

Finite element method for elliptic problems.

MA31-Numerical analysis of Partial Derivative Equations: Courses 11-12

Manuel Samuelides

Institut Supérieur de l'Aéronautique et de l'Espace
Toulouse, France

October 2013

Outline

- 1 Introduction
- 2 Approximation of Poisson equation
 - Poisson problem
 - Local approximation
 - Internal approximation
- 3 Implementing FEM
 - Linear system assembling

Objective of the course

- know the principle of finite element method approximation,
- be able to use it to solve a variational problem,
- be able to control approximation errors of the method.

Remark

This course is widely inspired from a similar course at "Ecole des Mines, ParisTech" from professor Michel Kern, senior searcher at INRIA who is quoted in reference.

Outline

- 1 Introduction
- 2 Approximation of Poisson equation
 - Poisson problem
 - Local approximation
 - Internal approximation
- 3 Implementing FEM
 - Linear system assembling

Poisson problem in linear elastic deformation

- Let ϵ be the strain tensor which is the gradient of the displacement field u : $\epsilon = \nabla(u)$
- The Hooke law is governing the **stress tensor** σ : $\sigma = C\epsilon$ where C is the stiffness tensor ,[1].
- The elastic equilibrium of a solid in the hypothesis of small strains is given by the linear law: $\vec{\nabla} \cdot \sigma + \rho f = 0$ [3]. So we get the Navier equation (vector version of Poisson equation) $\Delta u + \rho f = 0$
- The boundary conditions on $\partial\Omega$ may be of two kinds
 - Fixed position constraint gives Dirichlet boundary condition on $\partial\Omega_D$: $\forall x \in \partial\Omega_D, u(x) = 0$
 - Applied force governs the stress on $\partial\Omega_N$ and gives Neumann boundary conditions $\forall x \in \partial\Omega_D, \sigma(x) \cdot \vec{n} = g(x)$

Linear elastic deformation equation

We shall study here as a simple example the linear elastic deformation of a 2d structure with a normal force (traction or compression). It gives the following Poisson problem with mixed boundary conditions:

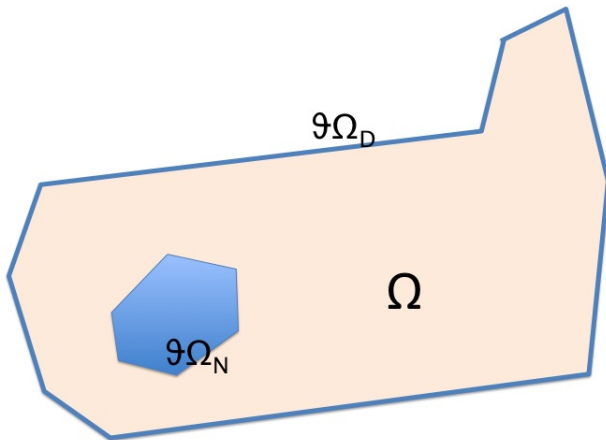
Problem

We give the external forces: the density one $f \in L^2(\Omega)$ and the Neumann boundary condition $g \in L^2(\partial\Omega_N)$

$$\begin{cases} \forall x \in \Omega, & -\vec{\nabla} \cdot \vec{\nabla}(ku)(x) = f \\ \forall x \in \partial\Omega_D, & u(x) = 0 \\ \forall x \in \partial\Omega_N, & k \frac{\partial u}{\partial \vec{n}}(x) = g(x) \end{cases}$$

A typical example

The polynomial boundary is given for further local approximation.



Variational formulation

Applying Stokes formula as in the last course, we get the variational formulation:

Problem

Let $\mathcal{H}_D = \{v \in H^1(\Omega) \text{ such that } v = 0 \text{ on } \partial\Omega_D\}$. Then the problem is to find $u \in \mathcal{H}_D$ such that

$$\forall v \in \mathcal{H}_D, a(u, v) = L(v)$$

with

- $a(u, v) = \int_{\Omega} k(x) \vec{\nabla} u(x) \cdot \vec{\nabla} v(x) dx$
- $L(v) = \int_{\Omega} f(x) v(x) dx + \int_{\partial\Omega_N} g(x) v(x) ds(x)$

Recall Lax-Milgram theorem

Theorem

Let a be a bilinear symmetric continuous and coercive form on a Hilbert space \mathcal{H} and L is a linear continuous form on \mathcal{H} :

$$\left\{ \begin{array}{ll} \forall (v, w) \in \mathcal{H} \times \mathcal{H}, & a(v, w) = a(w, v) \\ \forall (v, w) \in \mathcal{H} \times \mathcal{H}, \quad \exists M > 0, & |a(v, w)| \leq M \|v\| \|w\| \\ \forall v \in \mathcal{H}, \quad \exists \alpha > 0, & a(v, v) \geq \alpha \|v\|^2 \\ \forall v \in \mathcal{H}, \quad \exists C > 0, & |L(v)| \leq C \|v\| \end{array} \right.$$

Then $\exists ! u$ such that $\forall v \in \mathcal{H}, a(u, v) = L(v)$ and we have

$$\|u\| \leq \frac{MC}{\alpha}$$

Outline

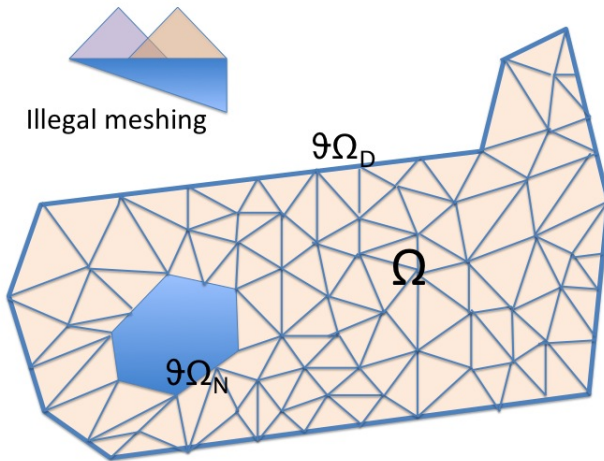
- 1 Introduction
- 2 Approximation of Poisson equation
 - Poisson problem
 - **Local approximation**
 - Internal approximation
- 3 Implementing FEM
 - Linear system assembling

Triangular meshing

- Like in FD approximation scheme, we shall build an approximation meshing.
- Since we use integral formulation no need of a rectangular meshing to compute partial derivative.
- In 2d equations polygonal meshings are used, in 3d equations polyedral meshings.
- We shall use here the simpler one: triangular meshing.
- Triangular meshing is subject to regularity conditions: Two triangles have either no intersection or a point intersection: one single node) or a line intersection: a whole edge and one only.
- Let \mathcal{T} be the set of triangles, $\mathcal{N} = \{x_j\}$ be the set of nodes and \mathcal{N}_D be the subset which is included in $\partial\Omega_D$.

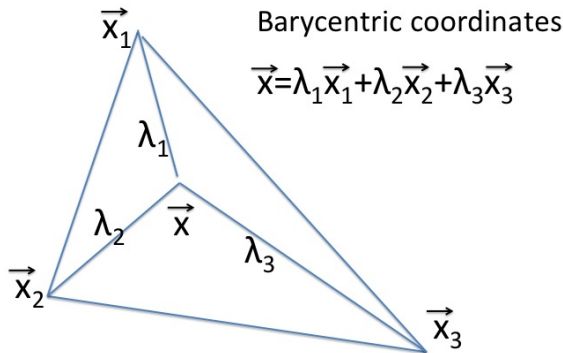
A typical meshing

The polynomial boundary is given for further local approximation.



Parametrization of a triangle

Each point of the convex triangle is defined by three positive barycentric coordinates $(\lambda_1, \lambda_2, \lambda_3)$ subject to $\lambda_1 + \lambda_2 + \lambda_3 = 1$



Piecewise affine functions

- If $T_j = (\vec{x}_1, \vec{x}_2, \vec{x}_3) \in \mathcal{T}$, we consider the 3d vector space \mathcal{F}_j of affine functions generated by the affine functions ϕ_1, ϕ_2, ϕ_3 such that $\phi_i(\vec{x}_j) = \delta_{ij}$.
- If $\vec{x} = \lambda_1 \vec{x}_1 + \lambda_2 \vec{x}_2 + \lambda_3 \vec{x}_3$ then

$$\forall \phi \in \mathcal{F}_j, \phi(\vec{x}) = \lambda_1 \phi_1(\vec{x}_1) + \lambda_2 \phi_2(\vec{x}_2) + \lambda_3 \phi_3(\vec{x}_3)$$

- We shall associate to the meshing the \mathcal{T} the subspace $\mathcal{H}_{\mathcal{T}} \in \mathcal{H}$ of continuous piecewise affine functions on the triangles of \mathcal{T} .
- We shall approximate elements of \mathcal{H} by elements of $\mathcal{H}_{\mathcal{T}}$ with a uniform convergence when the diameter h of the meshing goes to 0.

Local affine approximation

Consider a continuous function u on Ω and a triangle $\mathcal{T} = (\vec{x}_1, \vec{x}_2, \vec{x}_3) \in \mathcal{T}$, and the affine function on \mathcal{T} , $\phi \in \mathcal{F}$ defined by

$$\phi(\lambda_1 \vec{x}_1 + \lambda_2 \vec{x}_2 + \lambda_3 \vec{x}_3) = \lambda_1 u(\vec{x}_1) + \lambda_2 u(\vec{x}_2) + \lambda_3 u(\vec{x}_3)$$

Then

- ϕ is the unique affine interpolation of u on the three summits of \mathcal{T} .
- Moreover since u is uniformly continuous on the bounded domain Ω , when $h \rightarrow 0$, the affine interpolation ϕ of u converges uniformly towards u

Outline

- 1 Introduction
- 2 Approximation of Poisson equation
 - Poisson problem
 - Local approximation
 - Internal approximation
- 3 Implementing FEM
 - Linear system assembling

Space of internal approximation

Definition

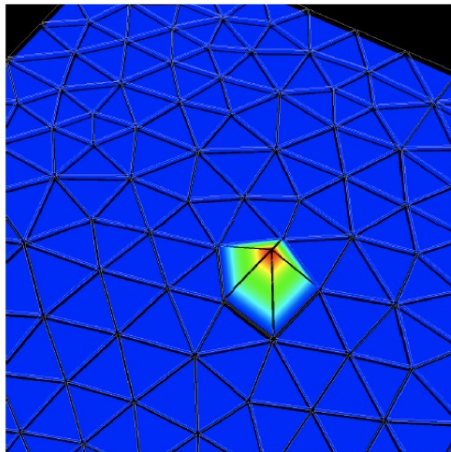
Let $(\mathcal{T}, \mathcal{N})$ be a triangular meshing of Ω with diameter h , we define by $\mathcal{H}_{\mathcal{T}}$ the subspace of piecewise affine functions and we set $\mathcal{H}_{\mathcal{T},D} = \mathcal{H}_{\mathcal{T}} \cap \mathcal{H}_D$.

Let $\vec{x}_i \in \mathcal{N}$, and let ϕ_i be the only function of $\mathcal{H}_{\mathcal{T}}$ such that $\forall \vec{x}_j \in \mathcal{N}, \phi_i(\vec{x}_j) = \delta_{ij}$.

The interpolation of u in \mathcal{H} is defined by the interpolation operator $I_{\mathcal{T}}$:
$$I_{\mathcal{T}}u = \sum_{i \in \mathcal{N}} u(\vec{x}_i) \phi_i$$

- The set of (ϕ_i) is clearly a basis of $\mathcal{H}_{\mathcal{T}}$ since a continuous piecewise affine function is equal to its interpolation.
- We have a uniform convergence of the interpolates $I_{\mathcal{T}}u$ when the diameter h of \mathcal{T} goes to 0.
- Let $\mathcal{N}_D^c = (\mathcal{N}_D)^c$, the set $(\phi_i)_{i \in \mathcal{N}_D^c}$ is a basis of $\mathcal{H}_{\mathcal{T},D}$

Representing one ϕ_i



Internal approximation of a variational formulation

Consider the variational problem on the Hilbert space \mathcal{H} with the bilinear symmetric continuous coercive form a and the linear continuous form L and a subspace \mathcal{H}_h .

We intend to approximate the solution u of

$$\text{Find } u \in \mathcal{H} \text{ such that } \forall v \in \mathcal{H}, a(u, v) = L(v)$$

by the solution u_h of an internal approximate problem

$$\text{Find } u_h \in \mathcal{H}_h \text{ such that } \forall v \in \mathcal{H}_h, a(u, v) = L(v)$$

We show the following projection theorem

Theorem

*The bilinear form a defines an equivalent norm on the Hilbert space \mathcal{H} (**the energy norm**). The approximate solution u_h is the orthogonal projection for this norm of \mathcal{H} onto \mathcal{H}_h .*

Proof of the projection theorem

- The properties of the energy product a gives

$$\forall v \in \mathcal{H}, \alpha \|v\|^2 \leq a(v, v) \leq M \|v\|^2$$

So the energy norm is equivalent to the original norm.

- We have

$$\forall v \in \mathcal{H}_h, a(u, v) = a(u_h, v) \Rightarrow a(u - u_h, v) = 0$$

It proves that u_h is the orthogonal projection of u onto \mathcal{H}_h .

- More over the projection property

$$a(u - u_h, u - h) = \inf_{v \in \mathcal{H}_h} a(u - v, u - v) \text{ implies}$$

$$\alpha \|u - u_h\|^2 \leq M \inf_{v \in \mathcal{H}_h} \|u - v\|^2$$

We use it to show the convergence of internal approximation.

Outline

- 1 Introduction
- 2 Approximation of Poisson equation
 - Poisson problem
 - Local approximation
 - Internal approximation
- 3 Implementing FEM
 - Linear system assembling

Linear system assembling

The solution of the approximated variational problem

Problem

Find $u \in \mathcal{H}_{\mathcal{T},D}$ such that $\forall v \in \mathcal{H}_{\mathcal{T},D}, a(u, v) = L(v)$

amounts to solve the linear system $\forall j \in \mathcal{N}, \sum_{i \in \mathcal{N}} a_{ji} u_i = b_j$
with $a_{i,j} = a(\phi_i, \phi_j)$ and $b_j = L(\phi_j)$

- If the nodes i and j are not the summits of one common element $a_{i,j} = 0$
- The computation is done element by element to avoid complexity.
- Each element adds its contribution to the $a_{i,j}$ and the b_j associated to the summits of that element.

Dirichlet boundary conditions

- The previous linear system does not take into account the Dirichlet boundary conditions. Actually the values of the functions $\phi_i \in \partial\Omega_D$ are constrained to 0 (in case of Dirichlet homogeneous boundary condition on $\partial\Omega_D$)
- If we separate from the other nodes the nodes of \mathcal{N}_D , the matrix decomposition of the previous system

$$\begin{pmatrix} A_{II} & A_{ID} \\ A_{DI} & A_{DD} \end{pmatrix} \begin{pmatrix} u_I \\ u_D \end{pmatrix} = \begin{pmatrix} b_I \\ b_D \end{pmatrix}$$

becomes

$$\begin{pmatrix} A_{II} & 0 \\ 0 & A_{DD} \end{pmatrix} \begin{pmatrix} u_I \\ u_D \end{pmatrix} = \begin{pmatrix} b_I - A_{ID}g_D \\ g_D \end{pmatrix}$$

- the matrix A is called the **stiffness matrix** of the meshing.

For Further Reading I



Yves Debard

Méthode des éléments finis: élasticité plane .

IUT du Mans, 2006-2007(free internet available).



Michel Kern

Introduction à la méthode des éléments finis .

Ecole Nationale Supérieure des Mines de Paris,
2004-2005(free internet available).



Christian Weilgosz, Bernard Peseux, Yves Lecointe

Formulations mathématiques et résolution numérique en mécanique .

M2R Université de Nantes, 2004(free internet available).

For Further Reading II



Stéphanie Salmon

Calcul scientifique .

M1 Université de Strasbourg, 2009 (free internet available).