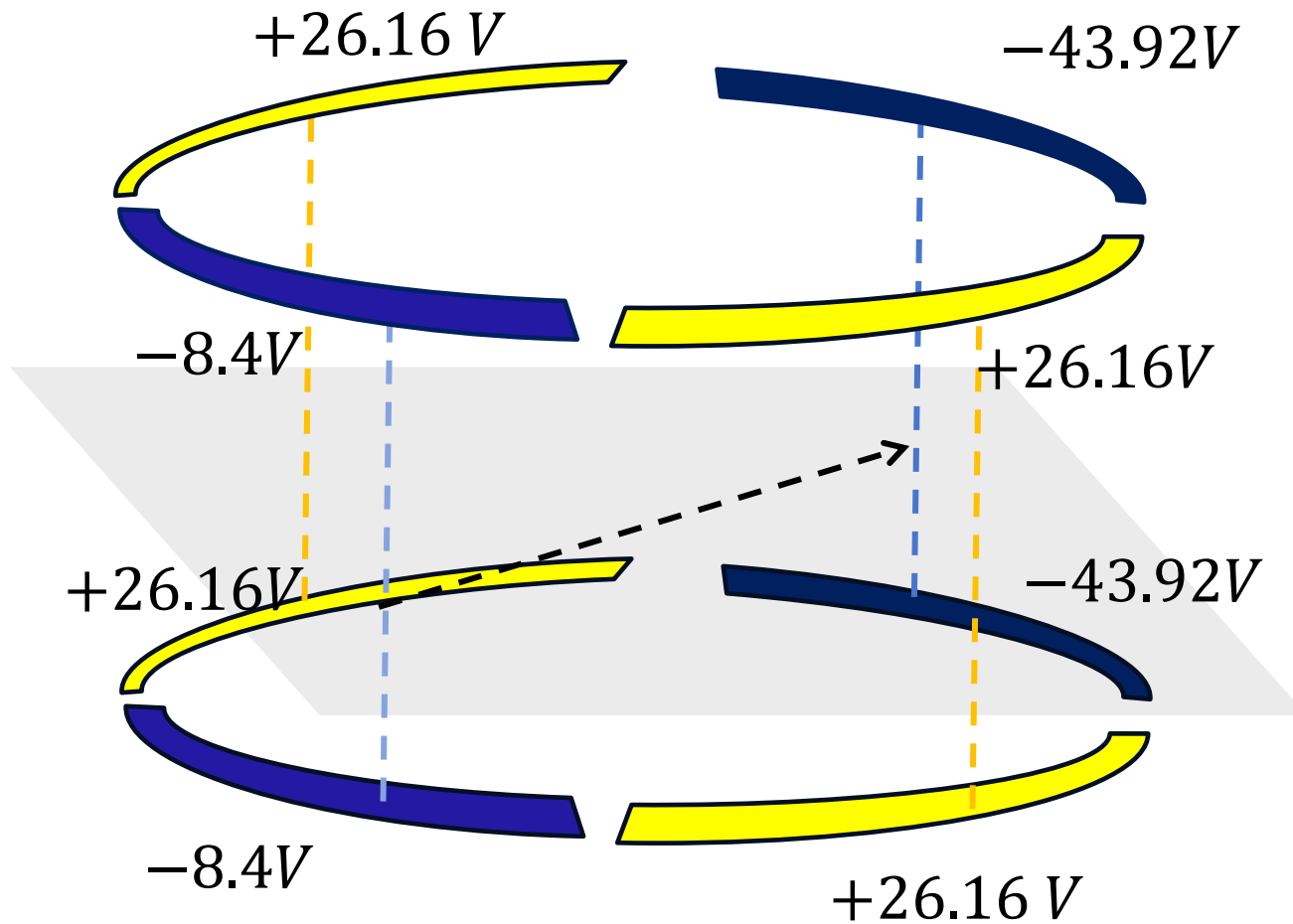
THE GRADIENT FIELDS GENERATED



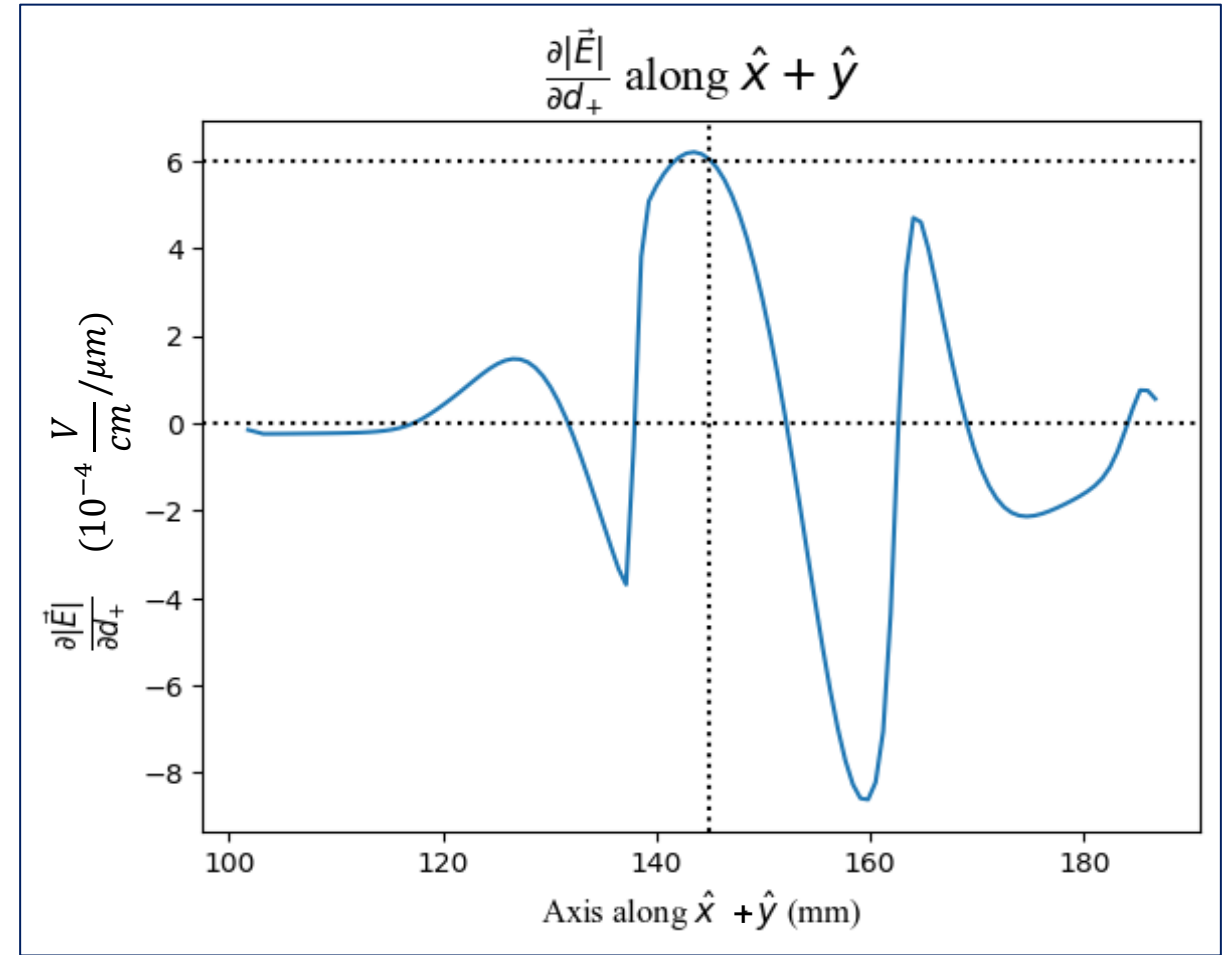
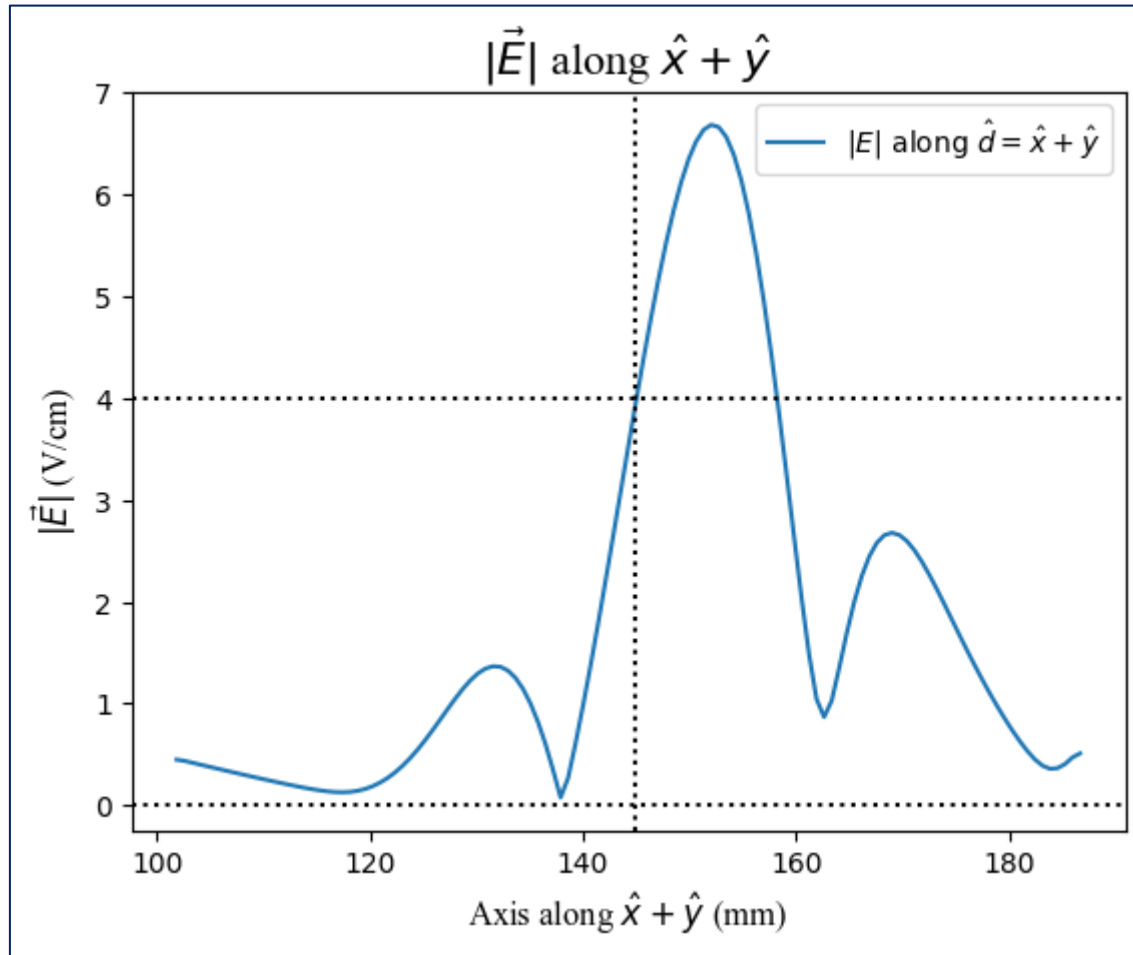
FINALIZING V_0



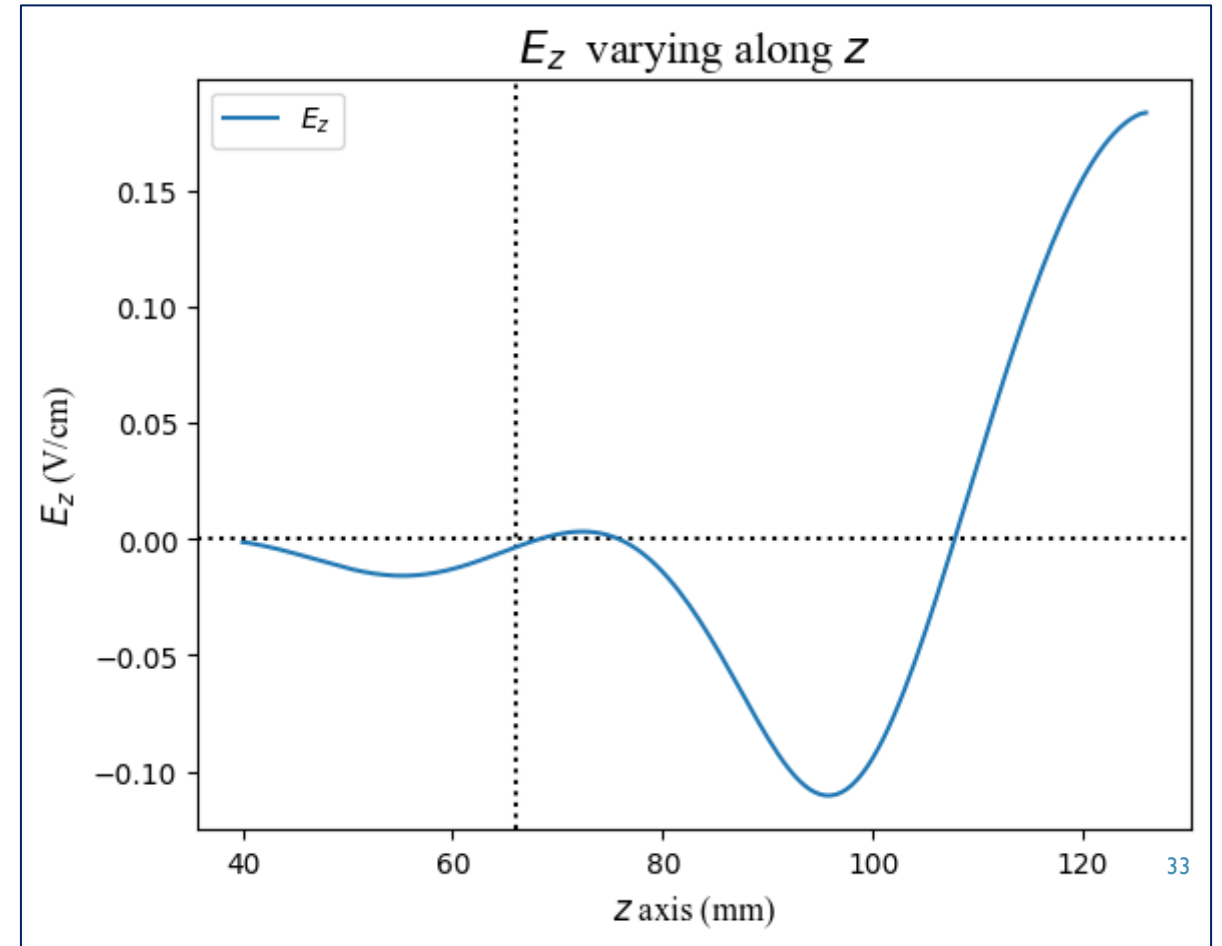
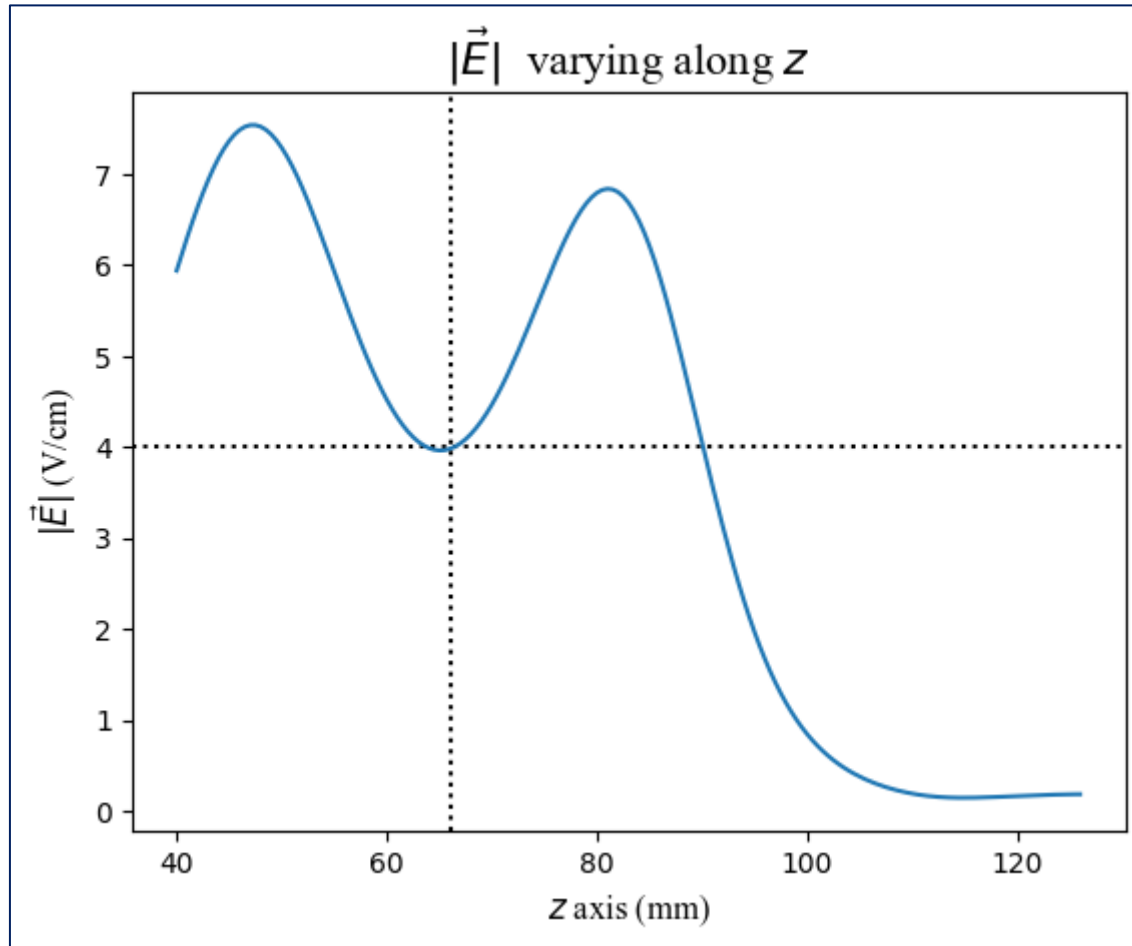
C. GENERATING OFFSET $|\vec{E}_{plane}|$ & $\frac{\partial |\vec{E}_{plane}|}{\partial r}$



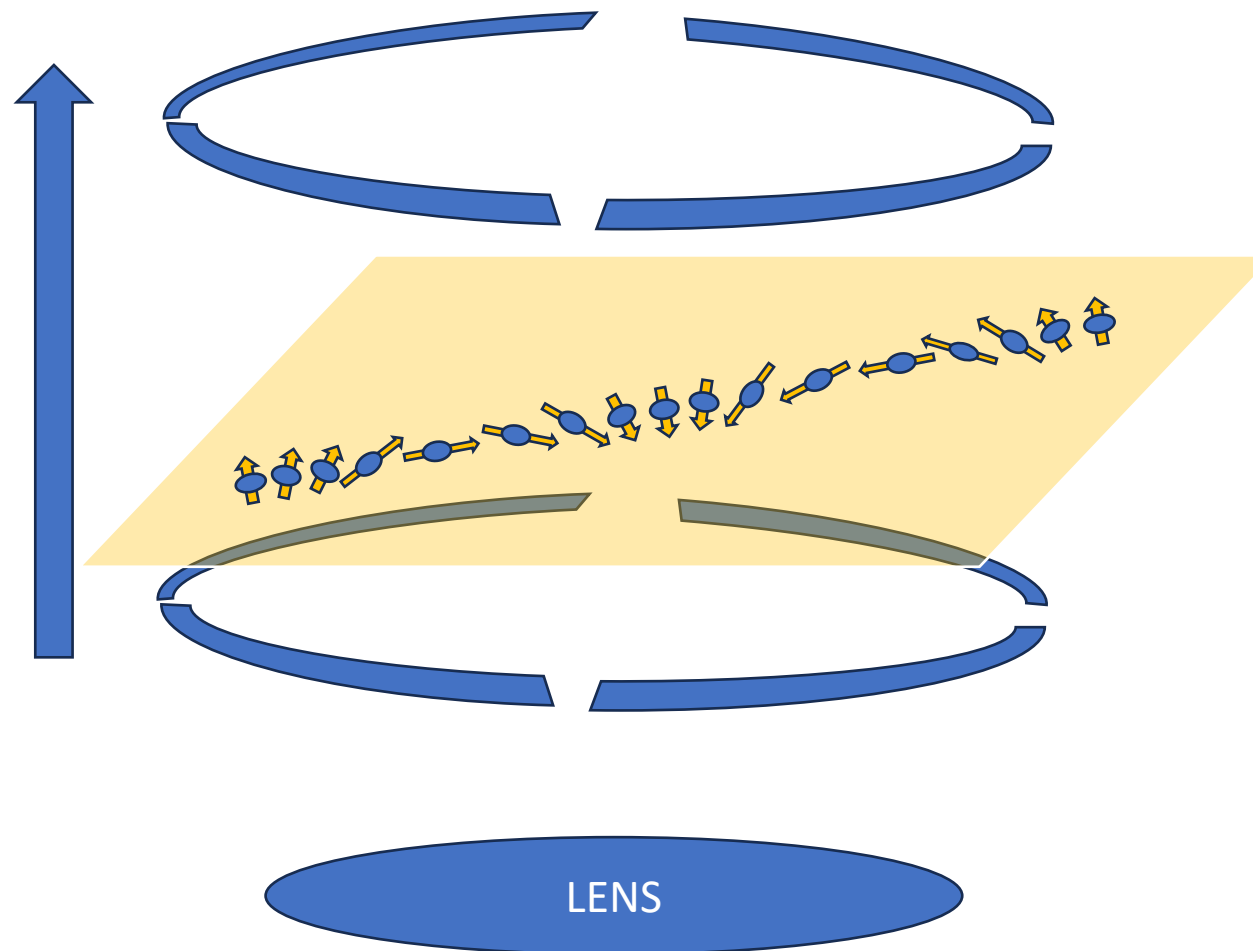
COMBINING THE OFFSET AND GRADIENT



CHARACTERIZING ALONG Z AXIS



FINAL DIRECTION OF SPIN SPIRAL

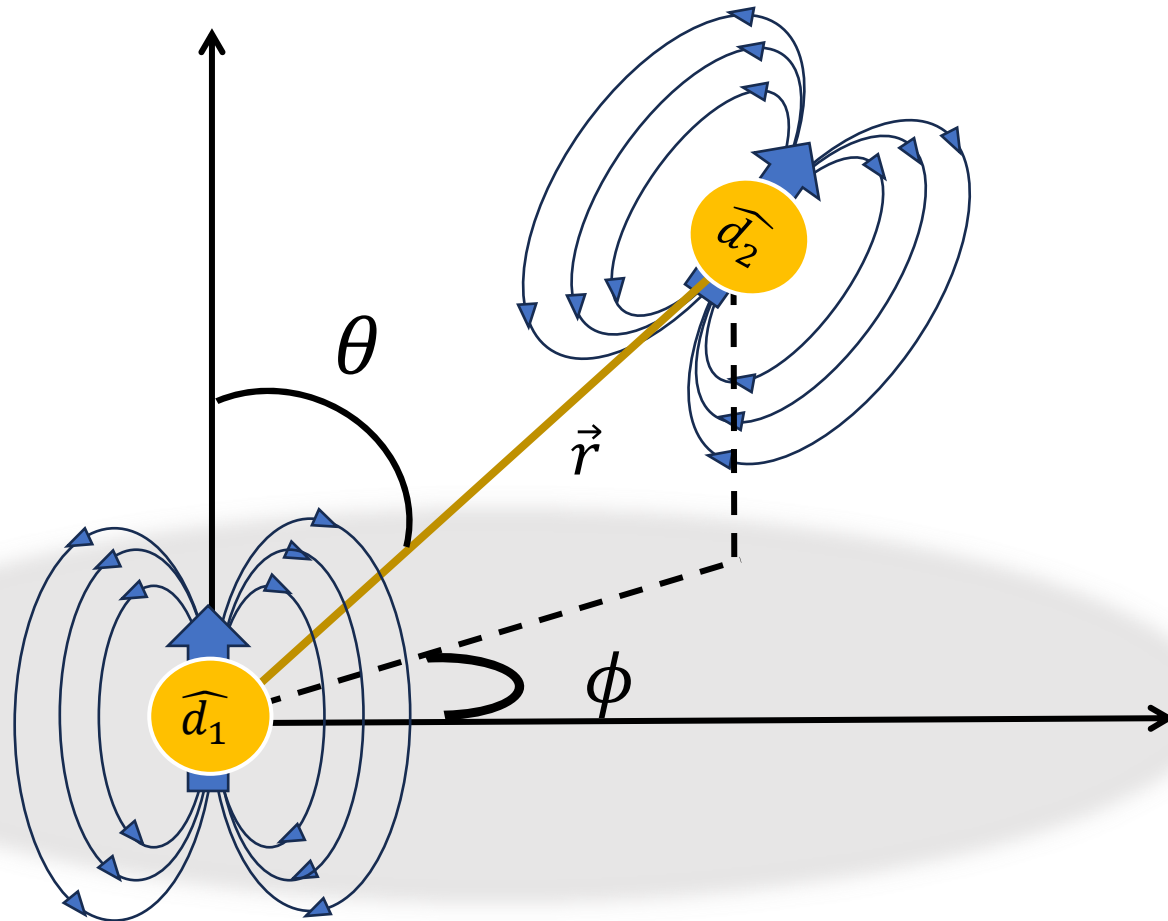




4. MAPPING TO A SPIN HAMILTONIAN

INTERACTION HAMILTONIAN FOR TWO ATOMS

To first order, as if dipoles are interacting



$$H_{Total} = H_1 \otimes \mathbb{I} + \mathbb{I} \otimes H_2 + H_{int}$$

$$H_{int} = \frac{1}{4\pi\epsilon_0} \frac{\widehat{d}_1 \cdot \widehat{d}_2 - 3(\widehat{d}_1 \cdot \hat{r})(\widehat{d}_2 \cdot \hat{r})}{|\vec{r}|^3}$$

HAMILTONIAN OF THE TWO ATOM SYSTEM

Using $\widehat{d}_{\pm} = \mp \frac{1}{\sqrt{2}} (\widehat{d}_x \pm i \widehat{d}_y)$ and $\widehat{d}_0 = \widehat{d}_z$

$$H_I(\theta, \phi, \vec{r}) = \frac{H_{dd}(\theta, \phi)}{4\pi\epsilon_0 |\vec{r}|^3}$$

$$\begin{aligned} & H_{dd}(\theta, \phi) \\ &= \frac{1 - 3 \cos^2(\theta)}{2} [2 \widehat{d}_{10} \cdot \widehat{d}_{20} + \widehat{d}_{1+} \cdot \widehat{d}_{2-} + \widehat{d}_{1-} \cdot \widehat{d}_{2+}] \quad \leftarrow \text{No azimuthal dependence} \\ &- \frac{3 \sin^2(\theta)}{2} [\widehat{d}_{1+} \cdot \widehat{d}_{2+} e^{-i\phi} + \widehat{d}_{1-} \cdot \widehat{d}_{2-} e^{i\phi}] \\ &- \frac{3 \sin^2(\theta)}{2} [(\widehat{d}_{1+} \cdot \widehat{d}_{20} + \widehat{d}_{10} \cdot \widehat{d}_{2+}) e^{-i\phi} + (\widehat{d}_{1-} \cdot \widehat{d}_{20} + \widehat{d}_{10} \cdot \widehat{d}_{2-}) e^{i\phi}] \end{aligned} \quad \left. \vphantom{\begin{aligned} &- \frac{3 \sin^2(\theta)}{2} [\widehat{d}_{1+} \cdot \widehat{d}_{2+} e^{-i\phi} + \widehat{d}_{1-} \cdot \widehat{d}_{2-} e^{i\phi}] \\ &- \frac{3 \sin^2(\theta)}{2} [(\widehat{d}_{1+} \cdot \widehat{d}_{20} + \widehat{d}_{10} \cdot \widehat{d}_{2+}) e^{-i\phi} + (\widehat{d}_{1-} \cdot \widehat{d}_{20} + \widehat{d}_{10} \cdot \widehat{d}_{2-}) e^{i\phi}] \right\}} \begin{array}{l} \text{Some terms could be} \\ \text{ignored for} \\ \text{unperturbed states} \end{array}$$

MAPPING TO A SPIN HAMILTONIAN

- |Chosen excited state⟩ = |↑⟩ & |Chosen ground state⟩ = |↓⟩
- Single Atom Hamiltonian = |↓⟩ ⟨↓| + ω|↑⟩ ⟨↑|
- Find H_{int} in the two spin basis : |↑↑⟩, |↑↓⟩, |↓↑⟩, |↓↓⟩
- Write H_{int} rotating frame of non-interacting Hamiltonians

$$\begin{pmatrix} a_{1,1} & a_{1,2} & a_{1,3} & a_{1,4} \\ a_{1,2}^* & a_{2,2} & a_{2,3} & a_{2,4} \\ a_{1,3}^* & a_{2,3}^* & a_{3,3} & a_{3,4} \\ a_{1,4}^* & a_{2,4}^* & a_{3,4}^* & a_{4,4} \end{pmatrix} \xrightarrow[\text{Rotating Wave Approximation}]{e^{i(H_1 \otimes \mathbb{I} + \mathbb{I} \otimes H_2)t}} \begin{pmatrix} a_1, & \begin{pmatrix} a_{1,1} & 0 & 0 & 0 \\ 0 & a_{2,2} & a_{2,3} & 0 \\ 0 & a_{2,3}^* & a_{3,3} & 0 \\ 0 & 0 & 0 & a_{4,4} \end{pmatrix} \begin{matrix} \omega t \\ \omega t \\ \omega t \end{matrix} \\ a_{1,2}^* e^{-i\omega t} \\ a_{1,3}^* e^{-i\omega t} \\ a_{1,4}^* e^{-i\omega t} \end{pmatrix}$$

MAPPING TO A SPIN HAMILTONIAN

$$H_{int} = \frac{1}{4\pi\epsilon_0|\vec{r}|^3} \times \begin{pmatrix} \langle \uparrow\uparrow | H_{dd} | \uparrow\uparrow \rangle & \langle \uparrow\uparrow | H_{dd} | \downarrow\uparrow \rangle & \langle \uparrow\uparrow | H_{dd} | \downarrow\downarrow \rangle & \langle \uparrow\uparrow | H_{dd} | \uparrow\downarrow \rangle \\ 0 & \langle \downarrow\uparrow | H_{dd} | \uparrow\uparrow \rangle & \langle \downarrow\uparrow | H_{dd} | \downarrow\uparrow \rangle & \langle \downarrow\uparrow | H_{dd} | \downarrow\downarrow \rangle \\ 0 & \langle \downarrow\uparrow | H_{dd} | \uparrow\downarrow \rangle & \langle \downarrow\uparrow | H_{dd} | \uparrow\uparrow \rangle & \langle \downarrow\uparrow | H_{dd} | \downarrow\downarrow \rangle \\ 0 & \langle \downarrow\downarrow | H_{dd} | \uparrow\uparrow \rangle & \langle \downarrow\downarrow | H_{dd} | \downarrow\uparrow \rangle & \langle \downarrow\downarrow | H_{dd} | \downarrow\downarrow \rangle \end{pmatrix}$$

$$\equiv J^\perp (S_1^x \otimes S_2^x + S_1^y \otimes S_2^y) + J^\parallel (S_1^z \otimes S_2^z) + h_z (S_1^z \otimes \mathbb{I} + \mathbb{I} \otimes S_2^z) + V \mathbb{I}$$

$$J^\parallel = E_{\downarrow\downarrow} + E_{\uparrow\uparrow} - 2E_{\downarrow\uparrow}$$

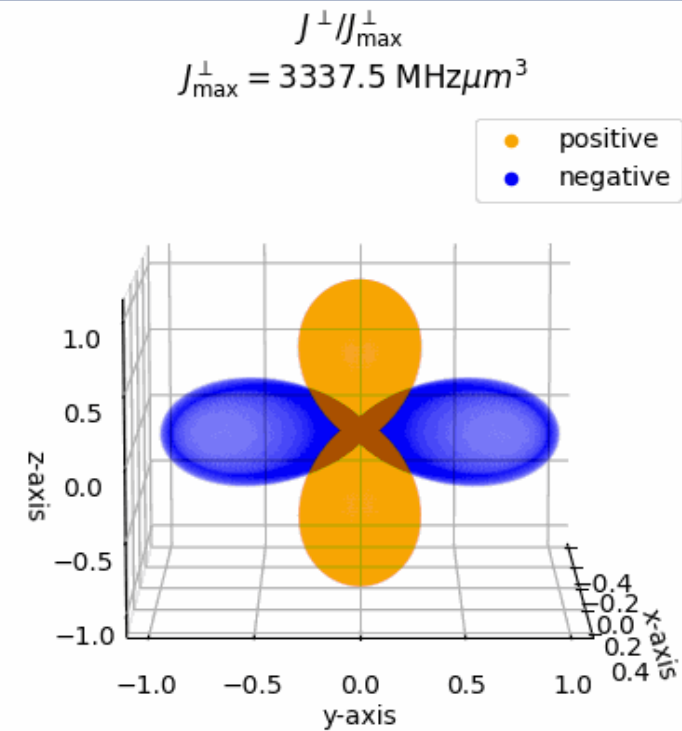
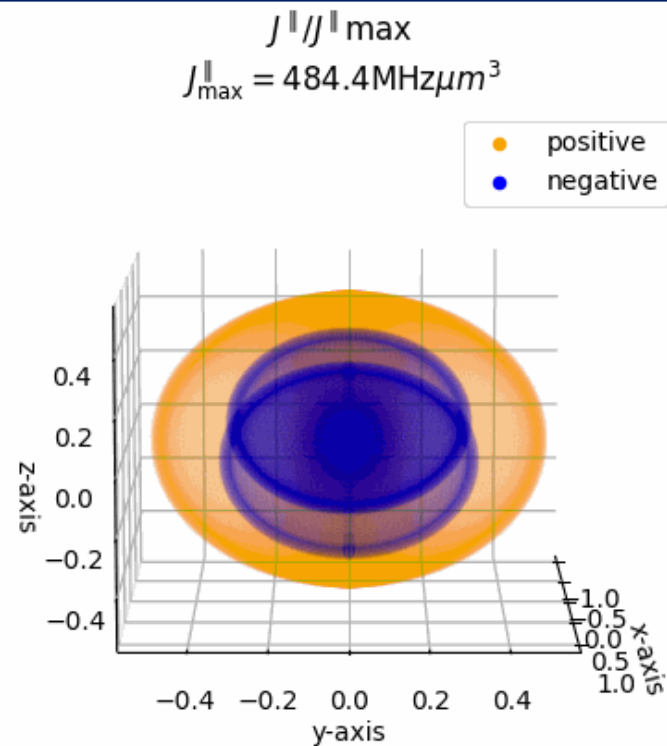
$$J^\perp = 2J$$

$$h_z = \frac{E_{\downarrow\downarrow} - E_{\uparrow\uparrow}}{2}$$

$$V = \frac{E_{\uparrow\uparrow} + E_{\downarrow\downarrow} + 2E_{\downarrow\uparrow}}{4}$$

ANGULAR DEPENDENCE OF J^{\parallel} AND J^{\perp}

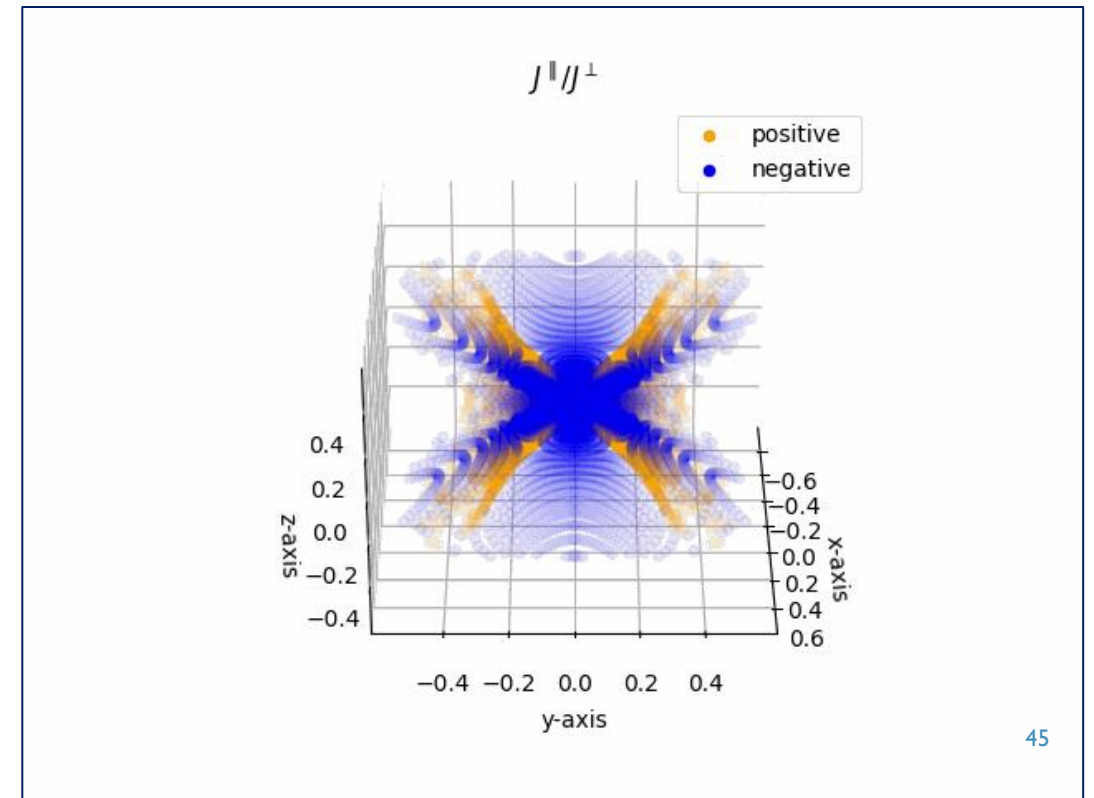
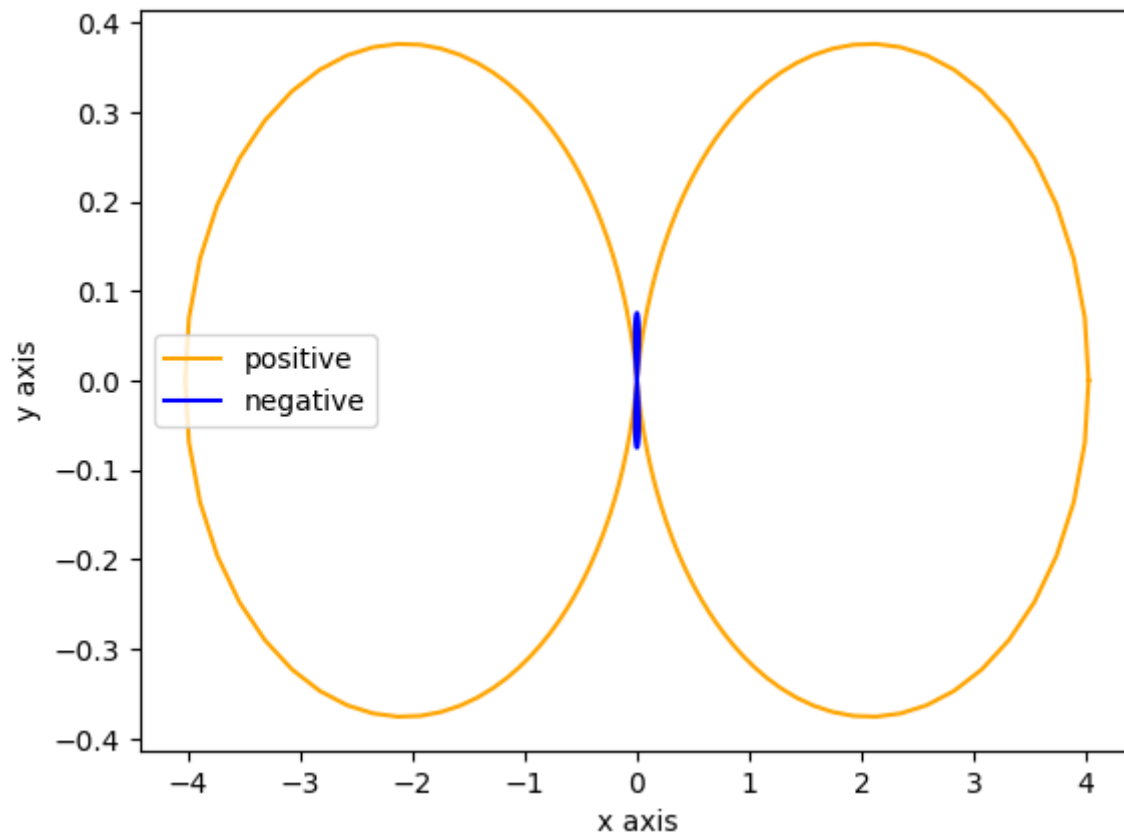
For $|48 s ; J = 0.5 ; m_J = 0.5 \rangle \rightarrow |48 p ; J = 1.5 ; m_J = 1.5 \rangle$



ANGULAR DEPENDENCE OF $J^{\parallel} / J^{\perp}$

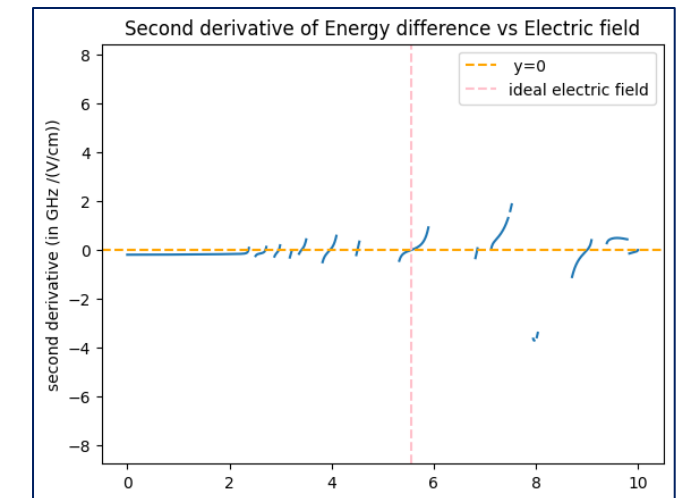
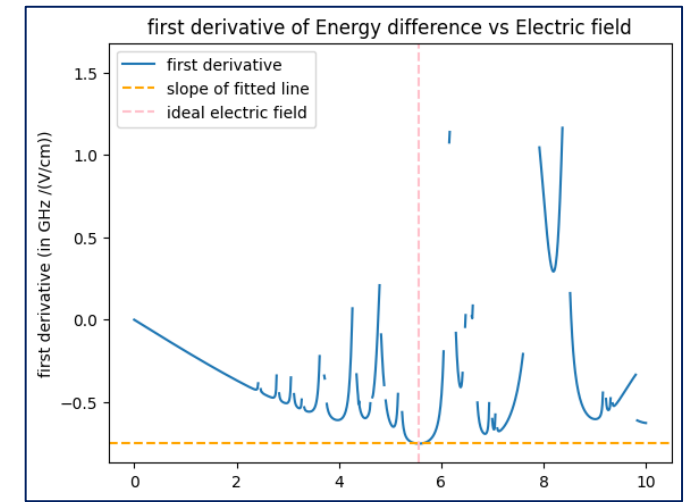
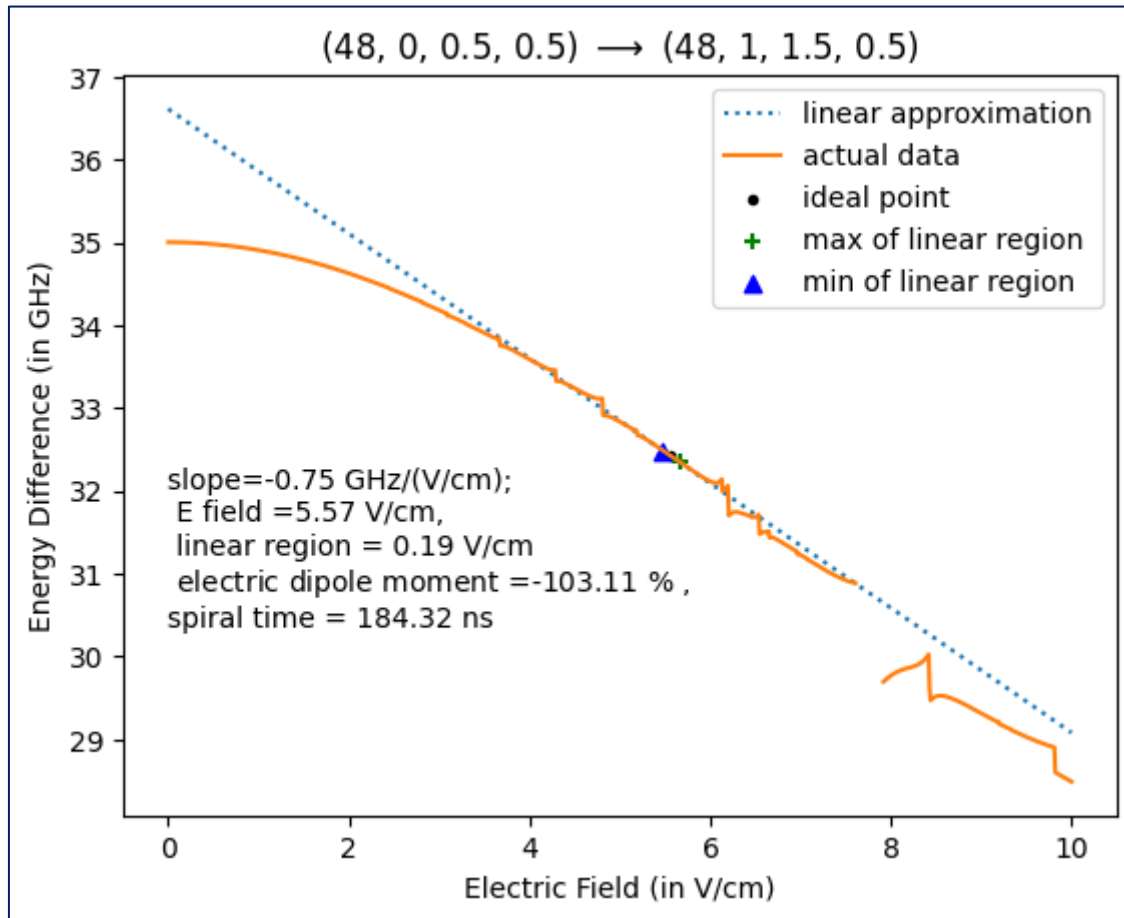
For $|48 s ; J = 0.5 ; m_J = 0.5 \rangle \rightarrow |48 p ; J = 1.5 ; m_J = 1.5 \rangle$

$J^{\parallel} / J^{\perp}$ in X-Y plane



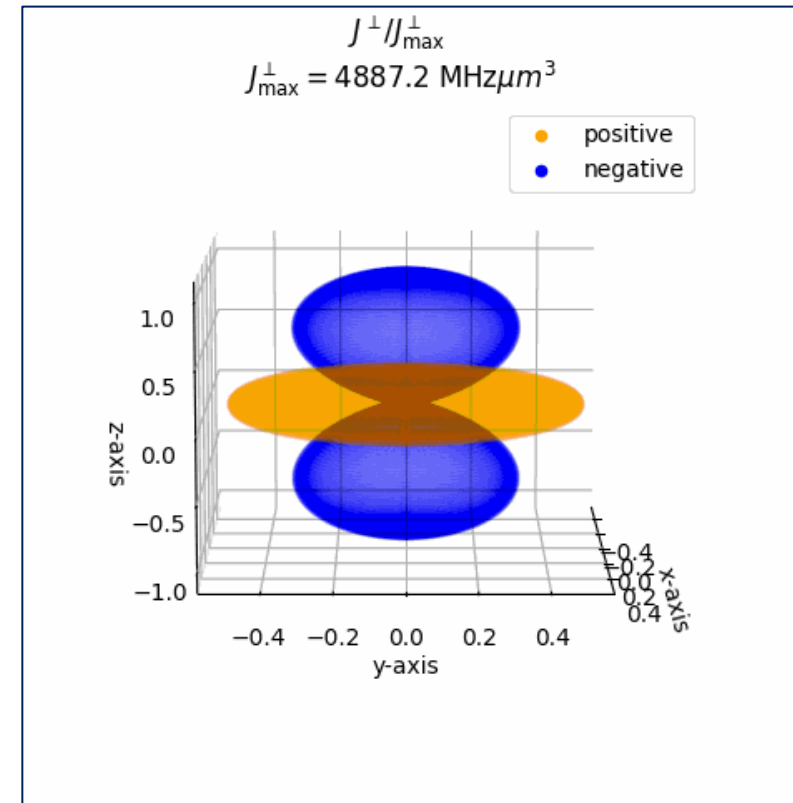
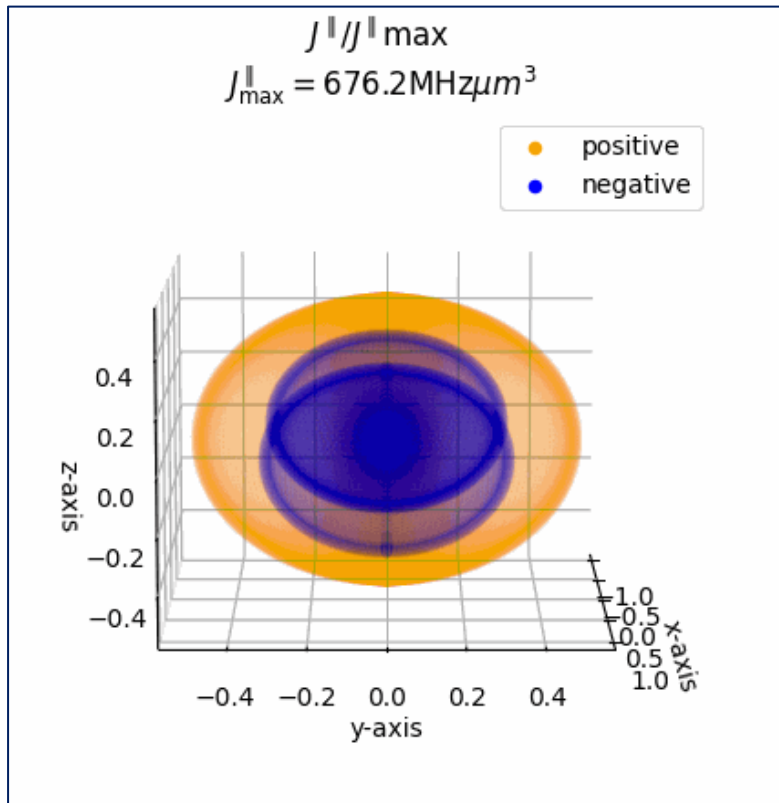
ANOTHER TRANSITION EXAMPLE

For $|48 s ; J = 0.5 ; m_J = 0.5 \rangle \rightarrow |48 p ; J = 1.5 ; m_J = 0.5 \rangle$



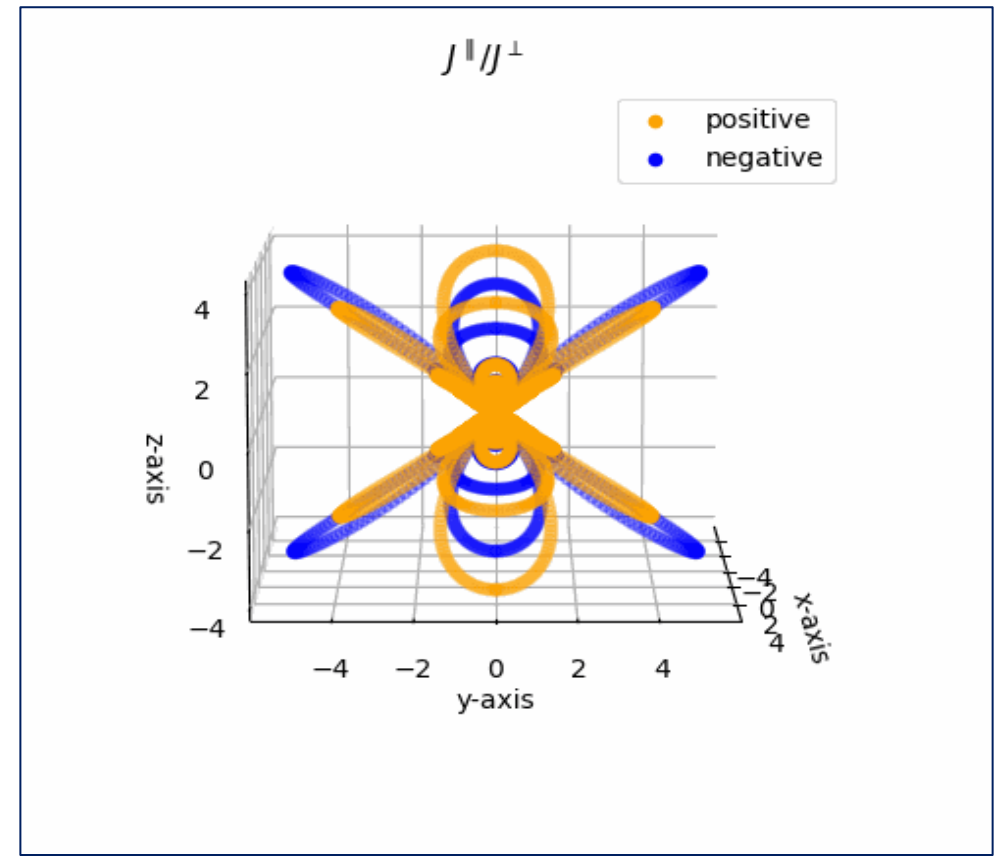
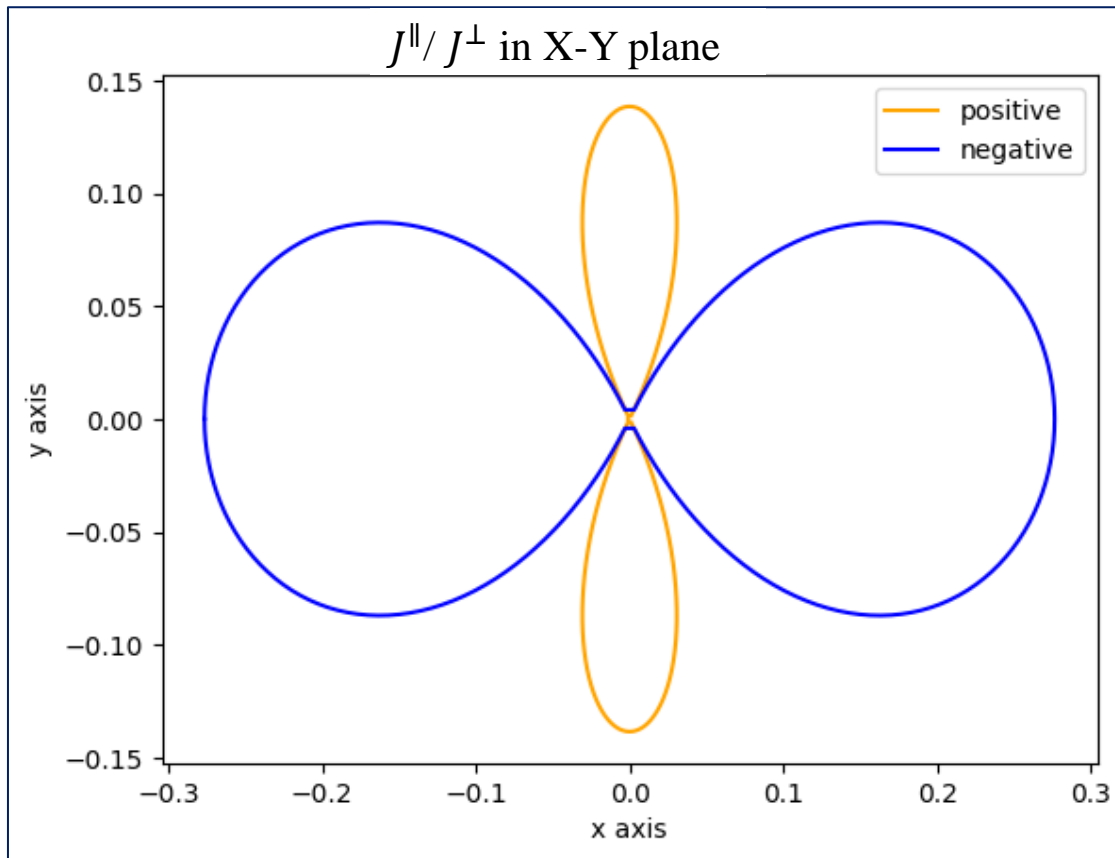
ANGULAR DEPENDENCE OF J^{\parallel} AND J^{\perp}

For $|48 s ; J = 0.5 ; m_J = 0.5 \rangle \rightarrow |48 p ; J = 1.5 ; m_J = 0.5 \rangle$



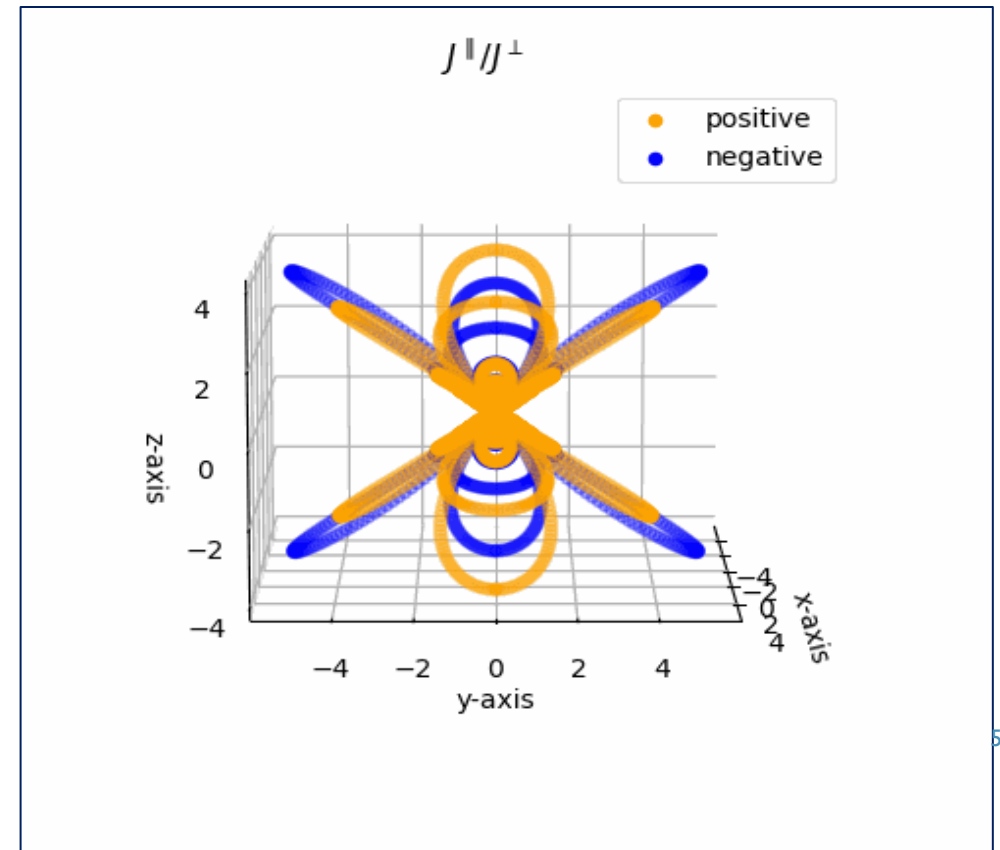
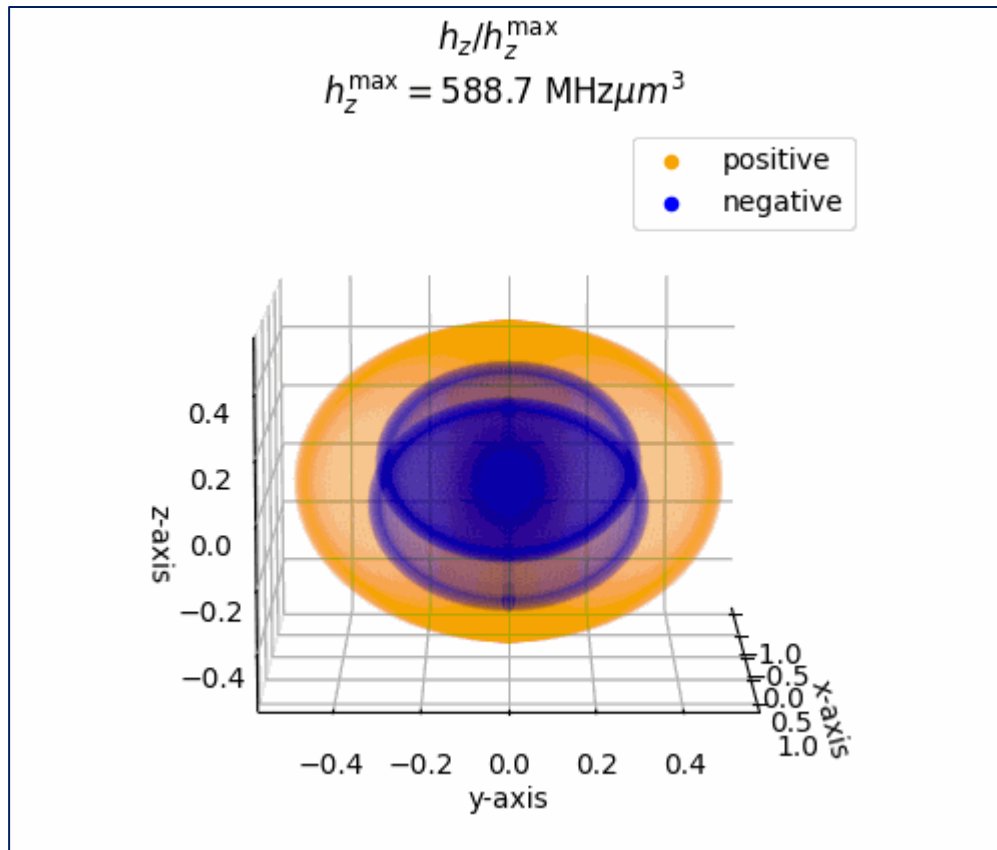
ANGULAR DEPENDENCE OF J^{\parallel}/J^{\perp}

For $|48 s ; J = 0.5 ; m_J = 0.5 \rangle \rightarrow |48 p ; J = 1.5 ; m_J = 0.5 \rangle$



ANGULAR DEPENDENCE OF h_z AND J^{\parallel}/J^{\perp}

For $|48 s ; J = 0.5 ; m_J = 0.5 \rangle \rightarrow |48 p ; J = 1.5 ; m_J = 0.5 \rangle$



OUTLOOK

- Can implement different Hamiltonians for different levels
- Analyze changes with principal quantum number
- Analyze effect of change in direction of fields on Spin model
- Look for Forster resonances

THANK YOU !



For making me a part of the Rydberg Team!