

SEARCH FOR RESONANT DOUBLE HIGGS PRODUCTION WITH  $b\bar{b}Z\bar{Z}$   
DECAYS IN THE  $b\bar{b}\ell\ell\nu\bar{\nu}$  FINAL STATE IN  $pp$  COLLISIONS AT  $\sqrt{s} = 13$  TeV

by

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Since the discovery of the Higgs boson in 2012 by the ATLAS and CMS experiments, most of the quantum mechanical properties that describe the long-awaited Higgs boson have been measured. Due to the outstanding work of the LHC, over a hundred of  $fb^{-1}$  of data have been delivered to both experiments. Finally, it became sensible for analyses teams to start working with a very low cross section processes, which made it possible to observe rare decay modes of the Higgs boson, e.g., a recent success in observing  $t\bar{t}H$  and  $VHbb$  processes. One of the main remaining untouched topics is a double Higgs boson production. However, additional hundred of  $fb^{-1}$  per year from the HL-LHC will not necessarily help us much with the SM double Higgs physics, the process may remain unseen even in the most optimistic scenarios. The solution is to work in parallel on new reconstruction and signal extraction methods as well as new analysis techniques to improve the sensitivity of measurements. This thesis is about both approaches: we have used the largest available dataset at the time the analysis has been performed and developed/used the most novel analysis methods. One of such methods is the new electron identification algorithm that we have developed at the CMS electron identification group, to which I have had a privilege to contribute during several years of my stay at CERN.

The majority of this thesis is devoted to techniques for the first search at the LHC for the double Higgs boson production mediated by a heavy narrow-width resonance

in the  $b\bar{b}ZZ$  channel:  $X \rightarrow HH \rightarrow b\bar{b}ZZ^* \rightarrow b\bar{b}\ell\ell\nu\bar{\nu}$ . The measurement searches for a resonant production of a Higgs boson pair in the range of masses of the resonant parent particle from 250 to 1000 GeV using  $35.9\text{ fb}^{-1}$  of data taken in 2016 at 13 TeV. Two spin scenarios of the resonance are considered: spin 0 and spin 2. In the absence of the evidence of the resonant double Higgs boson production from the previous searches, we proceed with setting the upper confidence limits.

*“... a place for a smart quote”*

*Lenin, 1922.*

## ACKNOWLEDGMENTS

This will be a long list!

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## CHAPTER 1

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# Physics Object Reconstruction

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Excellent spatial resolution of the CMS trackers, high granularity of the calorimeters, and almost  $4\pi$  coverage of the detector, allowed the CMS to introduce the particle flow (PF) algorithm [?] for a global event reconstruction. PF takes the input from all subdetectors, analyses the redundant information, removes the duplicate one, and forms the physics objects. PF procedure starts with identification of tracks and calorimeter clusters, then the reconstruction of the physics objects is performed, such as electrons, muons, jets, etc. In this section we will discuss the whole PF approach and all key elements of this reconstruction sequence.

## 1.1 Track Reconstruction

The reconstruction starts with the clusters of signals (“hits”) in the inner tracker. The information from these clusters in the Pixel and Strip subdetectors is aggregated based on their signal-to-noise ratios. The charge-weighted averaging is performed (for different particle charge hypothesis), as well as other corrections are further applied to identify the real hit positions.

The helix trajectory that the particle follows in the magnetic field inside of the detector is characterised by five parameters: the direction in  $\eta$ , the 3D position with respect to the reference point, which is the centre of the IP, and the curvature of the track with

the radius  $R$ . This information is enough to compute estimates of basic physics quantities, however, the this task is complicated by the presence of event high multiplicity (number of charged particles produced in the same event) and also by the physics aspect of the electron propagation in matter: an electron traversing the detector has nearly 85 % probability to emit a bremsstrahlung photon. Hadronic effects also need to be taken into account: a hadron has a 20 % probability to experience multiple scattering on the nuclei of the detector before reaching the HCAL.

To keep the track finding efficiency high, while maintaining low the efficiency of miss-identified tracks, track reconstruction is performed sequentially using the combinatorial track finder (CTF) [?]. First the "purest" tracks are reconstructed, they have high  $p_T$  and the hits point towards the primary vertex (PV). The term PV is used to refer to the vertex which is the actual point of origin of the produced particle, when several other hard scattering vertices are present in the event. Then these pure tracks are removed from the collection of tracks and another round of the track reconstruction starts. This procedure applied several times reduces the combinatorial factor and also simplifies the identification of tracks with the low  $p_T$  or those which do not point to the PV. During each iteration, the reconstruction follows these steps:

- Seed generation. Rough estimates of the particle trajectories ("seeds") are produced using either three hits or two hits and a PV constraint. Based on which iteration the algorithm is at, some additional constraints are applied, e.g., a minimal  $p_T$  requirement, the need for the seed to originate close to the beam spot, etc.
- Trajectory building. Initial seeds are projected towards the compatible hits in the next layers based on the Kalman Filter (KF) procedure [?]. The extrapolation is done until the outermost layer of the tracker is reached or when a "terminating condition" is satisfied, e.g., when the iteration accumulated the maximum number of invalid hits ("fake hits"). Each obtained trajectory is updated using a KF approach based on the compatibility of hits to form a better track candidate. The procedure is complicated

by the fact that the same initial seed can give rise to several track candidates or vice versa, the same track candidate may be compatible with different seeds. Additionally, the trajectory building step should take into account energy losses of the particle due to multiple scattering on the detector material, inhomogeneities of the tracker material, and the effects of the regions of non constant magnetic field.

- Track fitting. After the track candidate has been built, the track parameters are refitted by a KF and the "smoother". This step uses the full available information about the track and gives optimal estimates of the track parameters. To remove a large number of fake tracks, which are present due to a very complicated nature of the problem and the high track multiplicity in the event, a multivariate (MVA) selection is applied. MVA incorporates variables that discriminate real tracks from the fakes: the signed transverse curvature and impact parameters (with respect to the beam spot), the polar and azimuthal angles, number of missing hits, the fit quality variables, etc.

The CTF procedure runs for 10 iterations. For 2016 collisions with a mean pileup (additional hard scattering vertices) of 24, the CTF efficiency to identify real tracks varied from 80 to 95% with the mis-identification efficiency of 5 - 10 %.

### 1.1.1 Muon tracking

Muons are detected by the inner tracker and also by the outer (muon) tracker. This greatly improves muon track reconstruction and motivated the development of a dedicated muon reconstruction algorithms. The ninth and the tenth iterations of the CTF are focused on the muon reconstruction. These iterations are using three separate algorithms to identify:

- Standalone muons. This algorithm uses only muon tracker information: DTs, CSCs, RPCs. Hits from the inner chambers are used as seeds and projected to hits in the

outer chambers. Then, a standard KF procedure is used to identify track candidates, which are called standalone muons.

- Tracker muons. Only the inner tracking information is used to form tracks. Tracks are further projected to muon subsystems, where a compatibility with at least one muon hit is required. This algorithm works with low momentum muons: tracks with  $p_T$  above 0.5 GeV and a total momentum greater than 2.5 GeV.
- Global muons. Tracker tracks and standalone tracks are projected to the outermost layer of the muon system, checking the compatibility between two approaches. The resulting combined set of track hits is refitted to produce a global muon track. Mostly high momentum muons with  $p_T > ??$  200 GeV profit from this algorithm.

### 1.1.2 Electron tracking

Electrons are also detected by the inner tracker, however, their reconstruction is complicated by the fact that they emit bremsstrahlung photons and the trajectory becomes more complex. As a result, the clustering algorithms need also to identify the bremsstrahlung photons and account for the fact that the energy clusters corresponding to these photons may be located outside of the main electron trajectory, when extrapolated to the ECAL.

Since a KF approach assumes that energy losses are Gaussian, and this is not the case for electrons, a dedicated procedure is developed - a modified KF - the Gaussian Sum Filter (GSF) [144]. In this method the radiated energy losses are approximated by the sum of Gaussian distributions.

The electron seeds for the GSF are built using the ECAL information. Two different approaches are developed for the track reconstruction:

- Super Cluster based electrons. Cluster of energy in the ECAL and grouped together to form super clusters (SCs). Using the information of the energy spread among the

clusters, the curvature of the electrons is estimated and tracker seeds are formed, with the SC position as a constraint.

- Tracks based electrons. Tracks from the inner tracker are projected to ECAL clusters, checking the compatibility using quality variables, such as  $\chi^2$ , number of missing hits (absent hits along the path of the track), etc.

A typical momentum resolution for electrons in  $Z \rightarrow e^-e^+$  decays is approximately 1.7 - 4.5%.

Primary vertex finding Reconstructed tracks can be used to identify primary vertices, i.e. the locations of all proton-proton interactions in the event. These include the vertex corresponding to the ?hard? interaction, used for physics analysis and hereafter referred to as the primary vertex (PV), as well as additional parasitic interactions (pileup). The primary vertices are used to measure the position and size of the beam spot, i.e. the 3-D distribution of the luminous region, which feeds back into the track reconstruction sequence (both online and offline) as it is used as a constraint on track origin in some CTF iterations. Identifying pileup vertices is also of fundamental importance to mitigate the effect of pileup on object reconstruction performance, as will be described in Sec. 2.3.8. Lastly, knowledge of the position of the primary vertex is a crucial ingredient for b tagging (see Sec. 2.3.7). Vertex reconstruction consists of selecting a subset of tracks using quality criteria, and clustering them based on the z-coordinate of their PCA to the centre of the beam spot. To efficiently resolve close-by interactions while avoiding to split clusters corresponding to genuine vertices, the clustering is based on a deterministic annealing algorithm, seeking the global minimum of an analogue of free energy through step-wise reductions of the ?temperature? T. At infinite T, all tracks are assigned to a single vertex. As T is reduced, the vertices are allowed to split if the resulting configuration is more favoured. Once the minimum is reached, vertices containing at least two tracks are fitted to estimate their position. The resolution on vertex position varies from 10 ?m to 100 ?m, depending on the number and pT of the clustered tracks [195]. Finally, the beam spot can be measured by fitting the

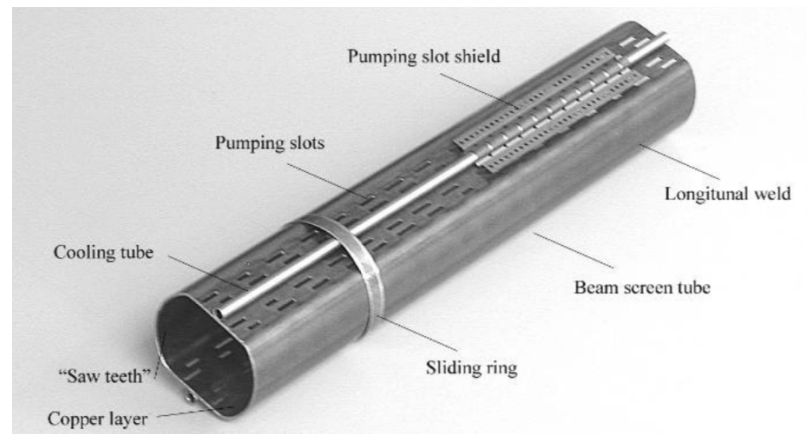
distribution of reconstructed vertices, a procedure repeated for every LS in the data. To pinpoint the vertex corresponding to the hard interaction in the event, objects are built, for each vertex candidate, using a jet clustering algorithm taking as input all the tracks associated to the vertex, as well as the vertex? missing transverse momentum. The ??2 vertex with the highest  $p_T$ , where the sum runs over the thus defined objects, is taken to be the hard interaction vertex.

**2.3.2. Calorimeter clusters and particle-flow links** In the PF algorithm, the calorimeters primarily serve to identify and measure neutral hadrons and photons, help in the reconstruction of electrons, and improve the energy measurement for charged hadrons (in particular at high  $p_T$ ). The clustering of energy deposits used in PF was specifically designed to resolve individual particles. Local cell energy maxima, above a certain threshold, define cluster seeds. Contiguous deposits are merged with the seeds, forming topological clusters. The procedure is carried out independently for each subdetector (ECAL, HCAL and HF) and partition (barrel and endcaps). Within a topological cluster, to take into account the overlap of energy deposits due to individual particles, clusters are identified using a Gaussian-mixture model, where the number of Gaussian energy deposits corresponds to the number of seeds, and the position and amplitude of each Gaussian are the cluster parameters. The single-particle PF response needs to be carefully calibrated, to correct i.a. for threshold effects and for the nonlinearity of the detector response. This calibration is carried out using simulated data, separately for photons using only the ECAL, and for neutral hadrons, which deposit

energy sometimes in both the ECAL and the HCAL, sometimes mostly in the HCAL. To identify individual particles in the events, the PF approach relies on information from the different subdetectors. To that end, the link algorithm combines PF elements (tracks, clusters) to create PF blocks, which form the basis for the different object reconstruction algorithms. For instance, an inner track is linked with a calorimeter cluster if it can be extrapolated to a position compatible with that of the cluster. Further, if tangents to a GSF track can be extrapolated to ECAL clusters, these clusters can be linked with the

track to recover bremsstrahlung photons. Pairs of tracks compatible with originating +? from a photon conversion into an  $e e$  pair also become linked. Similarly, groups of track are formed for reconstructed nuclear interactions within the tracker volume. Links between calorimeter clusters are only formed out of the tracker acceptance, when ECAL clusters are contained within the envelope of HCAL clusters. Lastly, as already hinted at in Sec. 2.3.1, links between inner tracks and muon tracks or segments establish global or tracker muons, respectively. With all the possible blocks in hand, object identification proceeds in steps. PF blocks corresponding to the candidates reconstructed at a given step are masked, and not considered further.

- somewhere a section on physics objects where you discuss that there are tracks, there are jets, there are b-jets, what it means an isolated lepton, what it means missing momentum, and so forth. For example, Ekaterina has a section on particle flow, if I remember right, where she discusses that. Or you can have a section that just discusses the proton-proton collisions environment and explains that many particles created in collisions are unstable and decay right away. What is observable in an experiment such as CMS is longer-lived or stable particles. These include electrons, muons, photons, and some ground state hadrons. At the same time, energetic quarks or gluons hadronize and can be seen in a particle detectors as jets of hadrons (so define jets here). Same for tau leptons, we detect them either as electrons/muons, or hadronic jets. Etc. If you have such a conceptual paragraph here, some time later you can have a particle flow discussion when it is time to be more specific. But discussion at this level above would allow you to then discuss the requirements on CMS design using terms like b-jets or isolated leptons.



**Figure 1.1:** Beam screen.

EGAMMA AT LEAST SAY THAT HAS BEEN WORKING AND SHOW SOME PLOTS?



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