

SEARCH FOR RESONANT DOUBLE HIGGS PRODUCTION WITH bbZZ  
DECAYS IN THE  $b\bar{b}\ell\ell\nu\bar{\nu}$  FINAL STATE IN pp COLLISIONS AT  $\sqrt{s} = 13$  TeV

by

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University of Nebraska, 2019

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Since the discovery of the Higgs boson in 2012 by the ATLAS and CMS experiments, most of the quantum mechanical properties that describe the long-awaited Higgs boson have been measured. Due to the outstanding work of the LHC, over a hundred of  $fb^{-1}$  of data have been delivered to both experiments. Finally, it became sensible for analyses teams to start working with a very low cross section processes, which made it possible to observe rare decay modes of the Higgs boson, e.g., a recent success in observing ttH and VHbb processes. One of the main remaining untouched topics is a double Higgs boson production. However, additional hundred of  $fb^{-1}$  per year from the HL-LHC will not necessarily help us much with the SM double Higgs physics, the process may remain unseen even in the most optimistic scenarios. The solution is to work in parallel on new reconstruction and signal extraction methods as well as new analysis techniques to improve the sensitivity of measurements. This thesis is about both approaches: we have used the largest available dataset at the time the analysis has been performed and developed/used the most novel analysis methods. One of such methods is the new electron identification algorithm that we have developed at the CMS electron identification group, to which I have had a privilege to contribute during several years of my stay at CERN.

The majority of this thesis is devoted to techniques for the first search at the LHC for the double Higgs boson production mediated by a heavy narrow-width resonance

in the  $b\bar{b}ZZ$  channel:  $X \rightarrow HH \rightarrow b\bar{b}ZZ^* \rightarrow b\bar{b}\ell\ell\nu\bar{\nu}$ . The measurement searches for a resonant production of a Higgs boson pair in the range of masses of the resonant parent particle from 250 to 1000 GeV using  $35.9\text{ }fb^{-1}$  of data taken in 2016 at 13 TeV. Two spin scenarios of the resonance are considered: spin 0 and spin 2. In the absence of the evidence of the resonant double Higgs boson production from the previous searches, we proceed with setting the upper confidence limits.

“... a place for a smart quote”

*Lenin, 1922.*

## ACKNOWLEDGMENTS

This will be a long list!

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## CHAPTER 1

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### Introduction

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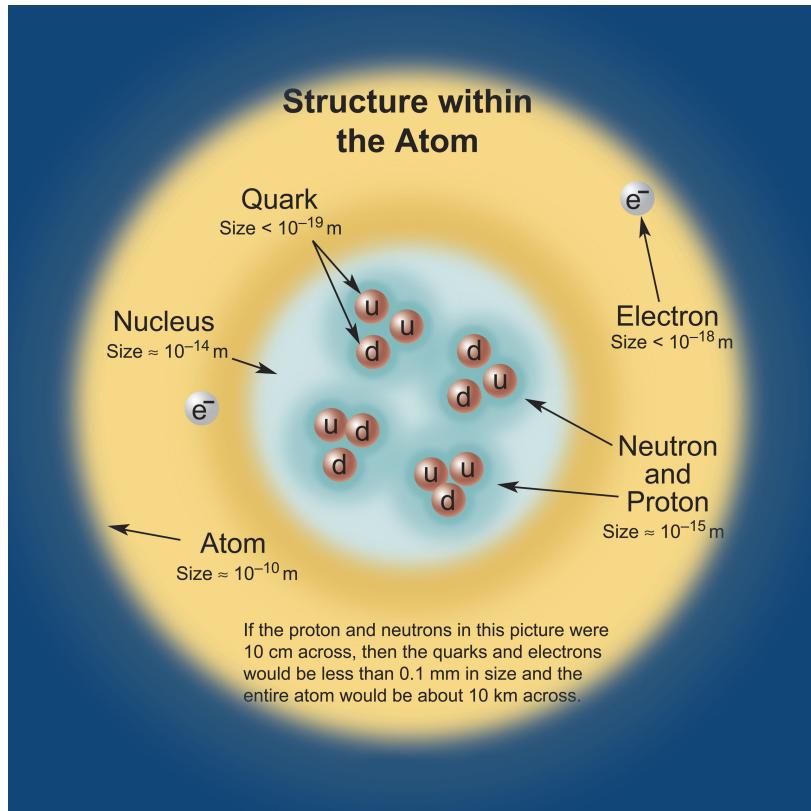
This chapter describes the development of particle physics following the historical order of how particles were discovered and parts of the Standard Model (SM) were constructed.

#### 1.1 "All things are made of atoms"

The SM is the theory of particles and their interactions that was built during many decades of intense experimental and theoretical work. Particles in this theory are elementary and have no size, meaning they cannot be divided further. This idea of the point-like particles is crucial since it reflects the goal of particle physics to find and describe the smallest and the most fundamental scale at which nature works. For the last several decades the SM is the most tested theory of elementary particles and forces and is presently generally accepted by the whole physics community. Formally, all SM elementary particles are split into two classes: fermions and bosons. Particles with a half-integer spin  $1/2$  (quarks and leptons) are called fermions since they obey Fermi-Dirac statistics [1]. The other class of particles is bosons. They are force carriers, have an integer spin, and are characterised by the Bose-Einstein statistics. A rigorous mathematical description of the SM will be given in the next chapter.

Our macroscopic world, from the smallest viruses to the biggest stars, is made of molecules and atoms. To show how deep and significant this simple idea is, let us quote Richard Feynman, a Physics Nobel Prize winner, who once summarised in a single phrase what he believed to be the most important fact about the world around us: "all things are made of atoms" [2]. Feynman himself was the father of quantum electrodynamics, and in this simple statement - delivered originally to Caltech students and now known to everyone through his series of physics books - he decided not to go into quantum mechanics principles and instead illustrated at the rather highly abstract level that everything is made of smaller particles. The quest to the smallest scale and the theory that would describe it were the key ideas that ultimately led to the development of the SM.

Today the physics community knows that molecules are made of atoms, which are not elementary particles either. Instead, atoms have heavy nuclei and light electrons "orbiting" around the nucleus on the electron shells. The nucleus is positively charged proportionally to the number of protons it contains. To provide the stability of the nuclei of the heavy atoms our world also needs neutrons, which have no electric charge. Going to an even smaller scale, it is now known that protons and neutrons are not elementary; rather, they are composed of point-like constituents that are called quarks (see Figure 1.1).



**Figure 1.1:** The structure of an atom. Approximate scale values are indicated.

Quarks were proposed by Gell-Mann and also by Zweig to explain periodicity in properties of observed subatomic particles [3]. Quarks are arranged in three families or three generations of doublets. A doublet is a mathematical construct that is used to explain a two-value system. For example, by design of Gell-Mann and Zweig, each quark doublet in their theory is a two quark system that has an "up" quark with the electric charge  $+1/3$  and a "down" quark with the charge  $-2/3$ . For antiquarks, the signs of the charges are reversed. The physics world before Gell-Mann and Zweig got used to the fact that particles have integer charges due to an enormous number of observations. The fact that the quark charge values were fractional was so revolutionary to Gell-Mann that he decided not to publish his article in a highly prestigious journal but, expecting a rejection, decided to go with a second tier one [3]. However,

with time out of all the theories that were trying to explain the difference in masses of observed hadrons it was the hypothesis of Gell-Mann and Zweig that turned out to be correct. Now, the quark theory is one of the key elements of the SM.

The SM includes six different types of quarks: up, down, charm, strange, top, and bottom. To distinguish one quark from another there is a "flavor" number assigned to them. For instance, a charm quark has +1 unit of "charmness", while a strange quark has -1 unit of "strangeness". All the other flavor fields are zero for quarks. This pattern is applied to all the other four quarks to fill the corresponding "quarkness" numbers. Another important characteristic of quarks was revealed at the  $e^+e^-$  colliders when physicists compared production rates of muons and hadrons. The theory was off by a factor of three. This was the motivation to introduce three quark colours: green, blue, and red.

Quarks were observed decades after electrons had been discovered. In particle physics, the electron belongs to a family of leptons. A lepton is an elementary particle with a spin 1/2 that participates in all but strong interactions, which will be discussed in more detail later in this chapter. The electron was discovered by Thompson [4] in 1897 when he was studying the properties of a cathode ray. Due to this discovery, that year may be considered the beginning of an era of a particle physics: dozens of particles were discovered in the next decades. In 1936, another lepton was observed, the muon [5], in an experiment of Anderson and Neddermeyer who studied cosmic radiation. In essence, a muon is almost a copy of an electron, but is 207 times heavier. No explanation for this mass difference exists in the SM. As a side note, according to Carl Bender, there is a story that Feynman was able to derive the mass of the muon starting with the mass of an electron, but the world has never seen that calculation published [6].

Leptons are also arranged in generations, analogously to quark families. Each

generation is a doublet that consists of a charged lepton (electron, muon or tau) with the charge  $-1$  and a neutral lepton (corresponding electron, muon, or a tau neutrino). Electron and muon neutrinos had been discovered in 1956 and 1962, respectively. The existence of the electron neutrino was deduced from the violation of the conservation of energy in a beta decay. The muon neutrino [7] was discovered by Schwartz, Lederman, and Steinberger during an experiment with a pion beam where leptons from the pion decays arrived to the aluminum spark chamber after passing a steel wall. Fifty-one events of interest had been observed after running the experiment for several months. Those events could not be initiated by electron neutrinos, since they will interact with the metal and produce electrons. The presence of narrow muon tracks in the chamber in each event, hence muons, was a clear indication that those neutrinos were of a different kind - they were muon neutrinos. Finally, a tau lepton and a tau neutrino were discovered in 1975 and 2000 correspondingly [8,9]. With that, all three families of the SM leptons were observed: a long-awaited tau neutrino, which decades ago was theoretically speculated to exist, was finally discovered experimentally. In a like manner to families of quarks, lepton masses grow with each generation, where a tau from the third generation is the heaviest lepton. To classify leptons of different families the lepton numbers were reserved: 1 unit of electron number to an electron and an electron neutrino, 1 unit of muon number to a muon and a muon neutrino, and 1 unit of tau number to a tau and a tau neutrino.

## 1.2 Fundamental forces

In nature there are four fundamental forces: gravitational, weak, electromagnetic, and strong forces. This thesis will classify all four forces [10] in terms of the relative strength, the range that they can cover, the spin of the mediator, and whether the

force's nature is attractive, repulsive, or both. This should be taken critically, since this is quite an ambiguous categorisation. It has a deep pedagogical meaning, though, because it helps to illustrate in which regime each of the forces is dominant. According to the world known mathematical physicist Carl Bender, this is of great importance since it is one of the main approaches to solving physics problems: to know which effects are the dominant and which are sub-dominant. This helps to justify what effect can be neglected and what approximation can be used. Thus, it allows the possibility to do calculations for problems where closed-form solutions do not exist, which is almost all the complex phenomena around us [6].

The first force on our list of forces is the gravitational force. This force governs the Universe at the macroscopic level: planets, solar systems, etc. The first theory of gravity was formulated by Newton [11]. Einstein later developed a new theory of gravity (GR). The key difference is that the Newtonian gravity had several "absolutes" that GR does not have: absolute time and space, a preferred separation of spacetime into time and spatial parts, absolute simultaneity, etc. [12]. Butterworth gives a good historical perspective in [13]. It is worth noting that the gravitational force is not included in the SM. Attempts are ongoing to expand the SM, e.g., adding the graviton as a mediator, but no real success so far has been achieved to create a renormalizable theory that would combine both the SM and gravity [14]. Surprisingly, gravity is the weakest force - the only reason why the motion of planets and galaxies is governed by gravity is because those are gigantic objects. Gravity effects become the dominant ones at the macroscopic scale because of an enormous number of particles involved in the interaction. If the strength of the strongest force, which is the strong force, is set to 1, then the strength of the gravity will be about  $10^{-41}$ . It is contemplated that the gravity mediator (the graviton), if it exists, would have a charge of zero, zero mass, spin 2, and should be a stable particle. The gravitational force is of the infinite range

and its nature is purely attractive, while the other three forces can exhibit both an attractive and a repulsive behaviour. Einstein's general relativity theory, though not a quantum theory, is the only working theory of gravity as of now.

The next force, the weak force, is mediated by a charged W (charge +1/-1) boson or a neutral Z boson, thus giving name to charged and neutral weak interactions correspondingly. All SM fermions, quarks and leptons, experience the weak force. The relative strength of the weak force is  $10^{-16}$  and the range of applicability is  $10^{-3}$  fm. All three weak bosons ( $W^+$ ,  $W^-$ , and Z) have spin 1 and are quite massive:  $m_{W^\pm} = 80.385$  GeV and  $m_Z = 91.189$  GeV. GeV is the unit of the "natural system of units", in which  $\hbar = c = 1$ . This system is very popular in the high-energy physics and is widely used in this thesis. Adoption of this system simplifies how many equations look and also makes a fine-structure constant  $\alpha \approx 1/137$  dimensionless. Using the natural system of units [15], masses, momenta, and energies are measured in electronvolts (eV), with GeV ( $10^9$  eV) and TeV ( $10^{12}$  eV) being the most popular units in a modern high-energy physics due to the energy regimes involved.

Charged weak interactions are interesting due to the fact that a primitive interaction vertex can be thought of as a point where a charged lepton is converted to a neutral lepton or vice versa. A good example is a muon decay, which is a conversion of the muon to a muon neutrino with the help of the W boson, which further decays to an electron and a corresponding electron antineutrino. It is worth noting that charged weak interactions do not conserve the flavor of quarks; e.g., members of doublets of the third and the second families can be converted into members of the lower family of quarks. This fact is reflected in the Cabibbo-Kobayashi-Maskawa (CKM) matrix [16]. This matrix describes the strength of the flavor-changing weak interactions. Since diagonal elements of this matrix are less than one and off-diagonal elements are non-zero, the CKM matrix represents a mismatch of quantum states of

quarks when they propagate freely and when they take part in the weak interactions. In other words, the CKM matrix with non-zero off-diagonal elements means cross-generation interactions are allowed and this is the information that the CKM matrix quantifies.

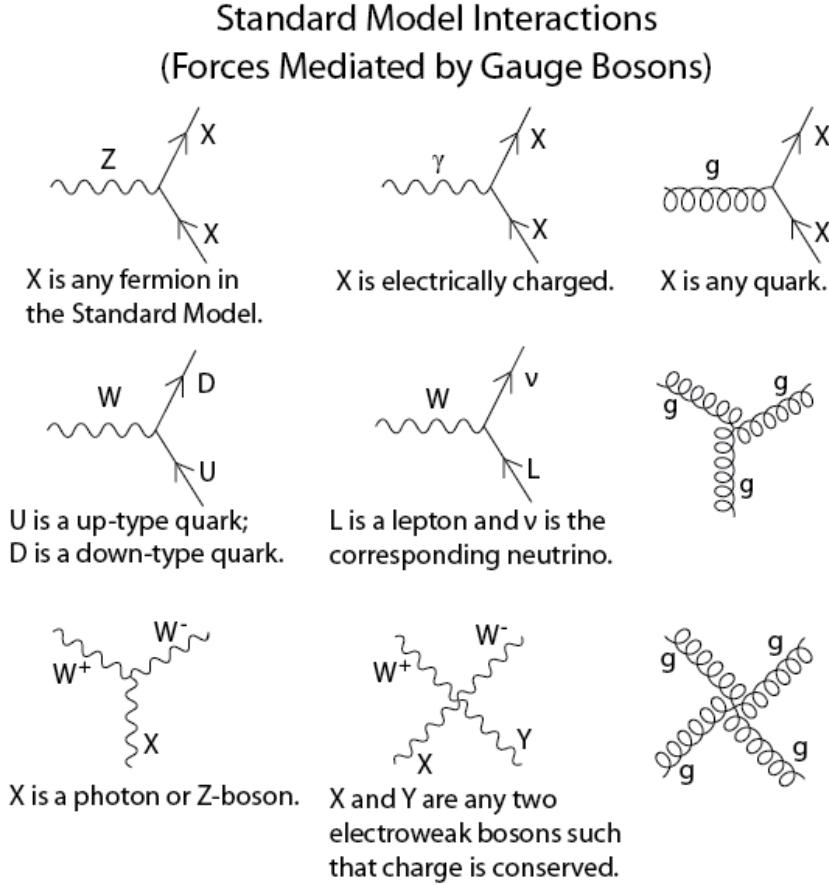
$$\begin{pmatrix} |V_{ud}| & |V_{us}| & |V_{ub}| \\ |V_{cd}| & |V_{cs}| & |V_{cb}| \\ |V_{td}| & |V_{ts}| & |V_{tb}| \end{pmatrix} = \begin{pmatrix} 0.97427 \pm 0.00015 & 0.22534 \pm 0.00065 & 0.00351^{+0.00015}_{-0.00014} \\ 0.22520 \pm 0.00065 & 0.97344 \pm 0.00016 & 0.0412^{+0.0011}_{-0.0005} \\ 0.00867^{+0.00029}_{-0.00031} & 0.0404^{+0.0011}_{-0.0005} & 0.999146^{+0.000021}_{-0.000046} \end{pmatrix}. \quad (1.1)$$

The third force, the electromagnetic (EM) force, is one of the main forces that we experience in our everyday life. The reason one can sit in the chair and does not fall further down due to gravity, is that electrons of the body repel electrons of the chair. The relative strength of the EM force is  $10^{-3}$  and the range of applicability is infinite. A photon, as the EM force's mediator, has zero mass and spin 1. The theory that describes photon interaction with leptons and quarks is called quantum electrodynamics (QED) and was developed in 1940s and 1950s by Tomonaga, Schwinger, Feynman, and Dyson [17]. Electric charge is conserved in EM interactions and no single photon-to-fermion vertex is possible; there are always two fermions that must be involved.

In the SM several multi-boson vertices are allowed. W and Z bosons that participate in weak interactions can couple to each other, so  $WWZ$ ,  $WWWW$ , and  $WWZZ$  vertices are possible in the SM. In addition, W bosons can couple to photons, so  $\gamma WW$ ,  $\gamma WWZ$ , and  $WW\gamma\gamma$  vertices are allowed too. Even though Z boson is massive and photon is a massless boson, Z boson has a neutral charge. This makes it possible that any interaction where the photon is a force carrier, can also be mediated by the Z boson.

The strong force, the fourth force of nature, is the strongest known force. Gluons are the carriers of this force and each gluon carries one unit of color and one unit of anticolor. There are nine types of gluons but, technically, the ninth gluon is a color invariant and would give rise to an infinite range of the strong force, which contradicts experiments. That is why modern physics assumes that in our world only eight gluons exist [3, 16]. Gluons carry color charge and, thus, can couple to each other. For several high order processes in quantum chromodynamics (QCD), 3- and 4-gluon vertices have to be introduced to restore gauge invariance and no higher order vertices are required [18].

To summarise the knowledge about the SM forces, one often refers to the Feynman diagram representation [19]. Fig. 1.2 shows all allowed SM particle interactions and corresponding simple vertices.



**Figure 1.2:** All SM interaction and simple vertices.

### 1.3 The Brout-Englert-Higgs mechanism

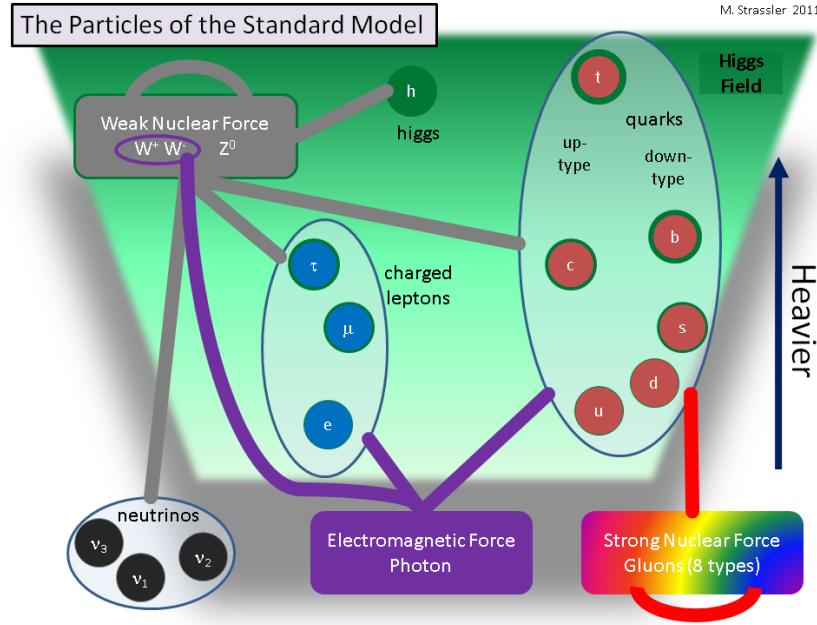
The description of the SM picture will not be complete without mentioning the main particle - the Higgs boson - that was predicted almost 60 years ago, but was not observed yet missing until 2012 (see Fig. 1.3 [20]). After the electroweak (EW) unification by Glashow, Salam, and Weinberg [21], it was still not clear what the origin of the mass of fundamental particles is. In 1964, Robert Brout and François Englert [22], Peter Higgs [23], Gerald Guralnik, C. Richard Hagen, and Tom Kibble [24] (BEHGHK authors), proposed the method by which the particles can acquire

mass. This technique consists of three stages and we will discuss them one-by-one:

1. The Brout-Englert-Higgs (BEH) mechanism
2. The BEH field
3. The Higgs boson.

The first stage, the BEH mechanism, is simply a spontaneous symmetry breaking (SSB) mechanism, which is a mathematical trick consisting of rewriting the original scalar fields in the EW Lagrangian, rearranging equations, and requiring that the fields are real. What does this lead to? The BEH authors started with a scalar complex field and a massless vector field and after SSB obtained a single real scalar field (Higgs boson) and a massive vector field. In terms of our physical world this it what gives mass to W and Z bosons.

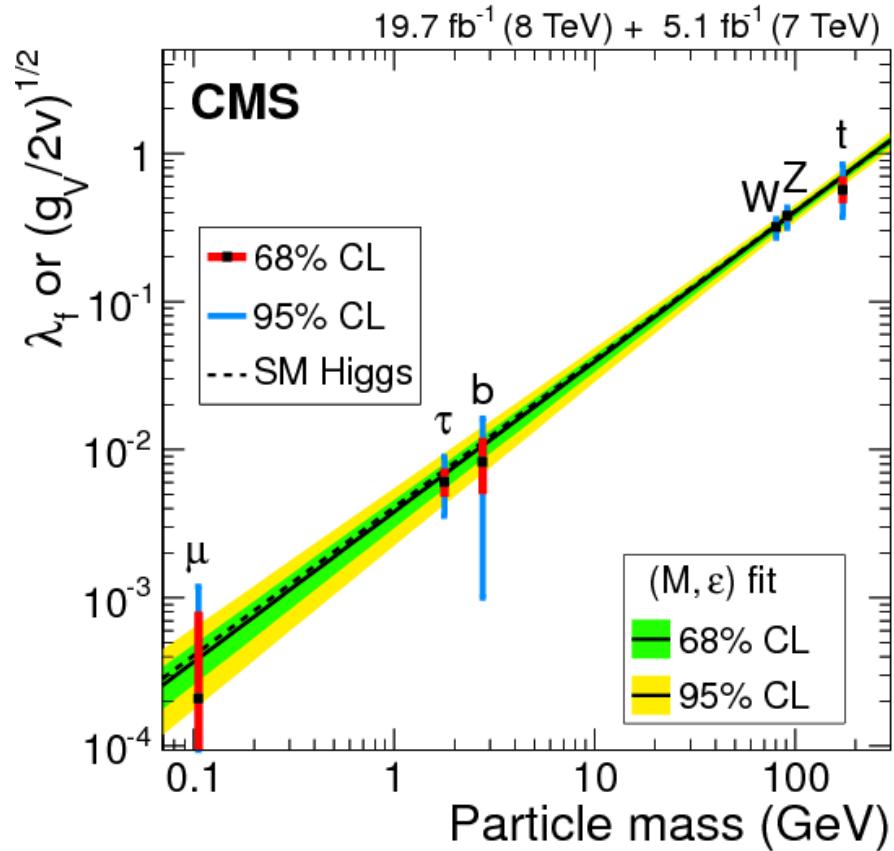
The second stage is the BEH field. It exists everywhere and has been present almost since the Big Bang [25]. It is a property of our world. All the fundamental particles that interact with the BEH field acquire mass. Those, which do not interact directly (at the tree level) have no mass and all their energy is in the form of the momentum. Such particles can travel with the speed of light. The more the particle interacts with the BEH field, the higher is the coupling to the Higgs boson or simply the higher is the mass of the particle. The coupling of the Higgs boson to fermions is proportional to the mass of the fermions, and for W and Z bosons it is proportional to the squared mass of bosons, making the top quark and the Z boson the most massive fermion and boson respectively (see Fig. 1.4 [26]).



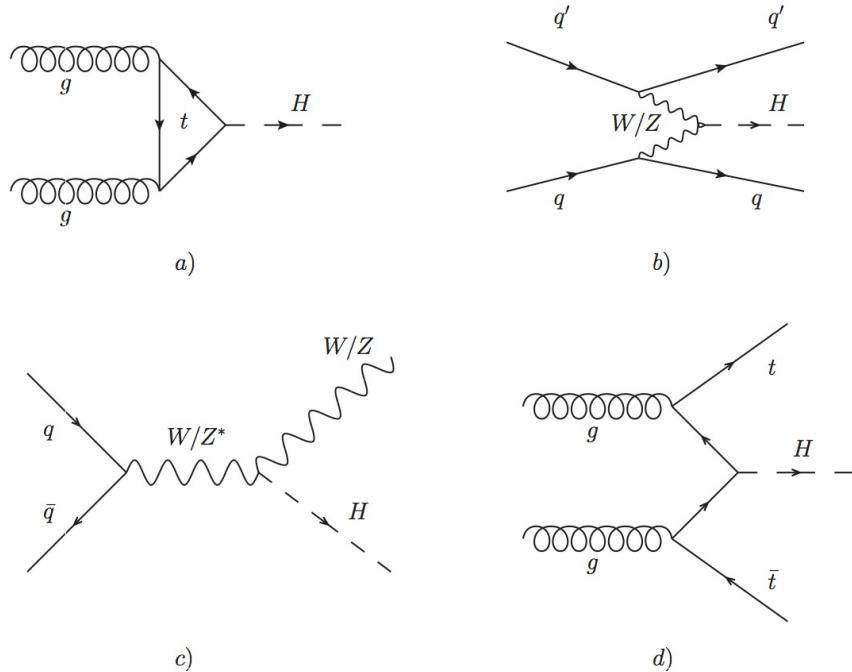
**Figure 1.3:** SM particles and force carriers. Self-interactions are also shown. The strength of the coupling to the Higgs boson increases from the bottom to the top, which is illustrated by the shades of the green color (the Higgs field).

The third and, arguably the most important stage, the Higgs boson. The Higgs boson is the excitation of the BEH field. Thus, the Higgs bosons can be produced at colliders by pumping more and more energy in a small space-time region exciting the BEH field to "produce" the Higgs bosons. In reality this happens through making the LHC beams more energetic and thus, during the collision, producing more energetic gluons (and also more energetic quarks). The main Higgs boson production mechanism is called a gluon fusion, when through the top quark loop a single Higgs boson is produced. This accounts for about 90% of the overall LHC Higgs production at the 13 TeV energy. The second mechanism is a vector boson fusion. The third mechanism is the associated production with a weak boson. The smallest contributor to the Higgs boson production is the  $t\bar{t}H$  process, which stands for the associated production of the Higgs boson with the top anti-top quark pair. All mentioned Higgs

boson production mechanisms are presented in the form of Feynman diagrams in Fig. 1.5.



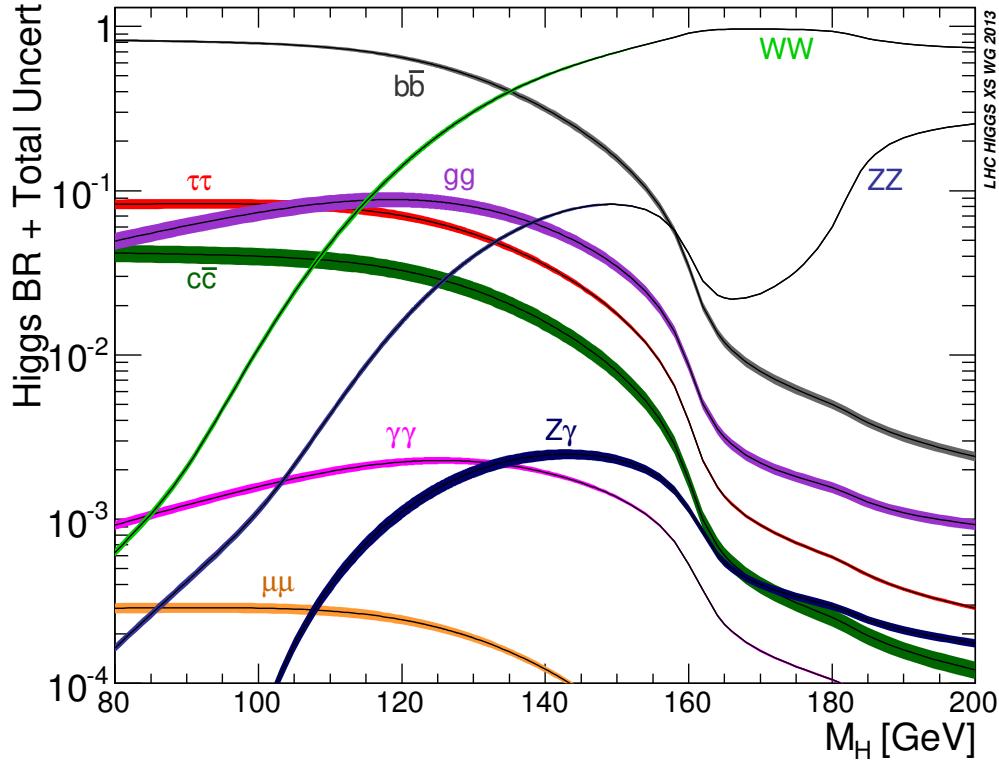
**Figure 1.4:** Coupling of particles to SM Higgs boson versus the mass of the particle, log-log scale is used.



**Figure 1.5:** SM Higgs boson production modes: a) a gluon fashion, b) a vector boson scattering, c) an associated production with a vector boson, d) an associated production with the top anti-top pair.

When describing Higgs bosons physics one cannot avoid mentioning the decay channels of the Higgs boson. In physics branching fraction term is reserved to quantify the probabilities with which a parent particle decays to daughter particles (see Fig. 1.6). The work of this thesis focuses on two specific Higgs boson decays,  $H \rightarrow b\bar{b}$  and  $H \rightarrow ZZ$ . The first one has the highest branching fraction, while the second one gives a clean signature when subsequent  $Z \rightarrow \ell\ell$  decays are selected. Before we conclude with the BEHGHK method, a little bit of history, an irony of life, actually. The BEH particle is called the Higgs boson, but Peter Higgs was not the first to publish the article on the BEH mechanism, in fact he was the last out of BEHGHK authors! His very first article was rejected since it contained no specific predictions or conclusions drawn from his calculations. This is why he was out-published by others. But this rejection made him write another article where he explicitly predicted an existence

of the new boson. And this is what has made all the difference, he was the first to predict a new boson, and this boson now is called the Higgs boson.



**Figure 1.6:** Higgs boson decay channels. At 125 GeV the dominant decay mode is  $H \rightarrow b\bar{b}$ .

Even though the facts above tell us about how great the SM is, the SM is still far from being perfect. Masses of elementary particles are the parameters in this theory, they do not come from SM predictions. It is hypothesised that the SM could be a part of the larger ultimate theory, the so-called "The Theory of Everything" (TOE), which is to be written (had been a lifelong journey of another genius, Einstein [27]). There is hope that the TOE will be able to explain many phenomena, such as the quark mass hierarchy, flavor mixing, etc. Also, in the SM all neutrinos are massless, however, it has been shown that they have a non-zero mass [28]. This fact is one of the main motivations for theorists to look for extensions of the SM.

## CHAPTER 2

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### Theory

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In the previous chapter we introduced the SM and discussed particles and their interactions that are described by this theory. In this chapter we will discuss first the general mathematical formalism of the SM and in the second part we will focus on the double Higgs boson physics in the Beyond the Standard Model (BSM) theory.

## 2.1 Lagrangian formalism of the Standard Model

The SM uses the Lagrangian mechanics as the mathematical approach to describe quantitatively the interactions of elementary particles and fields. The SM Lagrangian can be split into four main contributions [29]:

$$\mathcal{L}_{SM} = \mathcal{L}_{YM} + \mathcal{L}_{ferm} + \mathcal{L}_H + \mathcal{L}_{Yuk} \quad (2.1)$$

where

- $\mathcal{L}_{YM}$  represents gauge bosons and their **self**-interactions,
- $\mathcal{L}_{ferm}$  describes fermions and their interactions with the gauge bosons,

- $\mathcal{L}_H$  characterises Higgs boson, its self-interaction, and interaction with the gauge bosons to give them mass,
- $\mathcal{L}_{Yuk}$  gives details of fermions and their interactions with the Higgs boson, which through the Yukawa mechanism give mass to fermions.

The first term in the SM Lagrangian in full can be written as:

$$\mathcal{L}_{YM} = -\frac{1}{4}W_{\mu\nu}^i(x)W_i^{\mu\nu}(x) - \frac{1}{4}B_{\mu\nu}(x)B^{\mu\nu}(x) - \frac{1}{4}G_{\mu\nu}^a(x)G_a^{\mu\nu}(x) \quad (2.2)$$

where

$$B_{\mu\nu}(x) \equiv \partial_\mu B_\nu - \partial_\nu B_\mu \quad (2.3)$$

$$W_{\mu\nu}^i(x) \equiv \partial_\mu W_\nu^i(x) - \partial_\nu W_\mu^i(x) - g\varepsilon^{ijk}W_\mu^jW_\nu^k \quad (2.4)$$

$$G_{\mu\nu}^a(x) \equiv \partial_\mu G_\nu^a(x) - \partial_\nu G_\mu^a(x) - g_s f^{abc}G_\mu^bG_\nu^c \quad (2.5)$$

with  $\mu$  and  $\nu$  indices running from 0 to 3, SU(2) indexes  $i, j, k = 1, 2, 3$ , SU(3) indices given by  $a, b, c = 1, \dots, 8$ , and  $\partial_\mu$  and  $\partial_\nu$  represent four-vector covariant derivatives. According to the Noether's theorem, each symmetry is intrinsically connected to the conservation law [30]. The fields in the  $\mathcal{L}_{YM}$  are connected to their corresponding underlying symmetries in the following way:

- $B_{\mu\nu}$  corresponds to  $U(1)$  symmetry of the weak hypercharge  $Y_k$  with  $U(1)$  being a unitary one-by-one matrix (a scalar),
- $W_{\mu\nu}^i$  corresponds to  $SU(2)_I$  symmetry of the weak isospin  $I_w^i$ . Another common representation is  $SU(2)_L$ , since only left-handed SM fermions are transformed

under this symmetry.  $SU(2)_L$  is a unitary two-by-two matrix with the determinant equal to one.

- $G_{\mu\nu}^a$  corresponds to  $SU(3)_c$  symmetry of the QCD color charge with  $SU(3)_c$  being a unitary three-by-three matrix with the determinant equal to one.

The "B" field is a kinematic term, "W" and "G" terms describe interactions among the bosons,  $g$  and  $\varepsilon$  are  $SU(2)_L$  coupling and structure constants,  $g_s$  and  $f$  are coupling and structure constants for  $SU(3)_c$ .

The second term in the SM Lagrangian shows how fermions interact with the gauge bosons. Notice, that the mass terms are still absent:

$$\mathcal{L}_{ferm} = i\bar{\Psi}_L \not{D} \Psi_L + i\bar{\psi}_{l_R} \not{D} \psi_{l_R} + i\bar{\Psi}_Q \not{D} \Psi_Q + i\bar{\psi}_{u_R} \not{D} \psi_{u_R} + i\bar{\psi}_{d_R} \not{D} \psi_{d_R} \quad (2.6)$$

In the Eq. 2.6,  $\Psi$  represents a doublet of a charged lepton and a corresponding neutral lepton within the same lepton family of  $SU(2)_L$ , the subindex Q is reserved for a family of quarks, and  $\psi_R$  describes a right-handed leptonic singlet. Gauge boson interactions are present due to the derivative term:

$$D_\mu = \partial_\mu + igI_w^i W_\mu^i + ig'Y_w B_\mu + ig_s T_c^a G_\mu^a \quad (2.7)$$

Physical fields in this notation are represented by a linear combination of W and B fields:

$$\begin{aligned} A_\mu &= B_\mu \cos \theta_W + W_\mu^3 \sin \theta_W \\ Z_\mu &= -B_\mu \sin \theta_W + W_\mu^3 \cos \theta_W \end{aligned} \quad (2.8)$$

where  $\theta_W$  is known as the *Weinberg angle* [31].

With the first two terms of the SM Lagrangian one obtains a valid theory of fermions and bosons, however, these particles are massless in this theory [32], which evidently contradicts the reality. To solve this issue and to ensure that weak bosons are massive, one has to introduce a Higgs field. Higgs mechanism enters the SM Lagrangian through the corresponding Higgs Lagrangian term given by

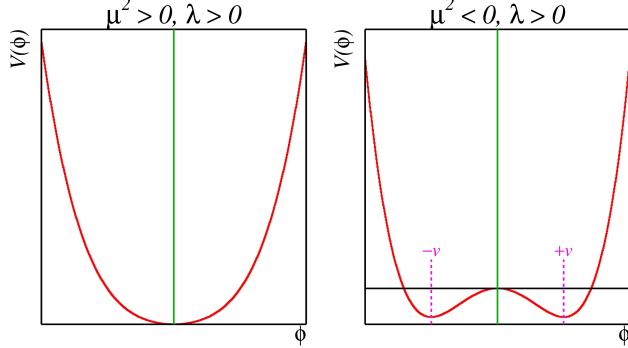
$$\mathcal{L}_H = (D_\mu \Phi)^\dagger (D^\mu \Phi) - V(\Phi), \quad V(\Phi) = -\mu^2 (\Phi^\dagger \Phi) + \frac{\lambda}{4} (\Phi^\dagger \Phi)^2 \quad (2.9)$$

where

$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 = (v + H + i\chi)/\sqrt{2} \end{pmatrix} \quad \text{with} \quad v = 2\sqrt{\frac{\mu^2}{\lambda}} \quad (2.10)$$

here  $\mu$  and  $\lambda$  are parameters of the Higgs potential. The Higgs field vacuum expectation value (*vev*)  $v$ , after the SSB, can be expressed in terms of  $\mu$  and  $\lambda$ . See Fig. 2.1 for the Higgs potential before and after the SSB. The importance of the  $\mathcal{L}_H$  in the SM Lagrangian is crucial: after rearranging terms the bosons finally have masses given by:

$$M_W = \frac{gv}{2}, \quad M_Z = \frac{M_W}{\cos \theta_W}, \quad M_H = \sqrt{2\mu^2} \quad (2.11)$$



**Figure 2.1:** Shape of the Higgs potential before and after the SSB that is determined at the leading orders by  $\mu$  and  $\lambda$  parameters [33].

The final contribution to the SM Lagrangian is the Yukawa term, with Yukawa Lagrangian given by:

$$\mathcal{L}_{Yuk} = -i\bar{\Psi}_L G_l \psi_{l_R} \Phi - i\bar{\Psi}_Q G_u \psi_{u_R} \tilde{\Phi} - i\bar{\Psi}_Q G_d \psi_{d_R} \Phi + h.c. \quad (2.12)$$

where  $\tilde{\Phi} = i\sigma^2 \Phi^*$ . The masses of fermions enter the equations through the  $3 \times 3$  matrices  $G$ , which are free parameters in the SM and have to be determined from the experiment. The mass of each fermion is proportional to the Yukawa coupling of the corresponding fermion to the Higgs boson, see Fig. 1.4.

The Higgs boson mass is proportional to the  $\mu$  parameter. In 2012, with precise single Higgs boson mass measurements from both ATLAS and CMS experiments, the value of  $\mu$  was determined. Since that time many analyses at CERN have been targeting the measurement of the  $\lambda$  parameter, because it is related to the shape of Higgs potential. In the SM having the "Mexican hat" Higgs potential, the simplest potential characterised by  $\mu$  and  $\lambda$  parameters, is sufficient to obtain the SSB phenomenon. This gives mass to fermions and bosons. However, the shape of the Higgs potential may be different and direct precise determination of the  $\mu$  and  $\lambda$  parameters

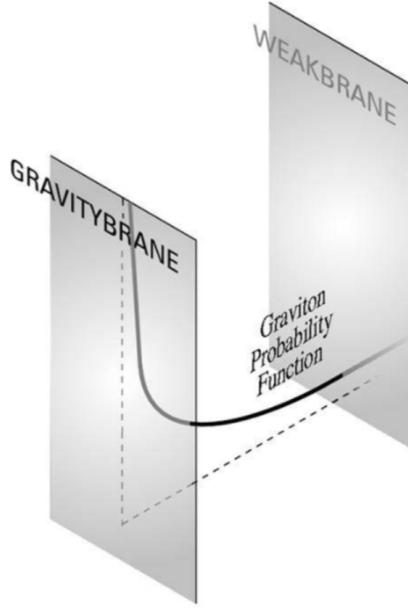
is a sensitive tool to test the limitations of the SM and may open doors to the BSM effects. The simplest interaction that is probing the Higgs potential directly is the one where two Higgs bosons (HH) are present. All this makes HH physics, the topic of this thesis, one of the main goals for the future High Luminosity LHC (HL-LHC).

**IN THE INTRO??**While the mass parameter has been measured fairly accurately,  $\lambda$  parameter requires even HL-LHC to run for many years to get enough statistics since HH processes are rare and are of almost three orders of magnitude lower rate than the single Higgs boson production. Technically, the amount of the HL-LHC data is not enough to reach the sensitivity of the SM for HH processes. At the same time, several BSM models predict resonant HH production to which even the current LHC data could be sensitive. In this theories, HH is produced through the decay of a heavy narrow width resonance, which is not a part of the SM. Thus, if such processes are found, this will open a new chapter in the HEP physics. In this thesis we focus on the resonant production of the HH system, which further decays to leptons and quarks. With the available CMS data, resonant HH analyses are starting to approach the needed sensitivity to rule out some BSM theories and test further the most promising ones.

## 2.2 Double Higgs in Beyond the Standard Model

Several BSM theories [34–36] predict a resonant production of the double Higgs boson events through a heavy resonance of a narrow width ( $\sim O(1 - 10)$  GeV) [37]. In this dissertation data is compared with respect to predictions from the Warped Extra Dimensions theory (WED) [38]. WED theory addresses the hierarchy problem by adding additional fifth dimension to the 4-dimensional (4D) space-time. In the framework that Randall and Sundrum (RS) [39] introduced, 4D space is an EFT

approximation of the higher dimensional space. The extra dimension exists between the gravity (Planck) and weak (TeV) flat 4D branes 2.2 and is called the "bulk". The bulk is described by the exponentially decaying metric.



**Figure 2.2:** 5D space in the RS model [40].

The free parameters of the RS model are the brane separation factor  $k$  and the size of the compactified dimension  $r_c$ . The curvature factor is given by  $k \approx \sqrt{\frac{\Lambda}{M_5^2}}$ , where  $\Lambda$  is the ultraviolet cutoff of the theory and  $M_5$  is the 5D Planck mass. The radius of the extra dimension  $r_c$  is proportional to the parameter  $1/k$  and the logarithm of  $1/vev$ . The hierarchy between the Planck scale and the electroweak scale is reproduced for  $k \cdot r_c \approx 11$ . In this case the RS model matches the observations of the Higgs boson being closer to the TeV brane and fermions having light mass (located near the Planck brane).

Since LHC had provided us with no evidence of the SM particles interacting with the RS particles, the RS model considered in this thesis hypothesises that SM particles are confined to branes. Another reason could be due to the fact that Kaluza-Klein

(KK) [41] partners of the SM particles are too massive to be produced at the LHC, but this scenario is not addressed in this dissertation. In the RS model under study two new particles appear. When the bulk is compactified, the WED theory predicts the existence of the KK excitations of the gravitational field, with the zero-th KK mode being a graviton, the mediator of the gravitational force. The graviton (spin 2) is the first WED particle predicted by the RS model. The graviton can propagate freely in the full higher-dimensional space of the 5D bulk. The other RS particle is a radion (spin 0). Its existence is required to stabilise the size of the extra dimension.

The theoretical arguments put forward by the authors [42] suggest the RS parameters  $k$  and  $\bar{M}_{Pl}$  to be constrained by the following range of values:  $0.01 \leq k/\bar{M}_{Pl} \leq 1$ . The parameter  $k$  is of the order of the Planck scale and  $\bar{M}_{Pl} = \sqrt{\frac{M_5^3}{k} \cdot (1 - e^{-2\pi kr_c})}$  is a reduced 4D  $M_{Pl}$ . Considered in this measurement graviton and radion are RS particles with a KK state mass of the order of TeV.

Let us denote a part of the KK 5D wave function, often called a profile, as  $f_X^{(n)}(\phi)$ , where  $n$  is referred to the  $n^{th}$  KK mode. Then the graviton can be decomposed as  $\sum_{n=0}^{\infty} h_{\mu\nu}^{(n)}(x_\mu) \cdot f_X^{(n)}(\phi)$ . Its zero-th mode corresponds to the massless graviton and the first mode corresponds to the lightest KK graviton (later graviton) which has the effective mass of the  $O(\text{TeV})$ . The profiles for all the matter fields are described by a combination of Bessel and exponential functions. The Lagrangian describing the interaction of the graviton with the SM fields is given then by

$$\mathcal{L}_{graviton} = -\frac{x_1 \tilde{k}}{m_G} h^{\mu\nu(1)} \times d_i T_{\mu\nu}^{(i)}, \quad (2.13)$$

where  $x_1 = 3.83$  is the first zero of the Bessel function for a given profile,  $\tilde{k} = k/\bar{M}_{Pl}$ ,  $h^{\mu\nu}$  is a symmetric tensor,  $m_G$  is the effective mass of the graviton of the order of TeV,  $d_i$  is an integral of the profiles of the SM fields and KK graviton, and  $T_{\mu\nu}^{(i)}$  is a

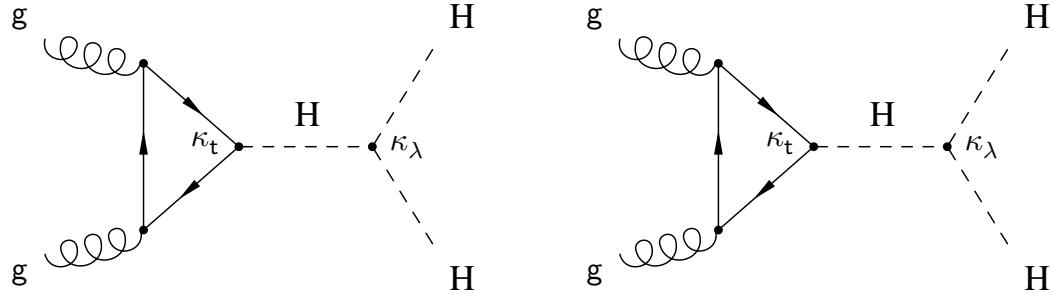
4D canonical energy-momentum tensor [43] for any SM field  $i$ . A free parameter  $\tilde{k}$  varies from 0.01 to 1 when  $m_G$  is varied from 100 to 1500 GeV.

For radion the Lagrangian is given by:

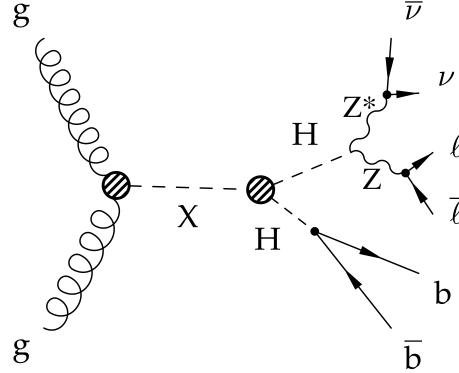
$$\mathcal{L}_{radion} = -\frac{r}{\Lambda_R} \times a_i T_\mu^{\mu(i)}, \quad (2.14)$$

where  $r$  is a 5D radion field,  $\Lambda_R$  is the scale parameter proportional to  $k \cdot \sqrt{(\frac{M_5}{k})^3}$ , and  $a_i$  is the coupling of the radion to the SM field  $i$ . In the studied RS model the profiles of the graviton and radion arise naturally as being localised at the TeV brane for the coupling of a radion and a graviton to the massive SM fields to be of the order of 1.

In the SM the HH system is produced predominantly via two diagrams shown on Fig. 2.3: the "box" and the "triangular" diagrams. They interfere destructively and the total cross section is thus lowered (Fig. 2.5 on the right). The box diagram dominates the double Higgs boson production and peaks near 400 GeV, as shown by the study performed by theorists, see [44]. In this measurement, though, the gravitons and radions in the search are expected to be produced by the BSM "contact interaction" Feynman diagram allowed by the WED scenario. These process is shown on Fig. 2.4. A graviton and a radion subsequent decays to HH system are thoroughly studied and the experimental results are compared to the theoretical predictions calculated for the WED model with the parameters  $\tilde{k} = 0.1$  and  $\Lambda_R = 3$  TeV.

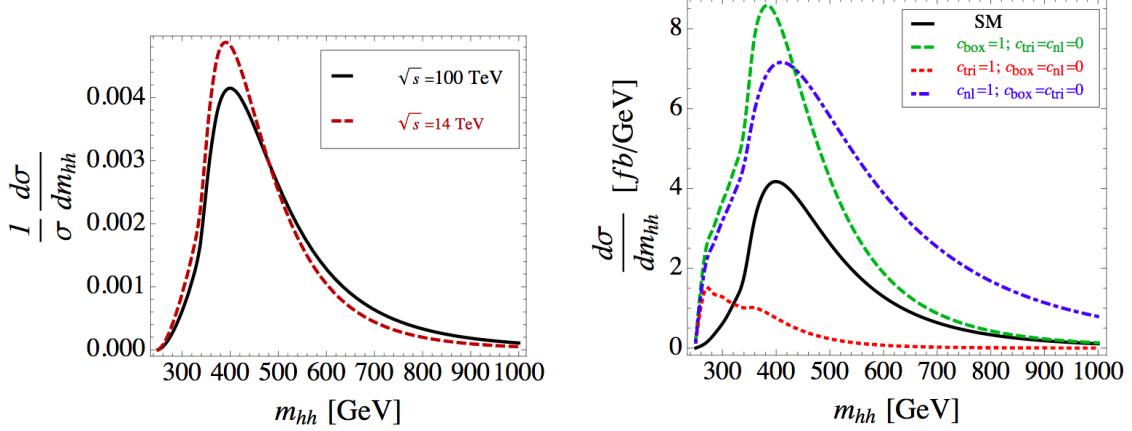


**Figure 2.3:** SM double Higgs boson production.



**Figure 2.4:** BSM Resonant double Higgs decay in the 2  $b$ , 2 lepton, and 2 neutrino final state.

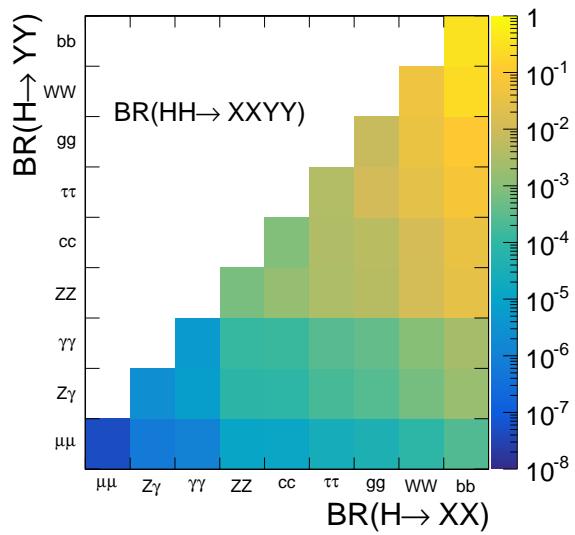
The kinematic distribution of the double Higgs mass remains to a high degree unchanged between 13-14 and 100 TeV (see Fig. 2.5 on the left), therefore, we can compare 100 TeV results produced by theorist to those analysed in this thesis that use the data delivered by the current 13 TeV LHC machine. Fig. 2.5 refers to the box and the triangular diagrams as "box" and "tri", and to the non-linear interaction as "nl" [45].



**Figure 2.5:** Left: comparison of the double Higgs boson mass distribution at the LO at  $\sqrt{14}$  and  $\sqrt{100}$  TeV center-of-mass energy. Right: the total SM HH cross section and the individual contributions [45].

This thesis separately addresses resonant graviton and radion decays into two SM Higgs bosons with the subsequent decays of one Higgs boson to a pair of b quarks, and the other Higgs boson to W or Z boson pairs. W bosons are allowed to decay only leptonically. For Z boson decays, the signature is characterised by the on-shell Z boson decaying into a pair of charged leptons and the off-shell Z boson decaying to neutrinos (see Fig. 2.4). The final state that this thesis focuses on consists of two b quarks, two charged leptons, and two neutrinos. This signature has a branching fraction of approximately 2.8%.

To finish this chapter, it is instructive to show all the decay channels of the double Higgs system to the SM particles, which is summarised in the Fig. 2.6. Both the horizontal and the vertical axes show decays of a single Higgs boson to two SM particles. In this representation, each square on the plot specifies a branching fraction of one of the double Higgs boson decays, with the probability of the decay given by the colour field on the right.



**Figure 2.6:** Double Higgs decay channels. The SM branching fractions are represented by the colour palette.

## CHAPTER 3

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### **LHC and the CMS experiment**

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CERN accelerator complex is a sequence of machines that produces and accelerates "bunches" of  $10^{11}$  protons to nearly the speed of light. In the Large Hadron Collider (LHC) the bunches collide at specific interaction points (IP), where the four main experiments are located: ALICE, ATLAS, CMS, and LHCb. We will start this section with the discussion of the LHC machine and then describe the CMS detector.

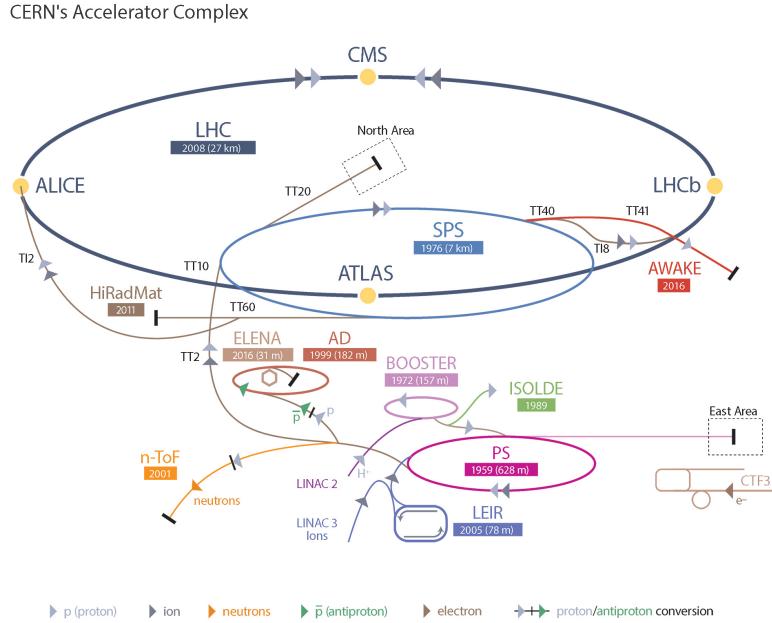
## **3.1 The Large Hadron Collider**

### **3.1.1 The history of the LHC**

The story of the LHC begins in 1977, when the CERN director general Sir John Adams suggested that the tunnel of the Large Electron-Positron Collider (LEP) can be reused to accommodate the future hadron collider of more than 3 TeV energies ???. At the 1984 ECFA-CERN workshop on a "Large Hadron Collider in the LEP Tunnel" ??, the physics goals of the LHC were stated: confirmation of the BEH mechanism, search for the Higgs Boson, and exploration of the origin of masses of W and Z bosons. The parameters of the proposed LHC were very ambitious: the centre-of-mass (COM) collision energy of 10 to 20 TeV, and a target instantaneous luminosity of  $10^{33-34} \frac{1}{cm^2 s}$ .

Large Hadron Collider (LHC) is the most powerful particle accelerator that has ever

been built. It is located at the border of France and Switzerland at a depth from 50 to 175 m underground. LHC ring is 26.7 km in circumference and it is the final stage in a sequence of accelerators. We will discuss the whole sequence of accelerators in the following section.



**Figure 3.1:** Schematic layout of the LHC.

### 3.1.2 The layout of the LHC

It is a complex process to start proton-proton collision in the LHC at 13 TeV and, therefore, the process consists of several stages (see Fig. 3.1). Everything begins with the bottle of hydrogen. The hydrogen atoms from the bottle are fed into the source chamber of the Linear Accelerator (Linac). In the chamber the hydrogen is heated up to the plasma state until electrons are stripped off of the hydrogen atoms. Then electrons are removed and remaining protons are directed to the first acceleration stage which increases the energy of protons to 50 MeV. After Linac, the beam of protons is injected into the Proton Synchrotron Booster (PSB). PSB contains four rings each accelerating a bunch of protons (a moving collection of protons of a narrow length) to 1.4 GeV. The third stage is the Proton Synchrotron

(PS), which splits the incoming beam into 72 bunches separated by 7.5 m. The energy of the protons is increased to 25 GeV. After that, the protons are sent to the Super Proton Synchrotron (SPS), where they are accelerated to 450 GeV. SPS then fills the LHC ring with two beams each consisting of 2808 bunches of protons with nearly  $10^{11}$  protons in total. It takes SPS about  $O(10)$  minutes to fill each LHC ring with bunches. In the LHC two beams are circulating in opposite directions in two separate beam pipes. During standard data taking beams circulate for  $O(10)$  hours.

### 3.1.3 LHC operations

The first LHC budget plan was finalised in 1996 and the final cost was approved just a few years later. The first proton beam entered the LHC ring in 2008. However, an incident intervened the LHC plans. It was caused by the mechanical damage of the tunnel equipment due to the release of the helium. Thus, the real data taking period (called LHC Run-1) had started only in 2010, lasted for two years, and 7-8 TeV COM energies were used. The recorded dataset contained enough Higgs bosons to claim a discovery of this rarely produced particle. After this achievement, the LHC was closed for the first long shutdown (LS1) that happened in 2012. During this time necessary upgrades of the main detectors and the LHC were performed. This was an unavoidable and essential step to prepare the LHC for more challenging environment of COM energies increased to 13 TeV.

If we denote the area of  $10^{-28} \text{ m}^2$  as barn (b), with the femtobarn ( $fb$ ) equal to  $10^{-43} \text{ m}^2$ , then in terms of these new units the LHC can theoretically produce  $80 - 120/\text{fb}$  (inverse femtobarns) of data a year. In practice numbers were lower, because LHC operated at the revolution frequency below the nominal, used fewer proton bunches in the beam, etc. All this resulted in lower than expected instantaneous luminosity, which is a very important term in collider physics and will be explained in the next section.

The LHC Run-2 has started in 2015 and the CMS collected  $4.2 \text{ fb}^{-1}$  of data that year. Over the course of the 2016 data taking, an integrated luminosity of  $35.9 \text{ fb}^{-1}$  was recorded.

This luminosity is the amount of data that has been collected by the CMS detector and later approved by the CMS physics coordination for the use in the physics analyses. The data set of proton-proton collisions collected in 2016 at 13 TeV COM energy is used in this thesis to analyse double Higgs boson decays. Together with the 2017 and 2018 data taking, almost  $150 \text{ fb}^{-1}$  have been delivered and recorded by the CMS detector during the whole Run-2 period of four years.

At the moment of writing this thesis, the LHC has entered the LS2. The next data taking will resume in 2020 and proton-proton collisions will continue for three years with the expected delivered integrated luminosity equal to nearly  $300 \text{ fb}^{-1}$ . This will conclude the LHC Phase-1 programme.

The new upgraded LHC, the High-Luminosity LHC (LHC) or the Phase-2, will start operations in 2026 and run until 2035. The COM energy will be increased to 14 TeV and one expects to record an unprecedented dataset of  $3000 \text{ fb}^{-1}$ .

### 3.1.4 Luminosity

The instantaneous luminosity  $\mathcal{L}$  is the coefficient which relates the cross section  $\sigma$  of the process to the number of events  $N_{events}$  produced during the interaction:  $N_{events} = \mathcal{L}\sigma$ . Luminosity is the parameter controlled by the machine and can be written as:

$$\mathcal{L} = \frac{N_b^2 f_{rev}}{4\pi\sigma_x\sigma_y}$$

where  $N_b$  is the number of particle in the colliding bunch,  $n_b$  is the number of colliding bunches in the beam,  $f_{rev}$  is the revolution frequency of the beam,  $\sigma_x$  and  $\sigma_y$  are the standard deviations of the beam density profile (BDP) in the transverse plane, where it is assumed that the BDP of both beams can be described by a Gaussian distribution.

To maximise the amount of collected data, the luminosity parameter should be as high as possible. It is worth noting that the luminosity is not constant and decays with time due to the degradation of the initial circulating beams. Theoretical decay time (the time to reach  $1/e$  level) is approximately 29 h. In practice, taking into account the decrease of protons in

the bunch due to collisions, contributions from the intrabeam scattering, scattering on the residual gas, etc., the real luminosity lifetime is about 15 h.

A useful variation of the luminosity parameter is a total integrated luminosity. This is the number normally quoted for the dataset collected over the period T:

$$L = \int_0^T \mathcal{L} dt.$$

In collider physics the "beam dump" is a process of burning off exhausted low luminosity beams by intentionally directing them towards the target made of concrete and steel. The time from the start of the collisions to the beam dump is usually called the "run".

We can calculate the amount of data delivered by the LHC during a single run period  $O(10)$  h. Performing the integration, we obtain:

$$L = \mathcal{L}_0 \tau_{\mathcal{L}} \left[ 1 - e^{-\frac{\tau_{\text{run}}}{\tau_{\mathcal{L}}}} \right],$$

where  $\mathcal{L}_0$  is the initial peak instantaneous luminosity at the start of the run,  $\tau_{\text{run}}$  is the total duration of a run, and  $\tau_{\mathcal{L}}$  is the luminosity lifetime. The optimum run time is 12 hours. During the runs, the LHC centre needs to dump the old beams, fill the rings with the new beams, and increase ("ramp") the energy of new beams to 13 TeV. After that a new run can be started. This restarting process normally takes two to six hours.

### 3.1.5 LHC infrastructure

The equipment of the LHC tunnel serves several purposes with the main objective to keep the colliding beams on the circular orbit. This requires a complex synchronised work of bending dipole magnets, cooling systems, accelerating radio frequency cavities, and vacuum insulation systems.

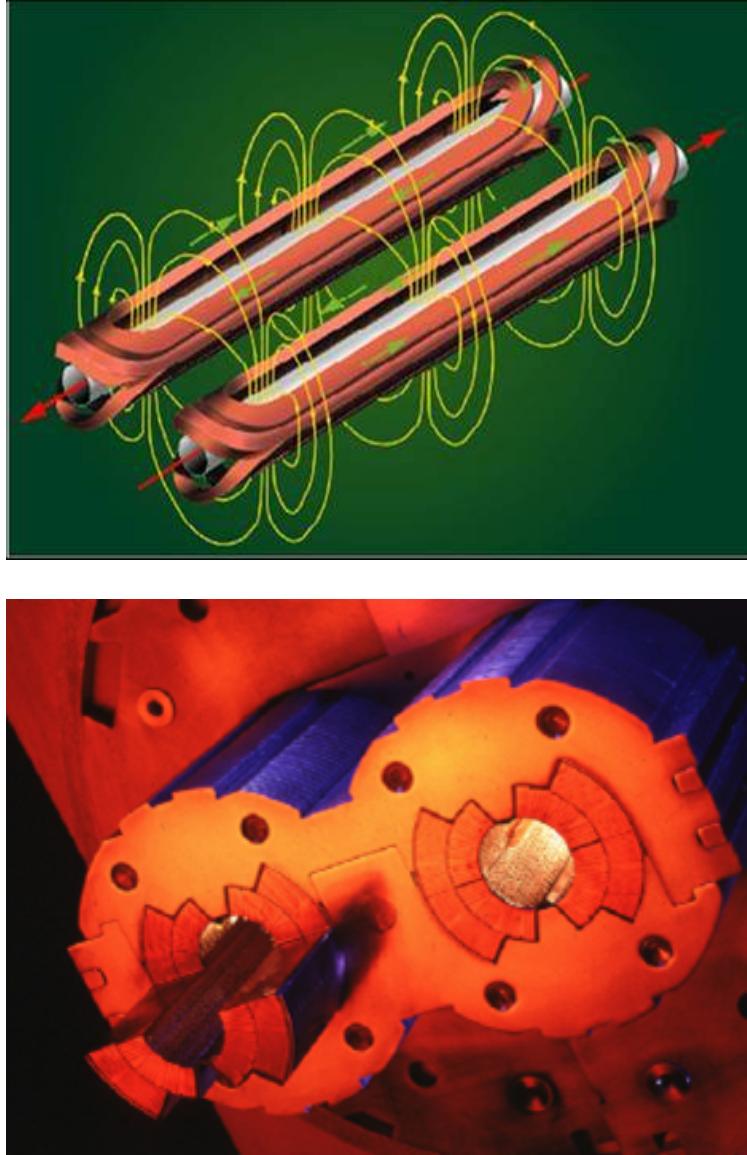
#### 3.1.5.1 Magnets

Most of the LHC circumference is used by 1232 superconducting magnets placed evenly around the tunnel to approximate the circular orbit. These are dipole magnets (see Fig. 3.2) that bend the beam and keep it on the circular orbit, that is why they are com-

monly called "Main Bends" (MB). The proven technology existed since Tevatron and relied on NbTi superconductors. This technology also satisfied the LHC cost and performance requirements, thus, it was decided to reuse the same choice of the alloy for the LHC superconducting dipole magnets that steer the proton beams.

The dipoles need to produce the magnetic field of 8.3T. Each dipole is 16.5 *m* (with ancillaries) long and 570 *mm* in diameter and is placed inside of the dipole cryostat which is called the "Helium bath".

This cryostat is a long cylindrical tube 914 *mm* in diameter made of low-carbon steel, where the dipole mass is cooled down to 1.9 *K*. Even though the inner structure of such cryostat is very complex and includes two beam pipes, two sets of coils for two beam pipes, vacuum pipes etc., one normally calls this compound object simply a dipole magnet. The name "dipole" is reserved for MBs since for each beam pipe the magnet consist of two "poles" that provide a vertical magnetic field similarly to a simple dipole system of magnets.

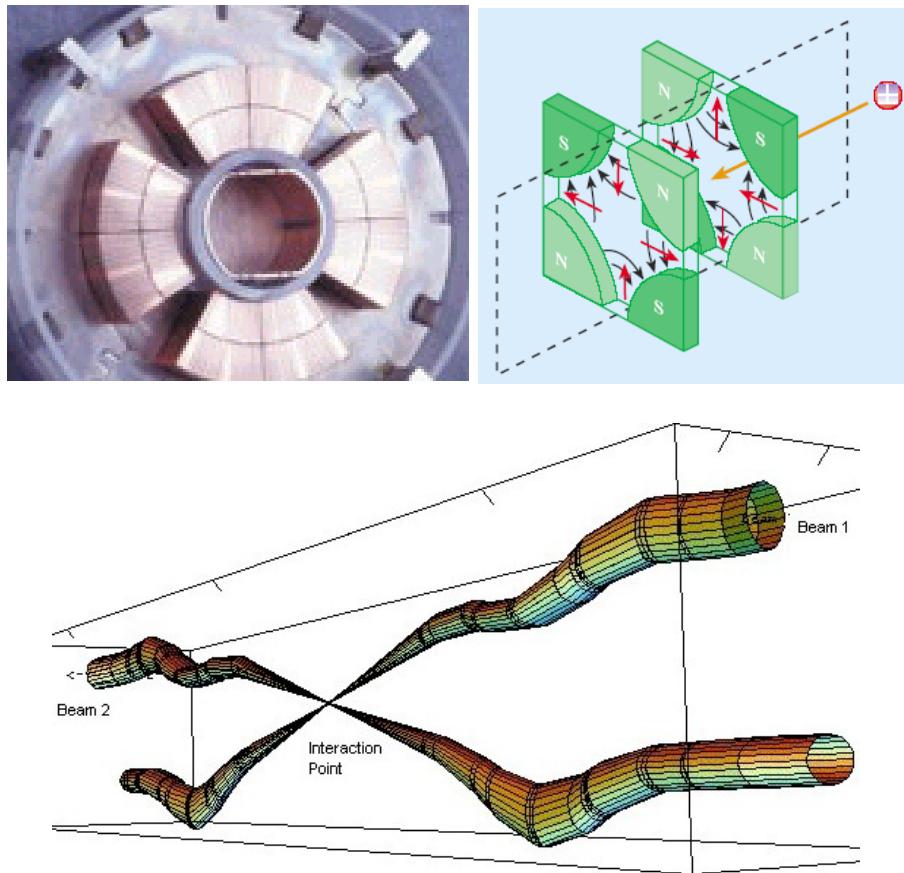


**Figure 3.2:** LHC dipole magnets. Top: two dipole coils and magnetic field lines. Bottom: two beam pipes with the coils inside of the dipole magnet.

A dipole magnet must be curved to help a chain of dipoles complete 360 degrees. The curvature is 5.1 *mrad* per dipole, which is equivalent to a sagitta of about 9 mm, corresponding to a radius of curvature of 2812.36 m.

The other important set of magnets is quadrupoles. They are used to ensure the proper beam dynamics. In total 392 quadrupole magnets ranging from 5 to 7 metres in length

are used to squeeze the beam in transverse direction and to keep it narrow during the run duration. Additional special quadrupole magnets (SQM) are installed right before the IPs to focus the beams even more. That increases the density of protons in the beam and guarantees the maximum luminosity. In addition, SQMs help to decrease the chance of the parasitic collisions when bunches from the same beam or bunches outside of the IP centre interact (see Fig. 3.3). To further correct the beam path (orbit), about 5000 higher order correcting magnets are used, which are evenly spaced around the circular trajectory of the LHC.



**Figure 3.3:** LHC quadrupoles. Top left: the coil of the quadrupole magnet. Top right: schematic view of the magnetic fields in the quadrupole. Bottom: two beams and the IP.

To power the LHC, 1612 electrical circuits are used. Mostly these circuits are needed to power the dipole and quadrupole magnets, which is done in eight evenly spaced location

of the LHC. A total of 3286 current leads are needed to connect all the circuits and power cables. More than a thousand of the leads operate between 600 A and 13 kA (see Fig. 3.4). The other leads operate in the range 60 to 120 A.



**Figure 3.4:** 13 kA high-temperature superconducting current lead.

### 3.1.5.2 Cooling System

To ensure that dipoles are in the superconducting state, they have to be cooled to 1.9 K using superfluid helium-4.

The cooling (cryogen) system is needed to keep superconducting LHC magnets at the appropriate temperature. The choice of the cooling gas depends on the magnet type and location. This dictates the required range of temperatures, which differs from system to system by 75 K. The cryogen system uses layered design with the temperature becoming progressively colder going from outside the dipoles closer to the beam pipe.

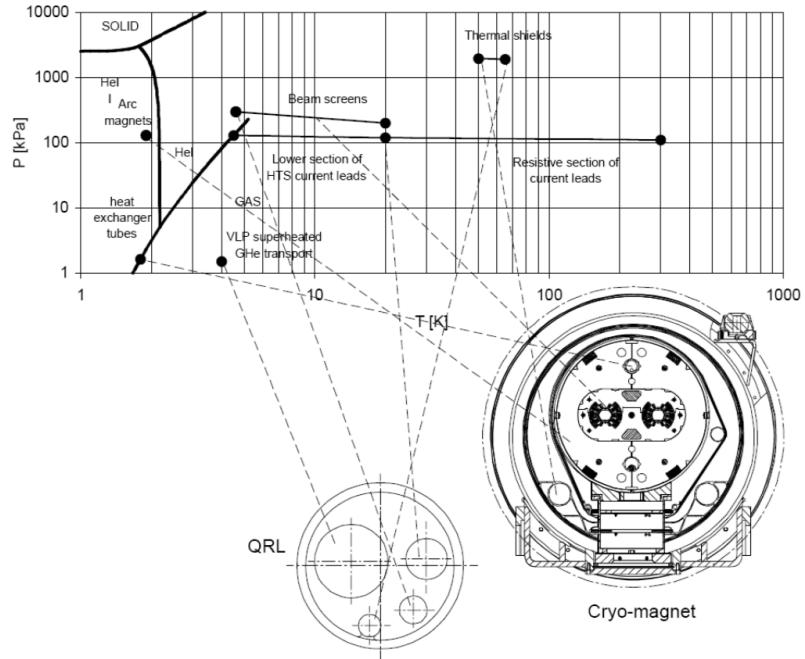
The "coldest" part of the cryogen system is designed for the inner part of the dipoles. This system (see Fig. 3.5) must cool down 37 Mkg of the LHC magnets within 15 days to

the required temperatures, which is done through the system of pipes that transports and directs the flow of the superfluid helium. The cryogen system must also be able to deal with the fast increases of the pressure flow and flow surges, as it is crucial for the LHC operation to keep dipoles constantly cooled and at the superconducting state.

The LHC tunnel is inclined in the horizontal plane by 1.41°. This translates to 120 m difference in the vertical location of two diametrically opposite points of the tunnel with respect to the surface level; and results in the additional hydrostatic pressure that can affect the flow of helium. This has been an important concern during the design of the cryogen system.

Since the cost to cool the LHC equipment to 1.8-1.9 K temperatures is high, several temperature levels are employed (see Fig. 3.5):

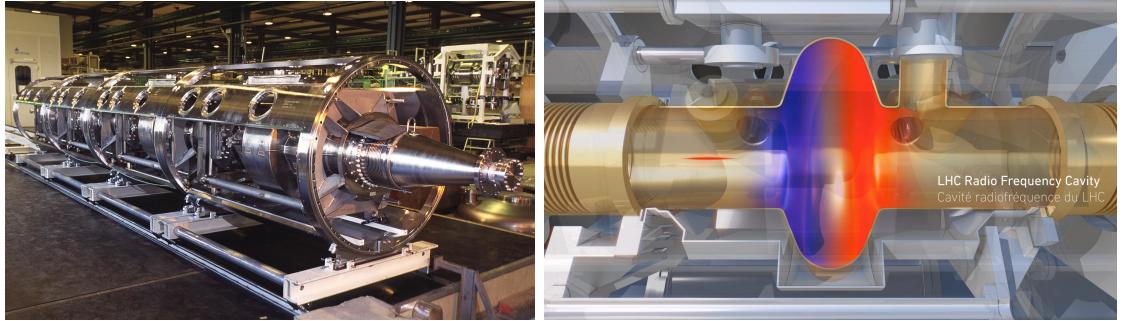
- 50 to 75 K for the thermal shielding used in the dipoles,
- 20 to 300 K for upper ("warm") sections of the high-temperature superconducting current leads,
- 4.6 to 20 K for lower temperature interception,
- 4.5 K for radio frequency cavities and lower ("cold") sections of the high-temperature superconducting current leads,
- 4 K for the transportation system that directs the 1.8 K helium to dipoles,
- 1.9 K for helium in the superfluid state to cool magnet masses.



**Figure 3.5:** LHC cryogenic states and the temperature scale.

### 3.1.5.3 Radio Frequency Cavities

Proton bunches need to be ramped to 7.5 TeV energies. To achieve this 13 TeV COM energy, eight superconducting radio-frequency cavities (RFC) are used per beam. They are located in front of the IPs of four experiments. Electromagnetic waves of 400 MHz with a peak field strength of 5.5 MV/m adjust the speed of protons in bunches. Each RFC (see Fig. 3.6) increases the energy of protons by 60 keV per revolution and it takes  $O(20)$  minutes to reach 6.5 TeV beam energy. The RFC frequencies are increased gradually by 1 kHz to match the speed up of protons in the bunch as they gain more energy. When the ramp is completed, the RFCs are used to compensate for small energy losses due to the synchrotron radiation (7 keV per revolution).



**Figure 3.6:** LHC RF cavities. Left: a cryomodule with four RF cavities. Right: a schematic drawing of a single RF cavity. The colour field is used to denote positive (red) and negative (blue) polarities. A narrow beam traversing the cavity is coming from the left and is shown in red.

### 3.1.5.4 Vacuum System

The work of the LHC depends on three vacuum systems [46]. Without them, dipoles will not be at the superfluid state, the beams will not be able to circulate, and no stable collisions would be taken. With a total of 104 kilometres of vacuum pipes, the LHC owns the largest vacuum system in the world. The main types of vacuum systems are:

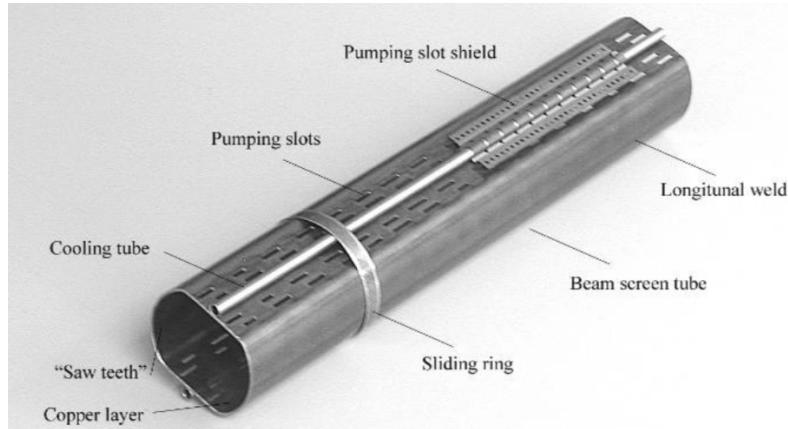
- insulation vacuum for cryomagnets,
- insulation vacuum for the helium distribution line,
- beam vacuum.

The insulation vacuum is needed to ensure the operations at both low temperatures of the magnets and the room temperatures in the tunnel. The insulation vacuum of  $10^{-6}$  mbar is used for a total of 15000 cubic metres. To build this vacuum system, the LHC used 250,000 welded joints and 18,000 vacuum seals.

The vacuum for the helium distribution lines is needed to protect from the heat the flow of the helium-4. This helium flow is used to cool down the dipole mass. Cryogenic distribution lines (QRL) of 3.3 km each are connected to eight cryogenic plants that pump the helium-4 into the LHC. The vacuum in these systems is at  $10^{-7} - 10^{-10}$  mbar level.

For the beam pipes the LHC uses ultra-high vacuum of  $10^{-10}$  mbar at cryogenic temperature of 5 K. The vacuum is getting progressively closer to  $10^{-11}$  mbar near the IPs, because in these locations collisions take place and any additional gas is highly undesirable. This vacuum is the emptiest space in the Solar System. This ultra-high vacuum is needed to reduce the beam degradation due to the beam-gas interactions in the pipe and parasitic collisions of bunches with the collimators near the IPs.

Vacuum system are affected by the heat produced from the synchrotron radiation emitted by the proton beams when they are bent. To reduce the amount of this heat and to narrow down the beam size in the transverse direction when the beam widens, the LHC uses "beam screens", which operate between 5 and 20 K.



**Figure 3.7:** Beam screen.

The beam screens are necessary to reduce the number of protons scattering on the residual gas of the beam pipes, which could lead to a magnet quench and even interrupt the machine operation.

The table below summarises the main heat sources that degrade the vacuum quality in the beam pipe, where the vacuum must exist at 1.9 K:

- synchrotron radiation ( $0.2 \text{ W/m}$  per beam),
- energy loss by nuclear scattering ( $30 \text{ mW/m}$  per beam),

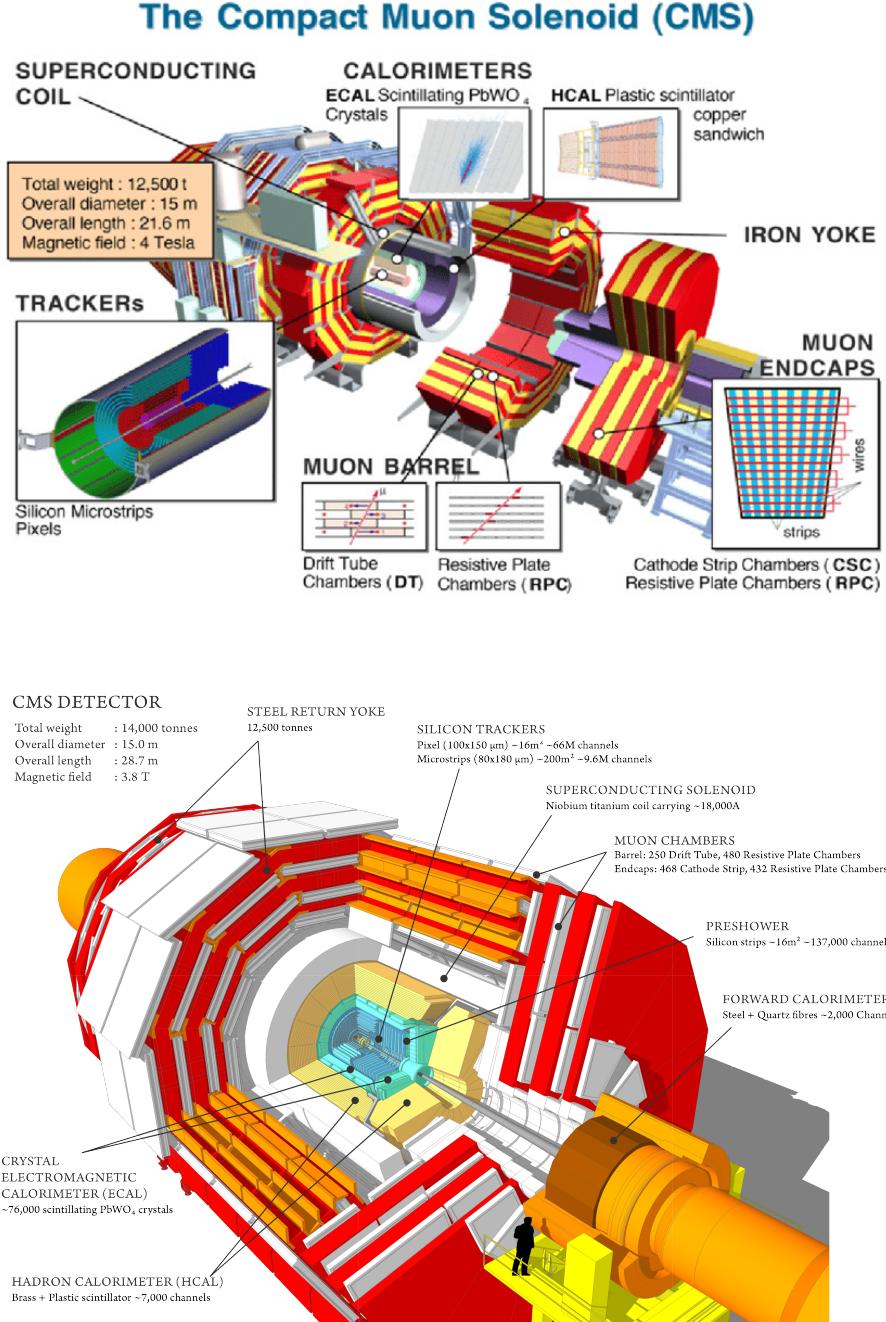
- image currents ( $0.2 \text{ W/m}$  per beam),
- electron cloud related effects (vary).

Now, that we discussed the LHC collider, we can continue with one of the main LHC detectors - the CMS detector - the one that was used to collect the data analysed in this thesis.

## 3.2 The CMS experiment

The Compact Muon Solenoid (CMS) is a multi-purpose particle detector built to study a variety of complex particle interactions produced by the LHC. CMS is located in the underground cavern at the "Point 5", which is one of the four main IPs of the LHC. The CMS detector with the additional computing infrastructure is able to detect the produced particles, measure their main physics parameters, and to send the related data to computing data centres for persistent storage.

The CMS detector has a cylindrical shape and consists of a central ("barrel") and two forward ("endcaps") sections (see Fig. 3.8). CMS is the heaviest detector ever built with the mass of nearly 12500 tons. The mass is explained by the amount of the used superconducting metal, which serves as the magnet. The CMS is 21.6 m long and 14.6 m high. The CMS has an onion-like structure of concentric layer of detectors around the IP. In addition, at the outer part it has a large superconducting solenoid to produce inside the detector a homogeneous magnetic field of 3.8 T.



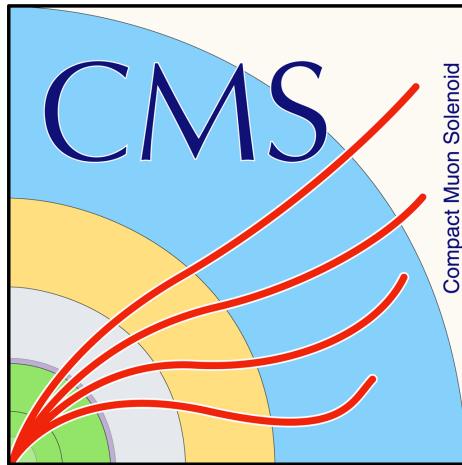
**Figure 3.8:** CMS experiment with the main sub-detectors.

All sub-detectors can be categorised into trackers and calorimeters [47]. As the particle passes through the material of the tracker, it leaves a "track", which is a path of the emerging particle. Trackers focus on the direction and the track curvature of the charged particles.

Tracking information allows the determination of the particle's momentum.

There are two trackers in CMS: an inner tracking system that encloses the IP and the outer tracking system that is located outside of the solenoid magnet. The first system contains the Pixel and the Strip trackers. The second tracking system is dedicated for the muon detection and is usually called a muon tracker or a muon system. This system is embedded within a steel yoke of the magnet.

The magnet yoke is made of five barrel wheels. Such an arrangement saves the CMS some space and also is used for the magnetic flux return. Additionally, it serves as a support for the embedded muon system, which is located outside of the ECAL and HCAL systems. Muons are energetic enough to traverse the ECAL and leave the detector. This muon system-magnet yoke structure provides a return field of the magnet of about 2 T and is used to measure the momentum of muons. This "two-directional" magnetic field with respect to the magnetic yoke, causes the muons trajectories to be bent in opposite directions in the inner tracker in contrast to the outer tracker. This important feature of the CMS detector is depicted in the CMS logo (see Fig. 3.9).



**Figure 3.9:** The logo of the CMS experiment that is showing curved trajectories of the emerging muons.

The CMS has two calorimeters: the electromagnetic and the hadronic calorimeters. They both rely on high density materials either to sample or to contain almost all the energy

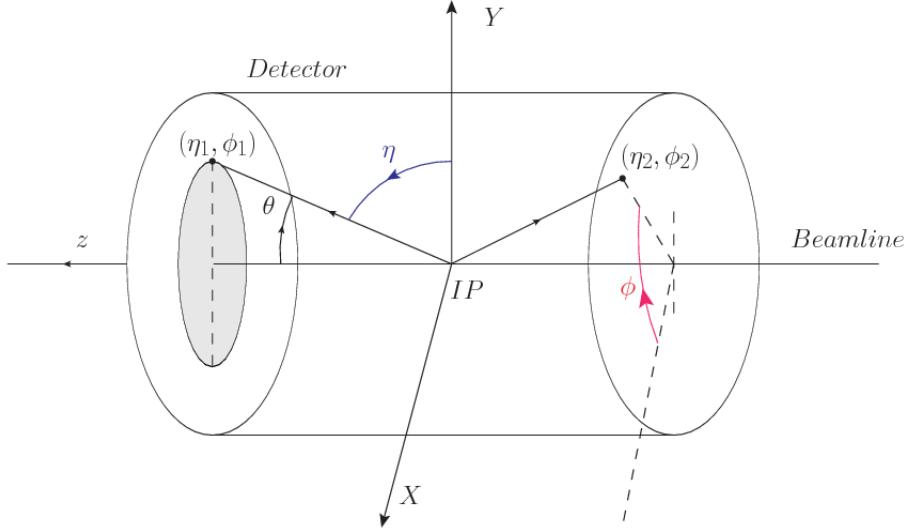
of the incoming particles with their secondary interaction products. However, these two systems focus on two different sets of particles. As will be discussed later, electromagnetic calorimeter (ECAL) is dedicated to measuring the energy of photons and electrons, while the hadronic calorimeter is targeting the measurement of the energy of hadrons.<sup>1!</sup> The rate of the incoming data at the LHC is 40 MHz. This corresponds to almost 70 TB produced every second! It is impossible to store that much data, and, most importantly, most of the information in this data is not interesting for future physics analyses. To reduce the data rate, the CMS uses a highly efficient system of triggers. The first one, the Level-1 (L1) trigger, reduces the nominal collision rate of 40 MHz to 100 kHz. The subsequent High-Level Trigger (HLT) further decreases the rate to 1 kHz. With the help of the trigger system, the original 40 TB per second rate is transformed into manageable 1 GB per second that is stored for offline analysis use.

### 3.2.1 The CMS coordinate system

The CMS uses a right-handed Cartesian coordinate system to define the axes of the colliding beams (see Fig. 3.10). The centre is located at the IP and the x axis points to the centre of the LHC ring. The y axis points upwards, and the z axis points along the proton beam direction. Since the CMS detector has a cylindrical shape, the polar system is used in the x-y plane: a standard set of the azimuthal angle  $\varphi$  and the radial coordinate  $r$ . The polar angle  $\theta$  is defined in the r-z plane and a widely used in this thesis angular variable  $\eta$  (called pseudorapidity) is defined as  $\eta = \ln \tan(\theta/2) = \ln(\frac{|\vec{p}|+p_z}{|\vec{p}|-p_z})$ . Additionally, a popular quantity in the collider physics - the rapidity - is given by  $y = 1/2 \ln(\frac{E+p_z}{E-p_z})$ . Rapidity is a function of the energy E and longitudinal momentum  $p_z$  of the particle (the projection of  $\vec{p}$  on the z axis. Note that  $\eta$  converges to  $y$  when the mass is negligible and the particle travels with the speed close to the speed of light. Most angular variables that are used currently in the modern high-energy physics (HEP) are defined in terms of  $\eta$  and  $\varphi$ :  $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\varphi)^2}$ , with  $\Delta\eta$  and  $\Delta\varphi$  being the absolute values of the relative

differences of  $\eta$ 's and  $\varphi$ 's of two particles.

Another extremely useful quantity is the projection of the momentum of a particle on the transverse plane and is called "transverse momentum"  $p_T$ . This variation of the momentum is independent of the z axis, hence, from the Lorentz boost. Similarly, the "transverse energy" of a particle is defined as  $E_T = \sqrt{m + p_T}$ .



**Figure 3.10:** Coordinate system of the CMS detector ???. Two particles (1 and 2) are shown with the corresponding angular variables  $(\Delta\eta_1, \Delta\varphi_1)$  for the first and  $(\Delta\eta_2, \Delta\varphi_2)$  for the second particle respectively.

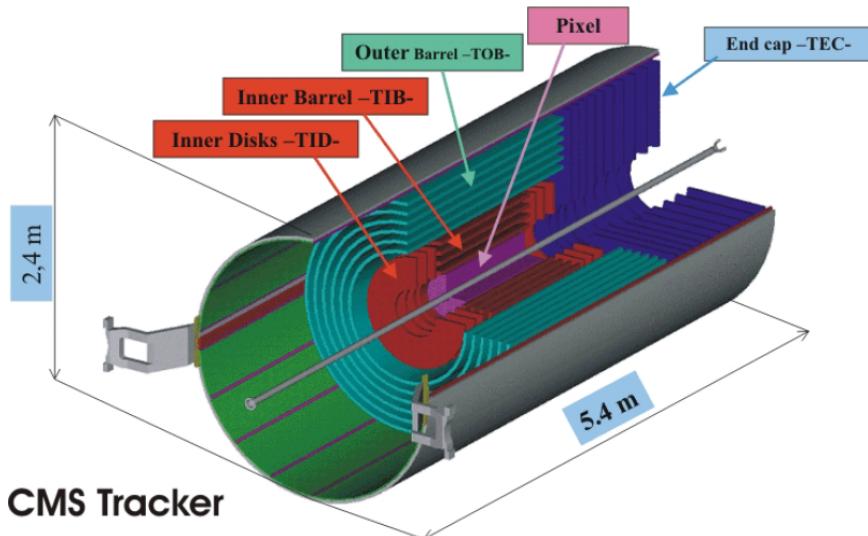
### 3.2.2 The Inner Tracker

The inner tracker [48] (see Fig. 3.11) is the closest subdetector to the IP. Using the tracker the experiment measures the trajectories of charged particles and reconstructs decay vertices. Since this system is constantly under the radiation coming from the interactions with the particle flux of nearly 100 MHz/cm at  $r = 4$  cm, the design of the tracker focused on two main requirements: high granularity for precise determination of the vertices and tracks, and robustness against the radiation-hard environment with the operational time of at least 10 years. As a solution to both challenges, the CMS relies on the silicon technology that provides the tracker with the large surface of thin but highly granular active detectors. The

tracking system has a diameter of 2.4 m and a length of 5.4 m covering the detector space of  $|\eta| < 2.5$ .

The inner most part of the tracker - the Pixel detector ("Pixel")- consists of three layers in the barrel at the radii of 4.4 cm, 7.3 cm, and 10.2 cm respectively. The Pixel also has two detector disks in forward regions. They are positioned 34.5 and 46.6 cm away from the IP. The Pixel is made of 1440 modules which contain 66 million pixel cells. Each cell is 100 by  $150 \mu\text{m}$  with  $285 \mu\text{m}$  thickness, which allows the determination of "hit" positions (the passage of the particle through the Pixel cells) in two directions  $z-\varphi$  in the barrel and  $r-\varphi$  in the endcaps.

The spatial resolution of each pixel about  $10 \mu\text{m}$  in the  $r-\varphi$  plane and  $20 \mu\text{m}$  along the  $z$  direction. The spatial information that comes from the tracker is used to determine the main interaction point of the hard scattering ("the primary vertex") and also additional interaction vertices ("pileup"). Tracker also helps to reconstruct the displaced vertices ("the secondary vertices") of the particles that decay relatively fast, e.g., b-jets, which will be discussed later in this chapter.



**Figure 3.11:** The inner tracker. Pixel and Strip detectors are shown.

The outer part of the inner tracker is the strip tracker. It contains several subsystems and

is made of almost 9.3 million strips arranged in different configurations in 15148 modules. The first subsystem is the tracker inner barrel (TIB), which consists of the four barrel layers of strip modules. The second subsystem is the tracker inner disks (TIDs), which is made of three disks of strip modules. Increasing the radius to about 60 cm, the tracker outer barrel (TOB) starts. TOB is made of six layers of strips. Finally, to cover high  $\eta$  regions, the tracker endcaps (TECs) are used, which are made of two sets of nine disks of strips.

Each strip is about  $O(20)$  cm long. Its thickness varies from  $320 \mu\text{m}$  for TIB and TID, to  $320 \mu\text{m} - 500 \mu\text{m}$  for TOB and TEC, respectively. Also width changes from  $80 \mu\text{m} - 141 \mu\text{m}$  for TIB and TID, to  $97 \mu\text{m} - 184 \mu\text{m}$  for TOB and TEC, correspondingly. The resolution on the single point in the radial direction is  $20 - 50 \mu\text{m}$ , and in the z direction it varries from 200 to 500  $\mu\text{m}$ , depending on the value of r.

All subsystems of the inner tracker have to be cooled down to about  $-20^\circ$ . This requirement is needed to minimise the damage of the tracker caused by the radiation from the collisions and to reduce overheating of the electronics.

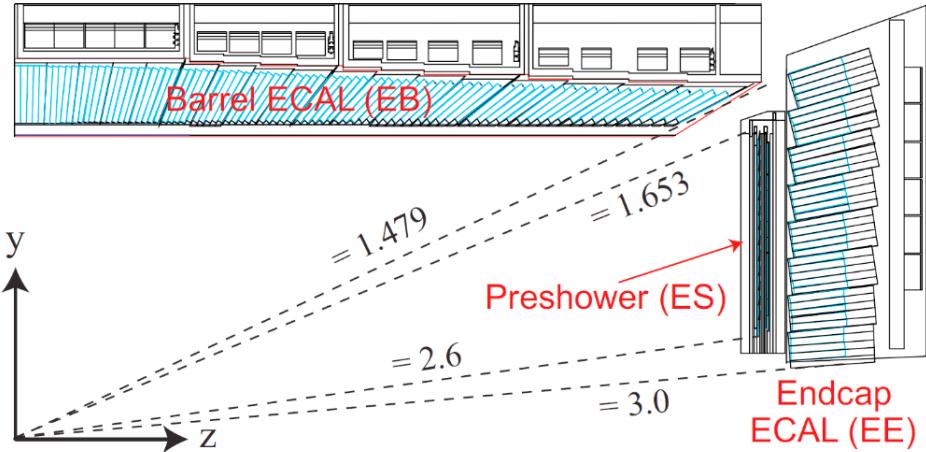
The material of the inner tracker has 0.4 to 1.8 radiation lengths ( $X_0$ ), which corresponds to 0.1 to 0.5 nuclear interaction lengths ( $\lambda_i$ ). Numbers vary with the  $\eta$ .

### 3.2.3 The ECAL

The inner tracker and the ECAL provide the detector with complementary measurements. The tracker focuses on the direction and the momentum of the particle and identifies only charged particles. The ECAL [49] (see Fig. 3.12), on the other hand, determines the energy of the particles and detects all particles that interact electro-magnetically, including photons and neutral pions. However, primarily the ECAL is designed to measure precisely the energy of electrons and photons.

The ECAL is a highly granular detector that relies on the lead tungstate crystal ( $\text{PbWO}_4$ ) technology. Electrons and photon passing through the crystal interact with its material and their energy is converted into the produced electromagnetic "shower".  $\text{PbWO}_4$  crystals are

known for being a popular choice of the scintillators: interactions with the crystal material produce the scintillation light that is further read out by the electronics. The  $\text{PbWO}_4$  crystals have a high density ( $8.28\text{g}/\text{cm}^3$ ), a small radiation length ( $X_0 = 0.89\text{ cm}$ ), a short Moliere radius ( $R = 2.2\text{ cm}$ ), and a fast response (80% of its scintillation light is produced within 25 ns). These characteristics are making  $\text{PbWO}_4$  crystals ideal candidates for the ECAL, since they guarantee an excellent containment of the electromagnetic shower within the crystals.



**Figure 3.12:** The ECAL and the Preshower detectors.

The ECAL has a barrel part (EB), covering the  $|\eta| < 1.479$ , and two endcaps (EE) covering  $1.479 < |\eta| < 3.0$ . In the barrel ECAL is made of 61 200 crystals. Each crystal is 22 by 22 mm with a length of 23 cm. In the endcaps ECAL has 7324 crystals. There each crystal is 28.62 by 28.62 mm with a length of 22 cm. The crystals' layout is following a quasi-geometric projection with axes of crystals slightly tilted to ensure particle trajectories are never aligned with the intercrystal cracks. This layout is optimised for the best particle shower containment with respect to the position of the interaction point.

The resolution of the ECAL is a function of energy of the incident particle  $E$  and can be decomposed into three terms. The first term is a stochastic term that is inversely proportional to the square root of the number  $N$  of scintillation photons produced in the

interaction. In the main formula N is replaced by E, since N is proportional to E. The second term is a "noise" term that describes the noise in the detector. The third term is related to detector imperfections and is represented by a constant C. The final dependence of the ECAL energy resolution  $\sigma$  on the particle energy E is given by:

$$\left(\frac{\sigma}{E}\right)^2 = \left(\frac{S}{\sqrt{E}}\right)^2 + \left(\frac{N}{E}\right)^2 + C^2 \quad (3.1)$$

From the dedicated calibration studies, the parameters in the formula above are found to be equal to: S = 2.8%, N = 12%, and C = 0.3%. As a "standard" procedure, the CMS often optimises the performance of the subdetectors for 45 GeV electrons, since they correspond to a classical Drell-Yan decay of Z boson to two electrons. In this case, a typical energy resolution for 45 GeV electrons is about 2% in EB and 2-5% for EE. Near the Z peak (91 GeV), the constant terms dominates the resolution.

The ECAL is operated at a temperature of 18 ° C and the "active width" of the ECAL material corresponds to 25 X<sub>0</sub>.

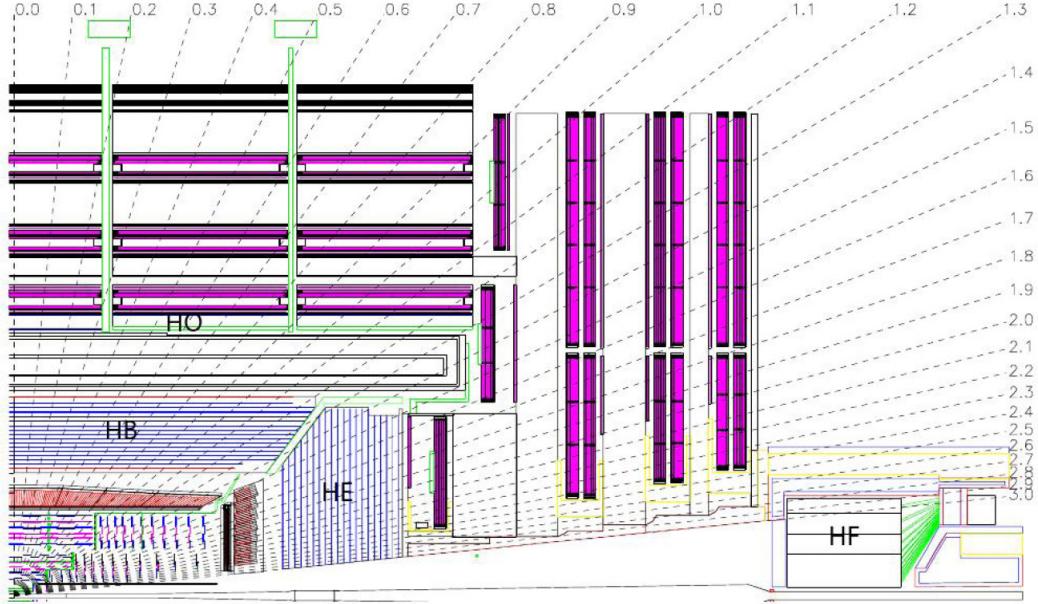
An additional subdetector, called the "Preshower", is installed right in front of the EE and covers  $1.653 < |\eta| < 2.6$ . The Preshower is designed to improve the discrimination of single photons from diphoton decays of neutral pions  $\pi^0 \rightarrow \gamma\gamma$ . This is a sampling calorimeter in which the material that produces the particle shower is distinct from the material that measures the deposited energy. Typically the two materials alternate. The Preshower has two lead layers which launch the electromagnetic showers. This "samples" the energy of the particles traversing the Preshower material. After these layers, 2 mm-wide silicon strips are placed. They measure the deposited energy and transverse profile of the shower shape initiated by the lead layers. The "thickness" of the Preshower material corresponds to 3 X<sub>0</sub>.

### 3.2.4 The HCAL

Hadrons normally go through the ECAL layers without being stopped. To absorb these particles, the HCAL [50] (see Fig. 3.13) is placed around the ECAL. The HCAL focuses on particles that hadronise. This is a process of the formation of hadrons out of quarks and gluons. The HCAL detects with the charged and neutral hadrons such as pions, kaons, protons, and neutrons. Hadrons also produce collimated streams of secondary particles (jets) and these jets are identified by the HCAL. Additionally, the HCAL is used to measure indirectly the transverse energy of neutrinos, by the momentum imbalance technique, which will be discussed later in this chapter.

The HCAL is split into HCAL barrel (HB) and HCAL endcap (HE) sections. They cover  $|\eta| < 1.3$  and  $1.3 < |\eta| < 3.0$  respectively. HB and HE are sampling calorimeters. They are made of a brass absorber and of active plastic scintillating tiles. The brass plates in HB have thickness of 56.5 mm and in HE the thickness if increased to 79 mm. The absorber material corresponds to  $5.82 \lambda_I$  at  $\eta = 0$  to almost  $10 \lambda_I$  at  $|\eta| < 1.3$ .

The gaps in the absorber of the HCAL are filled with an active medium of 70000 plastic scintillator tiles. The scintillation light is guided by wavelength shifting fibres (WLSs) to hybrid photodiodes (HPDs). The scintillator is quite fast with the 68 % of the light been collected within 25 ns.



**Figure 3.13:** The HCAL with the  $\eta$  coverage map.

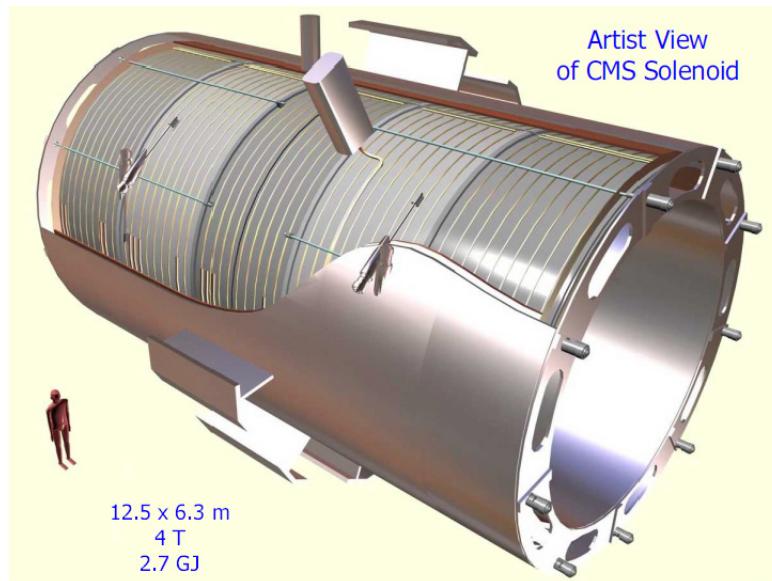
The CMS also has an outer calorimeter (HO) placed above the HB outside the solenoid. HO is called a tail catcher system and increases the total calorimeter thickness to  $11.8 \lambda_I$  in the barrel, with the magnet coil working as an extra absorption layer. The HO consists of five rings of scintillator tiles. A supplementary iron plate of 19.5 cm in thickness and a second layer of sensitive material are placed around  $\eta = 0$  to enhance the absorber depth there.

In the forward directions, two forward calorimeters (HF) extend the coverage to  $|\eta| = 5.2$ . The HF is composed of steel absorbers and quartz fibres that produce Cherenkov light when the particle in the material travels faster than the light in that medium. The light is further collected by photomultiplier tubes (PMTs).

Since the HCAL is located between the ECAL and the internal surface of the solenoid, the space allocated for the HCAL is not enough for the HCAL to fully absorb the hadronic showers and this imperfect containment of the hadronic shower limits the performance of the HCAL. Comparing with the formula 3.1, for the single pions the values are given by: by S = 115 %, N = 52 %, and C = 5.5 % [51].

### 3.2.5 The Superconducting Solenoid

The NbTi superconducting solenoid (see Fig. 3.14) of 6 m in diameter is the core of the CMS experiment. The magnet operates at a temperature of 4.5K. The bulk of the CMS detector weight (90 %) comes from the magnet steel return yoke and structural supports which together weigh 12500 tonnes.



**Figure 3.14:** The CMS superconducting solenoid. The person on the left is shown to emphasise the size of the magnet.

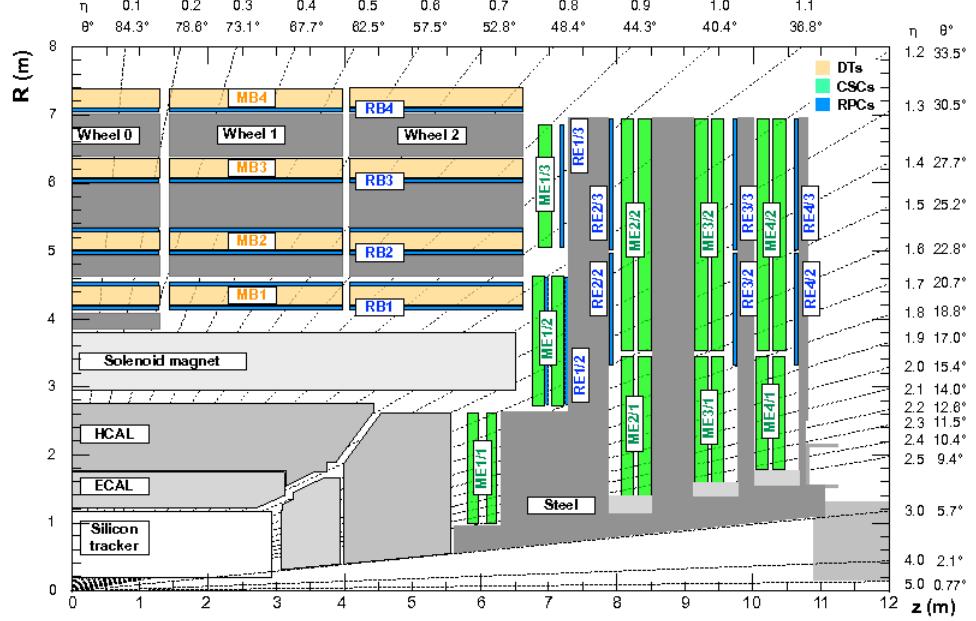
The solenoid is central part in the CMS detector design. The idea was to have a uniform magnetic field capable of bending the trajectories of charged particles as they traverse the detector. When a low energy particle is produced, it has a helical path and will be fully contained within the detector. On the other hand, when a highly energised particle is produced, the trajectory is seen as a "straight" incomplete arc. Both situations lead to imperfect measurement of the momentum. The primary measurements of the tracking system are presumed to be Gaussian distributed; but the momentum of the particle that the tracker measures is not Gaussian distributed. However, the sagitta is Gaussian distributed, and that is why widely used in the particle physics.

When the particle in the magnetic field passes thorough the material of the detector, the path deviates from the ideal circular line due to random fluctuations and multiple scattering. The sagitta term is used to quantify the depth of the circular arc and is equal to the distance from the centre of the arc to the centre of its base. Since the sagitta is following a Gaussian distribution, it may be approximated by simpler expressions in many calculations of the momentum resolution.

The magnetic field strength  $B$  and the length of the track  $L$  are dictated by the design of the detector. Since the momentum resolution is given by  $\sigma_p/p^2 \approx \sigma_x/BL^2$  (see [47]) and improves linearly with magnetic field  $B$ , this is the reason why the CMS decided to invest much of the detector space and budget in the magnet. For a track of the length of  $O(1)$  m in the magnetic field of  $O(3)$  T, the sagitta is equal to 1 mm, which can be measured very precisely.

### 3.2.6 The Muon Tracker

Many physics analyses in the CMS rely on a precise measurements of the muons in the detector. Although muons are detected by the inner tracker, that information cannot be used by the trigger (will be discussed in the following subsection). Therefore, CMS has an outer tracker or muon tracker [52] (see Fig. 3.15) located outside the calorimeters and the solenoid. Because of the typical muon energies, muons produced in collisions at the LHC traverse the detector material with the minimal energy losses. To measure energy of muons, the CMS uses the muon tracker, which relies on various gaseous detector technologies. The muon tracker is inserted into the gaps of the flux-return yoke. Tracks in the muon system are used to reconstruct standalone muons, and in combination with the inner tracker, to reconstruct the global muons.



**Figure 3.15:** The CMS muon tracker. DT, CSC, and RPC detectors are shown in yellow, green, and blue respectively.

CMS muon system has three subdetectors: the drift tubes (DTs), the cathode strip chambers detectors (CSCs), and the resistive plate chambers (RPCs). In the barrel region, the CMS is equipped with the DT system, which is 250 drift tubes arranged into five barrel section ("wheels"). Each wheel is made of four concentric rings of DT stations. The working elements of the DT system - cylindrical cells with the rectangular base of 4.2 by 1.3 cm<sup>2</sup> - are tubes with an anode wire in the mix of argon and CO<sub>2</sub> gases. DT cells are 2.4 m long and are organised in three groups of four elements (three "super-layers"). When the muon passes through super-layers, it ionises the gas in the cells and released electrons start moving to anodes. Using the time it takes for electrons to reach the anodes, the muon position and direction can be determined. DT resolution of a single-cell hit positions ranges from 200  $\mu$ m in the  $r\varphi$  plane to 200-600  $\mu$ m for forward directions.

CSCs are used in the forward direction to cover the region of  $0.9 < |\eta| < 2.4$ . CSCs chambers are multi-wire chambers made of cells that have a trapezoidal shape. Chambers contain radial copper cathode strips and, perpendicular to those, gold-plated tungsten anode

wires. Each cell is filled with the mix of Argon, CO<sub>2</sub>, and CF<sub>4</sub> gases. The strip cells have a single-layer resolution of 300-900  $\mu\text{m}$ . A CSC chamber provides a spatial resolution of 40 -150  $\mu\text{m}$ .

To improve the performance of DTs and CSCs, RPCs are used and are covering the barrel and endcaps in the range of  $|\eta| < 1.9$ . RPCs are double-gap chambers consisting of two resistive 2 mm in thickness Bakelite layers separated by a 2 mm layers filled with a mix of C<sub>2</sub>H<sub>2</sub>F<sub>4</sub>, *i*C<sub>4</sub>H<sub>10</sub>, and SF<sub>6</sub> gases.

RPCs operate in avalanche mode, producing an avalanche when the muon traverses the gas of the cell. RPCs have a spatial resolution of 0.8 - 1.2 cm, which is not as good as the ones provided by other muon subsystems, but RPCs have an advantage in terms of an excellent time resolution - just 3 ns. The barrel and the endcaps contain in total 10 RPC stations.

### 3.2.7 The Triggers and DAQ

The CMS trigger [53] is a system responsible for selecting events of interest and storing them for the offline analysis. The trigger has two stages: the L1 trigger (see Fig. 3.16), which reduces the event rate from 40 MHz to 100 kHz, and the HLT trigger, which further decreases the rate to nearly 1 kHz. The L1 trigger consists of the custom hardware that processes a part of the information from calorimeters and outer tracker systems. The HLT trigger is a part of the detector readout system (DRS) and uses the full detector information for event reconstruction. The HLT is a computing farm consisting of 22000 CPU cores that produce a decision on whether to save or to skip the event in an average time of about 220  $\mu\text{s}$ . DRS is integrated in the higher level data acquisition (DAQ) system [54]. The selected events are collected and sent by the DAQ to the tapes of the CERN Tier-0 for the persistent storage.

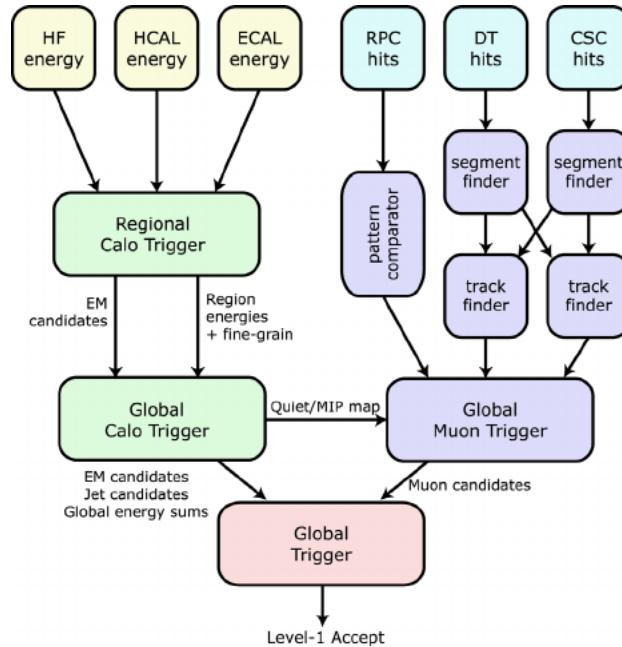
L1 and HLT systems have differences and similarities. They operate at the different time scales and the volumes of data they are processing are completely different. However,

the goals of these systems are the same - to identify and reconstruct physics objects and combine their properties to produce an acceptance/rejection decision for each event.

### 3.2.7.1 The L1 Trigger

L1 system [55] contains a "menu" of 500 algorithms or "seeds" designed to identify useful physics events. These menus include trigger criteria varying from basic single-object identification to complicated selections requiring some topological conditions to be met. Each seed has a set of assigned "Prescale" factors  $f$  that reduce the rate of events accepted by a particular trigger algorithm from 100% to  $100/f\%$ . Prescale factors are necessary since the luminosity level decreases during the run period. They adjust the trigger rate to keep it constant during the data taking time.

Since the processing time of the L1 system is very important for the whole CMS operation, the L1 is built using FPGAs and ASICs custom hardware. L1 produces decisions within  $3.8 \mu s$ . Data from all the calorimeters is first processed by the L1 regional calorimeter trigger (RCT) and then by a more selective global calorimeter trigger (GCT).



**Figure 3.16:** The CMS L1 trigger layout.

The RCT receives the information about energy deposits from all calorimeters and covers the range  $|\eta| < 5$ . The RCT processes this information in parallel and produces  $e/\gamma$  candidates as an output. The information from the inner tracker is not available, therefore, L1 identifies both electrons or photons, but cannot distinguish them. L1 also detects jets, taus, missing transverse energy (MET or  $\cancel{E}_T$  or  $\cancel{\not{p}}_T$ ), and muons. L1 RCT is responsible for determining the first estimates of the several main parameters of interest:  $p_T$ , isolation (described later in this chapter), etc.

First stage reconstruction uses particle "hits" in the muon detectors and analyses them using track finder algorithms. All three muon detectors of the CMS are used by the L1 muon trigger. Using DT and CSC systems, track segments from the hit information are identified. The pattern recognition algorithms are applied to these segments to reconstruct muon candidates and measure their momenta.

More complex but slower algorithms then re-use hits for a more precise particle identification using a global muon trigger (GMT). The hits from the RPCs are used directly by the pattern comparator trigger (PACT) that reconstructs muon candidates at the high radii. Then several regional track finder algorithms sort the identified muon candidates and send this information to the GMT. Each candidate contains  $p_T$  and angular information.

The GMT then combines the muon information from different subsystems to avoid duplicating the candidates. The GMT also performs more tight track quality checks and may discard a portion of the input candidates it received.

Finally, the information from the GCT and GMT is combined by a global trigger (GT). The GCT sorts the created  $e/\gamma$  candidates, identifies jets, and calculates  $\cancel{E}_T$ . The final decision of the GT is to store or to skip the event. If the event satisfies the acceptance requirements and is going to be kept, the L1 accept signal (LAS) is generated and propagated by the trigger control and distribution system (TCDS) to all subdetectors.

The GT is the final step of the CMS L1 trigger system and implements a menu of triggers. The output of this system is used as an input to the HLT algorithms.

### 3.2.7.2 The HLT Trigger

The selection done by the HLT mimics the offline analysis - for all reconstructed objects in the event: electrons, muons, and jets - the identification criteria is applied to select only events of interest. What these events are each offline analysis defines in a different way, but to name a few, almost all analyses need true prompt leptons, well reconstructed jets, or some other commonly used objects.

The HLT computing farm has the event filter farm, which consists of filter-builder units (FBU). In the FBU the parts of the events and information from different detector subsystems is combined to produce "complete" events. Then the filter unit unfolds the raw detector data into experiment specific data structure and performs the event reconstruction and trigger filtering.

The whole event processing procedure of the HLT is centered around the HLT path. The HLT path is a set of algorithmic instructions that in a sequential manner reconstructs physics objects and performs the object selection. The complexity of the steps in the path sequence increases and the quality of the physics objects (the probability to have a correct label) improves too. After this step is completed, selected events are sent to another software processing farm. In this storage manager farm, the date is archived, stored locally on disk, and later sent to the CMS Tier-0 computing center for offline use.

Most data enters the queue for processing and is ready to be sent to Tier-0 very soon. In some cases the special data, the "parked" data, may be collected and kept until the run is finished. In this situations the CMS tape is used and the data has a high-priority for "parking". This may include "hot topic" analyses such as vector boson fusion or parton distribution studies such as Drell-Yan process.

The output of the HLT is limited by capacities of the Tier-0. This includes the bandwidth of the data transfer as well as the amount of available tape. This complicates the work of the DAQ, since in addition to physics data streams, the calibration streams also need to be stored. These streams, though, use information only from few subdetectors.

### 3.2.7.3 The DAQ system

The DAQ systems in the modern high energy physics are responsible for any tasks. The challenges are well known: high data rates and volumes, limited tape space and transfer bandwidth. CMS DAQ is based on the homogeneous architecture, scales well with the different beam energy regimes and data rates, and stable performance in a variety of operating conditions.

To illustrate an example of a complex computing task that is elegantly solved by the DAQ system, let us discuss in more details the aforementioned FBUs system of the HLT. FBU relies on a single multi-core machine and the communications with other units are done via shared memory. The data from the full detector is used for the filtering process. Complicated offline-like reconstruction algorithms are then used for the full precision event selection. With more CPU cores available for the Run-2, the per-event time budget is increased to "comfortable" 175 ms per event, which is a long enough time to run most of the CMS reconstruction algorithms.

The current CMS DAQ was developed to address these core requirements:

- The data from one or several data transfer lines is available for other lines,
- the event building is done in parallel profiting from multiple processing units,
- almost real-time process monitoring,

Proper design patterns are used in the software for DAQ, which decouple the user interface from the implementation. The design also allows for the remote control. The software system can be run on a number of different operating systems and hardware platforms. The memory management tools of the underlying system are not linked directly to the applications, it is done using a dedicated abstract addressing scheme.

### 3.2.8 The CMS design

Now that we discussed all CMS subdetectors and DAQ, we can look back and summarise in one list what were the requirements on the design of the CMS detector to successfully complete its physics program. Here we refer to the CMS Technical Design Report [55]:

- good muon momentum resolution over the momentum scale covering almost a TeV range, good dimuon resolution (mostly  $Z \rightarrow \mu\mu$  and  $H \rightarrow \mu\mu$ ) at the O(100) GeV. The capability to determine correctly the charge of the highly energetic muon all the way up to 1 TeV,
- good momentum resolution of all charged particles in the inner tracker,
- good diphoton mass resolution with the focus on the  $H \rightarrow \gamma\gamma$  discovery channel. Also, the ability to reject  $\pi^0 \rightarrow \gamma\gamma$ , which is one of the main background processes to many physics analyses. This requirement mostly concern the performance of the ECAL,
- good resolution of the missing transverse energy (discussed in the section below) and of the mass of the two-jet system. This task depends heavily on the performance of the HCAL.

In the next chapter we will dive deep into the physics analysis of the data produced by the LHC and collected with the CMS detector that has been used to perform the measurement of the double Higgs boson decays.

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