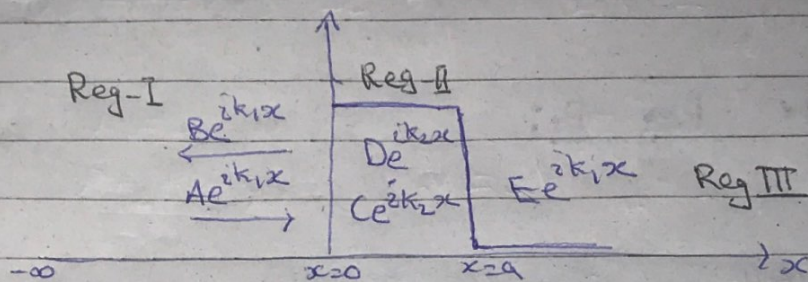


Potential Barrier:-



Case 1 $E > V_0$:-

Beam of Particles of mass 'm'

$$V(x) = \begin{cases} 0 & x < 0 \\ V_0 & 0 < x \leq a \\ 0 & x > a \end{cases}$$

For Reg I:-

From Schrodinger wave eq:

$$V(x) = 0 \quad x < 0$$

$$\frac{d^2 \psi_1}{dx^2} + \frac{2m}{\hbar^2} (E - V) \psi_1 = 0$$

$$|V=0|$$

$$\frac{d^2 \psi_1}{dx^2} + \frac{2m}{\hbar^2} (E - 0) \psi_1 = 0$$

where,

$$\frac{2mE}{\hbar^2} = K_1^2$$

$$\frac{d^2 \psi_1}{dx^2} + K_1^2 \psi_1 = 0 \quad \longrightarrow (1)$$

By solving this we get:

$$\psi_1 = Ae^{2K_1x} + Be^{-2K_1x} \longrightarrow \textcircled{A}$$

For Reg 2:

$$V(x) = V_0 \quad 0 < x \leq a$$

$$\frac{d^2\psi_2}{dx^2} + \frac{2m(E-V_0)}{\hbar^2}\psi_2 = 0$$

where;

$$\frac{2m(E-V_0)}{\hbar^2} = K_2^2$$

where;

$$\frac{2m(E-V_0)}{\hbar^2} = K_2^2$$

$$\frac{d^2\psi_2}{dx^2} + K_2^2\psi_2 = 0 \longrightarrow \textcircled{ii}$$

By solving this we have:

$$\psi_2 = Ce^{2K_2x} + De^{-2K_2x} \longrightarrow \textcircled{B}$$

For Reg 3:

$$V(x) = 0 \quad x > a$$

$$\frac{d^2\psi_3}{dx^2} + \frac{2m(E-V)}{\hbar^2}\psi_3 = 0$$

$$\frac{d^2\psi_3}{dx^2} + \frac{2m(E-0)}{\hbar^2}\psi_3 = 0$$

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$$\frac{d^2\psi_3}{dx^2} + \frac{2mE}{\hbar^2} \psi_3 = 0$$

where;

$$\frac{2mE}{\hbar^2} = K_1^2$$

So,

$$\frac{d^2\psi_3}{dx^2} + K_1^2 \psi_3 = 0 \rightarrow \text{(iii)}$$

By solving this we have:

$$\psi_3 = E e^{iK_1 x} + F e^{-iK_1 x} \rightarrow \text{(C)}$$

* where in eq (A) 'A' gave the information of incident wave in Reg-I, and 'B' gave the information of reflected wave in Reg-I.

* In eq (B) 'C' gave the information of transmitted wave in Reg-II, and 'D' gave information of reflected wave in Reg-II.

* In eq (C) 'E' gave information of transmitted wave in Reg-III and 'F' gave information of reflected wave in Reg-III. In Reg III Here is no reflected wave so, $F e^{-iK_1 x} = 0$

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Find the values of Constant:

To find A, B, C, D, E we will apply Boundary conditions:

at $x=0$

$$\psi_1 = \psi_2$$

at $x=a$

$$\psi_2 = \psi_3$$

$$\frac{d\psi_1}{dx} = \frac{d\psi_2}{dx}$$

$$\frac{d\psi_2}{dx} = \frac{d\psi_3}{dx}$$

For eq A, B at $x=0$

when:

$$\psi_1 = \psi_2$$

By Comparing:

$$Ae^{2k_1x} + Be^{-2k_1x} = Ce^{2k_2x} + De^{-2k_2x}$$

$$Ae^{2k_1(0)} + Be^{-2k_1(0)} = Ce^{2k_2(0)} + De^{-2k_2(0)}$$

$$A + B = C + D \rightarrow \textcircled{D}$$

when: $\frac{d\psi_1}{dx} = \frac{d\psi_2}{dx}$

Taking Differentiate of eq A & B

$$\psi_1 = Ae^{2k_1x} + Be^{-2k_1x}$$

$$\frac{d\psi_1}{dx} = 2k_1Ae^{2k_1x} - 2k_1Be^{-2k_1x} \rightarrow \textcircled{E}$$

$$\psi_2 = Ce^{2k_2x} + De^{-2k_2x}$$

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$$\frac{d\psi_2}{dx} = ik_2 C e^{ik_2 x} - ik_2 D e^{-ik_2 x} \rightarrow (F)$$

Comparing eq (E) & (F)

$$ik_1 A e^{ik_1 x} - ik_1 B e^{-ik_1 x} = ik_2 C e^{ik_2 x} - ik_2 D e^{-ik_2 x}$$

$$ik_1 A e^{ik_1(a)} - ik_1 B e^{-ik_1(a)} = ik_2 C e^{ik_2(a)} - ik_2 D e^{-ik_2(a)}$$

$$ik_1 A - ik_1 B = ik_2 C - ik_2 D \rightarrow (G)$$

Now For eq 'B' and 'C' at
 $x=a$

when:

$$\psi_2 = \psi_3$$

By Comparing:

$$C e^{ik_2 x} + D e^{-ik_2 x} = E e^{ik_1 x}$$

$$C e^{ik_2(a)} + D e^{-ik_2(a)} = E e^{ik_1(a)} \rightarrow (H)$$

when: $\frac{d\psi_2}{dx} = \frac{d\psi_3}{dx}$

Taking differentiation of eq 'B' & 'C'

$$\psi_2 = C e^{ik_2 x} + D e^{-ik_2 x}$$

$$\frac{d\psi_2}{dx} = ik_2 C e^{ik_2 x} - ik_2 D e^{-ik_2 x} \rightarrow (I)$$

$$\psi_3 = E e^{ik_1 x}$$

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$$\frac{d\psi_3}{dx} = ik_1 E e^{ik_1 x}$$

$$\frac{d\psi_3}{dx} = ik_1 E e^{ik_1 (a)} \rightarrow \textcircled{J}$$

Compare I & J

$$ik_2 C e^{ik_2 a} - ik_2 D e^{-ik_2 a} = ik_1 E e^{ik_1 a} \rightarrow \textcircled{K}$$

$$k_2 (i C e^{ik_2 a} - i D e^{-ik_2 a}) = k_1 i E e^{ik_1 a}$$

$$i C e^{ik_2 a} - i D e^{-ik_2 a} = \frac{k_1}{k_2} i E e^{ik_1 a}$$

$$i (C e^{ik_2 a} - D e^{-ik_2 a}) = \frac{k_1}{k_2} i E e^{ik_1 a}$$

$$C e^{ik_2 a} - D e^{-ik_2 a} = \frac{k_1}{k_2} E e^{ik_1 a} \rightarrow \textcircled{L}$$

Now:

$$C e^{ik_2 a} + D e^{-ik_2 a} = E e^{ik_1 a} \rightarrow \textcircled{H}$$

$$C e^{ik_2 a} - D e^{-ik_2 a} = \frac{k_1}{k_2} E e^{ik_1 a} \rightarrow \textcircled{L}$$

For transmission coefficient:

$$T = \frac{k_1}{k_2} \frac{|E|^2}{|A|^2}$$

$$T = \left[\frac{1 + \frac{V_0^2}{4E(E-V_0)} \sin^2 \left(a \sqrt{\frac{2mV_0}{\hbar^2}} \sqrt{\frac{E}{V_0} - 1} \right)}{1} \right]^{-1}$$

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$$\gamma = a \sqrt{\frac{BmV_0}{\hbar^2}} \quad \frac{E}{V_0} = \epsilon$$

$$T = \left[1 + \frac{\sin^2 \lambda \sqrt{\epsilon - 1}}{4\epsilon(\epsilon - 1)} \right]^{-1}$$

$$T = \left[\frac{1 + \sin^2 \lambda \sqrt{\epsilon - 1}}{4\epsilon(\epsilon - 1)} \right]^{-1}$$

$$R + T = 1$$

$$R = 1 - T$$

$$R = \left[1 + \frac{4\epsilon(\epsilon - 1)}{\sin^2 \lambda \sqrt{\epsilon - 1}} \right]^{-1}$$

Case-I

$$E > V_0$$

$$E \gg V_0 \quad \epsilon \gg 1$$

$$T = 1$$

$$R = 0$$

Case-II

$$T = 1 \quad \text{when:}$$

$$\sin^2 \lambda \sqrt{\epsilon - 1} = 0$$

$$\lambda = \sqrt{\epsilon - 1} = n\bar{\lambda}$$

$$\frac{\sqrt{\epsilon - 1}}{\lambda} = n\bar{\lambda}$$

$$\epsilon = \frac{n^2 \bar{\lambda}^2}{\lambda^2} - 1$$

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$$E = \frac{n^2 \hbar^2 V_0}{a^2 2m V_0} + V_0$$

$$E_n = \frac{n^2 \hbar^2 \pi^2}{2a^2 m} + V_0$$

$$n = 1, 2, 3, \dots$$

Total Transmission is possible when energy of incident particles in integral multiple of $\left(\frac{\hbar^2 \pi^2}{2a^2 m} + V_0 \right)$