Ab initio study of toroidal carbon nanotubes with encapsulated atomic metal loops

Mark T. Lusk and Nathaniel Hamm*

Department of Physics, Colorado School of Mines, Golden, Colorado 80401, USA (Received 6 June 2007; revised manuscript received 10 June 2007; published 19 September 2007)

Toroidal carbon nanotubes can serve as hosts for encapsulated loops of atomic metal wires. Such composite structures have been analyzed using density functional theory for a semiconducting C_{120} torus encapsulating chains of Fe, Au, and Cu atoms. The sheathed metal necklaces form a zigzag structure and drop the highest occupied molecular orbital and/or lowest unoccupied molecular orbital band gap to less than 0.1 eV. The iron composite is ferromagnetic with a magnetic moment essentially the same as that of bcc iron. The azimuthal symmetry of these toroidal composites suggests that they may offer novel electromagnetic properties not associated with straight, metal-encapsulated carbon nanotubes.

DOI: 10.1103/PhysRevB.76.125422 PACS number(s): 73.22.-f, 71.15.Mb, 71.20.Be, 73.21.Hb

I. INTRODUCTION

Fullerene-based toroids combine the novel characteristics of carbon nanotubes with electromagnetic and mechanical properties unique to closed-ring geometries. An intriguing possibility, examined in this paper and shown in Fig. 1, is that atomic chains of metal can be encapsulated within such toroidal carbon nanotubes (TCNs). The carbon torus would serve as a stabilizing sheath for the atomic metal necklaces resulting in rigid, stable structures at finite temperatures with novel paramagnetic or ferromagnetic properties. Beyond applications that exploit the properties of independent rings, assemblies of such ferromagnetic tori might be expected to self-assemble into columnar structures with unusual macroscopic behavior. The electronic structure and mechanical stability of these robust atomic necklaces are investigated using density functional theory (DFT) with the objective of establishing ground state geometries, binding energies, and magnetic moments.

At least four types of single-wall fullerene toroids have been considered since the idea was first proposed 15 years ago. Two variations of type 1 TCNs exploit the curvature induced by pentagonal and septagonal defects to generate polygonal graphitic tori^{1,2} and rings as small as C_{120} have been considered.³ The symmetric type 2 TCNs are toroidal polyhexes⁴ created by joining the ends of single-wall carbon nanotubes.⁵ Perhaps the most technologically promising category of toroid is derived from C₆₀ molecules that are joined with pentagonal and septagonal rings and looped into undulated rings.^{6,7} These are referred to here as type 3 toroids. Finally, Haeckelite tubules, based on a two-dimensional bravais lattice with five-, six-, and seven-member rings, can also be joined at the ends to make the type 4 TCNs.^{6,7} The first two types of TCNs are theoretical constructs that are thought to be mechanically stable. It has been estimated that TCNs of type 1 are, in fact, more stable than C₆₀ molecules,³ while type 2 TCNs are expected to be stable provided they have a sufficiently large toroidal diameter.^{5,9} Carbon toroids of types 3 and 4, on the other hand, have been realized experimentally and are demonstrably stable.^{6,7}

For all toroids, delocalized π electrons can be induced into ballistic, persistent motion by the application of a magnetic field perpendicular to the plane of the torus. The resulting orbital magnetic response of tori can be substantial and has been investigated using tight binding theory with the

magnetic vector potential accounted for using London theory. ^{10,11} Type 1 toroids are predicted to exhibit a large diamagnetic response, ^{5,12,13} while a class of conducting type 2 TCNs are expected to be able to generate colossal paramagnetic moments. ^{5,14,15} More recently, carbon toroids of types 3 and 4 TCNs have been shown experimentally to have paramagnetic behavior consistent with tight binding model predictions. ⁶ The septagonal and pentagonal rings in these structures are predicted to result in ferromagnetic behavior at sufficiently low temperatures in the absence of any flaws. All types of TCNs are expected to exhibit Aharanov-Bohm effects. ¹³

Paralleling activity associated with toroidal fullerene structures have been advanced in the fabrication and analysis of isolated atomic chains of metal atoms. Metal contacts created with a scanning tunneling microscope can be slowly pulled apart to extrude an atomic chain of Au atoms. 16,17 Stable monatomic chains of Au atoms can support currents of 80 μA —clear evidence of ballistic transport in this alternate manifestation of one-dimensional conductors 16 —and the quantum conductance in these metallic atomic chains has also been measured. 17 A subsequent DFT analysis of atomic chains of Au atoms, intended to address observed anomalies, indicated that zigzag structures are more stable than linear

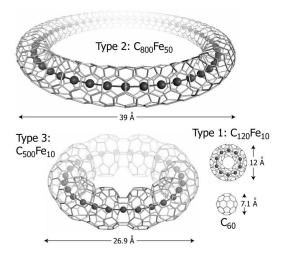


FIG. 1. Atomic chains of Fe atoms encapsulated within three types of toroidal carbon nanotubes. A bucky ball is shown at lower right for size comparison. All structures are drawn to scale.

chains as a result of reduced transverse kinetic energy of electrons. ¹⁸ Such configurations are most likely metastable since previous investigations have identified polyhedral ground state structures for small clusters of Fe and Au atoms. ^{19,20}

Iron structures sheathed within straight single-wall carbon nanotubes have been studied both computationally and experimentally. 21,22 A DFT analysis of atomic chains of Fe within single-wall carbon nanotubes concluded that the spacing between Fe atoms and their nearest C neighbors determined the character of the Fe wire. 21 The Fe–C bond length of 1.9 Å associated with (3,3) tubes resulted in a composite with semiconductor character and no magnetic moment. On the other hand, an Fe–C bond length of 2.1 Å associated with (4,4) tubes resulted in a composite with the metallic character and magnetic moment of $2.6\mu_B$ per atom. 21 The geometry of such tubes precludes the generation of axial persistent currents and a colossal paramagnetic response, but they might be used to help predict the conduction and ferromagnetic nature of wires encapsulated within TCNs.

The present work brings together the unusual behavior of toroidal carbon nanotubes and atomic metallic chains within a simple setting: a type 1 C_{120} torus with D_{5d} symmetry² encapsulating closed, monatomic chains of Au, Fe, and Cu atoms. The iron necklace is of particular importance because of the possibility that it has a ferromagnetic nature at finite temperatures. The gold necklace was also considered,

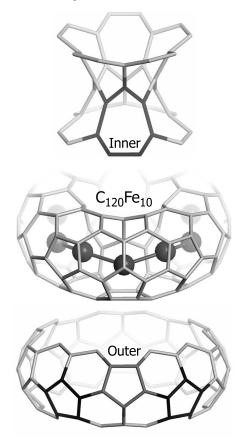


FIG. 2. Break-out view of the inner and outer structures of the $C_{120}Fe_{10}$ molecule with shading used to highlight polygonal defect types.

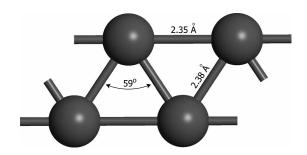


FIG. 3. Ground state structure for a periodic chain of Fe atoms.

though, because of the focus that Au has received in atomic metal wires. 16,17 Analysis of Cu atoms was thought to offer an intermediate composite since it is a noble metal, such as Au, and a period 4 metal, such as Fe. The type 1 TCN chosen for analysis is a polygonal C_{120} torus with D_{5d} symmetry chosen for its predicted stability and minimal size.² The C₁₂₀ torus is composed of an inner surface of 10 septagons and an outer surface of 10 pentagons patched with 20 hexagons, as shown in Fig. 2. Encapsulated atomic loops composed of ten metal atoms were considered in order to exploit the computational expedient of D_{5d} symmetry while creating M-Mbond lengths on the order of those calculated for isolated metal chains. DFT was used to determine the ground state configurations and to quantify M-M bond length, minimum C-M bond length, magnetic moment, density of states, conduction character, and binding energy of the encapsulated metal atoms.

II. METHOD

All calculations were performed with the real-space, numerical atomic orbital, DFT code, DMOL. 23 A norm conserving, spin unrestricted, semicore pseudopotential approach was employed with electron correlation accounted for using the Perdew-Wang generalized gradient approximation. 24 Double numerical basis sets were used with 4p polarization functions for Fe atoms and 3d polarization functions for C

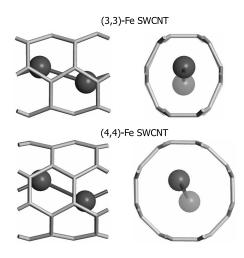


FIG. 4. Ground state structures of (3,3) and (4,4) SWCNTs embedded with atomic Fe wire.

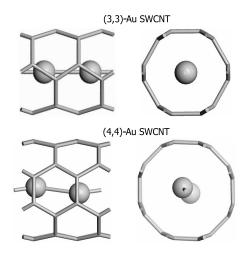


FIG. 5. Ground state structures of (3,3) and (4,4) SWCNTs embedded with atomic Au wire.

atoms. A smearing value of 0.272 eV was used for all structures

As a check on the method, an Fe dimer was determined to have a $(^7\Delta_u)$ collinear ferromagnetic ground state with a bond length of 2.00 Å and magnetic moment of $7.0\mu_B$ by Hirshfeld analysis—consistent with both experimental and other DFT results. S As a further check, the atomic spacing, binding energy, and magnetic moment of bcc Fe were evaluated to be 2.48 Å, -4.89 eV, and $2.22\mu_B$ —reasonable matches to the experimental values of 2.47 Å, -4.32 eV, and $2.22\mu_B$ per atom. S DFT values of 2.76 eV and $2.08\mu_B$ per atom were obtained elsewhere.

III. RESULTS

A. Freestanding metal chains

Periodic structures of pairs of freestanding Fe, Au, and Cu atoms were analyzed in order to understand the character of metallic chains prior to encapsulating them in carbon toroids. The ground state configuration for the Au chain was found to be a zigzag structure with a bond length of 2.83 Å and an angle between adjacent atoms of 57°. There is also a local equilibrium of higher energy at 134°, and this is close to that obtained by Sánchez-Portal *et al.* who identified an equilibrium angle of 131°. Analogous results, summarized in Tables II and III, were obtained for Fe and Cu chains, and all

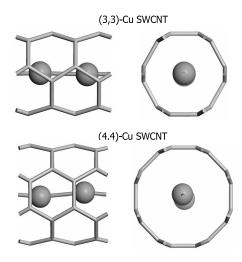


FIG. 6. Ground state structures of (3,3) and (4,4) SWCNTs embedded with atomic Cu wire.

three chains exhibit ground state bending angles of approximately 60°; an affinity for symmetric, three-atom bonding is optimum within the geometric constraints imposed. As shown in Fig. 3, such structures amount to two-dimensional, triangular lattice structures rather than linear chains. Key property data are summarized in the tables. All positive eigenvalues of the Hessian matrices associated with periodic cells of four atoms indicate at least linearized stability at 0 K for all three chains. A Jahn-Teller distortion was noted for metallic strands of eight atoms or more, and this is consistent with previous *ab initio* investigations.³⁰

B. Atomic metal wires within single-wall carbon nanotubes

Single-strand atomic wires of metal atoms encapsulated within single-wall carbon nanotubes (SWCNTs) are of interest on their own and can be used to anticipate the electronic structure and stability of wires encapsulated within carbon tori. The tube diameter of the C_{120} torus is between that of (3,3) and (4,4) SWCNTs, and so these fullerenes were chosen for more detailed analysis with metal wires. Previous DFT studies considered atomic wires of Fe within SWCNTs. ^{21,31} The analysis for (n,n) tubes considered only one metal atom per unit cell, though, and therefore did not allow for the possibility of the metal chain forming a zigzag configuration. Since the toroidal structure should allow for

TABLE I. Summary of properties for Fe structures. The band gap is HOMO-LUMO for the toroidal molecule.

Structure	<i>M−M</i> (Å)	$M-C$ $(\mathring{\mathrm{A}}_{\min})$	M-M angle (deg)	Gap (eV)	Mom. (μ_B/M)	E_B (eV/M)
Fe bcc	2.49			0.00	2.22	4.89
Fe chain	2.35		59	0.00	3.25	2.87
(3,3)-Fe	2.59	1.94	144	0.48	0.00	1.75
(4,4)-Fe	2.77	2.22	126	0.00	0.00	1.95
C ₁₂₀ -Fe ₁₀	2.49	2.04	139	< 0.1	2.28	1.82

Structure	<i>M−M</i> (Å)	$M-C$ ($\mathring{\mathrm{A}}_{\mathrm{min}}$)	angle (deg)	Gap (eV)	Mom. (μ_B/M)	E_B (eV/M)
Au fcc	2.88			0.00	0.00	3.05
Au chain	2.83		57	0.00	0.00	2.07
(3,3)-Au	2.46	2.34	180	0.00	0.00	-6.80
(4,4)-Au	2.52	2.56	156	0.00	0.00	-0.12
C ₁₂₀ -Au ₁₀	2.43	2.00	143	< 0.1	0.00	-5.54

TABLE II. Summary of properties for Au structures. The band gap is HOMO-LUMO for the toroidal molecule.

such corrugation, the SWCNT analysis was performed for double-cell, periodic structures with both Fe, Au, and Cu wires. The resulting equilibrium configurations are shown in Figs. 4–6 with property data summarized in Tables I–III. As expected, the encapsulated wires do exhibit a zigzag structure although the carbon sheath flattens the wires relative to their freestanding configuration. For the (4,4) tube, the Fe wire causes a diametral tube strain of 0.94%, while the Au wire causes a larger diametral expansion of 2.44%. Reassuringly, the Cu wire causes a dilation between these two— 2.1%. In the case of the smaller (3,3) tube, the diametral expansion of the Au tube is 8.7%, while the Fe tube shows a nonuniform distortion with bilateral strains of 8.2% and -5.2%. The Cu tube is uniformly dilated, like the Au case, with a strain of 3.4%. The allowance for zigzagging within the tubes significantly impacts the electronic structure of the composite. If a periodic cell half as long is considered, so that there is only one Fe atom per period, the band gap is 0.87 eV instead of 0.48 eV. The first number is consistent with the 0.88 eV obtained previously.³¹ The magnetic moments of both straight and zigzagging chains of Fe atoms were extinguished within (4,4) CNTs. An earlier analysis on straight, sheathed chains reported a magnetic moment of $2.6\mu_B$ per Fe atom.³¹

C. C₁₂₀ torus

The C_{120} torus was first analyzed in the absence of atomic wire, and the energy per C atom was found to be less than that of C_{60} by 0.09 eV (smearing) and 0.16 eV (no smearing). This is in contrast to an earlier tight binding calculation

which predicted that the toroidal state would be of lower energy. The difference is small, though, and the C_{120} torus should be very stable. A reactive molecular dynamics code was used to estimate that the structure is stable to at least 2000 K. The torus has no intrinsic magnetic moment and exhibits a band gap of 0.91 eV between its highest occupied and lowest unoccupied molecular orbitals (HOMO/LUMO). For comparison, the HOMO/LUMO band gap of C_{60} was found to be 1.5 eV—close to the experimental measurement of 1.6 eV.³²

D. Atomic metal necklaces within a C_{120} torus

The ground state was obtained for M_{10} loops encapsulated within the C₁₂₀ torus with the resulting structures shown in Figs. 7–9. The associated properties are summarized in Tables I–III. Consistent with the ground state morphology of freestanding metallic chains, the encapsulated necklace of Fe atoms forms a zigzag structure which is highlighted in the bottom of the figure. The minimum Fe-C bond length was found to be between that of the (3,3) and (4,4) SWNTs. The toroidal composite has a HOMO and/or LUMO band gap of less than 0.1 eV and has a magnetic moment of $2.28\mu_B$ per atom—approximately the same as that associated with bcc Fe. The spatial distribution of this magnetic moment is shown in Fig. 7 through an isosurface plot of the spin excess. At zero temperature, the Fe atoms have aligned spin axes and so the composite is ferromagnetic. In contrast, neither of the straight nanotubes with Fe showed any intrinsic moment because of stronger Fe-C bonding for the tubes.

Encapsulation of a Au_{10} necklace was found to be energetically unfavorable, while the Cu_{10} encapsulation is an en-

TABLE III. Summary of properties for Cu structures. The band gap is HOMO-LUMO for the toroidal molecule.

Structure	<i>M</i> - <i>M</i> (Å)	M-C (Å _{min})	M-M angle (deg)	Gap (eV)	Mom. (μ_B/M)	E_B (eV/M)
Cu fcc	2.56			0.00	0.00	3.62
Cu chain	2.42		59	0.00	0.00	2.17
(3,3)-Cu	2.46	2.24	177	0.00	0.00	-0.83
(4,4)-Cu	2.47	2.69	168	0.00	0.00	1.09
C ₁₂₀ -Cu ₁₀	2.39	1.92	144	< 0.1	0.00	0.01

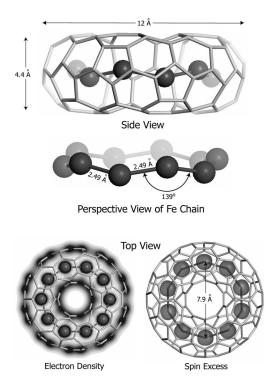


FIG. 7. Ground state structure of $C_{120}Fe_{10}$. The axial positions of the Fe atoms alternate between ± 0.21 Å. Also shown are the electron density cross section (center) and spin excess isosurface (bottom right).

ergetically neutral process. As indicated in Tables II and III, these results are consistent with those associated with CNT encapsulations.

The DFT investigation indicates that the C_{120} torus with all three types of metal necklaces is structurally stable at zero temperature. The TCN has a strong affinity for Fe atoms (1.82 eV/Fe) and a neutral affinity for Cu atoms, but the placement of a Au necklace within the torus would require an energy input of 5.54 eV per Au atom. This precludes the study of Au wires within this setting. The (3,3) and (4,4) CNTs can be used to estimate the affinity of C₁₂₀ toroids for metal atoms; the binding energy was found to be midway between the CNT values for each of the three metal chains. This allowed the prediction, in advance, that C_{120} would exhibit a neutral affinity for the Cu necklace. The presence of the metal atoms causes the TCN composite to have a HOMO/LUMO band gap of less than 0.1 eV in all three cases. When encapsulated with an Fe necklace, the torus exhibits ferromagnetic properties close to that of bcc iron, and it will be useful to determine to what degree, if any, this persists at finite temperatures. It is not clear what sort of chain structure and composite properties would result from the encapsulation of chains with symmetry at odds with that of the host torus.

IV. SUMMARY

The azimuthal symmetry of encapsulated atomic metallic necklaces within toroidal carbon sheaths offers the prospect of synthesizing composites with novel electromagnetic prop-

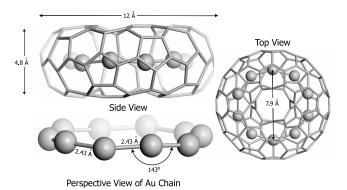
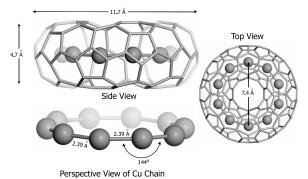


FIG. 8. Ground state structure of $C_{120}Au_{10}$. The axial positions of the Au atoms alternate between ± 0.076 Å.

erties that are not associated with straight, metalencapsulated CNTs. As a first step in this direction, a particularly simple encapsulation was analyzed using density functional theory to determine its ground state geometry, structural stability, binding energy, electrical conductivity, and magnetic moment. Of the four known classes of TCNs, a type 1 C₁₂₀ molecule was chosen because it results in a computationally tractable composite, and only ten metal atoms were considered in order to take full computational advantage of the D_{5d} symmetry of the torus. The TCN exhibits positive, neutral, and negative affinities for Fe, Cu, and Au atoms, respectively, and metal atom encapsulation resulted in linearly stable structures in all three cases. In contrast to Fe chains within (3,3) and (4,4) CNTs, Fe atoms within the C₁₂₀-Fe₁₀ molecule exhibit a net magnetic moment comparable to bcc Fe.

All toroidal nanotubes should be capable of serving as sheathing hosts for atomic metal necklaces with novel electromagnetic properties. The possibility of encapsulating metal chains within the C₆₀-coalesced tori of types 3 and 4 is of particular interest because these are structures that have already been synthesized.^{6,7} For type 3 TCNs, a reasonable strategy might be to encapsulate Fe atoms within the C₆₀ molecules prior to coalescing them into a toroid. Sufficiently long CNTs with embedded metal atoms could also be deformed to synthesize type 2 TCN composites. Both approaches would admit the construction of alloyed necklaces as well—the one-dimensional, atomic analog of powder metallurgy.



1

FIG. 9. Ground state structure of $C_{120}Cu_{10}$. The axial positions of the Cu atoms alternate between ± 0.025 Å.

- *mlusk@mines.edu
- ¹B. I. Dunlap, Phys. Rev. B **46**, 1933 (1992).
- ²S. Ihara, S. Itoh, and J.-I. Kitakami, Phys. Rev. B **47**, 12908 (1993).
- ³S. Itoh, S. Ihara, and J.-I. Kitakami, Phys. Rev. B **47**, 1703 (1993).
- ⁴E. Kurby, R. B. Mallion, and P. Pollak, J. Chem. Soc., Faraday Trans. **89**, 1945 (1993).
- ⁵ V. Meunier, P. Lambin, and A. A. Lucas, Phys. Rev. B **57**, 14886 (1998).
- ⁶J. A. Rodríguez-Manzo, F. López-Urías, M. Terrones, and H. Terrones, Nano Lett. 4, 2179 (2004).
- ⁷H. Terrones, F. Lopez-Urias, E. Munoz-Sandoval, J. Rodriguez-Manzo, A. Zamudio, A. Elias, and M. Terrones, Solid State Sci. 8, 303 (2006).
- ⁸ H. Terrones, M. Terrones, E. Hernández, N. Grobert, J. C. Charlier, and P. M. Ajayan, Phys. Rev. Lett. **84**, 1716 (2000).
- ⁹N. Chen, M. T. Lusk, A. C. T. van Duin, and W. A. Goddard, Phys. Rev. B **72**, 085416 (2005).
- ¹⁰F. London, J. Phys. Radium **8**, 397 (1937).
- ¹¹ J. P. Lu, Phys. Rev. Lett. **74**, 1123 (1995).
- ¹²R. Haddon, L. Schneemeyer, J. Waszczak, S. Glarum, R. Tycko, G. Dabbagh, A. Kortan, A. Muller, A. Mujsce, M. Rosseinsky, S. Zahurak, A. Makhija, F. Thiel, K. Raghavachari, E. Cockayne, and V. Elser, Nature (London) 350, 46 (1991).
- ¹³R. Haddon, Nature (London) **388**, 31 (1997).
- ¹⁴M. F. Lin and D. S. Chuu, Phys. Rev. B **57**, 6731 (1998).
- ¹⁵L. Liu, G. Y. Guo, C. S. Jayanthi, and S. Y. Wu, Phys. Rev. Lett. 88, 217206 (2002).
- ¹⁶ A. I. Yanson, G. R. Bollinger, H. E. van den Brom, N. Agra, and

- J. M. van Ruitenbeek, Nature (London) 395, 783 (1998).
- ¹⁷H. Ohnishi, Y. Kondo, and K. Takayanagi, Nature (London) 395, 780 (1998).
- ¹⁸D. Sanchez-Portal, E. Artacho, J. Junquera, P. Ordejon, A. Garcia, and J. M. Soler, Phys. Rev. Lett. 83, 3884 (1999).
- ¹⁹G. Rollmann, P. Entel, and S. Sahoo, Comput. Mater. Sci. **35**, 275 (2006).
- ²⁰P. Jain, Struct. Chem. **16**, 421 (2005).
- ²¹M. M. Rahman, M. Kisaku, T. Kishi, D. Matsunaka, W. Diño, H. Nakanishi, and H. Kasal, J. Phys.: Condens. Matter 16, S5755 (2004).
- ²²M. Weissmann, G. García, M. Kiwi, R. Ramírez, and C. C. Fu, Phys. Rev. B **73**, 125435 (2006).
- ²³B. Delley, J. Chem. Phys. **113**, 7756 (2000).
- ²⁴J. P. Perdew and Y. Wang, Phys. Rev. B **45**, 13244 (1992).
- ²⁵B. V. Reddy and S. N. Khanna, Phys. Rev. Lett. **93**, 068301 (2004).
- ²⁶G. Ackland, D. J. Bacon, A. Calder, and T. Harry, Philos. Mag. A 75, 713 (1997).
- ²⁷C. Kittel, *Introduction to Solid State Physics*, 7th Ed. (Wiley, New York, 1996).
- ²⁸M. Acet, H. Zähres, E. F. Wassermann, and W. Pepperhoff, Phys. Rev. B **49**, 6012 (1994).
- ²⁹V. Garcia-Suarez, C. Newman, C. Lambert, J. Prunedal, and J. Ferrer, Eur. Phys. J. B **40**, 371 (2006).
- ³⁰ A. Yanson, Ph.D. Thesis, Leiden Institute of Physics, 1975.
- ³¹ M. Kisaku, M. Rahman, T. Kishi, D. Matsunaka, T. Roman, W. Dino, H. Nakanishi, and H. Kasai, Jpn. J. Appl. Phys., Part 1 44, 882 (2005).
- ³² A. Skumanich, Chem. Phys. Lett. **182**, 486 (1991).