
Tricks and Techniques

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I KNEW A LOT OF THINGS. OVER TIME, I FORGOT MANY THINGS. NOW, I ONLY KNOW A FEW THINGS. THEREFORE I TOOK THE WISE DECISION OF NOTING STUFF DOWN. TO START, I TRY TO SUMMARIZE THE MOST ELEMENTARY THINGS I KNEW, AND THE HOPE IS THAT I'LL KEEP ON ADDING NEW THINGS THAT I LEARN. FOR NOW, IT'S ALMOST ALL THE UNDERGRAD STUFF WHICH I AM TRYING TO SUMMARIZE AS QUICKLY AS POSSIBLE.

Chapter 1

Magic of Infinitesimals

1.1 Continuity

1.1.1 Any mid value

Statement: If function $f(x)$ is continuous in the interval $[a, b]$ with $f(a) = A$ and $f(b) = B$, then every value in the interval $[A, B]$ is achieved. Namely, $\forall N \in [A, B], \exists \xi \in [a, b]$ such that $f(\xi) = N$

Proof: Not required.

1.1.2 Max and Min value

Statement: If the function is continuous in a closed interval than it has a largest and a smallest value for atleast one x in that interval. (Other way of saying that it doesn't become infinite).

Use: In rolle's theorem, for the existence of highest value.

1.2 Differentiability

1.2.1 Rolle's Theorem

Statement : If the function is differentiable in the interval $[a, b]$ and $f(a) = f(b)$, $\exists \xi \in (a, b)$ such that $f'(\xi) = 0$.

Proof : Use sign change/existence argument at maximum.

Use : Finding remainder term in taylor series. Proving L'Hopital rule and various other places where we construct a new function on which we can apply Rolle to prove certain thing about a given function.

1.2.2 Mean Value theorem

Statement : For a differentiable function in a given interval $[a, b]$, $\exists \xi \in (a, b)$ such that $f'(\xi) = \frac{f(b)-f(a)}{b-a}$

Proof : Using Rolle's theorem.

Use : $f(b) = f(a) + (b-a)f'(a + \theta(b-a))$, $\theta \in [0, 1]$.

1.2.3 Taylor series

A function which can be differentiated arbitrary number of times can always be expanded as

$$f(b) = f(a) + (b-a)f'(a) + \frac{(b-a)^2}{2!}f''(a) + \dots + \frac{(b-a)^n}{n!}f^{(n)}(a) + R_n$$

The remainder term can be beautifully calculated using Rolle's theorem as

$$R_n = \frac{(b-a)^{n+1}}{(n+1)!}f^{(n+1)}(\xi)$$

Clearly, this tends to zero as n tends to infinity if the derivative is bounded (which is, because we assumed it exists). A clever trick for finding out Taylor series for indefinite integrals is to find the Taylor series of integrand and integrate them.

1.3 Infinitesimals

Summing up (infinite times), or dividing is same as summing up (infinite times) or dividing only the principal part.

Three properties are important for geometric visualization of calculus:

Property 1 : Any general arc differs from its chord by an infinitesimal of a higher order. Using this, we can imagine arcs as chords (and same extension in higher dimensions, curved surface as planes). We can find the curve length by ignoring this higher order infinitesimal difference between arc and chord because infinite sum of infinitesimals is same as infinite sum of their principal parts.

Property 2 : Perpendicular distance from one end of the infinitesimal arc to the tangent at the other end is an infinitesimal of higher order than arc, however the length of the tangent from the foot of perpendicular to the point of tangency is of the same order. You can prove this by using Maclaurine series in transformed coordinates or by simple geometric construction.

Property 3 : Same trigonometric laws are obeyed. This is because the chord becomes tangent to the arc as their length approaches zero. The rules of chord extend to the arcs.

1.4 Differentials

Note that for an independent variable $d^2x = 0$, not otherwise. For independent variable x , differential dx is same as infinitesimal increment Δx to x . For dependent variable y , the differential dy is the principal part of the infinitesimal change Δy in y . The rules for differentials can easily be derived.

$$\lim_{\Delta x \rightarrow 0} \frac{\Delta y}{\Delta x} = f'(x)$$

$$\Delta y = f'(x)\Delta x + \epsilon\Delta x$$

$$dy = f'(x)dx$$

Knowing usual rules for differentials, we can use them as algebraic quantities, no need to think of differential operators.

Note the difference between the context of the symbol $\frac{d^2y}{dx^2}$. Which has context independent meaning only when x is independent variable.

Implicit Function Theorem : If $F(x, y) = 0$ and $|F_x| < \infty$ with $F_y \neq 0$, then $y(x)$ exists and is differentiable. Can be easily extended to higher dimensions and/or involving more than one function using jacobians.

Chapter 2

Sequences

I list all the theorems which you can prove in order and digest.

Defn : A sequence (x_n) is said to converge to x if $\forall \epsilon > 0, \exists N \in \mathbb{N}$ such that $|x_n - x| < \epsilon, \forall n > N$.

Theorem 1 : If sequence (x_n) converges, it has a unique limit point.

Theorem 2 : If sequence (x_n) converges, it is bounded. (Bounded : $\exists M \in \mathbb{R}$ such that $\forall n, |x_n| < M$).

Theorem 3 : If a **monotone** sequence (x_n) is bounded, then it converges. The limit is $\sup\{x_n : n \in \mathbb{N}\}$

Theorem 4 : Every subsequence of a convergent sequence converges to the limit of sequence.

Theorem 5 : Every sequence has a monotone subsequence.

Theorem 6 : (Bolzano) Every bounded sequence has a convergent subsequence. (Just a consequence of prev theorem)

Defn : Limsup/ Liminf of sequence (a_n) are defined by

$$\limsup_{n \rightarrow \infty} a_n = \lim_{n \rightarrow \infty} \sup(\{a_k : k \geq n\})$$
$$\liminf_{n \rightarrow \infty} a_n = \lim_{n \rightarrow \infty} \inf(\{a_k : k \geq n\})$$

Theorem 7 : limsup/liminf of a bounded sequence always exists.

Theorem 7 : Let (x_n) be a bounded sequence. Then \exists subsequences (x_{n_k}) and (x_{m_k}) s.t.

$$\lim_{n \rightarrow \infty} x_{n_k} = \limsup x_n$$

$$\lim_{n \rightarrow \infty} x_{m_k} = \liminf x_n$$

Theorem 8 : If every convergent subsequence of (x_n) has the same limit x , then (x_n) converges to x .

Cauchy Sequence : A sequence (x_n) is cauchy sequence, if $\forall \epsilon > 0, \exists N \in \mathbb{N}$ such that $|x_n - x_m| < \epsilon, \forall m, n > N$

Theorem 9 : A sequence (x_n) converges iff it's a Cauchy sequence.

Chapter 3

Series

Same presentation as chapter 2.

Defn : If (a_n) is a sequence of real numbers then the series $\sum_{n=1}^{\infty} a_n$ converges to $S \in \mathbb{R}$ if the sequence (S_n) of partial sum converges to S .

Trick : For geometric and telescopic series, you can find explicit formula for n^{th} term of the sequence (s_n) .

Theorem 1 : A series $\sum a_n$ of positive term converges iff its partial sums are bounded.

Theorem 2 : (Cauchy Condition) The series $\sum a_n$ converges iff $\forall \epsilon > 0, \exists N \in \mathbb{N}$ such that $|S_n - S_m| < \epsilon, \forall n > m > N$

Defn : Series $\sum a_n$ converges absolutely if $\sum |a_n|$ converges. It converges conditionally if $\sum a_n$ converges but $\sum |a_n|$ diverges.

Theorem 4 : Absolutely convergent series $\sum a_n$ converges. Moreover, $\sum a_n$ converges absolutely, iff series $\sum a_n^+$ and $\sum a_n^-$ converges.

Following tests can all be proved simply, using cauchy criterion.

Test for divergence : If series $\sum a_n$ converges, then $\lim_{n \rightarrow \infty} a_n = 0$

Comparison test : If $b_n \geq 0, \sum b_n$ converges, then $\sum a_n$ converges absolutely if $|a_n| \leq b_n$

Ratio test : If (a_n) is a sequence of nonzero real numbers, let

$$r = \lim_{n \rightarrow \infty} \left| \frac{a_{n+1}}{a_n} \right|$$

Series Σa_n converges absolutely if $0 \leq r < 1$, diverges if $1 < r \leq \infty$

Root test : For (a_n)

$$r = \lim_{n \rightarrow \infty} \sup |a_n|^{\frac{1}{n}}$$

The series converges absolutely if $0 \leq r < 1$ and diverges if $1 < r < \infty$.

Alternating series test : If (a_n) is a decreasing sequence of nonnegative real numbers, such that $\lim_{n \rightarrow \infty} a_n = 0$, then the alternating series $\Sigma(-1)^{n+1}a_n$ converges.

Limit Comparison test : If Σa_n and Σb_n are two series and

$$r := \lim_{n \rightarrow \infty} (a_n/b_n)$$

If $r = 0$: If Σb_n converges, then Σa_n converges.

If $r = \infty$: If Σb_n diverges, then Σa_n diverges.

If r is finite, either both converge or both diverge.

Theorem 4 : Every rearrangement of an absolutely convergent series converges to the same sum.

3.0.1 Power series :

There's only one power series for a given function (Coefficients are obviously uniquely defined by differentiation).

- If a power series converges for some $x = x_1$, it converges absolutely and uniformly for all $|x_2| < |x_1|$. This can be proven using comparison test.

Following are the consequences of uniform convergence :

1. The function defined by power series is continuous inside the region of convergence.
 2. The power series can be integrated term by term in the region of convergence.
 3. The power series can be differentiated term by term. This is a more restrictive condition : the continuity of all individual derivatives as well as the uniform convergence of the series formed by derivative is required.
- On the boundary of region of convergence, there is atleast one singularity.
 - Laurent Series and Analytic function.
 - Using residue calculus to sum the series.

Chapter 4

Integration

- Any upper sum will be greater than any lower sum. Infinite integrals (in integrands and in limits) are defined by taking appropriate limits. If function is discontinuous at only finitely many points, both limits approach to the same value, which is defined as the integral
- Differentiating under the integral sign. Uniform convergence in α is required for infinite limits.
- Integrating under the integral sign is rarely appreciated. Nice trick.
- A nice understanding of how lengths change under change of coordinates. The jacobian can easily be derived by parallelepiped volume formula.
- Green's and Stokes theorem (easy proof).

Chapter 5

Complex Variables

- Cauchy reimann conditions for analyticity are sufficient when first order derivatives are also continuous. Analyticity is defined as differentiability in a region.
- As an example $|z|^2$ is differentiable at $z = 0$ but not analytic anywhere.
- The corresponding real functions of analytic function are harmonic and the conjugate functions form an orthogonal family. Moving along one would give maximum change in the other. The mapping is conformal at point when $f'(z) \neq 0$ or $f'(z) \neq \infty$.

5.1 Cauchy's Theorem and Integral Formula:

If $f(z)$ is analytic in a simply connected region and C is a closed contour in that region then

$$\oint_C f(z)dz = 0$$

$$\frac{1}{2\pi i} \oint_C \frac{f(z)}{z - z_0} dz = f(z_0)$$

1. Taylor and laurent series are simple consequence of this.
2. Residue is very important for performing various integrals and summing series(as that is the only term that will remain after integration).
3. Apart from poles and essential singularity, branch points are also a type of singularity.
4. for $z^{\frac{1}{2}}$, the branch point is at $z = 0$, order is defined as the number of paths around that point to return to same value. In this case it's 2.

5.2 Analytic Continuation

In real analysis, a function defined on a given range can be smoothly extended to outer regions in many ways. There's only a unique way to do so for analytic functions. This great property makes complex analysis almost **GODLY**.

If two analytic functions coincide in any region, or on any finite line segment, they are the same function.

Chapter 6

Differential Equations

I won't even mention ODE's with constant coefficients.

6.1 First Order ODE

- One can always find integrating factor for first order ODE in principle. In practice, we only know direct formula for **linear** first order ODE. This also means that every differential constraint (in two dimensions) is holonomic.
- List of famous non-linear first order equations that can be solved

1. Bernoulli equation

$$\frac{dy}{dx} + P(x)y = Q(x)y^n$$

2. Clairaut's equation

$$y = x \frac{dy}{dx} + f\left(\frac{dy}{dx}\right)$$

3. D'Alembert's equation

$$y = xf\left(\frac{dy}{dx}\right) + g\left(\frac{dy}{dx}\right)$$

6.2 Second Order Linear ODE

Most important differential equations arising in physics are almost always second order and linear (ordinary or partial). Let's consider only the homogenous part right now.

$$y'' + P(x)y' + Q(x)y = 0$$

Ordinary Point : x_0 for which $P(x_0)$ and $Q(x_0)$ are finite.

Regular Singular Point : Singular x_0 for which $(x - x_0)P(x)$ and $(x - x_0)^2Q(x)$ remain finite.

Irregular Singular Point : Singular x_0 for which one of $(x - x_0)P(x)$ or $(x - x_0)^2Q(x)$ diverge.

- The only trick we know is Frobenius! Always check if the final solution converges. You have to expand about a clever point. Expanding about essential singularity won't work.
- Expressing ODE as $L(x)y(x) = 0$. If $L(x) = L(-x)$, then if $y(x)$ is a solution, so is $y(-x)$. So you can express general solution as a linear combination of odd and even.

Wronskian : To test for linear independence of functions, we obtain n equations by differentiating the relation

$$\sum_{i=0}^n a_i \phi_i(x) = 0$$

This is the origin of Wronskian! Also used in some ninja-technical way to find particular solution for inhomogenous linear equation.

6.2.1 Sturm-Liouville Theory

In physics, we don't really require general solutions. We require solutions with specific boundary conditions. Specific boundary conditions impose conditions on the parameters of the equation. For instance, for a string clamped at $x = 0$ and $x = l$, in the differential equation $\frac{d^2\psi}{dx^2} + k^2\psi(x) = 0$, k will have certain restrictions as you know. Here, k^2 is the eigenvalue.

Characterization of general features of eigenproblems arising **from linear second-order differential equations** is known as Sturm-Liouville theory.

Consider $\mathcal{L}\psi(x) = \lambda\psi(x)$, where

$$\mathcal{L}(x) = p_0(x)\frac{d^2\psi}{dx^2} + p_1(x)\frac{d\psi}{dx} + p_2(x)$$

Now \mathcal{L} is known as self adjoint if $p_0'(x) = p_1(x)$, which enables us to write

$$\mathcal{L}(x) = \frac{d}{dx} \left[p_0(x) \frac{d}{dx} \right] + p_2(x)$$

We can show that

$$\int_a^b v^*(x) \mathcal{L}u(x) dx = [v^* p_0 u' - (v^*)' p_0 u]_a^b + \int_a^b (\mathcal{L}v)^* u dx$$

- Dirichlet Boundary : u and v both vanish at endpoints.
- Neumann Boundary : u' and v' both vanish at endpoints.
- Any second order linear operator can be turned into Sturm-Liouville form by changing the scalar product to include weight factor.

If u and v are eigenfunctions of \mathcal{L} with respective eigenvalues λ_u and λ_v , then we can see that they are orthogonal. **This is great because we have just proven orthogonality of Trigonometric, Bessel, Legendre, Hermite, Laguerre, Chebyshev. Further, completeness of eigenfunctions is the origin of fourier series!**

Namely, if $f(x)$ satisfy boundary conditions of Sturm-Liouville and is continuous and piecewise differentiable on $[a, b]$ then eigenfunction expansion of f converges uniformly to f on $[a, b]$. If it's piecewise differentiable, then it converges to $[f(x_+) + f(x_-)]/2$ on $[a, b]$.

6.2.2 Variational Method

Suppose that the spectrum of the hamiltonian H has a finite greatest lower bound. In such cases, we have a simple yet very powerful method to guess the ground state energy and wavefunction.

$$\langle \Psi | H | \Psi \rangle \geq E_0$$

This allows us to guess a parameterized wavefunction and choose the parameter which minimizes the expectation value. The result can only get better because the ground state has lowest energy among all such ψ

6.3 Partial Differential Equations

Second order (mostly linear, but even non-linear) partial differential equations are everywhere. It's all of physics.

6.3.1 Characteristic

Understanding boundary conditions is important. Consider for example a simple first order PDE

$$\mathcal{L}\phi = a \frac{\partial \phi}{\partial x} + b \frac{\partial \phi}{\partial y} = 0$$

We generally try to change variables in such a manner that the equation reduces to something which only contains derivative w.r.t one term. Here $s = ax + by, t = bx - ay$ does the job. Reducing it to

$$\frac{\partial \phi}{\partial s} = 0 \quad \phi(x, y) = f(bx - ay)$$

Observe that the solution is constant along constant t . It is known as the characteristic curve. It's not necessary that solution remains constant on characteristic curve, but it can be solved using ODE methods on the characteristic curve. If $\phi(x, y)$ is specified on a curve segment, one can deduce its value on all the characteristics that intersect it.

$$\phi(x, y) = \phi(x_0, y_0) + \frac{\partial\phi(x_0, y_0)}{\partial x}(x - x_0) + \frac{\partial\phi(x_0, y_0)}{\partial y}(y - y_0) + \dots$$

We can obtain the values of partial derivatives using the equations determining the curve and the given partial differential equation (Arfken p.405).

If the boundary curve is along a characteristic, then there may be inconsistencies. If the boundary intersects same characteristic on more than one point, then it may lead to inconsistencies as well.

Exact general analysis of boundary conditions is complicated. Look up when you need it. For reference of three simple yet important PDE's

Table 9.1 Relation between PDE and Boundary Conditions

Boundary Conditions	Class of Partial Differential Equation		
	Elliptic	Hyperbolic	Parabolic
	Laplace, Poisson in (x, y)	Wave equation in (x, t)	Diffusion equation in (x, t)
Cauchy			
Open surface	Unphysical results (instability)	Unique, stable solution	Too restrictive
Closed surface	Too restrictive	Too restrictive	Too restrictive
Dirichlet			
Open surface	Insufficient	Insufficient	Unique, stable solution in one direction
Closed surface	Unique, stable solution	Solution not unique	Too restrictive
Neumann			
Open surface	Insufficient	Insufficient	Unique, stable solution in one direction
Closed surface	Unique, stable solution	Solution not unique	Too restrictive

Chapter 7

Asymptotics and Perturbation Methods

$$\lim_{x \rightarrow x_0} \frac{f(x)}{g(x)} = 1 := f(x) \sim g(x) \quad \text{as } x \rightarrow x_0$$

$$\lim_{x \rightarrow x_0} \frac{f(x)}{g(x)} = 0 := f(x) \ll g(x) \quad \text{as } x \rightarrow x_0$$

Given an asymptotic sequence (each successive term having lower order) $\{\phi_i\}$
The coefficients of the asymptotic expansion

$$f(x) \sim a_1 \phi_1(x) + a_2 \phi_2(x) + \dots$$

are uniquely determined.

If the remainder term is of lower order than the last term of the expansion for all n , the expansion is called asymptotic.

Caution1 : Two given functions differing by transcendently small terms (exponentials as compared to powers) can have same asymptotic expansions. Hence, these terms are ignored in asymptotic expansions (in powers of x). TST are related to essential singularities.

Caution2 : Need to keep higher order terms in exponentials (including \sin, \cos, \sinh, \cosh etc). Not doing so, can give wrong prefactor.

Caution3 : Differentiating can cause trouble too. Look at "Tauberian Theorems" (BO - p.127)

7.1 For Integrals

7.1.1 Integration by Parts

In general, if we have

$$I(x) = \int_a^b f(t) e^{x\phi(t)} dt \quad \text{as } x \rightarrow \infty$$

$$I(x) = \int_a^b \frac{f(t)}{x\phi'(t)} d(e^{x\phi(t)})$$

$$I(x) = \frac{1}{x} \left[\frac{f(t)}{\phi(t)} e^{x\phi(t)} \right]_{t=a}^{t=b} - \frac{1}{x} \int_a^b \frac{d}{dt} \left[\frac{f}{\phi'} \right] e^{x\phi(t)} dt$$

We hope that the second integral has lower order than the first term and that $\phi'(t)$ is not zero anywhere in the interval. This will fail when asymptotic expansion involves *log* or fractional powers of x .

7.1.2 Laplace's Method

For "sharply-peaked" integrands, where the dominant contribution comes from the nbd of a single point, we can expand the function involved in exponent around that point.

Example1 :

$$I(x) = \int_{-10}^{10} e^{-xt^2} dt \quad \text{as } x \rightarrow \infty$$

Integration by parts won't work because $\phi'(0) = 0$. Dominant contribution comes for $t = 0$.

$$I(x) = \int_{-\infty}^{\infty} e^{-xt^2} dt - 2 \int_{10}^{\infty} e^{-xt^2} dt$$

First part is simple gaussian integral, we can estimate the second integral by replacing t^2 by $10t$. Which will show that the second term is transcendentally smaller than the first term.

$$I(x) \sim \sqrt{\frac{\pi}{x}}$$

Stirling's Approximation :

$$\Gamma(x+1) = \int_0^{\infty} t^x e^{-t} dt$$

Where's the peak of this integrand ?

$$\Gamma(x+1) = \int_0^\infty e^{x \log(t) - t} dt$$

Maximum occurs when the exponent is maximized, which gives $x = t$. Maximum moves with x . We change coordinates where it doesn't change. Define $s = \frac{t}{x}$

$$\Gamma(x+1) = x^{x+1} \int_0^\infty e^{x \log(s) - s} ds$$

Expanding log about $s = 1$.

$$\Gamma(x+1) \sim x^{x+1} \int_0^\infty e^{-x \left(1 + \frac{(1-s)^2}{2}\right)} ds$$

$$\Gamma(x+1) \sim x^{x+1} e^{-x} \int_0^\infty e^{-x \frac{(1-s)^2}{2}} ds$$

We can't do this integral. But if we add TST, we have usual gaussian integral!

$$\Gamma(x+1) \sim x^{x+1} e^{-x} \sqrt{\frac{2\pi}{x}}$$

7.1.3 Stationary Phase

For an integral of the function $f(t)$ multiplied with an oscillating function, we get cancellations. If the frequency of oscillation is sufficiently big, we expect the integral to vanish. Rigorously speaking :

$$I(x) = \int_a^b f(t) e^{ix\psi(t)} dt$$

As $x \rightarrow \infty$, we try integration by parts to get

$$I(x) = \left[\frac{f(t) e^{ix\psi(t)}}{ix\psi'(t)} \right]_{t=a}^{t=b} - \frac{1}{ix} \int_a^b \frac{d}{dt} \left[\frac{f}{\psi'} \right] e^{ix\psi(t)}$$

Again, just hoping for usual conditions on $\psi'(t)$ and the second integral being of lower order, we see that

$$I(x) \sim O\left(\frac{1}{x}\right)$$

At points when $\psi'(t) = 0$, the first order oscillation of ψ vanishes, meaning it oscillates much more slowly (in second order of t) near this point. It gives a contribution of higher order than $O(\frac{1}{x})$. Therefore, we only integrate near neighbourhood of that point of stationary phase. After expanding around that point, a standard trick would be then to extend this region to infinity, if corresponding integral is easy to do. Since this extension would contribute to higher order terms only.

Expanding around the stationary phase ($\psi'(c) = 0$), and using Fresnel's integral (really just use Gamma function formula), we can show that if

$$\text{if } f^{(i)}(c) = 0 \quad \forall i < n \quad \text{then} \quad I(x) \sim O\left(\frac{1}{x^{\frac{1}{n}}}\right)$$

A useful way to think about both the laplace and stationary phase method in a coherent manner is to see that the function multiplying $f(t)$ has an exponential order. Meaning they change much faster than $f(t)$ as long as $f(t)$ doesn't have exponential order. Note that complex representation of sin and cosine make this evident for this oscillatory functions as well. In some sense then, the integral is wholly dominated by this exponential ordered terms (including sine, cosine). Therefore we only care about the regions where this exponential ordered terms are slowly varying (given that they decay eventually in case of laplace method). In laplace method, it was the peak. In stationary phase method, it is the stationary phase. The reason we only care about regions where the exponential term is slowly varying can easily be digested by seeing how well the integral $\int_0^\infty e^{-x} dx$ is approximated by $\int_0^a e^{-x} dx$. First order change in a would give exponentially better approximation. In this case, 0 was the point of slower change.

Bessel function

$$J_0(x) = \frac{1}{\pi} \int_0^\infty \cos(x \sin(t)) dt$$

$\psi'(t) = 0$ at $t = \frac{\pi}{2}$. Expanding $\sin(t)$ about $t = \frac{\pi}{2}$

$$\sin(t) = \sin\left(\frac{\pi}{2}\right) + \cos\left(\frac{\pi}{2}\right)\left(t - \frac{\pi}{2}\right) - \frac{1}{2}\sin\left(\frac{\pi}{2}\right)\left(t - \frac{\pi}{2}\right)^2 + \dots$$

$$\begin{aligned} \text{As } x \rightarrow \infty \quad & \frac{1}{\pi} \int_0^\pi e^{ix \sin(t)} dt \\ & \sim \frac{1}{\pi} \int_{\frac{\pi}{2}-\epsilon}^{\frac{\pi}{2}+\epsilon} e^{ix \left[1 - \frac{1}{2}\left(t - \frac{\pi}{2}\right)^2 + \dots\right]} dt \end{aligned}$$

As done again and again, expanding the limits of integration because if only affects higher order terms, we get

$$J_0(x) \sim \sqrt{\frac{2}{\pi x}} \cos\left(x - \frac{\pi}{4}\right) \quad \text{as } x \rightarrow \infty$$

7.1.4 Steepest Descent

This method makes previous two methods a consequence. Consider the integral

$$I(x) = \int_C h(t) e^{x\phi(t)} e^{ix\psi(t)}$$

If we choose a contour such that the $\psi(t)$ is constant on this contour, we can bring it out of the integral and use laplace's method. This will enable us to find higher order terms that we missed in stationary phase method. Please check the video by Professor Strogatz for example.

Chapter 8

General Relativity

- An observer is a worldline where the basis vector he chooses at each point of the worldline reduces the metric to η_{ab} . In effect, he's someone who uses flat metric at all points of the worldline. (Schuller Lect 13. 41:26)
- The cause of your confusion was that expanding the metric around some point (which you can flatten using gaussian normal coordinates at any point one at a time) is not the same as taking limit of the metric near some region. Taking this limit, we may not get flat metric all the time, it was a mere coincidence in case of Near Horizon geometry of Schwarzschild Blackhole. Remember AdS2 geometry in RN blackhole?
- All we can do in reality is measure proper distance, and use it to determine the coordinates that we use.

$$1. \Gamma^t_{mn} = \frac{1}{2}g^{st}[\partial_n g_{sm} + \partial_m g_{sn} - \partial_s g_{mn}]$$