

# Jets at the LHC: a fixed order perspective

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# Outline

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→ Three jet observables at NNLO QCD

R32 ratios

Event-shapes

→ Flavoured jets

Infrared safe definition of jet flavour?

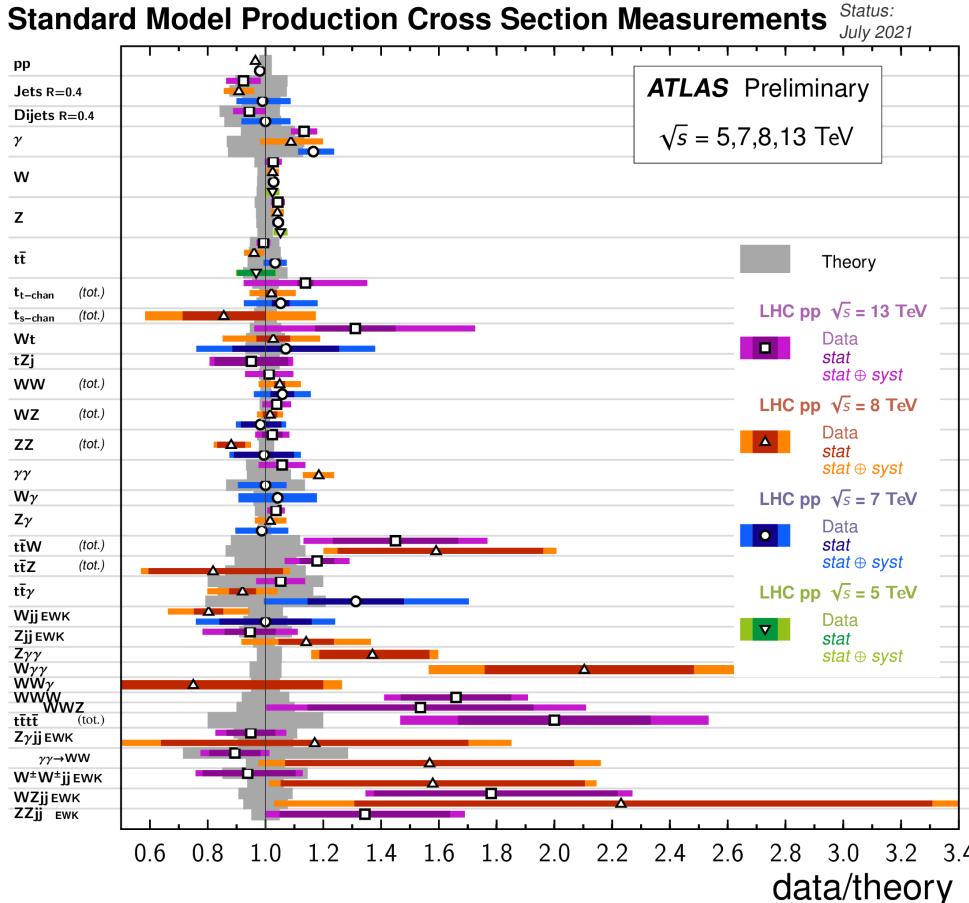
→ New proposal for a flavour safe algorithm.

→ Wrap-up and outlook

# SM measurements at the LHC

## Standard Model Production Cross Section Measurements

Status:  
July 2021



	$\int \mathcal{L} dt [fb^{-1}]$	Reference
pp	50x10 <sup>-3</sup>	PLB 761 (2016) 158
Jets R=0.4	8x10 <sup>-3</sup>	Nucl. Phys. B 848:548 (2014)
Dijets R=0.4	20.2	JHEP 09 (2017) 020
$\gamma$	20.2	JHEP 09 (2017) 020
W	20.2	JHEP 09 (2017) 020
z	20.2	JHEP 05 (2014) 050
tt	20.2	PLB 89 (2017) 077
$t_1$ -chan (tot.)	0.081	PRD 89 (2014) 073005
$t_2$ -chan (tot.)	0.081	PLB 759 (2016) 601
Wt	20.2	EPLC 79 (2019) 30
tZj	0.025	EPLC 79 (2019) 128
WW (tot.)	3.25	JHEP 02 (2017) 117
WZ (tot.)	20.2	JHEP 02 (2017) 117
ZZ (tot.)	0.025	EPLC 79 (2019) 128
$\gamma\gamma$	36.1	EPLC 80 (2020) 528
W $\gamma$	20.2	EPLC 74 (2014) 3109
Z $\gamma$	0.3	ATLAS-CONE-2021-003
ttW (tot.)	20.2	JHEP 04 (2017) 026
ttZ (tot.)	20.2	PRD 90, 112006 (2014)
ttg	20.3	PLB 756, 228-246 (2016)
Wjj EWK	20.3	JHEP 01 (2018) 63
Zjj EWK	20.3	JHEP 01 (2018) 66
Z $\gamma\gamma$	20.3	PLB 716, 142-159 (2012)
WW $\gamma$	139	JHEP 07 (2020) 124
WWZ	36.1	EPLC 79 (2019) 884
WW $\gamma\gamma$	20.3	EPJ C 763, 114 (2016)
WW $\gamma\gamma$	4.6	EPLC 76 (2019) 56
WWZ $\gamma$	20.3	PRD 93, 092004 (2016)
WWZ $\gamma$	4.6	EPLC 72 (2012) 2173
WWZ $\gamma$	38.3	JHEP 01 (2018) 020
WWZ $\gamma$	20.3	JHEP 03 (2018) 128
WWZ $\gamma$	1.6	arXiv:2107.09330 [hep-ex]
WWZ $\gamma$	139	JHEP 01 (2018) 086
WWZ $\gamma$	20.3	PRD 87, 112003 (2013)
WWZ $\gamma$	36.1	JHEP 03 (2020) 054
WWZ $\gamma$	20.3	PRD 93, 112002 (2016)
WWZ $\gamma$	4.6	PRD 93, 072003 (2019)
WWZ $\gamma$	20.3	JHEP 11, 173 (2015)
WWZ $\gamma$	139	arXiv:2103.12693
WWZ $\gamma$	20.3	JHEP 01 (2015) 015
WWZ $\gamma$	4.6	EPLC 79 (2019) 382
WWZ $\gamma$	20.2	JHEP 11 (2021) 086
WWZ $\gamma$	4.6	EPLC 91, 020007 (2015)
WWZ $\gamma$	20.2	EPLC 77 (2017) 474
WWZ $\gamma$	139	EPLC 81 (2021) 163
WWZ $\gamma$	20.3	JHEP 03 (2021) 203
WWZ $\gamma$	20.3	PRD 93, 112002 (2016)
WWZ $\gamma$	20.3	PRL 115, 031802 (2015)
WWZ $\gamma$	20.2	EPLC 77 (2017) 646
WWZ $\gamma$	78.3	ATLAS-CONE-2021-003
WWZ $\gamma$	78.3	arXiv:2106.11683
WWZ $\gamma$	139	ATLAS-CONE-2021-038
WWZ $\gamma$	20.3	JHEP 07 (2017) 077
WWZ $\gamma$	20.3	PRD 91, 063019 (2019)
WWZ $\gamma$	20.3	PRD 96, 032003 (2019)
WWZ $\gamma$	36.1	PRD 93, 161801 (2019)
WWZ $\gamma$	20.3	PRD 96, 092007 (2017)
WWZ $\gamma$	20.3	PRD 93, 092004 (2016)
WWZ $\gamma$	139	arXiv:2004.10612 [hep-ex]

New physics around the corner?

Precise measurements

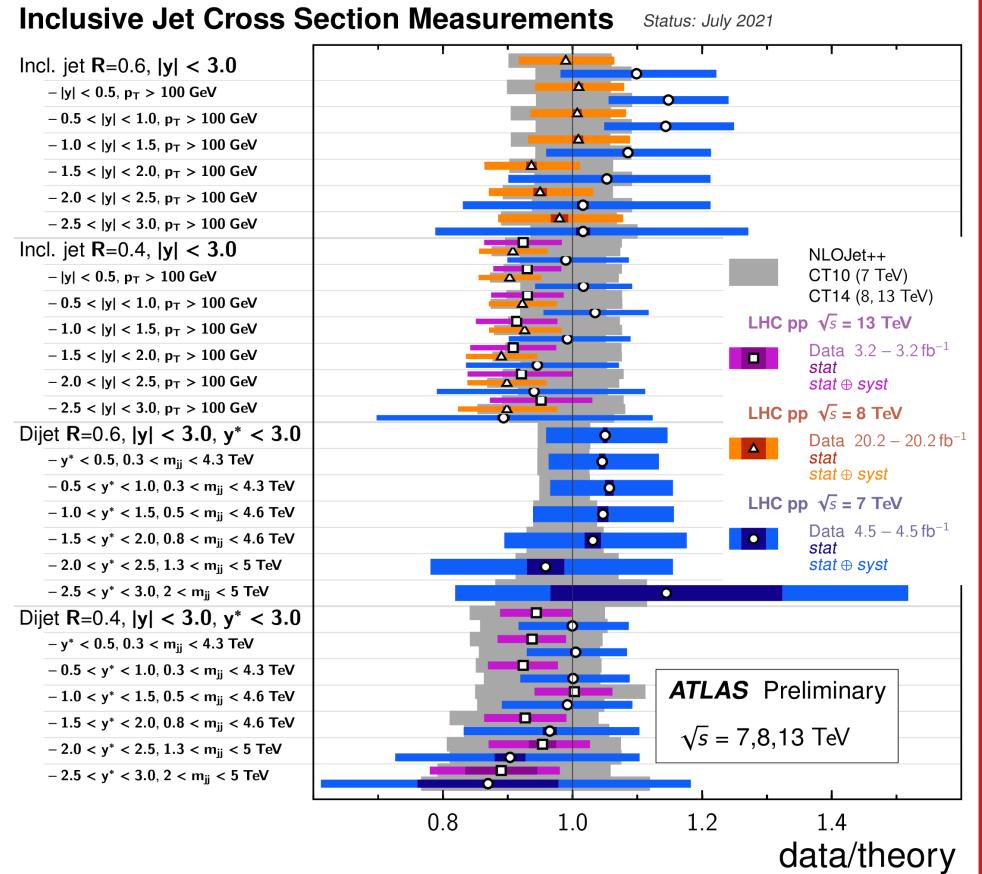
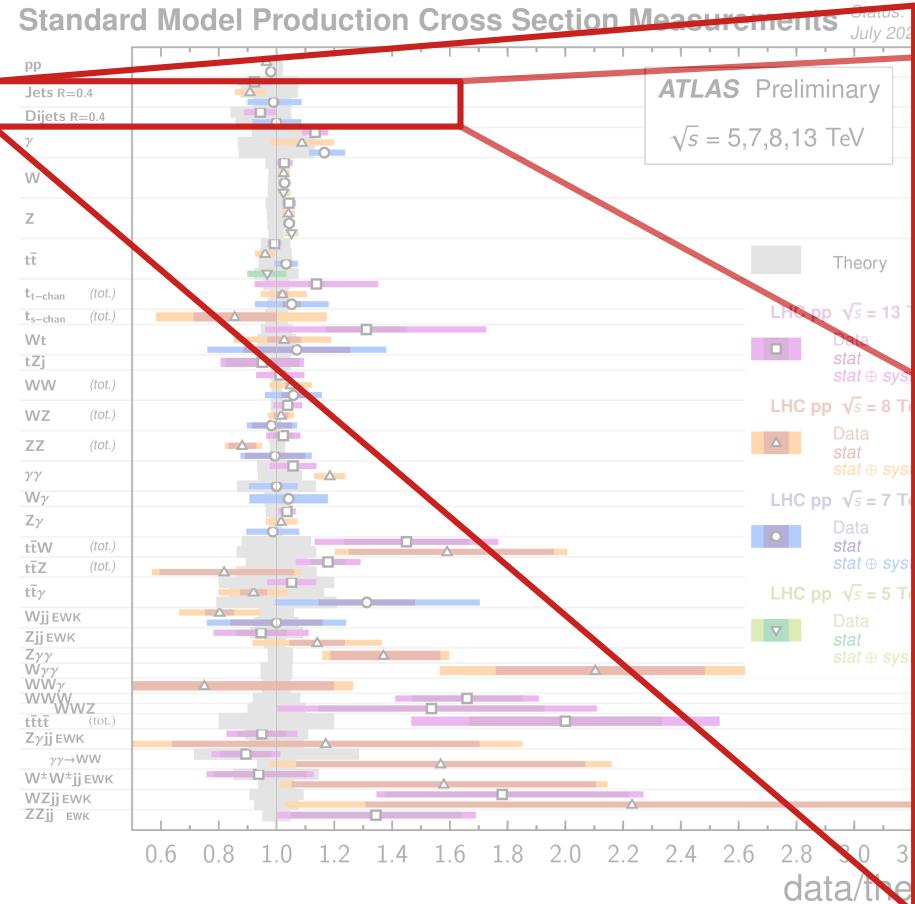
<->

Precise theory

Win-Win situation

- improved SM understanding
- possible indirect BSM signals

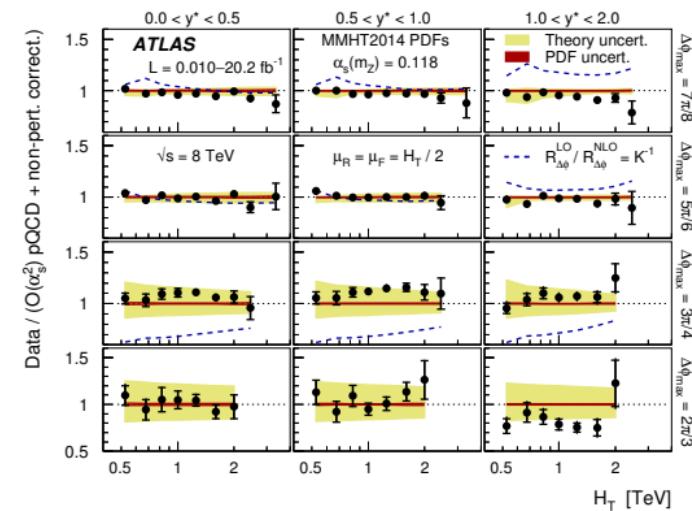
# SM measurements at the LHC



# Jet observables at the LHC

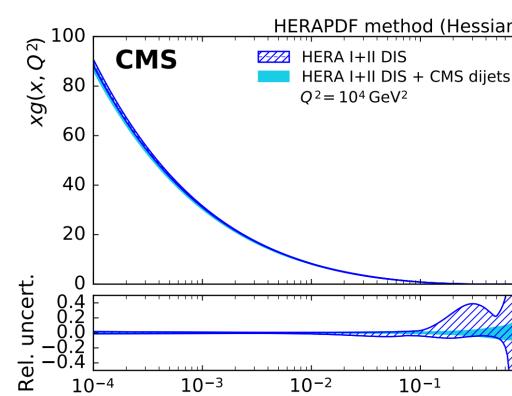
The LHC produces jets abundantly → many phenomenological applications

Tests of pQCD,  $\alpha_s$  extraction:  
R32 ratios, event-shapes



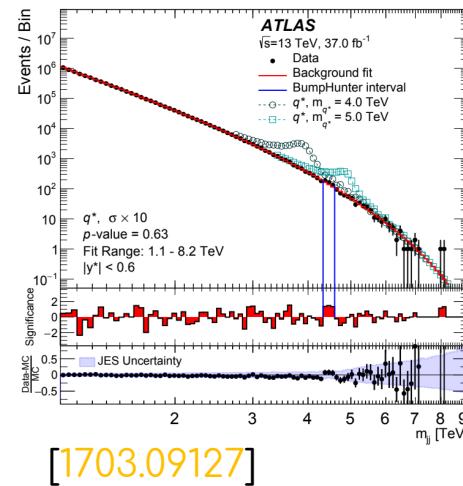
[1805.04691]

PDF determination:  
Single inclusive,  
Multi-differential dijet



[1705.02628]

BSM searches:  
dijet mass



[1703.09127]

Precision theory required!

Data driven

# Multi-jet observables at the LHC

Multi-jet final states:

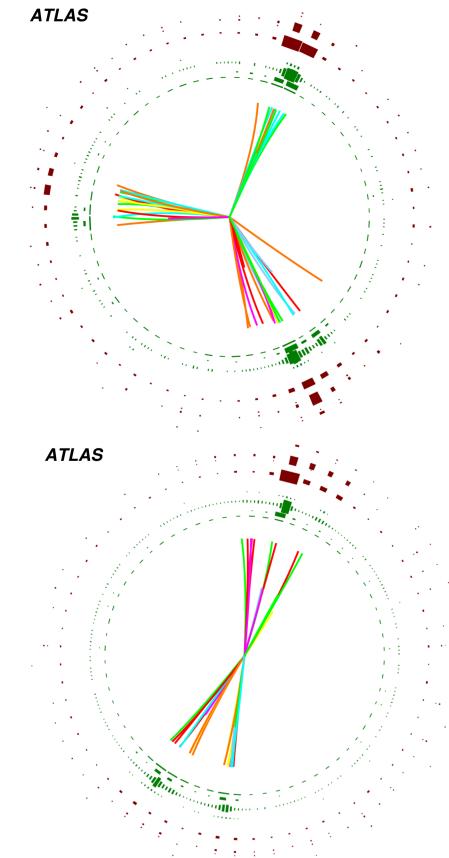
- Tests of pQCD at high energy
- Tests of MC modelling of LHC events
- Search for new physics

Study of perturbative QCD:

- R32 ratios

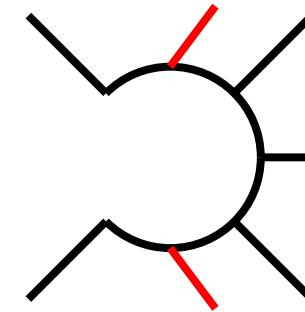
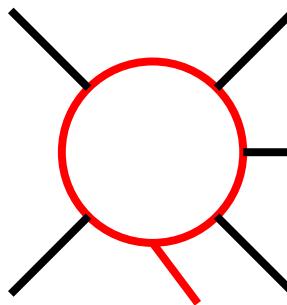
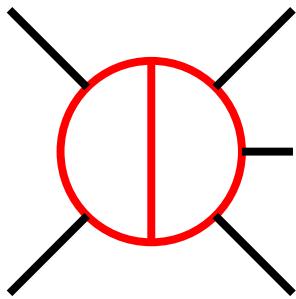
$$R_{3/2}(X, \mu_R, \mu_F) = \frac{d\sigma_3(\mu_R, \mu_F)/dX}{d\sigma_2(\mu_R, \mu_F)/dX} \sim \alpha_s$$

- Extraction of the strong coupling constant
- Transverse Energy-Energy Correlator
- Event shapes



Credits: [ATLAS:2007.12600]

# NNLO QCD prediction beyond $2 \rightarrow 2$



$2 \rightarrow 3$  Two-loop amplitudes:

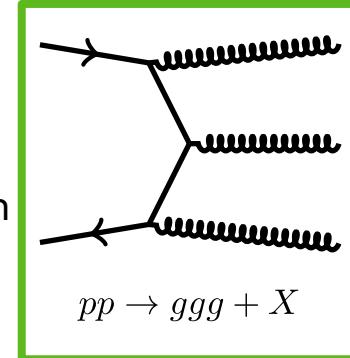
- Advances in amplitude techniques:  
IBPs, amplitude reconstruction and master integrals
- (Non-) planar 5 point massless amplitudes  
[Chawdry'19'20'21, Abreu'20'21, Agarwal'21, Badger'21]  
→ triggered by efficient MI representation  
[Chicherin'20]

Cross-sections → Combination with real radiation

- Various NNLO subtraction schemes available:  
qT-slicing [Catain'07], N-jettiness slicing  
[Gaunt'15/Boughezal'15], Antenna [Gehrmann'05-'08],  
Colorful [DelDuca'05-'15], Projectction [Cacciari'15],  
Geometric [Herzog'18], Unsubtraction [Aguilera-Verdugo'19], Nested collinear [Caola'17],  
Sector-improved residue subtraction [Czakon'10-'14,'19]

# Three-jet production

- Sector-improved residue subtraction [Czakon'10'14'19]
  - Efficient c++ implementation → STRIPPER
  - Highly automated to deal with enormous amount of channels in three-jet production  
→ O(1k) sectors → O(1M) individual MC integrals
  - Still computationally very challenging! → O(1M CPUh)
- Many-leg, IR stable one-loop amplitudes → OpenLoops [Buccioni'19]
- Double virtual amplitudes in leading-colour approximation [Abreu'21]
  - Sub-leading colour corrections expected to be small
  - Analytical expressions challenging
  - Fast numerical evaluation → very small contribution to computational cost



**Only** Approximation made:

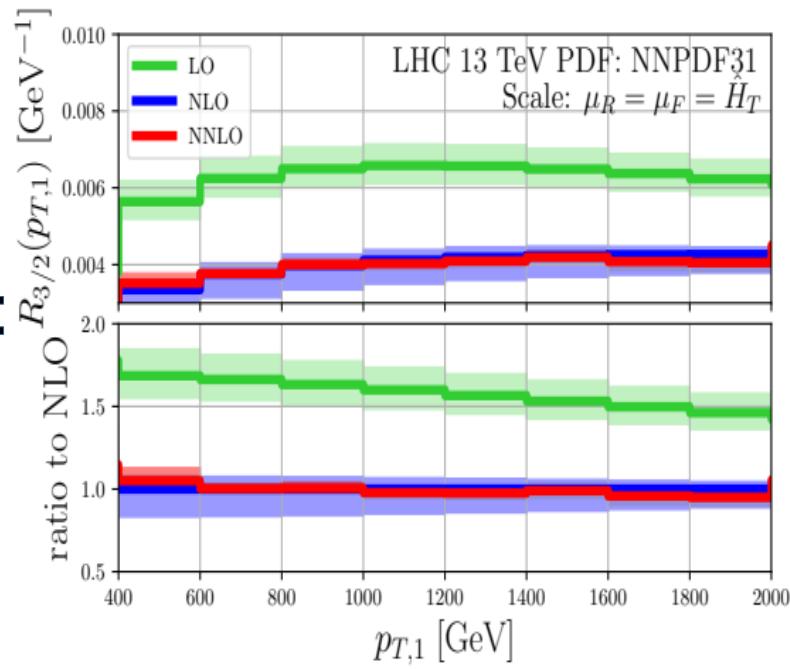
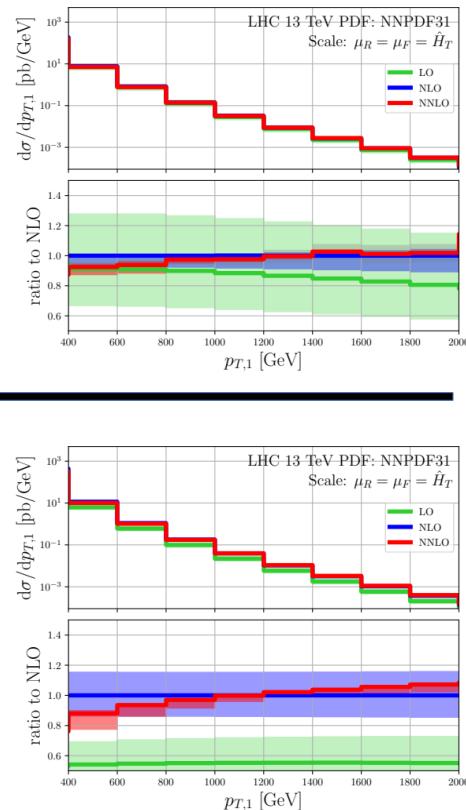
$$\mathcal{R}^{(2)}(\mu_R^2) = 2 \operatorname{Re} \left[ \mathcal{M}^{\dagger(0)} \mathcal{F}^{(2)} \right] (\mu_R^2) + |\mathcal{F}^{(1)}|^2(\mu_R^2) \equiv \mathcal{R}^{(2)}(s_{12}) + \sum_{i=1}^4 c_i \ln^i \left( \frac{\mu_R^2}{s_{12}} \right)$$

$$\mathcal{R}^{(2)}(s_{12}) \approx \mathcal{R}^{(2)l.c.}(s_{12})$$

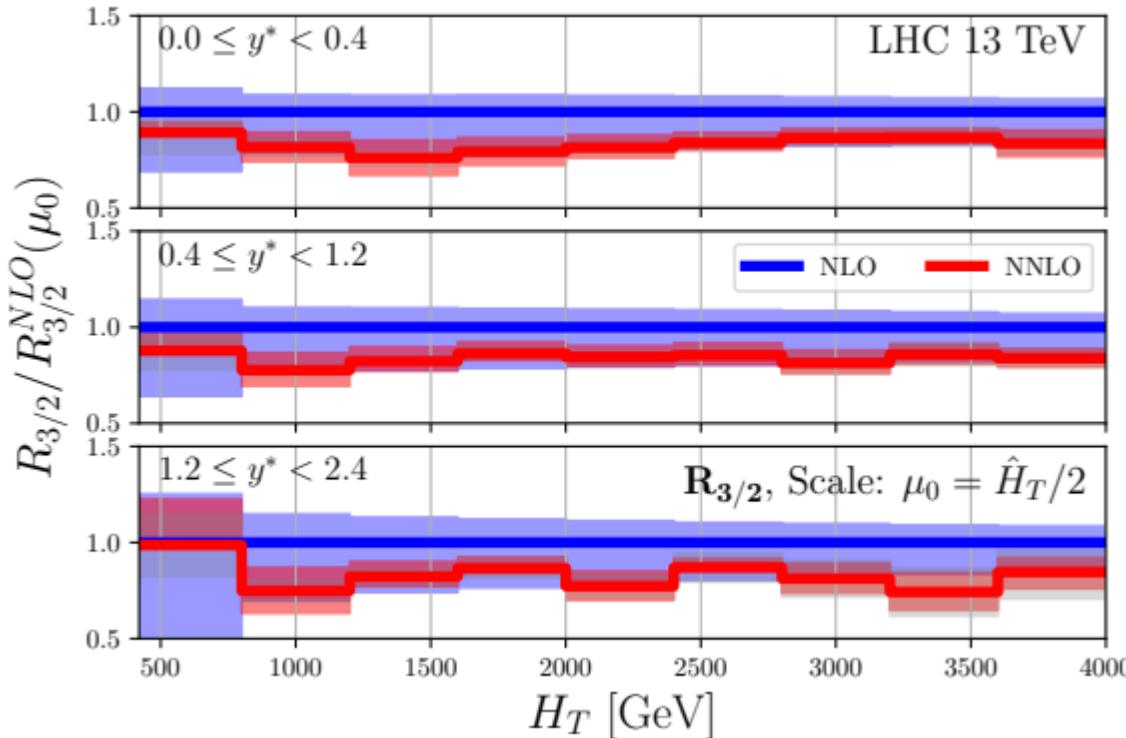
# Three-jet production - R32( $p_T$ 1)

- LHC @ 13 TeV, NNPDF31
- Require at least three (two) jets:
  - $p_T(j) > 60$  GeV and  $|y(j)| < 4.4$
  - $H_{T,2} = p_T(j_1) + p_T(j_2) > 250$  GeV
- Scales:

$$\mu_R = \mu_F = \hat{H}_T = \sum_{\text{partons}} p_T$$



# Three-jet production – R<sub>3/2</sub>(H<sub>T</sub>,y\*)



Double differential w.r.t.  $H_T = \sum_{\text{jets}} p_T$  and  $y^* = |y(j_1) - y(j_2)|/2$

Central scale choice:  $\hat{H}_T/2$

# Three-jet production – azimuthal decorrelation

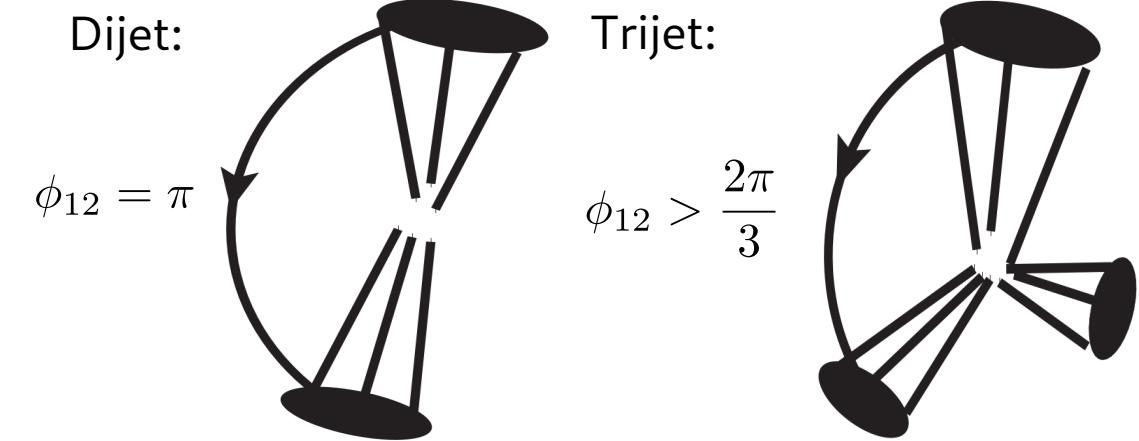
Kinematic constraints on the azimuthal separation between the two leading jets ( $\phi_{12}$ )

$\phi_{12}$  sensitive to the jet multiplicity:

$$2j: \phi_{12} = \pi$$

$$3j: \phi_{12} > \frac{2\pi}{3}$$

4j: unconstrained

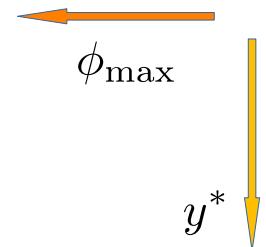


Study of the ratio:

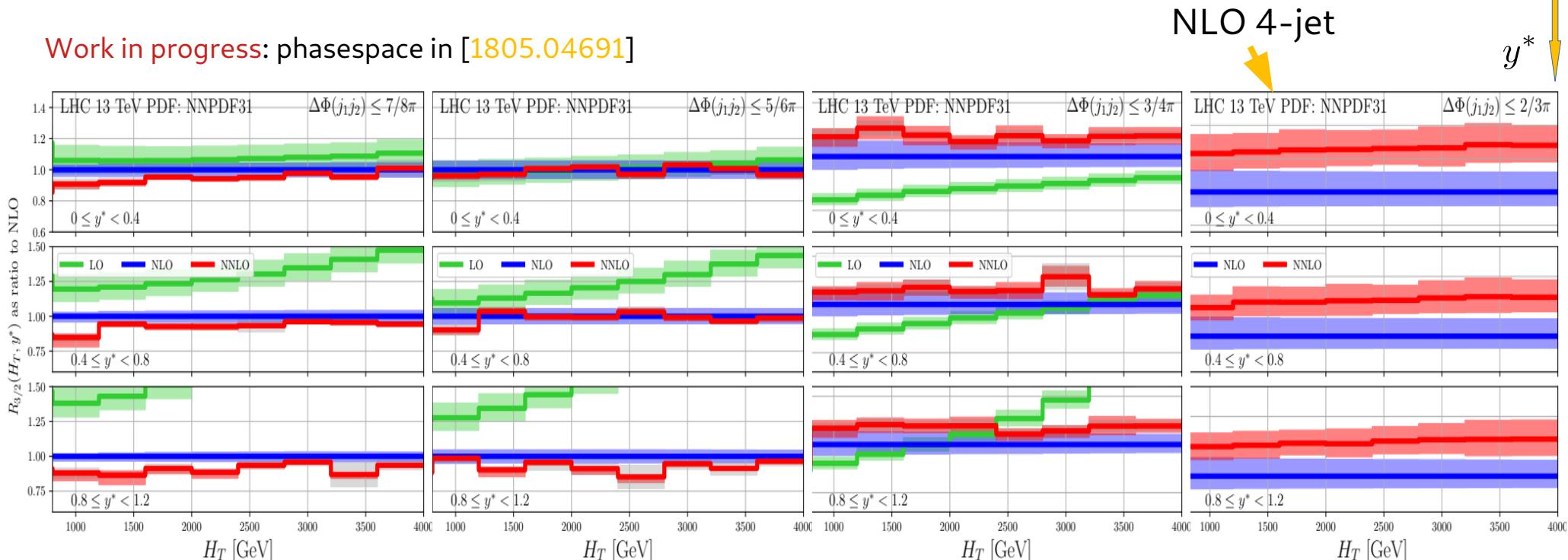
$$R_{32}(H_T, y^*, \phi_{\max}) = \frac{d\sigma_3(H_T, y^*, \phi_{12} < \phi_{\max})}{d\sigma_2(H_T, y^*)}$$

# Three-jet production - azimuthal decorrelation

NNLO/NLO K-factor smaller than NLO/LO  
Scale dependence is reduced



Work in progress: phasespace in [1805.04691]



# Outlook: Extraction of the strong coupling constant from multi-jet events at the LHC

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- Transverse Energy-Energy Correlator TEEC
- Event shapes

# Transverse Energy-Energy Correlator @ LHC

## TEEC: Transverse Energy-Energy Correlation

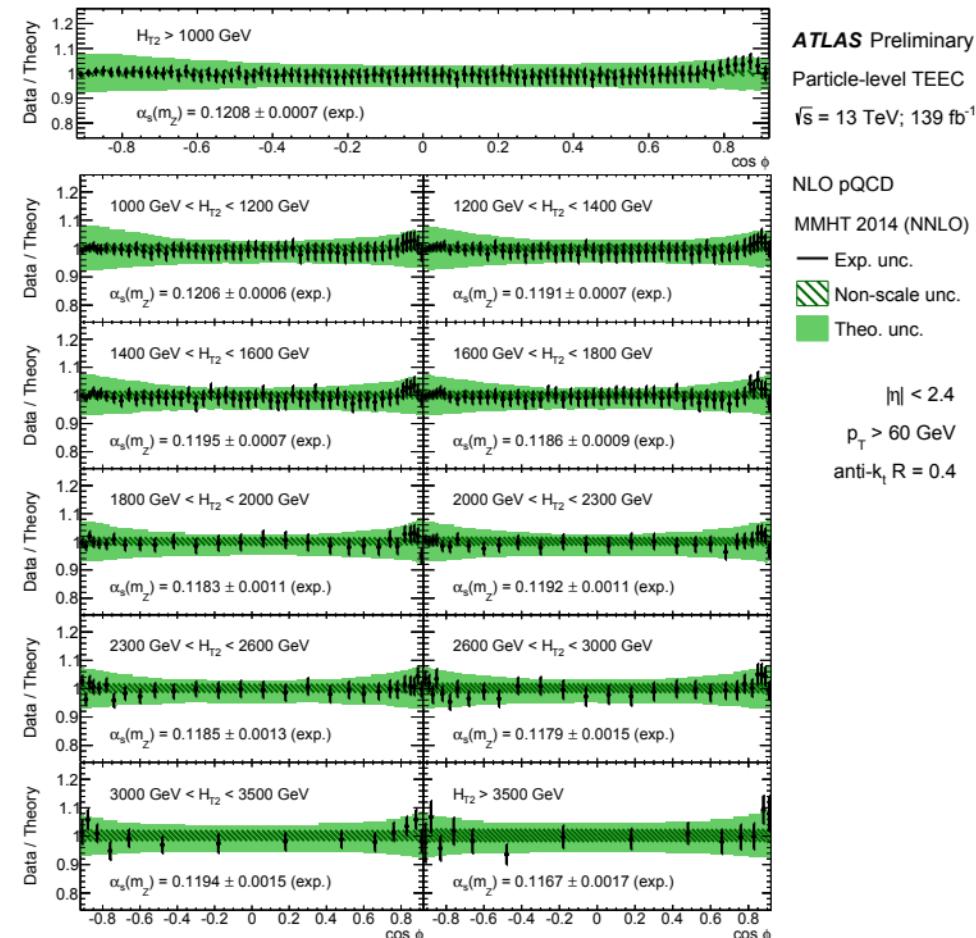
$$\frac{1}{\sigma} \frac{d\Sigma}{d \cos \phi} = \frac{1}{N} \sum_{A=1}^N \sum_{ij} \frac{E_{\perp,i}^A E_{\perp,j}^A}{\left( \sum_k E_{T,k}^A \right)^2} \delta(\cos \phi - \cos \phi_{ij})$$

ATLAS measurement of the TEEC and ATEEC:

- @ 8 TeV [[ATLAS:1707.02562](#)]
- @ 13 TeV [[ATLAS-CONF-2020-025](#)]

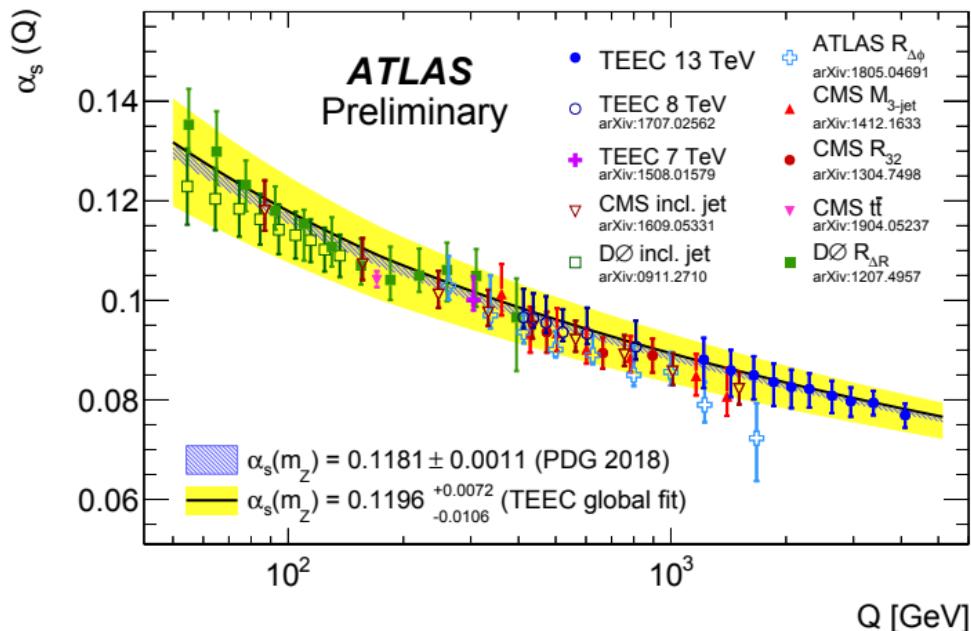
TEEC in HT2 bins:

- from 1000 GeV to 3500 GeV and above
- sensitivity to different energy scales



# Transverse Energy-Energy Correlator @ LHC

Extraction of alphas in different HT bins → test of SM running



$\langle Q \rangle$ [GeV]	$\alpha_s(m_Z)$ value (MMHT 2014)
Global	$0.1195 \pm 0.0002$ (stat.) $\pm 0.0006$ (syst.) $+0.0084_{-0.0106}$ (scale) $\pm 0.0009$ (PDF) $\pm 0.0003$ (NP)
Inclusive	$0.1198 \pm 0.0002$ (stat.) $\pm 0.0006$ (syst.) $+0.0078_{-0.0095}$ (scale) $\pm 0.0010$ (PDF) $\pm 0.0002$ (NP)
1219	$0.1202 \pm 0.0003$ (stat.) $\pm 0.0006$ (syst.) $+0.0079_{-0.0098}$ (scale) $\pm 0.0010$ (PDF) $\pm 0.0002$ (NP)
1434	$0.1184 \pm 0.0003$ (stat.) $\pm 0.0007$ (syst.) $+0.0078_{-0.0098}$ (scale) $\pm 0.0011$ (PDF) $\pm 0.0002$ (NP)
1647	$0.1188 \pm 0.0004$ (stat.) $\pm 0.0007$ (syst.) $+0.0073_{-0.0087}$ (scale) $\pm 0.0012$ (PDF) $\pm 0.0001$ (NP)
1856	$0.1177 \pm 0.0006$ (stat.) $\pm 0.0008$ (syst.) $+0.0072_{-0.0083}$ (scale) $\pm 0.0013$ (PDF) $\pm 0.0006$ (NP)
2064	$0.1174 \pm 0.0008$ (stat.) $\pm 0.0009$ (syst.) $+0.0069_{-0.0078}$ (scale) $\pm 0.0013$ (PDF) $\pm 0.0007$ (NP)
2300	$0.1185 \pm 0.0009$ (stat.) $\pm 0.0010$ (syst.) $+0.0063_{-0.0067}$ (scale) $\pm 0.0014$ (PDF) $\pm 0.0005$ (NP)
2636	$0.1166 \pm 0.0016$ (stat.) $\pm 0.0012$ (syst.) $+0.0062_{-0.0066}$ (scale) $\pm 0.0015$ (PDF) $\pm 0.0000$ (NP)
2952	$0.1141 \pm 0.0029$ (stat.) $\pm 0.0013$ (syst.) $+0.0062_{-0.0069}$ (scale) $\pm 0.0018$ (PDF) $\pm 0.0003$ (NP)
3383	$0.1164 \pm 0.0043$ (stat.) $\pm 0.0015$ (syst.) $+0.0050_{-0.0044}$ (scale) $\pm 0.0017$ (PDF) $\pm 0.0001$ (NP)
4095	$0.1029 \pm 0.0163$ (stat.) $\pm 0.0014$ (syst.) $+0.0066_{-0.0012}$ (scale) $\pm 0.0010$ (PDF) $\pm 0.0003$ (NP)



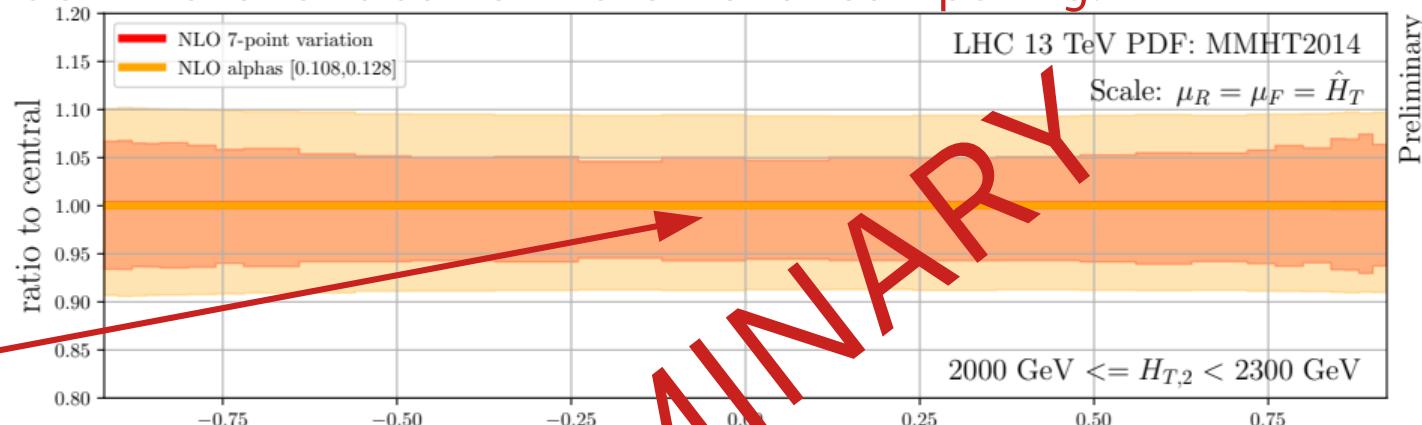
FO scale uncertainty limiting factor!

# NNLO QCD corrections to TEEC @ LHC

Massive thanks to Manuel Alvarez and Javier Llorente for computing!

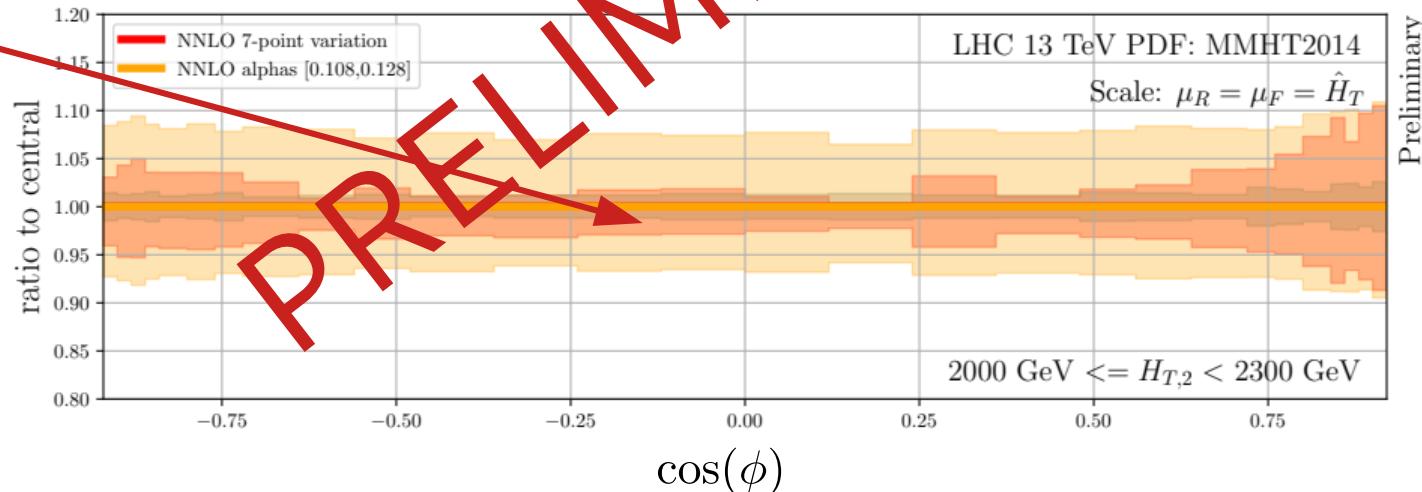
NLO

Reduction in  
scale dependence  
by factor 2-3



NNLO

PRELIMINARY



# Event shapes at the LHC

ATLAS measurement of event shapes @ 13 TeV using multi-jet events (139fb-1) in HT2 bins and high pT jets (> 100 GeV): [ATLAS:2007.12600]

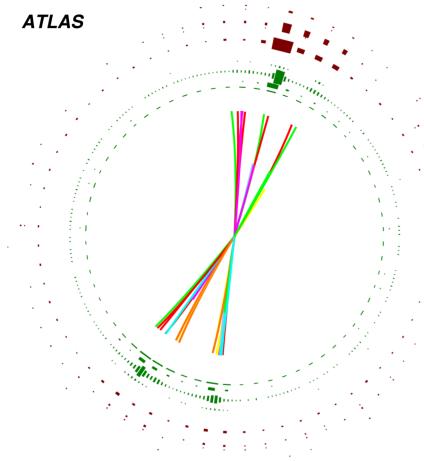
Transverse Thrust:  $\tau_T = 1 - \frac{\sum_i^{\text{jets}} |\vec{p}_{T,i} \cdot \hat{n}|}{\sum_i^{\text{jets}} |\vec{p}_{T,i}|}$

Thrust Minor:  $T_m = \frac{\sum_i^{\text{jets}} |\vec{p}_{T,i} \times \hat{n}|}{\sum_i^{\text{jets}} |\vec{p}_{T,i}|}$

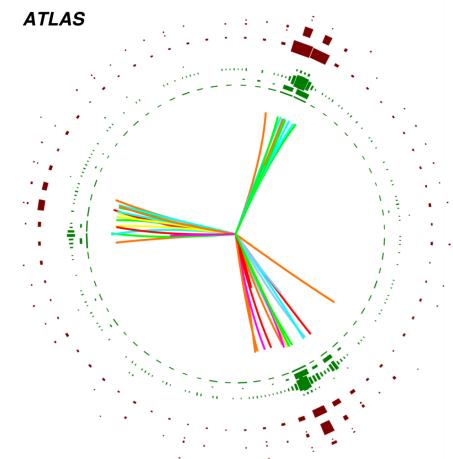
More quantities based on eigenvalues of (transverse) linearised sphericity tensor:

$$\mathcal{M}_{xyz} = \frac{1}{\sum_i^{\text{jets}} |\vec{p}_i|} \sum_i^{\text{jets}} \frac{1}{|\vec{p}_i|} \begin{pmatrix} p_{x,i}^2 & p_{x,i}p_{y,i} & p_{x,i}p_{z,i} \\ p_{y,i}p_{x,i} & p_{y,i}^2 & p_{y,i}p_{z,i} \\ p_{z,i}p_{x,i} & p_{z,i}p_{y,i} & p_{z,i}^2 \end{pmatrix}$$

Back-to-Back

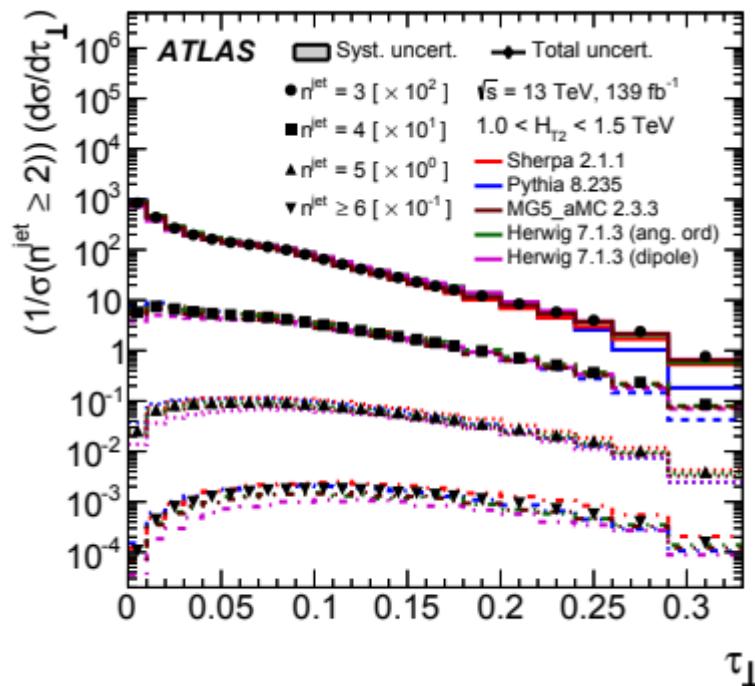


Spherical

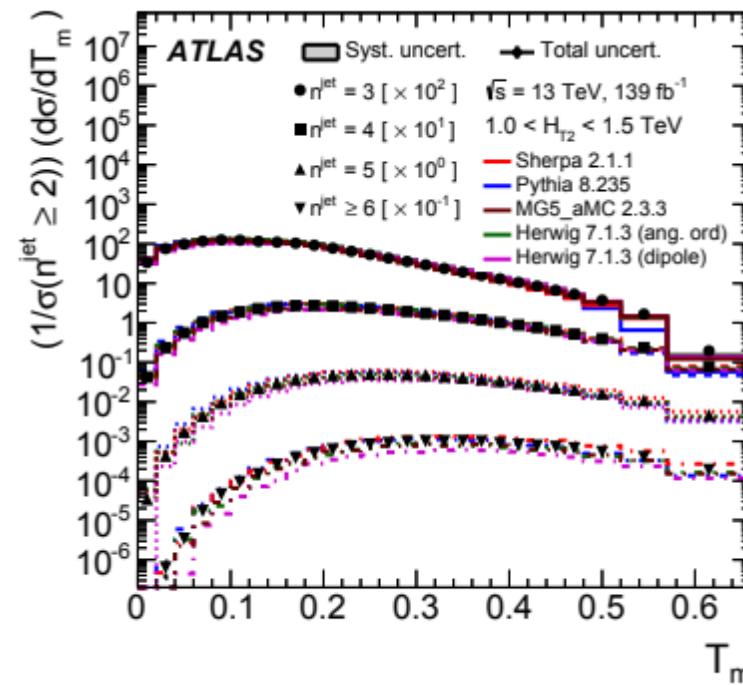


# Event shapes at the LHC

Transverse thrust:



Transverse thrust minor:

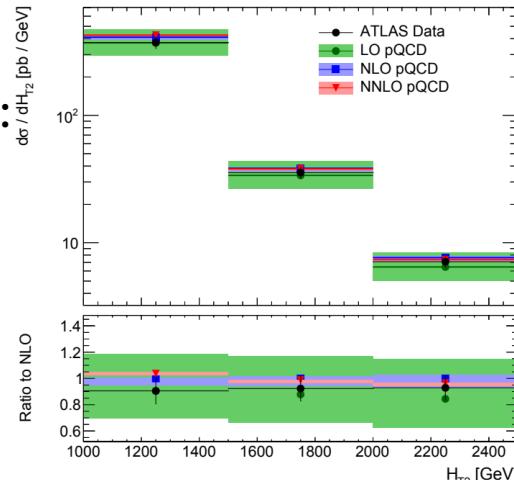


[ATLAS:2007.12600]

# NNLO QCD corrections to event shapes

Comparison of public data from HEPdata

HT 2 denominator:

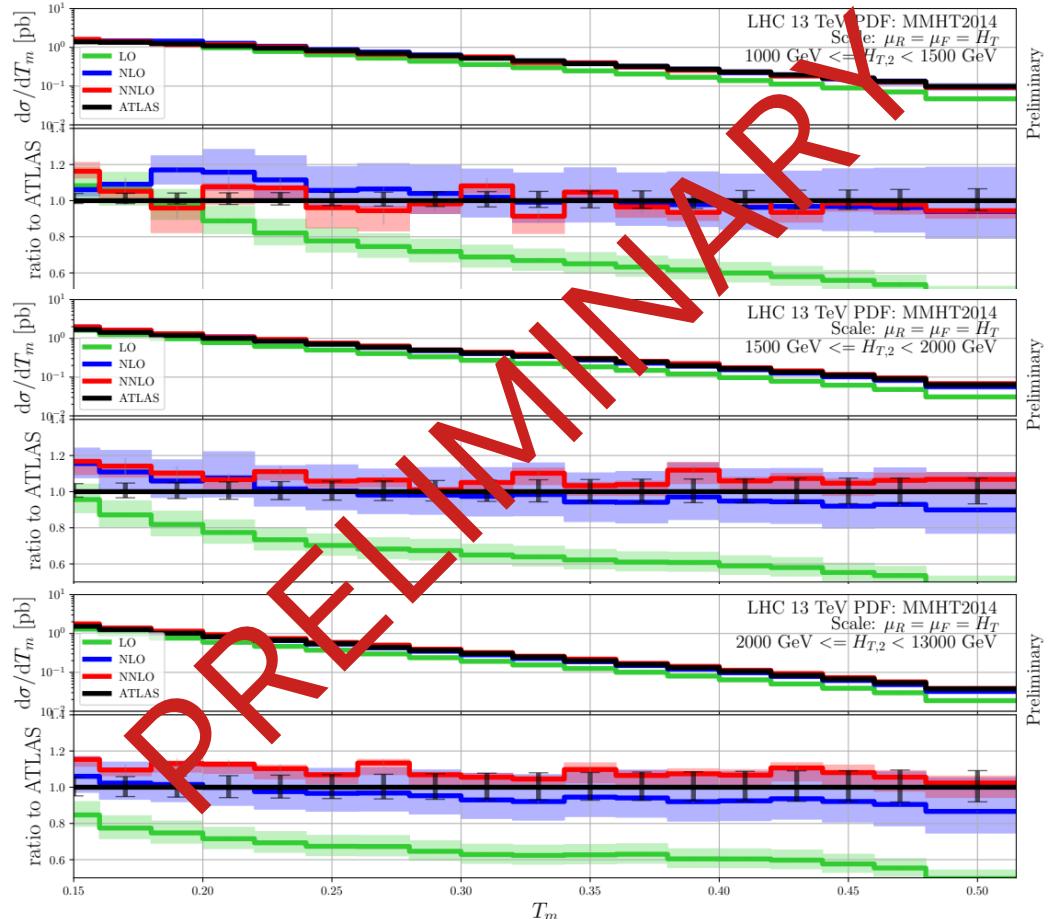


Credits:

Javier Llorente!

Example Thrust-Minor:

- Beautiful perturbative convergence
- Significant reduction of perturbative corrections



# Flavoured Jets

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# Flavoured jets

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- Jets are a tool to connect QCD of quarks&gluons to actually strongly interacting particles, i.e. hadrons.
- They are defined by a suitable algorithm: experimentally and theoretically
- Jet-substructure reveals additional information:
  - Separation of quark and gluon initiated jets
  - Jets of definite flavour:

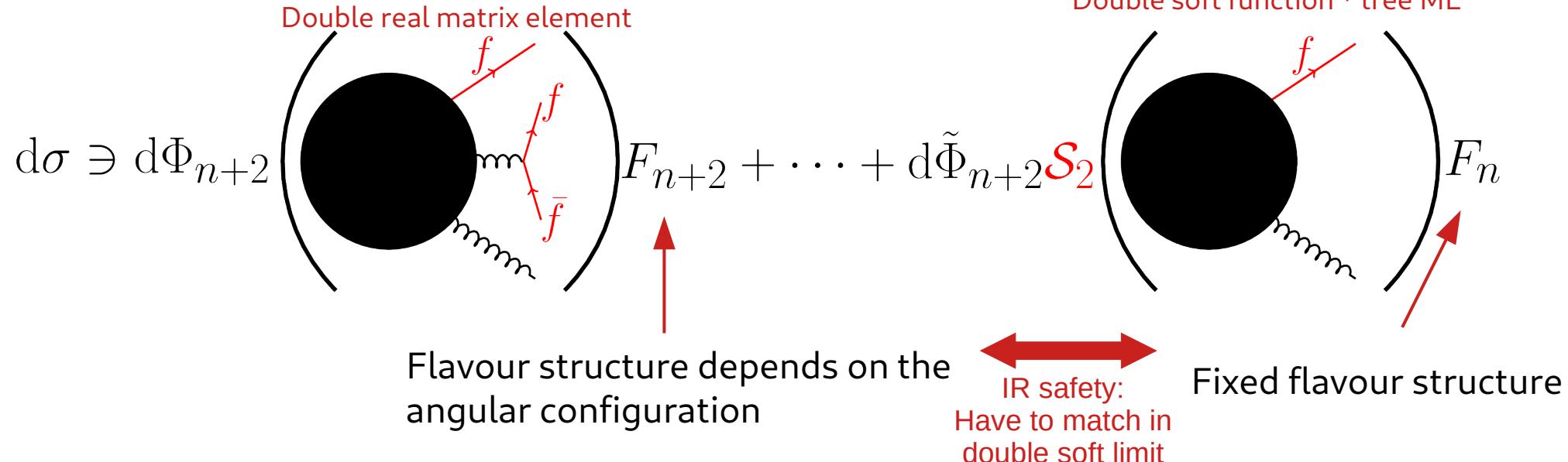
Experimentally	<i>Displayed vertices</i> of heavy intermediate particles: D/B mesons
MC Event Simulation	Similar objects due to hadronization and detector simulations
Partonic computations	<ul style="list-style-type: none"><li>• Impose relation between quarks and hadrons (quark model)</li><li>• Massless quarks: emission of soft flavoured pairs → gluons → <b>Implications for IR safety in FO computations beyond NLO</b></li></ul>

- Why are partonic computations for flavoured jets interesting?
  - Higher order perturbation theory (not necessarily available matched to PS)
  - Extraction of SM parameters or PDFs

# Fixed order flavoured jets beyond NLO

What is the problem with FO flavoured jets?

Example NNLO: double real radiation and subtraction



- If  $F(n+2)$  does not treat the flavour pair appropriately:
  - double soft singularity not subtracted
  - **Implies correlated treatment of kinematics and flavour information**

# Solution: Modified jet algorithms

- Implies correlated treatment of kinematics and flavour information

Standard kT algorithm [Ellis'93]:

Pair distance:

$$d_{ij} = \min(k_{T,i}^2, k_{T,j}^2) R_{ij}^2$$

$$R_{ij}^2 = (\Delta\phi_{ij}^2 + \Delta\eta_{ij}^2)/R^2$$

"Beam" distance for determination condition:

$$d_i = k_{T,i}^2$$

Flavour kT algorithm [Banfi'06]:

Pair distance:

$$d_{ij} = R_{ij}^2 \begin{cases} \max(k_{T,i}, k_{T,j})^\alpha \min(k_{T,i}, k_{T,j})^{2-\alpha} & \text{softer of } i,j \text{ is flavoured} \\ \min(k_{T,i}, k_{T,j})^\alpha & \text{else} \end{cases}$$

Beam distance:

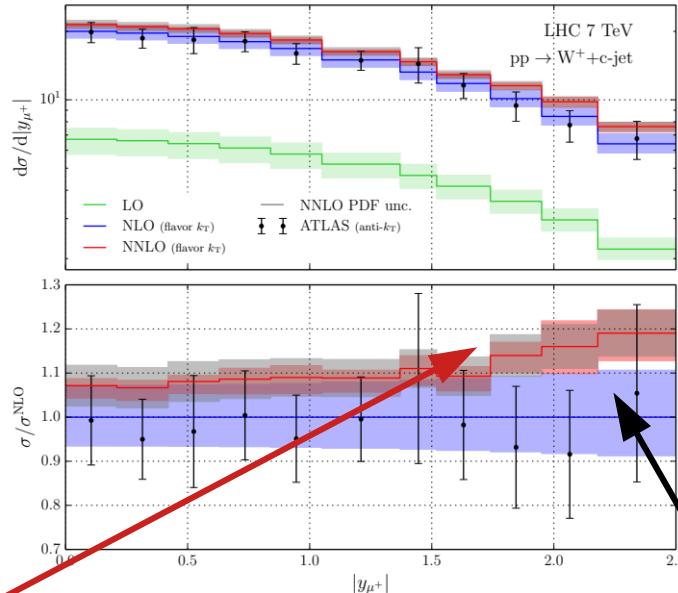
$$d_{i,B} = \begin{cases} \max(k_{T,i}, k_{T,B}(y_i))^\alpha \min(k_{T,i}, k_{T,B}(y_i))^{2-\alpha} & i \text{ is flavoured} \\ \min(k_{T,i}, k_{T,B}(y_i))^\alpha & \text{else} \end{cases}$$

$$d_B(\eta) = \sum_i k_{T,i} (\theta(\eta_i - \eta) + \theta(\eta - \eta_i)) e^{\eta_i - \eta}$$

$$d_{\bar{B}}(\eta) = \sum_i k_{T,i} (\theta(\eta - \eta_i) + \theta(\eta_i - \eta)) e^{\eta - \eta_i}$$

# Problem solved, isn't it?

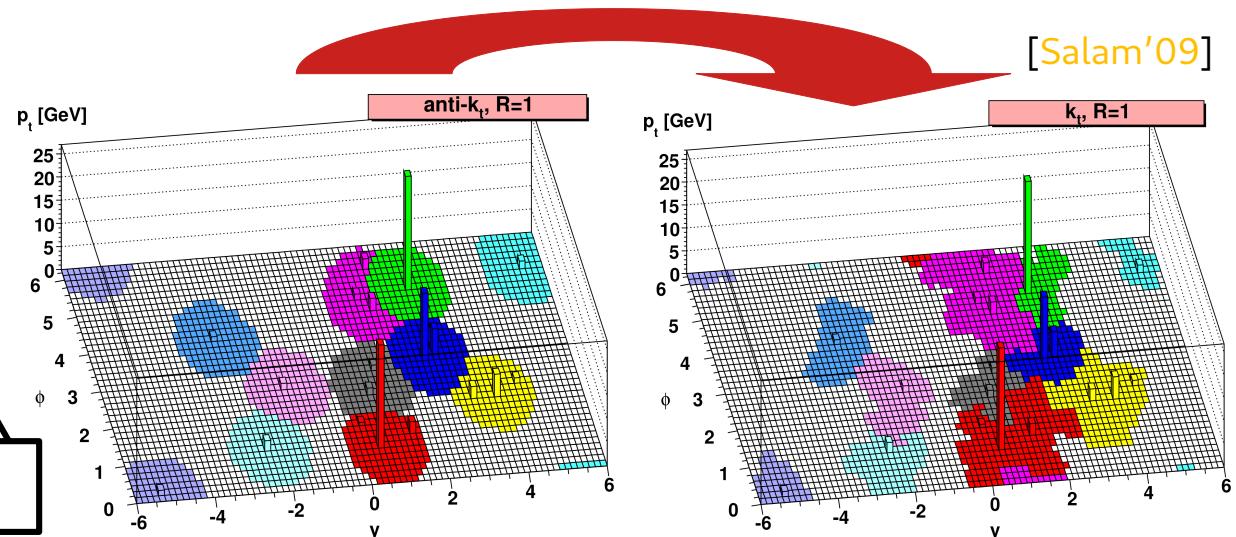
Real world example: W+c-jet at NNLO QCD with flavour-kT [Czakon'20]



NNLO QCD with flavour kT

ATLAS data with standard anti- $k_T$

A proper comparison would require to  
unfold experimental data  
→ (flavour-)  $k_T$  and anti- $k_T$  cluster partonic jets  
differently → Non-trivial procedure.



# What about flavour anti-kT?

Anti-kT:  $d_{ij} = \min(k_{T,i}^{-2}, k_{T,j}^{-2})R_{ij}^2 \quad d_i = k_{T,i}^{-2}$

The energy ordering in anti-kT prevents correct recombination of flavoured pairs in the double soft limit.

Proposed modification:

A soft term designed to modify the distance of flavoured pairs.

$$d_{ij}^{(F)} = d_{ij} \begin{cases} \mathcal{S}_{ij} & i,j \text{ is flavoured pair} \\ 1 & \text{else} \end{cases}$$

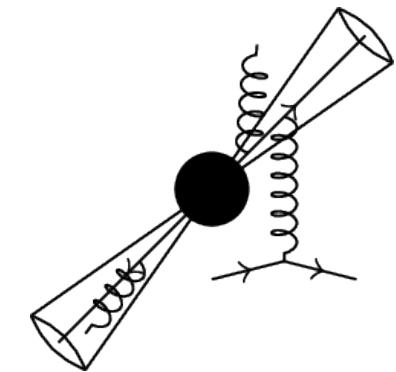
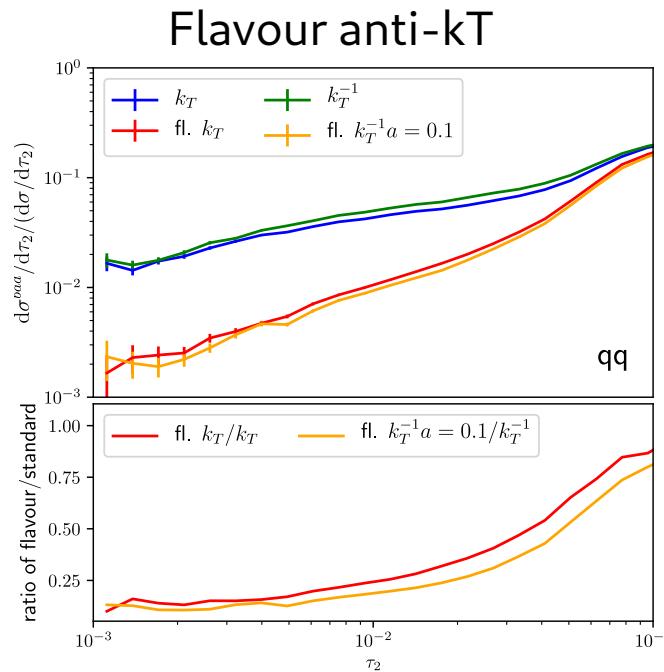
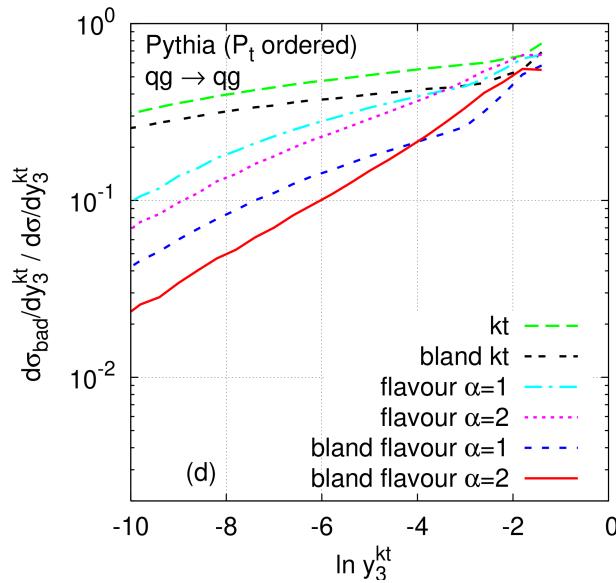
$$\mathcal{S}_{ij} = 1 - \theta(1 - x) \cos\left(\frac{\pi}{2}x\right) \quad \text{with} \quad x = \frac{k_{T,i}^2 + k_{T,j}^2}{2ak_{T,\max}^2}$$

# Tests of IR safety with parton showers

Dress tree-level di-jet events (definite flavour structure: “qq”, “qg” or “gg”) with radiation and study jet flavour (q or g) as function of kinematics.

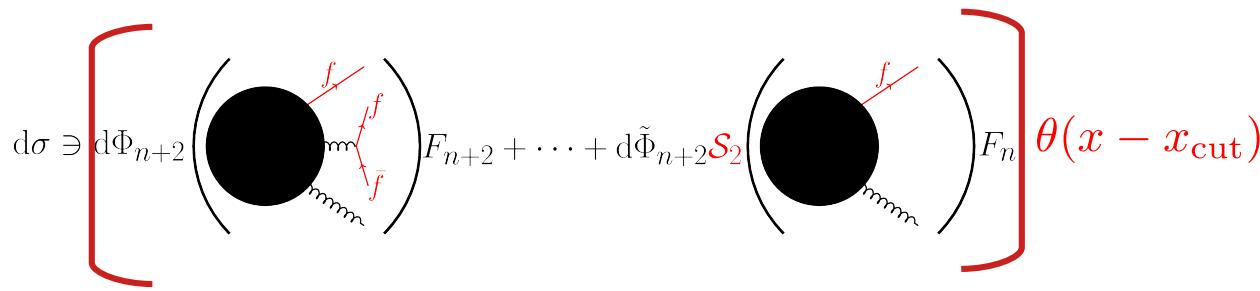
In the di-jet limit the flavour needs to correspond to tree level flavours  
→ misidentification rate needs to vanish in dijet back-to-back limit

Flavour  $k_T$  vs.  $k_T$  [Banfi'06]:



# Tests of IR safety with NNLO FO computations

IR sensitivity of jet cross sections on (technical) IR regulating parameter  $x$



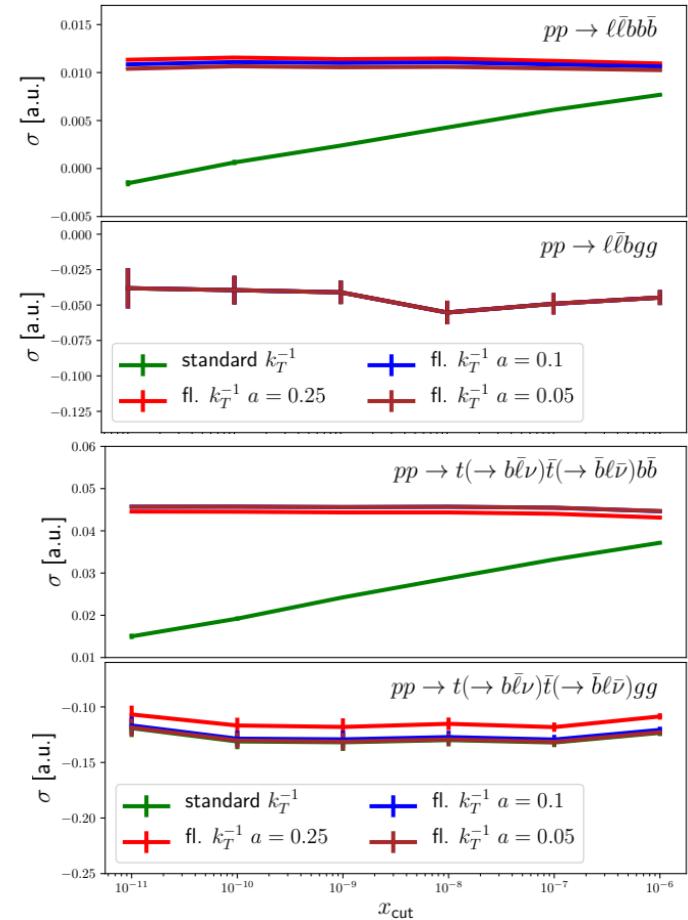
In the limit  $x_{cut} \rightarrow 0$ :

IR safe jet flavour

→ no dependence on  $x_{cut}$

IR non-safe jet flavour

→ logarithmic divergent



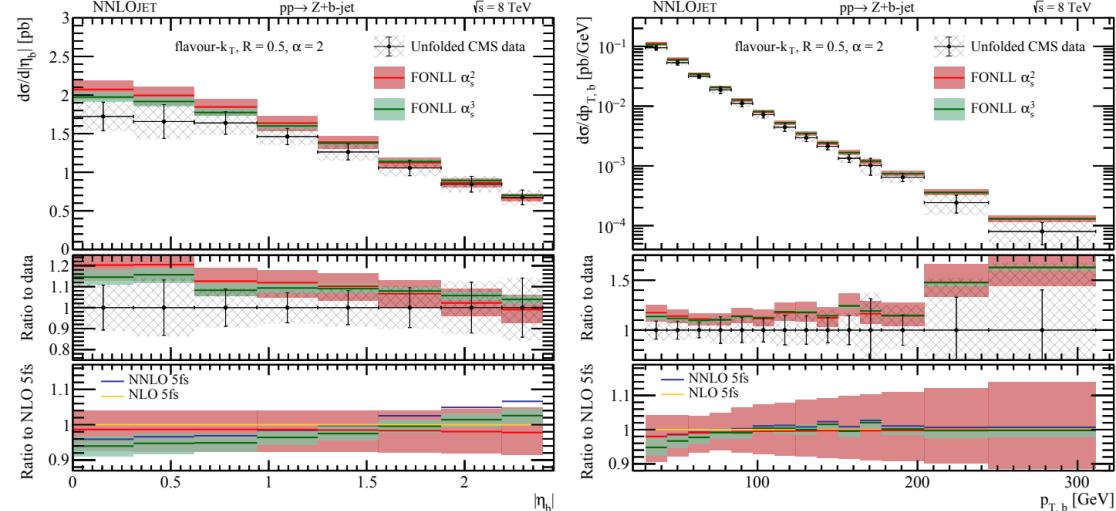
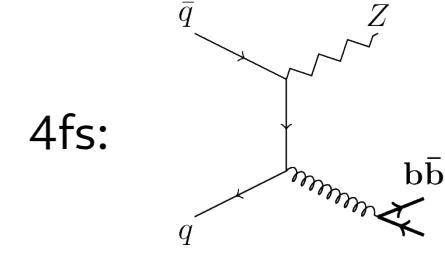
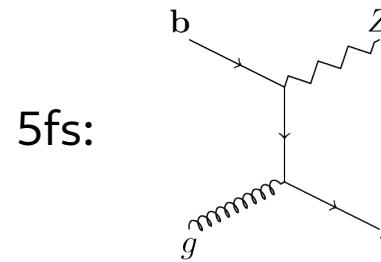
# Phenomenology: Z+b-jet

Benchmark process:

Well studied up to  $\mathcal{O}(\alpha_s^3)$  [Gauld'20]:

- Defined with flavour-kT algorithm
  - Unfolding of experimental data (RooUnfold, bin-by-bin unfolding)
  - Matching between four- and five-flavour schemes (FONLL) [Gauld'21]
- $$d\sigma^{\text{FONLL}} = d\sigma^{5\text{fs}} + (d\sigma_{m_b}^{4\text{fs}} - d\sigma_{m_b \rightarrow 0}^{4\text{fs}})$$
- CMS measurement @ 8 TeV [CMS 1611.06507]

$pp \rightarrow Z(l\bar{l}) + b\text{-jet}$



# Phenomenology: Tunable parameter

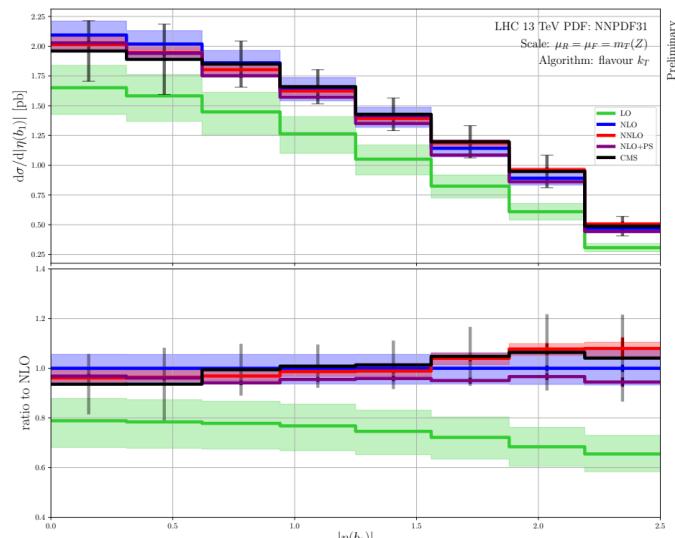
Benchmark process:  $pp \rightarrow Z(l\bar{l}) + b\text{-jet}$

Preliminary

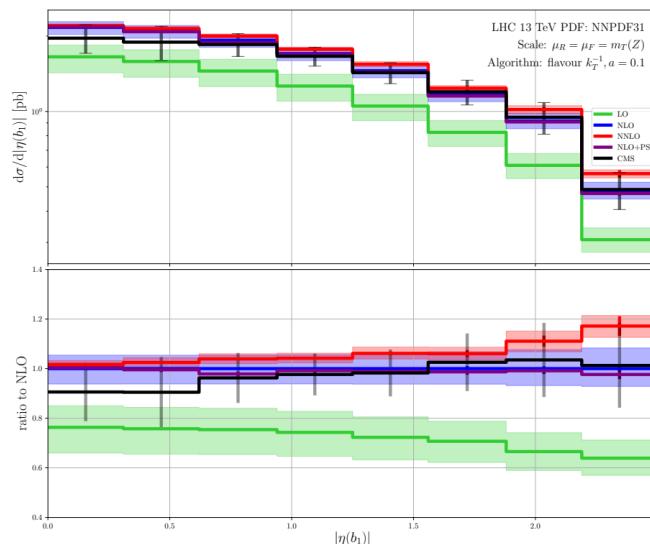
Tunable parameter a:

- Limit  $a \rightarrow 0 \Leftrightarrow$  original anti-kT (IR unsafe)
- Large  $a \Leftrightarrow$  large modification of cluster sequence

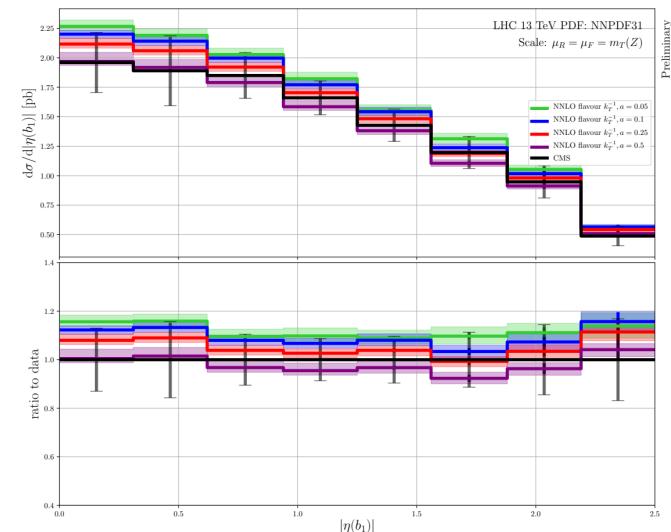
Flavour kT:



Flavour anti-kT:  $a = 0.1$



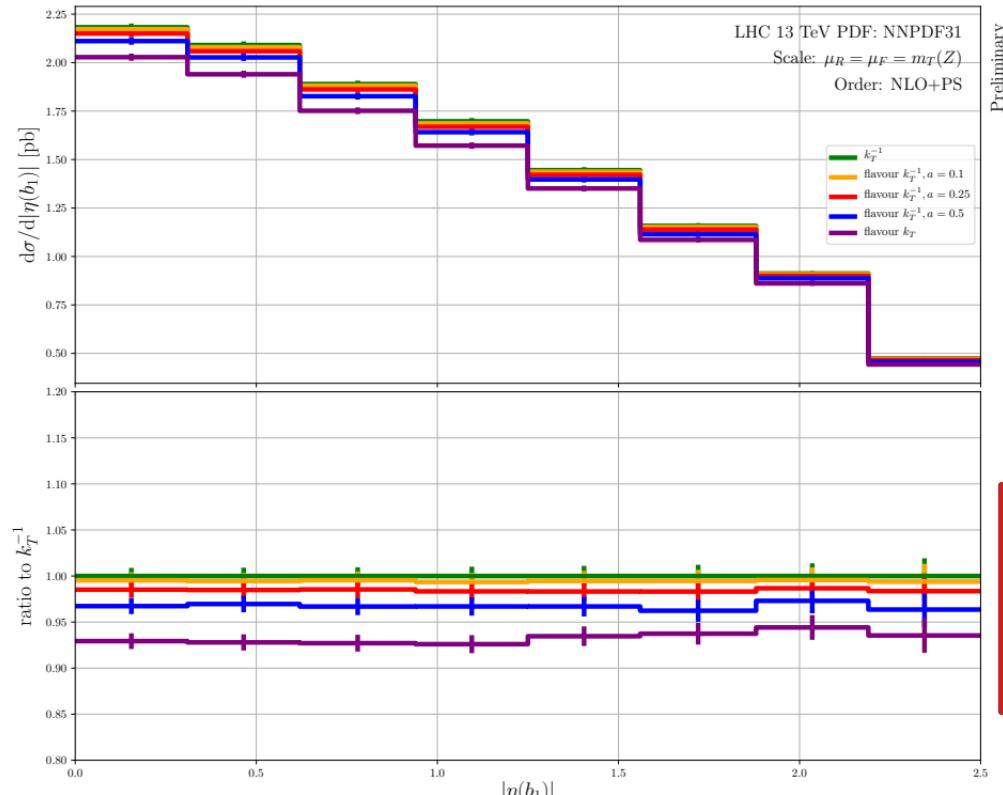
Comparison of different parameter a to data:



# Phenomenology: Tunable parameter II

Preliminary

What happens in the presence of many flavoured partons? → NLO PS



Tunable parameter a:

- Small a: Flavour anti- $k_T$  results are more similar to standard anti- $k_T$  → **small unfolding factors**
- Larger a: Larger modification of clustering

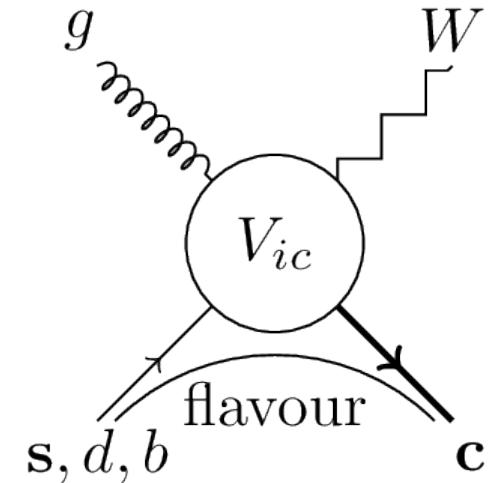
Good FO perturbative convergence +  
Small difference to standard anti- $k_T$   
→ a~0.1 is a good candidate

# W+c-jet

Idea: Identify final state c-quarks to access s-quark PDFs.

- Reduction of PDF uncertainties
- Shed light on  $s\bar{s}$  asymmetry

- Non-diagonal CKM and  $g \rightarrow \bar{c}c$  reduce s-PDF sensitivity
- Large NLO corrections → higher order corrections?
- Theoretical treatment:
  - Massive c (3-flavour scheme):
    - Resummation of mass logs at high pT → PS
    - Higher order predictions?
  - Massless c:
    - c-quark part of the PDFs
    - NNLO QCD available
    - **Jet definition?**



$$V_{sc} > V_{dc} \gg V_{bc}$$

# W+c-jet with flavour kT at NNLO QCD

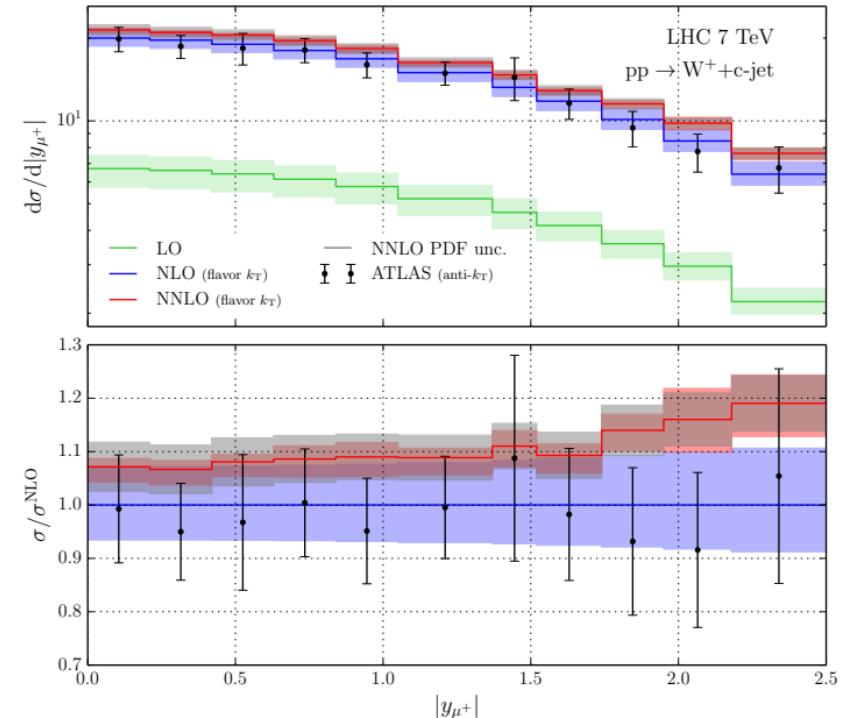
In collaboration with: Czakon, Mitov, Pellen

NNLO QCD 7 TeV results [2011.01011]:

- Full NNLO corrections for V<sub>cs</sub> contribution
- Off-diagonal CKM only LO QCD
- Comparison flv. kT results vs. ATLAS [1402.6263]

Update for 13 TeV measurement:

- Full CKM through NNLO QCD
- Study of different jet-algorithms:
  - Impact of beam-function d\_ib in flv kT
  - New anti-kT algorithm
- Study of different flavour tag definitions/setups:
  - Modulus vs. absolute flv tag definition
  - OS minus SS
  - “Inclusive c-jet” rates



# W+c-jet with flavour anti-kT

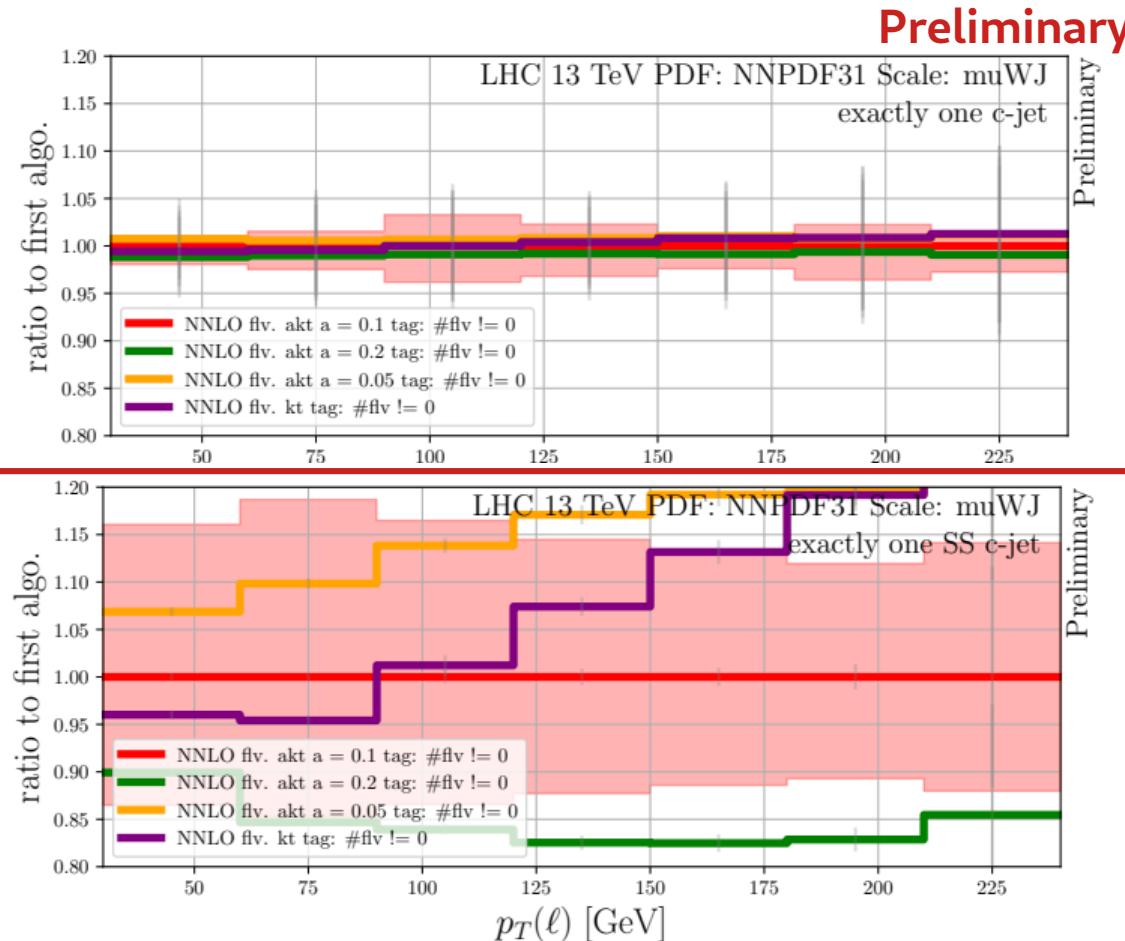
In collaboration with: Czakon, Mitov, Pellen

Exactly one c-jet requirement:

- Comparison of parameters  $a$ :  
→ small dependence < 2%
- Comparison to flv kT:  
→ small dependence @ NNLO < 2%

ONLY large effect in SS contribution

- Exactly one c-jet of SS type:  
Larger dependence ~15%  
(roughly size of NNLO scale band)
- BUT: SS contribution ~2-5%
- => OS ~0.2-0.5% dependence

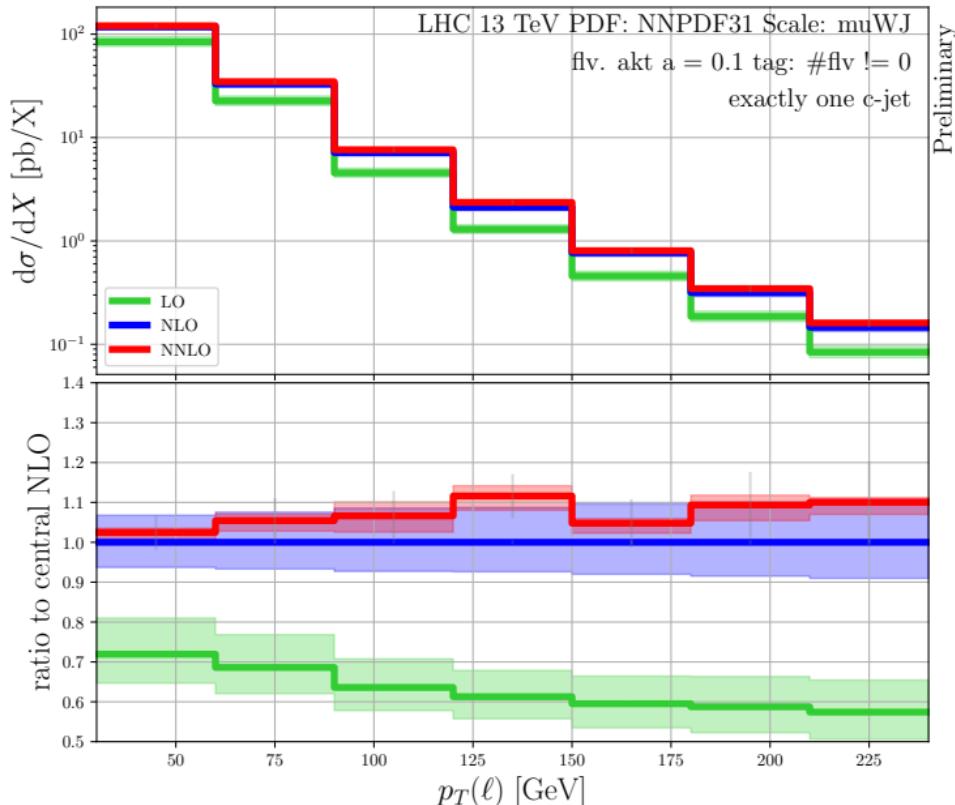


# Flavour tags: OS - SS

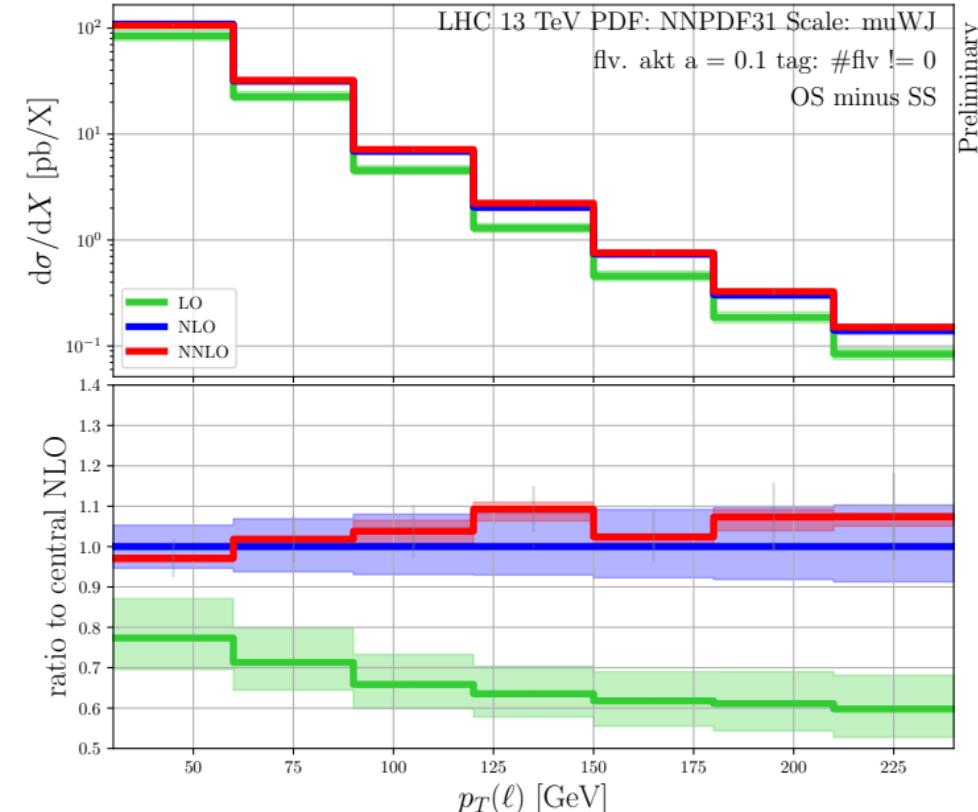
In collaboration with: Czakon, Mitov, Pellen

Preliminary

Exactly 1 c-jet:



OS-SS:



# Some final remarks

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- What is that  $kT_{\text{max}}$  parameter?

Some scale to define what **soft** means.

Examples:

1.  $pT$  of hardest pseudo jet or lepton

at a clustering step

2. Some fixed dynamical scale, e.g.  $pT(Z)$ ,  $pT(\text{lep})$ , ...

3. Some fixed hard scale:  $m_{\text{top}}$ ,  $m_Z$  etc.

→ The choice impacts the clustering.

- Besides c/b jets: What about gluon/quark jet identification?

Conceptually not a problem. Not yet studied in detail.

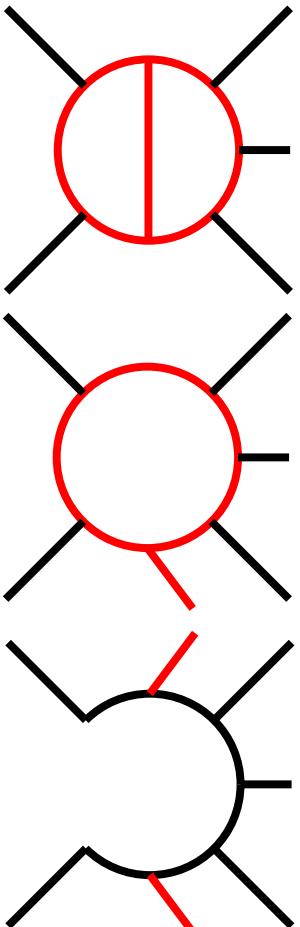
But might introduce some more sensitivity to actual form of  $S_{ij}$  ??

- More complicated examples:  $\text{pp} \rightarrow W \bar{b} \bar{b}$  ! LO sensitivity to flv jet algorithm

$$d_{ij}^{(F)} = d_{ij} \begin{cases} S_{ij} & i,j \text{ is flavoured pair} \\ 1 & \text{else} \end{cases}$$

$$S_{ij} = 1 - \theta(1-x) \cos\left(\frac{\pi}{2}x\right) \quad \text{with} \quad x = \frac{k_{T,i}^2 + k_{T,j}^2}{2ak_{T,\text{max}}^2}$$

# Summary and Outlook



Precision jet observables allow for many pheno applications!

- First NNLO QCD phenomenology results for three jet production  
R32 ratios, azimuthal decorrelation, event-shapes
- Future application to alphaS extraction

Flavoured jet observables

- New proposed flavour safe version of anti- $k_T$
- Phenomenological applications to  $Z+b$ -jet,  $W+c$ -jet, top-quark pairs
- Many more applications ahead: open- $b$ 's,...

# Backup

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# b-jets in top-pair production&decay

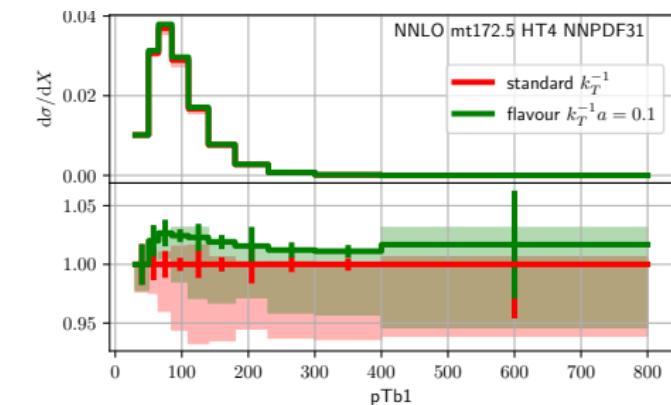
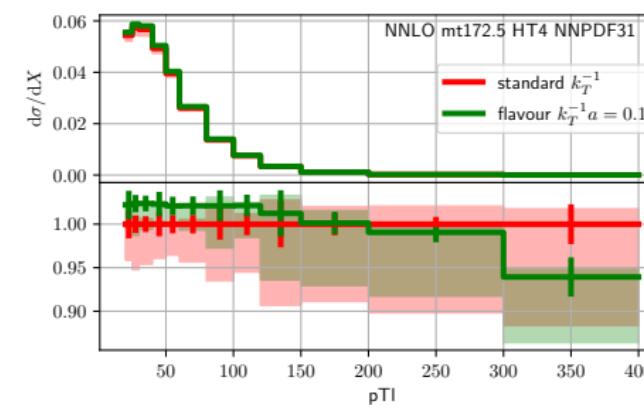
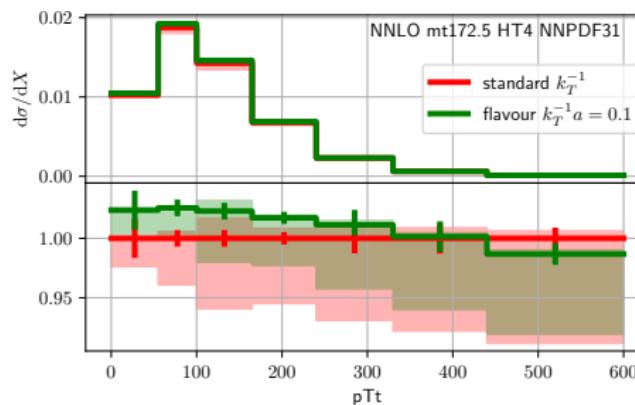
NNLO QCD corrections [Czakon'20] to:

$$pp \rightarrow t(\rightarrow b\bar{l}\nu)\bar{t}(\rightarrow \bar{b}\ell\bar{\nu}) + X$$

Flavour sensitive channels like:

$$pp \rightarrow t\bar{t}b\bar{b} \rightarrow \bar{\ell}\nu\ell\bar{\nu} \boxed{b\bar{b}b\bar{b}}$$

Small numerical impact from extra bbar emissions  
in  $pp \rightarrow b\bar{b}$  [Catani'20] and single-top production [Berger '17'18, Campbell '20]  
→ naive treatment via cut-off procedure



Naive 'cut-off' treatment vs. proposed IR safe flavour anti- $k_T$