

Precision Measurements - Theory

Andreas von Manteuffel
Rene Poncelet
Raoul Röntsch



THE HENRYK NIEWODNICZAŃSKI
INSTITUTE OF NUCLEAR PHYSICS
POLISH ACADEMY OF SCIENCES

Current issues

NNLO + PS

- generality
- accuracy

Mixed EW x QCD corrections

- amplitudes
- subtraction

FO subtraction schemes at NNLO

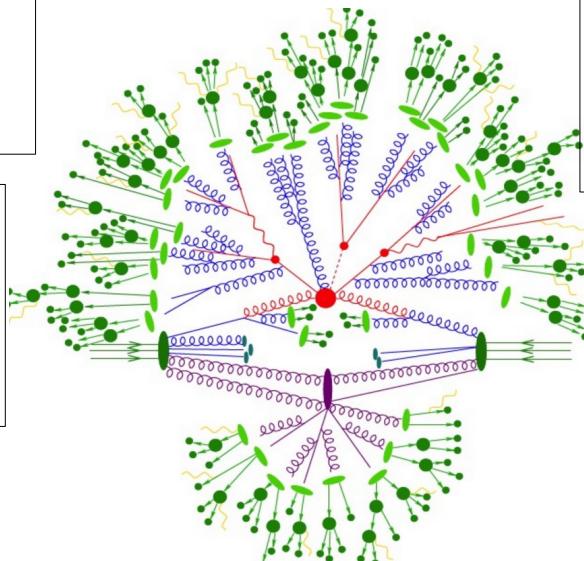
- generality & efficiency
- numerical integration

Multi-loop amplitudes

- Methods
- Automation

Theory uncertainties

- fixed-order
- resummation
- PDFs (aN3LO/flavour schemes)



Flavoured jets

- flavoured algorithms
- experimental implementation

Toponium/top-pair threshold

- bound state
- resummation

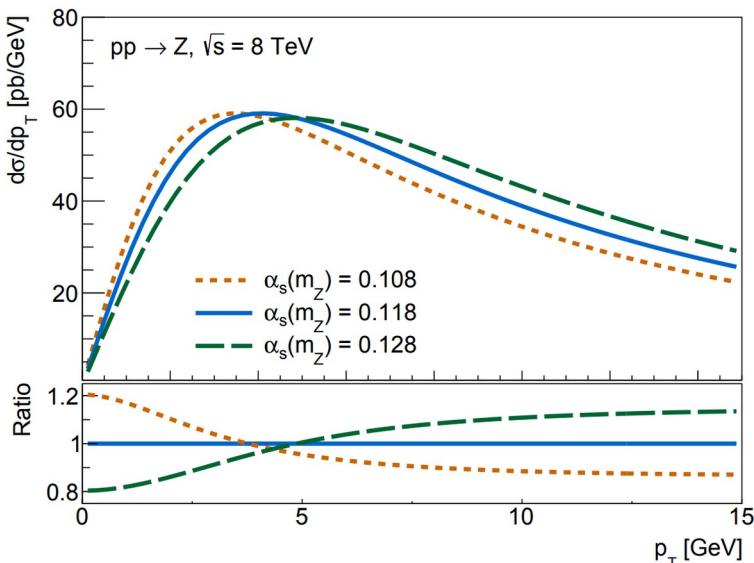
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Theory uncertainties

Precision example: strong coupling from pT(Z)

[ATLAS 2309.12986]

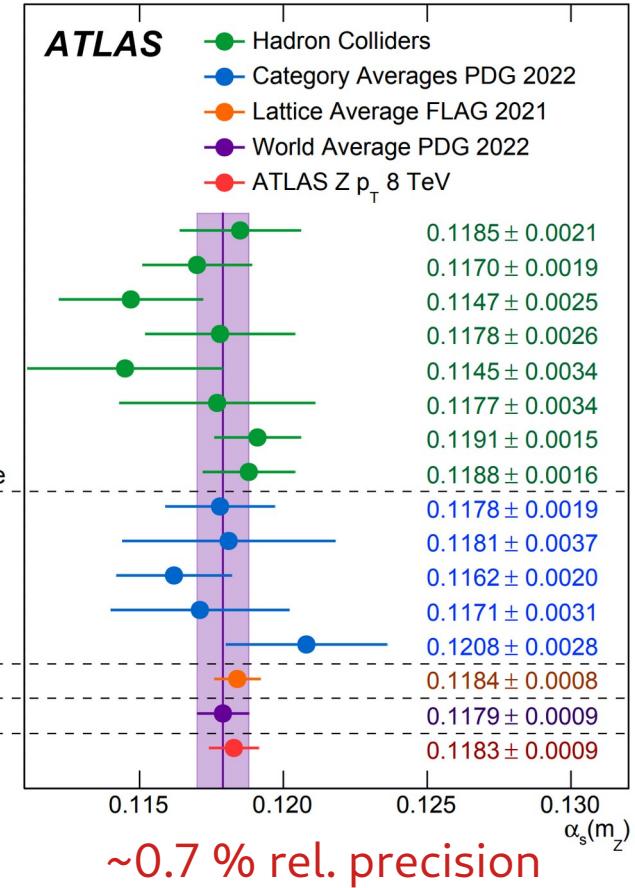
Sensitivity of Z-boson's recoil to the strong coupling constant:



→ at low pT resummation regime!

→ theory uncertainty?

ATLAS ATEEC
CMS jets
H1 jets
HERA jets
CMS $t\bar{t}$ inclusive
Tevatron+LHC $t\bar{t}$ inclusive
CDF Z p_T
Tevatron+LHC W, Z inclusive
 τ decays and low Q^2
QQ bound states
PDF fits
 e^+e^- jets and shapes
Electroweak fit
Lattice
World average
ATLAS Z p_T 8 TeV



Theory uncertainties in Higgs production

[LHC(H)XS)WG YR4' 16]

N4LO approximation
[Das, Moch, Vogt '20]

aN3LO PDFs
[MSHT'22]
[NNPDF'24]

Exact top-mass dependence
through NNLO QCD
[Czakon, Harlander, Klappert, Niggetiedt'21]

Input parameters

\sqrt{S}	13 TeV
m_h	125 GeV
PDF	PDF4LHC15_nnlo_100
$\alpha_s(m_Z)$	0.118
$m_t(m_t)$	162.7 GeV ($\overline{\text{MS}}$)
$m_b(m_b)$	4.18 GeV ($\overline{\text{MS}}$)
$m_c(3\text{GeV})$	0.986 GeV ($\overline{\text{MS}}$)
$\mu = \mu_R = \mu_F$	62.5 GeV ($= m_H/2$)

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
+0.10 pb	± 0.18 pb	± 0.56 pb	± 0.49 pb	± 0.40 pb	± 0.49 pb
-1.15 pb					
+0.21%	$\pm 0.37\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$
-2.37%					

N3LO HEFT
[Mistlberger'18]

Improved QCD-EW predictions
[Bonetti, Melnikov, Trancredi'18] [Anastasiou et al '19]
[Bonetti et al. '20][Bechetti et al. '21] [Bonetti, Panzer, Trancredi '22]

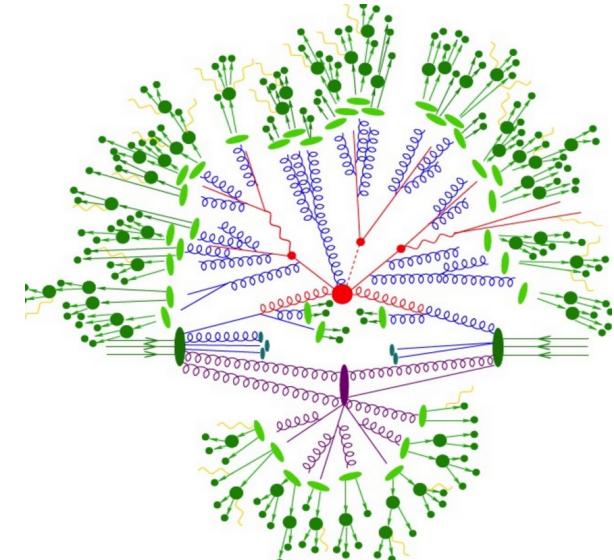
Bottom-top-interference
[Czakon, Eschment, Niggetiedt,
Poncelet, Schellenberger '23, '24]

Uncertainties in precision phenomenology

Experiments are getting more precise → theory uncertainties matter!

Sources of theory uncertainties:

- parametric (values of coupling parameters etc.)
→ variation of parameters within their uncertainties
- parton distribution functions (PDFs)
→ different error propagation methods (fit parameter, replicas,...)
- non-perturbative parameters in Monte Carlo simulations.
→ needs data constraints by definition. Problematic if dominant effect...
- **missing higher orders in fixed-order and resummed predictions (MHOU)**
→ tricky because we are trying to estimate the unknown....



[Credit: SHERPA]

Missing higher orders

Notation from: [Tackmann 2411.18606]

Beyond Scale Variations: Perturbative Theory Uncertainties from Nuisance Parameters

Generic perturbative expansion:

$$f(\alpha) = f_0 + f_1\alpha + f_2\alpha^2 + f_3\alpha^3 + \dots$$

f_i : the coefficient of the series, potentially unknown

We can compute the truncated series: \hat{f}_i : the true value, i.e. a value we actually computed

$$f^{\text{LO}}(\alpha) = \hat{f}_0 \quad f^{\text{NLO}}(\alpha) = \hat{f}_0 + \hat{f}_1\alpha \quad f^{\text{NNLO}}(\alpha) = \hat{f}_0 + \hat{f}_1\alpha + \hat{f}_2\alpha^2$$

The missing terms are the source of uncertainty.

(assume convergence \rightarrow the first missing is the dominant one)

$$f^{\text{LO+1}}(\alpha) = \hat{f}_0 + f_1\alpha \quad f^{\text{NLO+1}}(\alpha) = \hat{f}_0 + \hat{f}_1\alpha + f_2\alpha^2 \quad f^{\text{NNLO+1}}(\alpha) = \hat{f}_0 + \hat{f}_1\alpha + \hat{f}_2\alpha^2 + f_3\alpha^3$$

Challenge: how to estimate f_1, f_2, f_3, \dots without computing them?

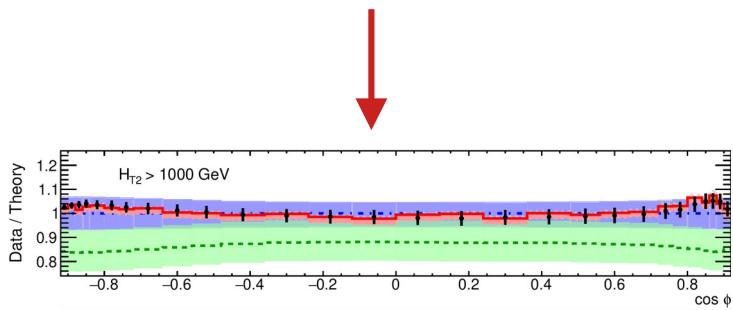
Theory uncertainties from scale variations

Lets focus on QCD as an example: $\alpha = \alpha_s(\mu_0)$

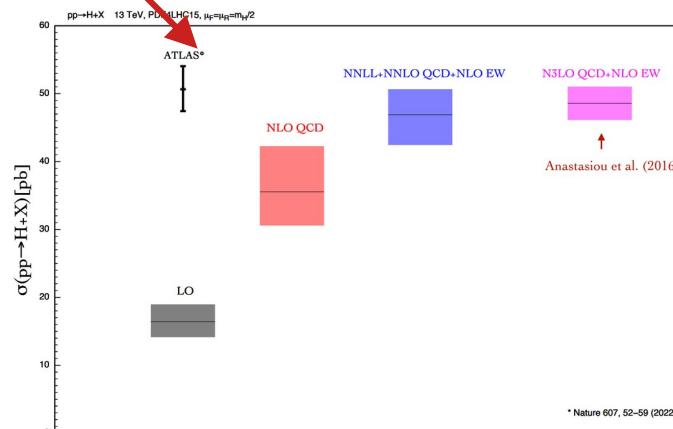
$$d\sigma = d\sigma^{(0)} + \alpha_s d\sigma^{(1)} + \alpha_s^2 d\sigma^{(2)} + \dots \xrightarrow{\text{RGE}} \mu \frac{d\sigma^{(n)}}{d\mu} = \mathcal{O}\left((n_0+n+1)\right)$$

Of same order as the next dominant term
→ exploiting this to estimate size of $d\sigma^{(n+1)}$

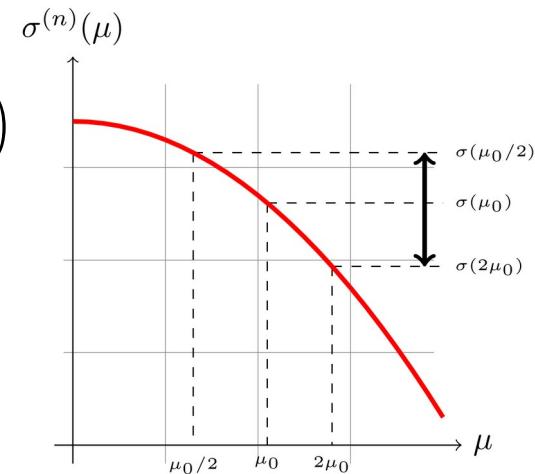
Sometimes it does works ... sometimes not



[ATLAS 2301.09351]



[Grazzini, talk]



Short comings of scale variations

- **not always reliable ...** however in most cases issues are understood/expected:
new channels, phase space constraints, etc. → often we can design workarounds
- however, some issues are more fundamental:
 - how to choose the **central scale?** → **not a physical parameter**, no 'true' value
(Principle of fastest apparent convergence, principle of minimal sensitivity,...)
 - how to propagate the estimated uncertainty, **no statistical interpretation!**
 - what about **correlations**? Based on 'fixed form' of the lower orders and RGE.

(At the moment) two alternative approaches under investigation:

"Bayesian"

[Cacciari,Houdeau 1105.5152]

[Bonvini 2006.16293]

[Duhr, Huss, Mazeliauskas, Szafron 2106.04585]

"Theory Nuisance Parameter"

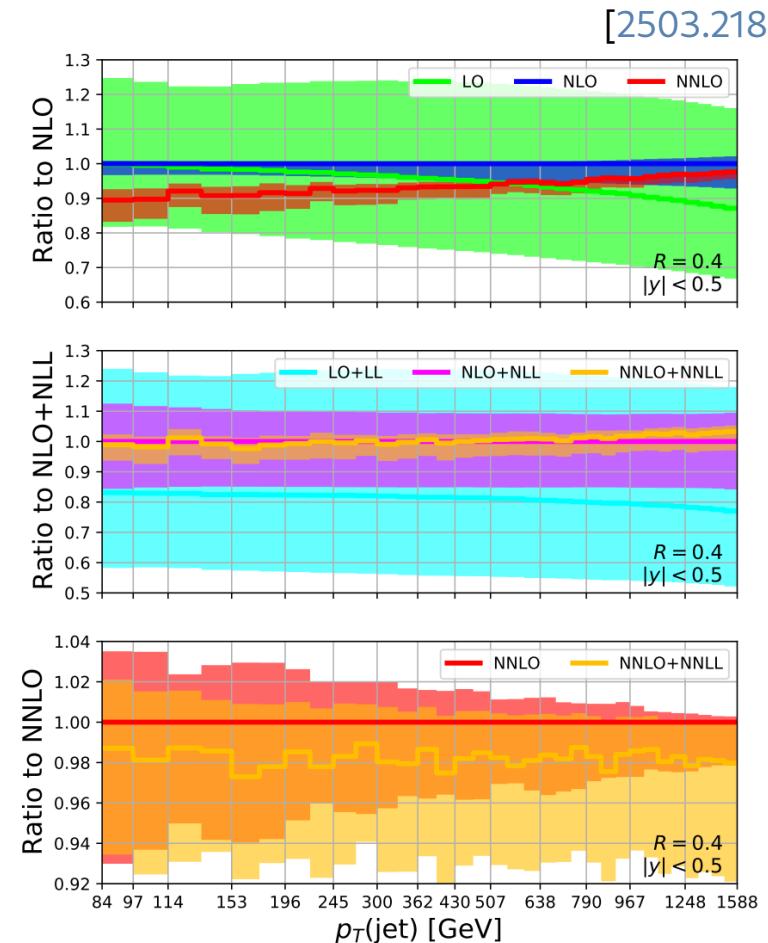
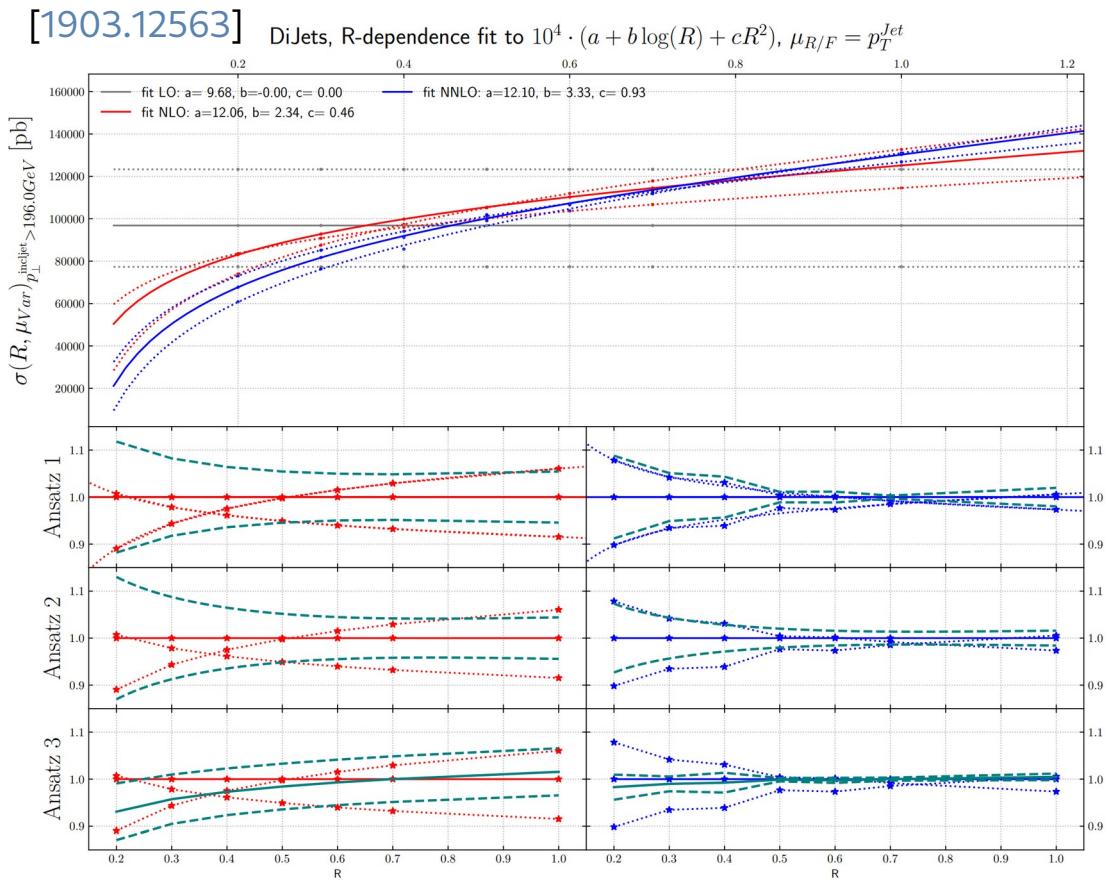
[Tackmann 2411.18606]

→ W mass extraction: [CMS 2412.13872]

[Cal,Lim, Scott,Tackmann Waalewijn 2408.13301]

[Lim, Poncelet, 2412.14910]

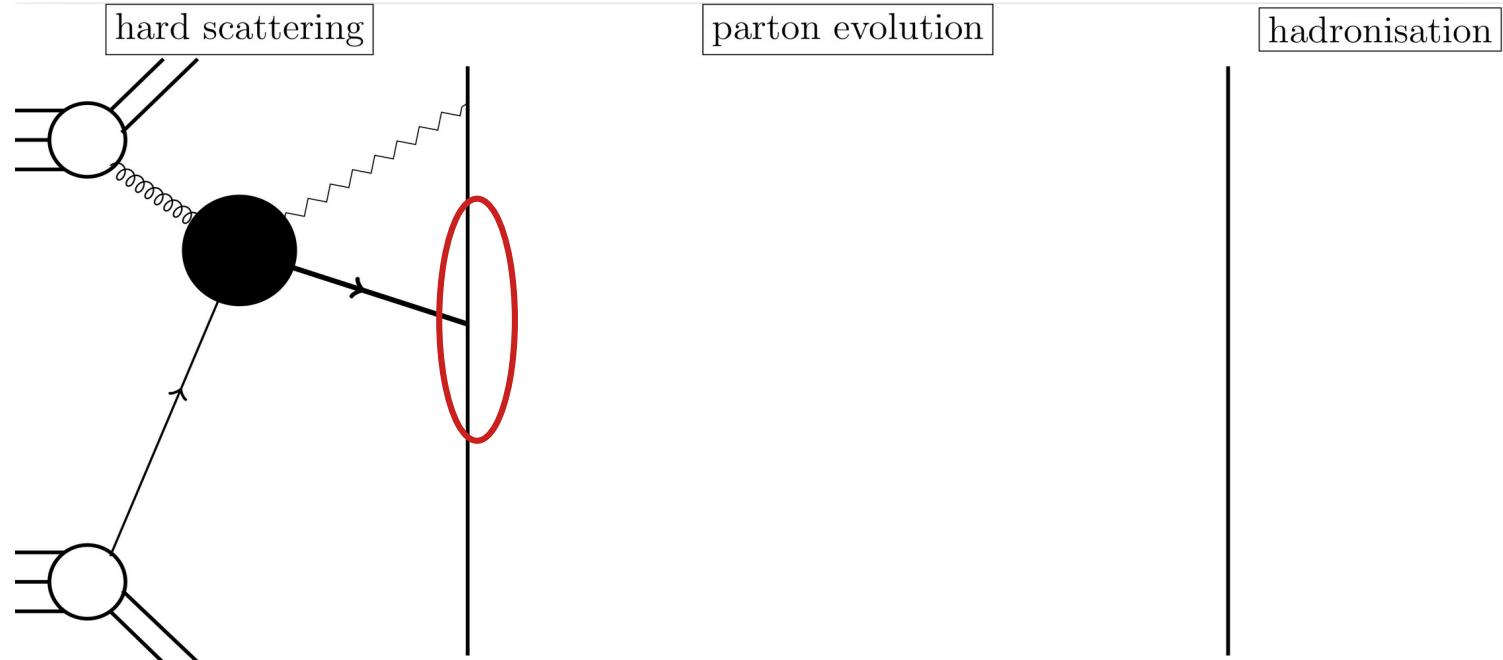
Example: uncertainties in jet cross sections



Flavoured jets

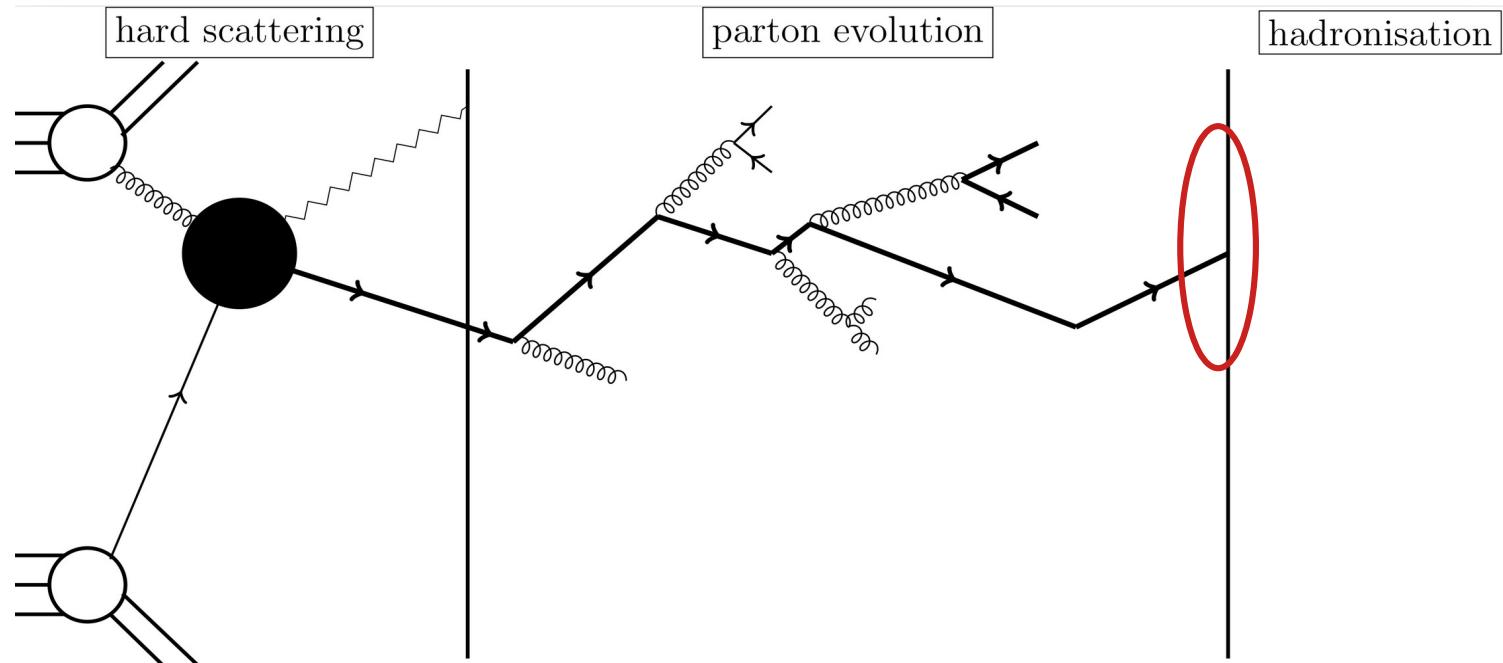
Flavoured jets

jet that initiated from a "hard scatter" product of specific flavour:
bottom, charm , "quark/gluon"



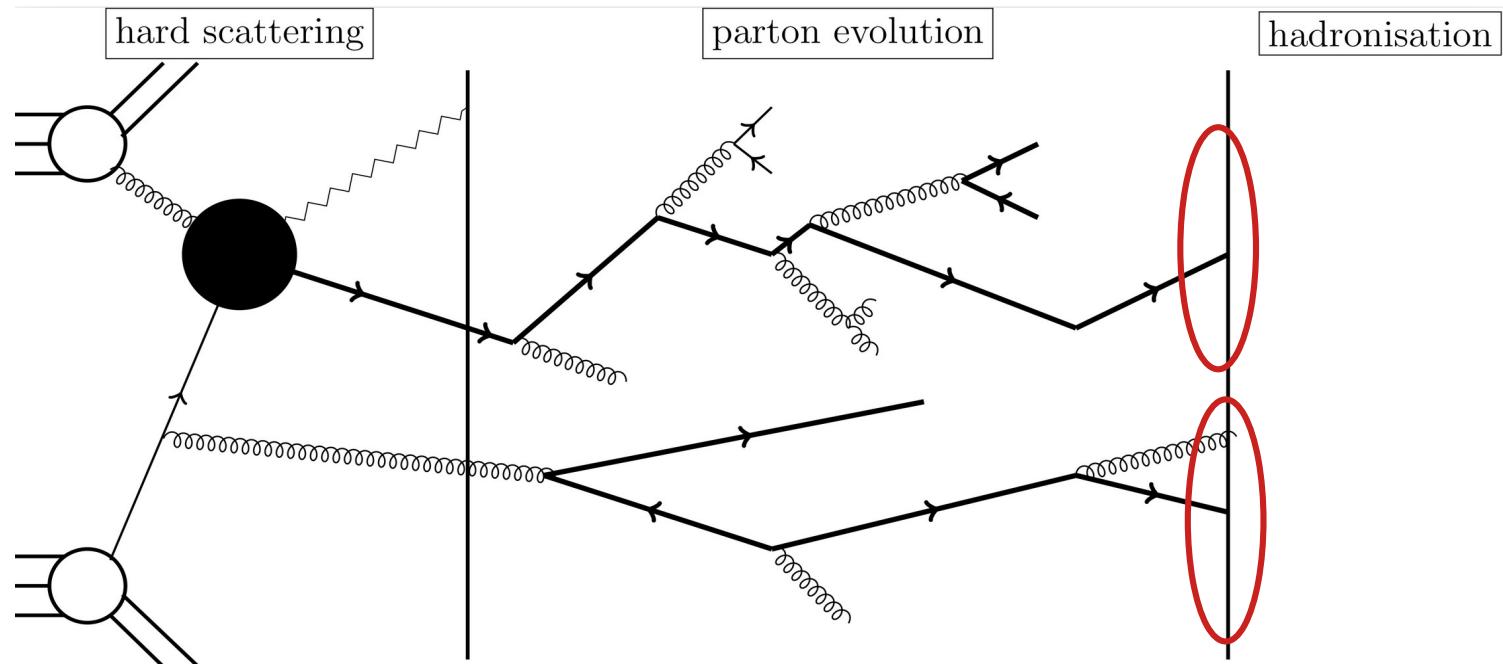
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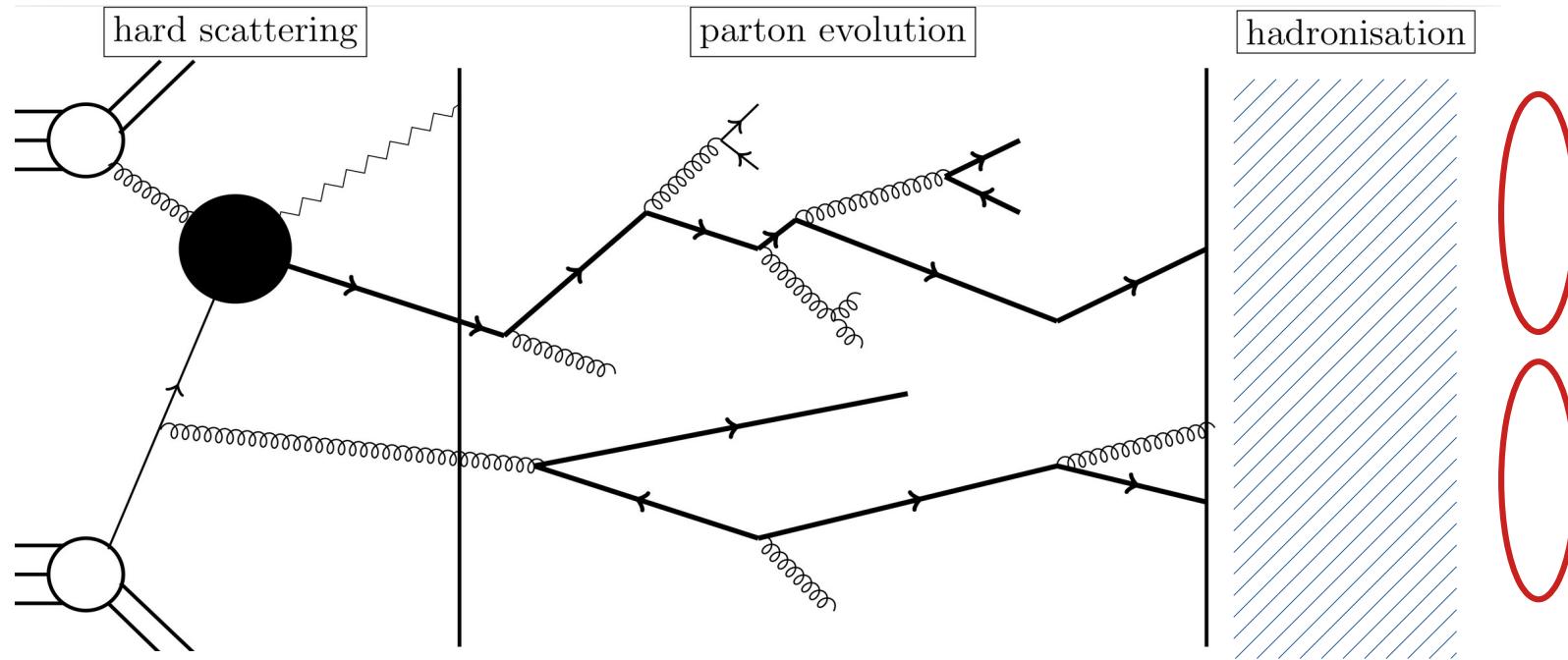
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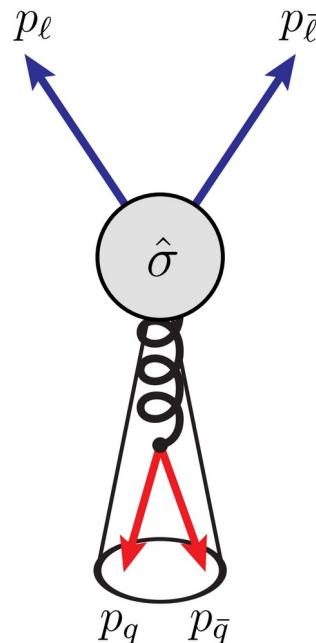


Flavoured jets

jet that initiated from a "hard scatter" product of specific flavour:
bottom, charm , "quark/gluon"



Infrared safety issues with flavoured jets I



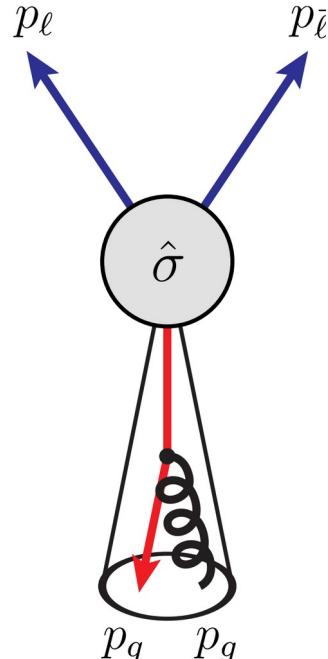
- IRC unsafe due to $g \rightarrow$ quark-anti-quark splitting
 - Quarks massless: cross-section not defined
 - Quarks massive: logarithmic sensitivity to quark mass
- Can be resolved by proper flavour recombination schemes:

<i>jet contents scheme</i>	b	$b + \bar{b}$	$b + b$	
"any flavour"	b	b	b	simplest experimentally (but collinear unsafe for $m_b \rightarrow 0$)
net flavour	b	g	$2b$	theoretically "ideal" definition; but not robust wrt B-Bbar oscillations
flavour modulo 2	b	g	g	theoretically OK; robust wrt B-Bbar oscillations

Picture from
[Gauld et al. 2302.12844]

[Salam]

Infrared safety issues with flavoured jets II

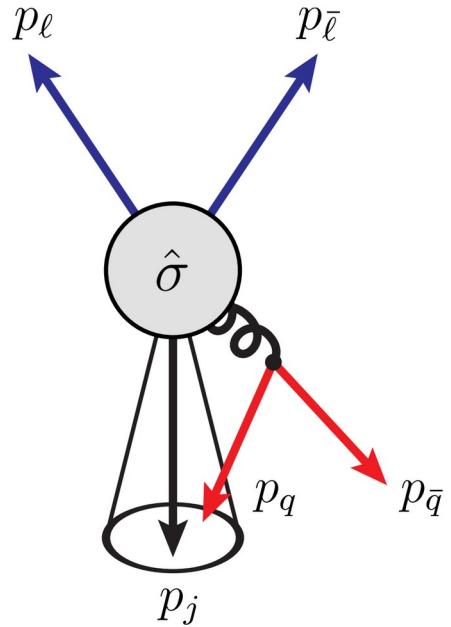


- Collinear unsafe if pT requirement on the quark is present
- Not implementable in pQCD with massless quarks
→ proper treatment needs fragmentation functions
→ NNLO QCD example:

B-hadron production in NNLO QCD: application to LHC ttbar events with leptonic decays,
Czakon, Generet, Mitov and Poncelet, 2102.08267

Picture from
[Gauld et al. 2302.12844]

Infrared safety issues with flavoured jets III



- Starting at NNLO QCD:
→ Soft singularity from quark pairs
- Massless quarks → cross section not defined
- Massive quarks → logarithmic IRC sensitivity $\ln\left(\frac{m}{p_T}\right)$
- **Needs modified jet algorithms!**

Picture from
[Gauld et al. 2302.12844]

New proposals for flavour-safe anti-kT jets

- Flavour with Soft-drop

Practical Jet Flavour Through NNLO

Caletti, Larkoski, Marzani, Reichelt 2205.01109

- Flavour anti-kT

Infrared-safe flavoured anti-kT jets,

Czakon, Mitov, Poncelet 2205.11879

- Fragmentation approach

A Fragmentation Approach to Jet Flavor

Caletti, Larkoski, Marzani, Reichelt 2205.01117

B-hadron production in NNLO QCD: application to LHC ttbar events with leptonic decays,

Czakon, Generet, Mitov and Poncelet, 2102.08267

- Flavour dressing → standard anti-kT + flavour assignment

QCD-aware partonic jet clustering for truth-jet flavour labelling

Buckley, Pollard 1507.00508

A dress of flavour to suit any jet

Gauld, Huss, Stagnitto 2208.11138

- Interleaved flavour neutralisation

Flavoured jets with exact anti-kT kinematics and tests of infrared and collinear safety

Caola, Grabarczyk, Hutt, Salam, Scyboz, Thaler 2306.07314

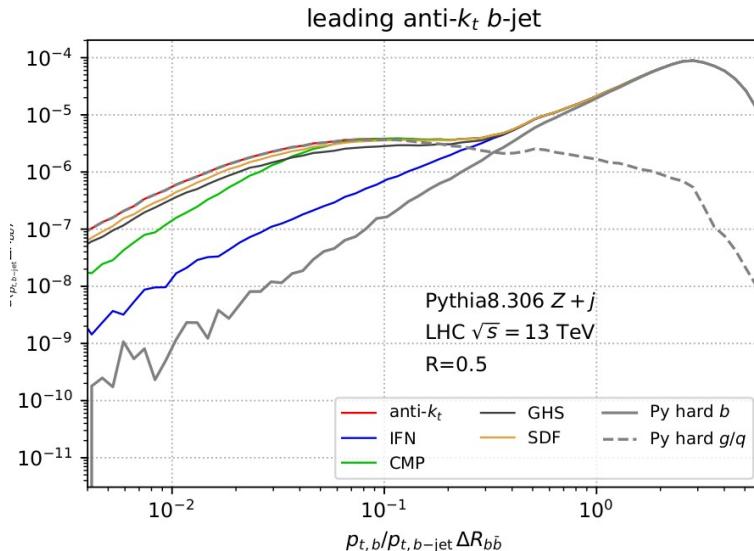
- TBC...

Outcome LH23

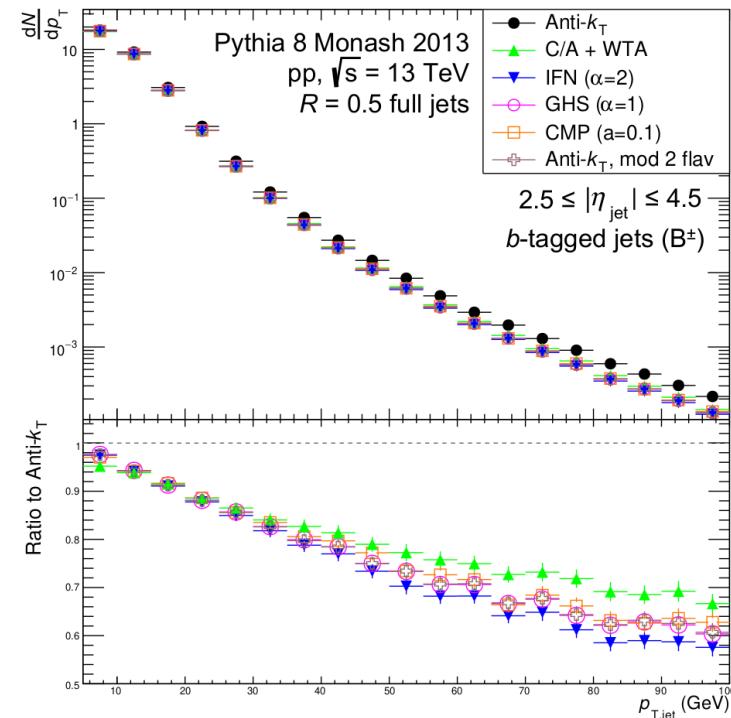
Today on the arxiv

Flavoured jet algorithms: a comparative study

Arnd Behring,¹ Simone Caletti,² Francesco Giuli,³ Radosław Grabarczyk,^{4,5} Andreas Hinzmann,⁶ Alexander Huss,⁷ Joey Huston,⁸ Ezra D. Lesser,⁷ Simone Marzani,⁹ Davide Napoletano,¹⁰ René Poncelet,¹¹ Daniel Reichelt,^{7,12} Alberto Rescia,⁹ Gavin P. Salam,¹³ Ludovic Scyboz,¹⁵ Federico Sforza,⁹ Andrzej Siódtek,^{7,16} Giovanni Stagnitto,¹⁰ James Whitehead,¹⁶ Ruide Xu¹⁷

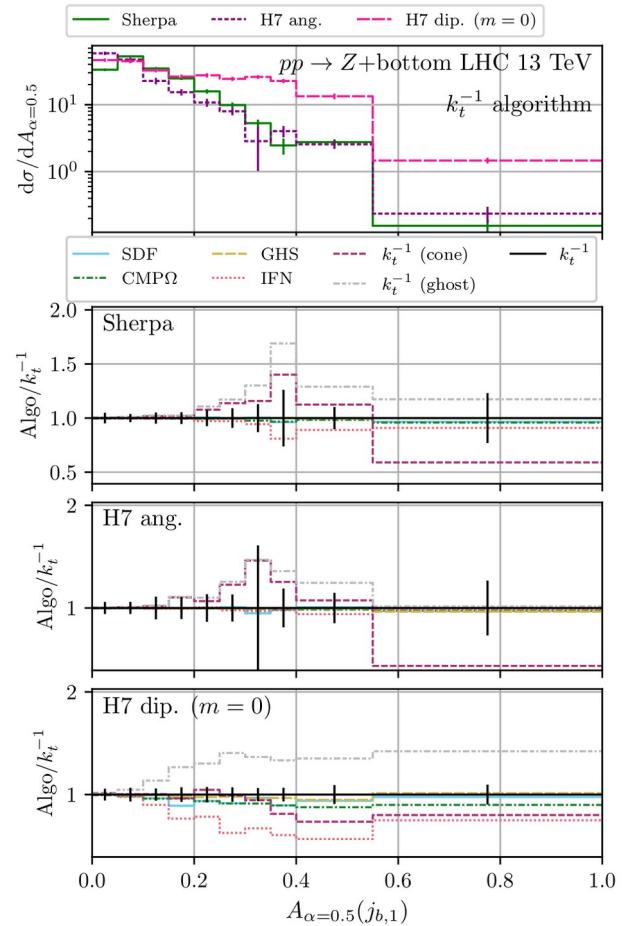
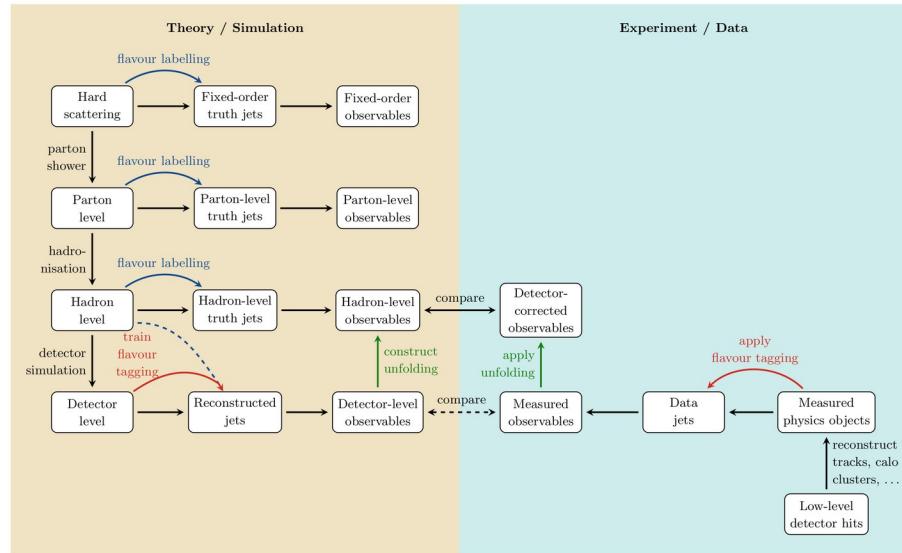


Many analyses from theory and experiment groups → improved understanding on all sides



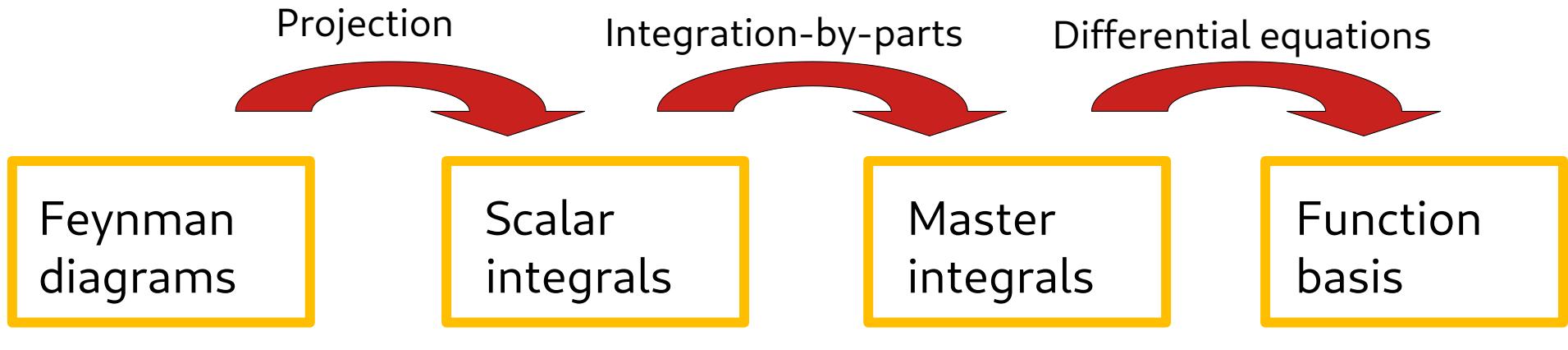
Open Questions

- impact of hadronisation?
- impact on jet substructure?
- how to incorporate in experiment?
- revisit the truth level labels in the training of ML taggers?
- can we get an experimental handle on $g \rightarrow f\bar{f}$?



Amplitudes

Typical multi-loop calculation work-flow



Finite Field methodology to circumvent
excessively large intermediate expressions

Integration-By-Parts reduction

$$a_i^{(L),p} = \sum_i c_{j,i}(\{p\}, \epsilon) \mathcal{I}(\{p\}, \epsilon) \quad \rightarrow \text{prohibitively large number of integrals}$$

$$\mathcal{I}_i(\{p\}, \epsilon) \equiv \mathcal{I}(\vec{n}_i, \{p\}, \epsilon) = \int \frac{d^d k_1}{(2\pi)^d} \frac{d^d k_2}{(2\pi)^d} \prod_{k=1}^{11} D_k^{-n_{i,k}}(\{p\}, \{k\})$$

Integration-By-Parts identities connect different integrals \rightarrow system of equations
 \rightarrow only a small number of independent “master” integrals

$$0 = \int \frac{d^d k_1}{(2\pi)^d} \frac{d^d k_2}{(2\pi)^d} l_\mu \frac{\partial}{\partial l^\mu} \prod_{k=1}^{11} D_k^{-n_{i,k}}(\{p\}, \{k\}) \quad \text{with} \quad l \in \{p\} \cap \{k\}$$



$$a_i^{(L),p} = \sum_i d_{j,i}(\{p\}, \epsilon) \text{MI}(\{p\}, \epsilon)$$

Master integrals

Differential Equations: $d\vec{MI} = dA(\{p\}, \epsilon)\vec{MI}$

[Remiddi, 97]

[Gehrmann, Remiddi, 99]

[Henn, 13]

Canonical basis: $d\vec{MI} = \epsilon d\tilde{A}(\{p\})\vec{MI}$

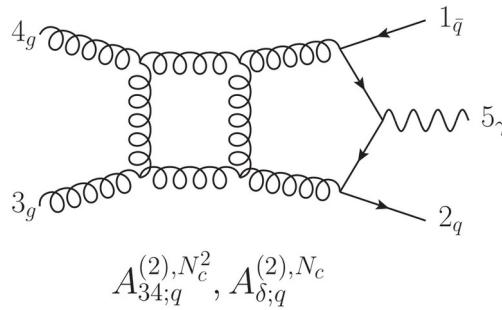
Simple iterative solution



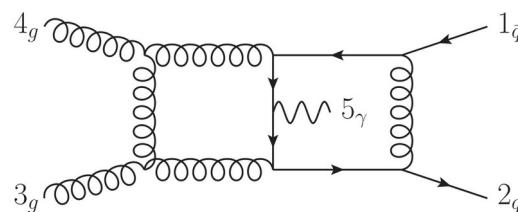
$$MI_i = \sum_w \epsilon^w \tilde{MI}_i^w \quad \text{with} \quad \tilde{MI}_i^w = \sum_j c_{i,j} m_j$$

Chen-iterated integrals
"Pentagon"-functions
[Chicherin, Sotnikov, 20]
[Chicherin, Sotnikov, Zoia, 21]

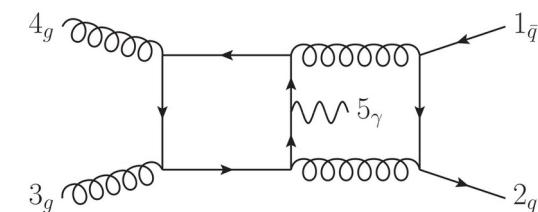
Reconstruction of Amplitudes



$$A_{34;q}^{(2),N_c^2}, A_{\delta;q}^{(2),N_c}$$



$$A_{34;q}^{(2),1}, A_{\delta;q}^{(2),N_c}, A_{\delta;q}^{(2),1/N_c}$$



$$A_{34;l}^{(2),N_c}, A_{34;l}^{(2),1/N_c}, A_{\delta;l}^{(2),1/N_c^2}$$

[Badger et al '23]

Many possible optimizations

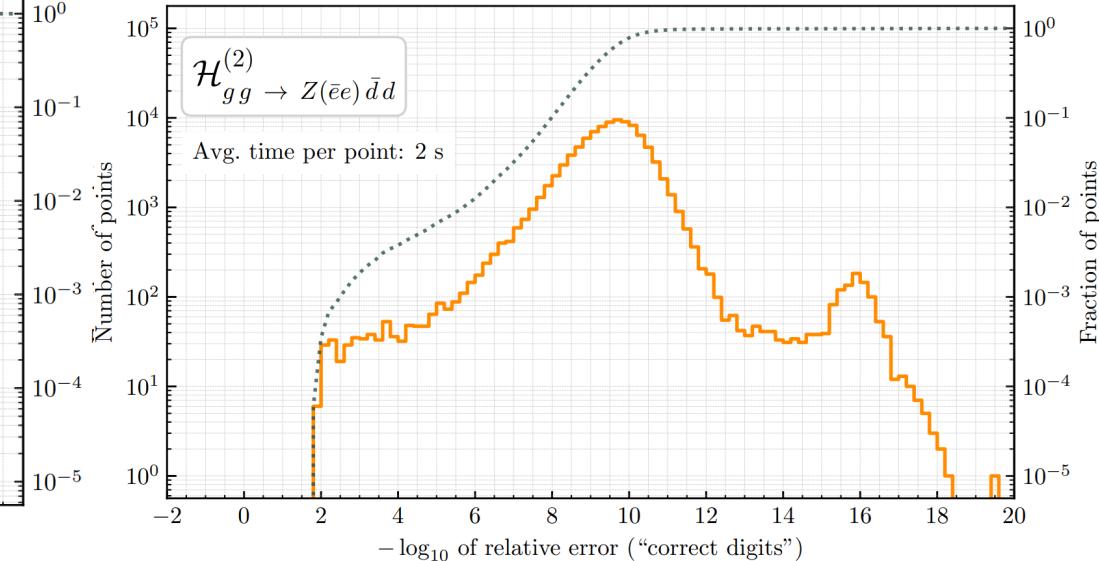
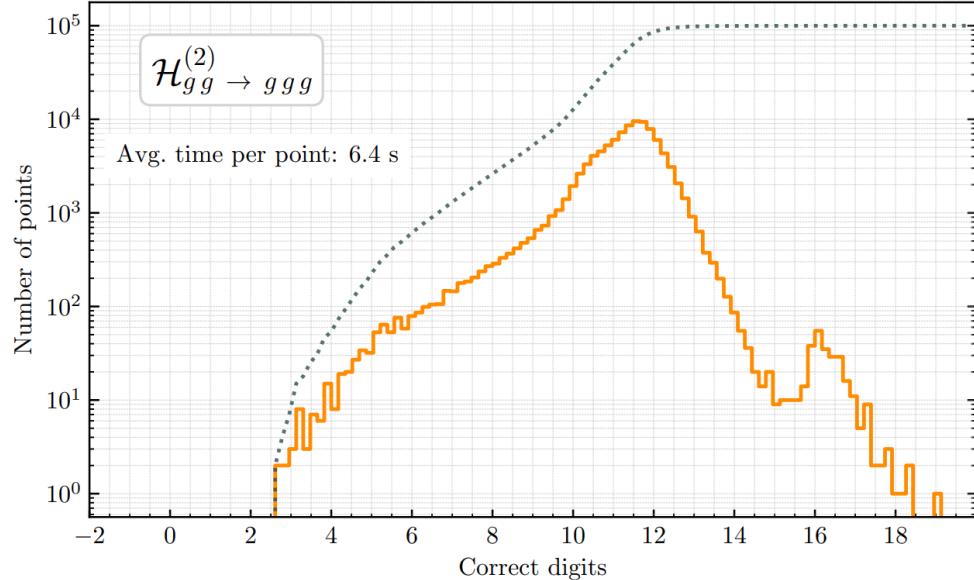
- Syzygy's to simplify IBPs
- Exploitation of Q-linear relations
- Denominator Ansaetze
- On-the-fly partial fractioning
- ...

amplitude	helicity	original	stage 1	stage 2	stage 3	stage 4
$A_{34;q}^{(2),1}$	- + + - +	94/91	74/71	74/0	22/18	22/0
$A_{34;q}^{(2),1}$	- + - + +	93/89	90/86	90/0	24/14	18/0
$A_{34;q}^{(2),1/N_c^2}$	- + + - +	90/88	73/71	73/0	23/18	22/0
$A_{34;q}^{(2),1/N_c^2}$	- + - + +	90/86	86/82	86/0	24/14	19/0
$A_{34;l}^{(2),1/N_c}$	- + - + +	89/82	74/67	73/0	27/14	20/0
$A_{34;l}^{(2),1/N_c}$	- + + - +	85/81	61/58	60/0	27/18	20/0
$A_{34;q}^{(2),N_c^2}$	- + - + +	58/55	54/51	53/0	20/16	20/0

Reduction of complexity

Numerical evaluation of two-loop matrix elements

[De Laurentis, Ita, Sotnikov'23][De Laurentis, Ita, Page, Sotnikov'25]



Efficient function basis representation power the speed.
If too expensive \rightarrow numerical interpolation grids (used a lot for 2D problems)?
Methods for high dimensions?

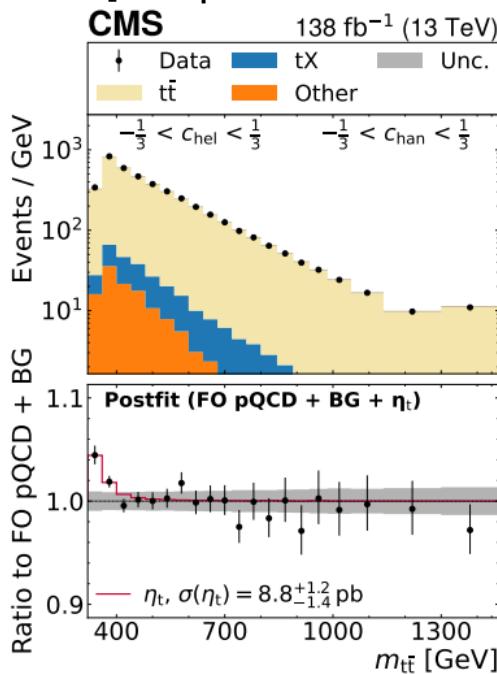
Topics for amplitude discussions

- Geometry, elliptic functions, Calabi-Yau, ...
- Dimensional regularization and gamma5
- Automation @ two-loops
- Standardisation of interfaces
- Efficiency:
 - ML
 - **Interpolation**
 - partial fractioning
 - singularities

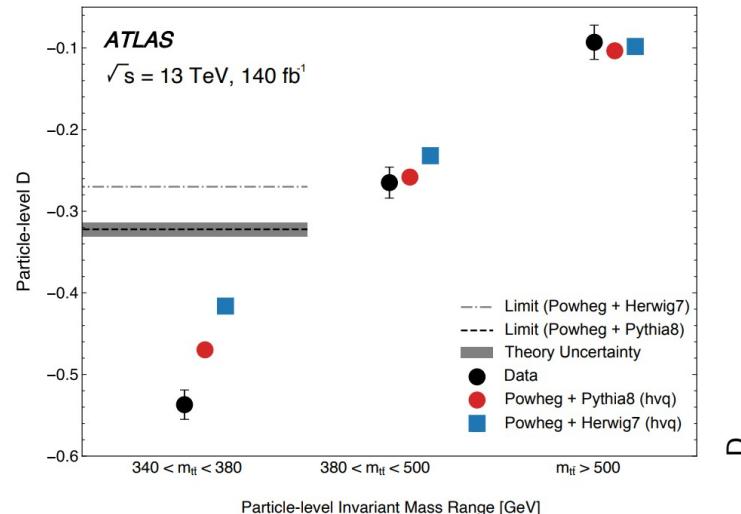
Toponium / top-pair threshold

Surprises at the threshold

[CMS'25] Toponium

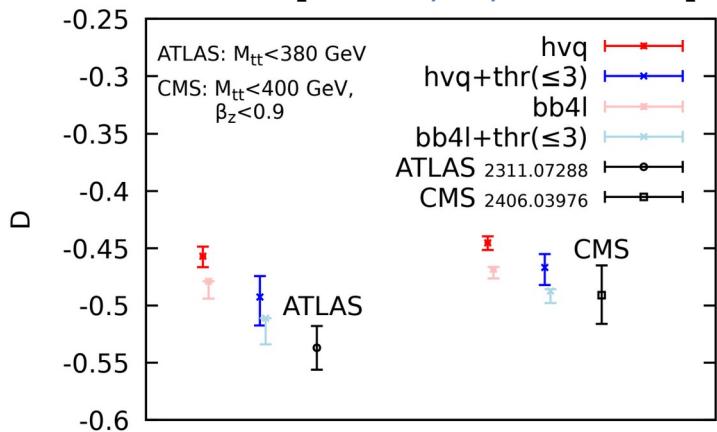


[ATLAS'23] Entanglement



Threshold QCD enhancements

[Nason,Re,Rottoli'25]



Sounds like super topic for a discussion :)

Numerical integration and non-positive integrands or negatively weighted events

The problem

- Numerical integration of highly dimensional integrands → Monte Carlo Sampling

Integral

$$I = \int_{\mathbf{x} \in \Omega} d\mathbf{x} f(\mathbf{x})$$

MC estimate

$$\hat{I} = \frac{1}{N} \sum_{i=1}^N f(\mathbf{x}_i), \quad \delta \hat{I} = \sqrt{\frac{1}{N-1} \left(\frac{1}{N} \sum_{i=1}^N f^2(\mathbf{x}_i) - \hat{I}^2 \right)}$$

MC error estimate

- Variance reduction techniques improve performance, mapping $\mathbf{H} : \Omega \rightarrow \Omega, \mathbf{x} \mapsto \mathbf{H}(\mathbf{x})$

$$I = \int_{\mathbf{H}(\mathbf{x}) \in \Omega} d\mathbf{H} \frac{f(\mathbf{x})}{h(\mathbf{x})}$$

$$h(\mathbf{x}) = \left| \det \left(\frac{\partial \mathbf{H}(\mathbf{x})}{\partial \mathbf{x}} \right) \right|$$



Find with $h(\mathbf{x})$ adaptive MC techniques: VEGAS [Lepage'78], Parni [Hameren'14],
ML techniques: Normalising Flows Iflow [Bothmann'20] Madnis [Heimel'22], ...

Non-positive definite integrands

- Non-definite integrands introduce computational challenges
→ cancellation between +/- parts increase the variance
- Consider extreme case: $|f(x)|/h(x) = w = \text{const.}$

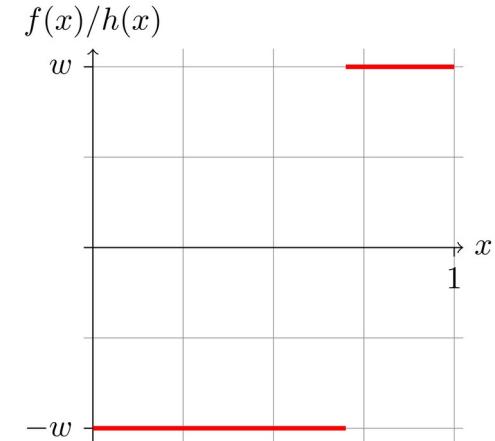
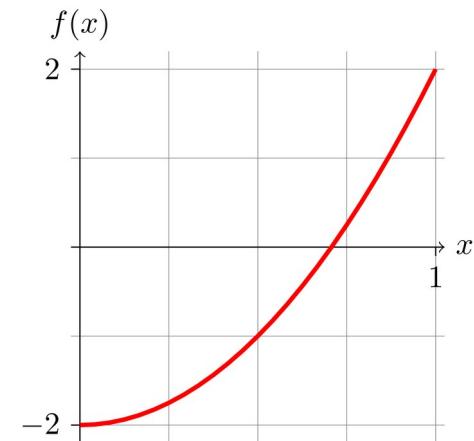
MC estimate: $\hat{I} = w \frac{N_+ - N_-}{N} \equiv w(2\alpha - 1) \quad \alpha = N_+/N$

- Lower bound on variance:

$$\text{Var}(\hat{I}) = w^2 - w^2(2\alpha - 1)^2 = w^2(4\alpha(1 - \alpha))$$

$$\rightarrow \text{relative uncertainty: } \frac{\delta \hat{I}}{\hat{I}} = \frac{1}{\sqrt{N-1}} \frac{\sqrt{\alpha(1-\alpha)}}{\alpha - \frac{1}{2}}$$

Rephrased: at some point it doesn't matter any more how good your adaptive MC is...

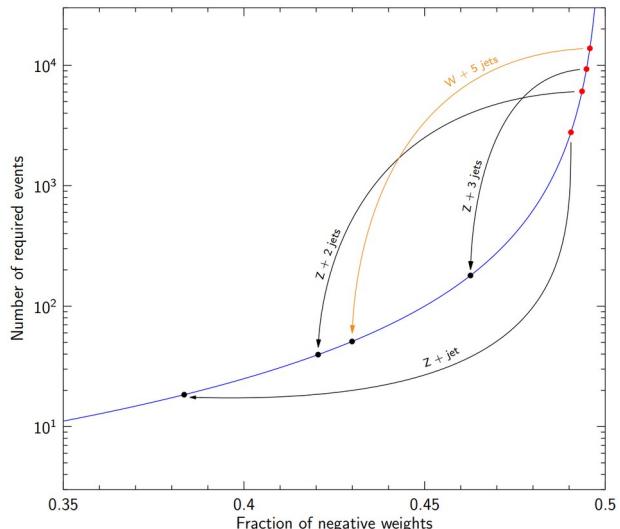


Negative Weights

How to improve convergence
of non-positive definite MC integrals?
→ phase space splitting

[Janssen, Schumann, Poncelet'25]

[Andersen, Maier, Maître, 2023]



How to avoid negative weighted events?
→ Design a positive matching
for example ESME, KrKNLO, ...

[van Beekveld, Ravasio, Helliwell, Karlberg, Salam,
Scyboz, Soto-Ontoso, Soyez, Zanolí'25]
[Whitehead, Sarmah, Siódmod'24]

I have negative weights
→ How to get rid of them?
Cell-resampling [Anderson, Maier'21]

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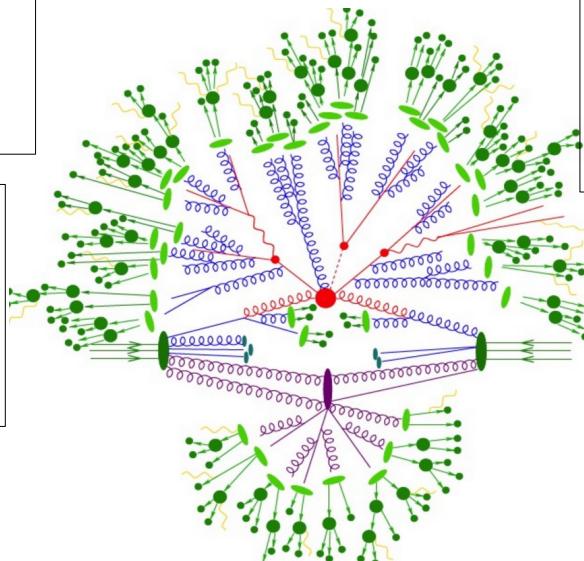
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Summary

Looking forward to a productive week with lots of interesting discussions!

