

# Quantum interfaces

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## Abstract

Rational agents acting as observers use “knowables” to construct a vision of the outside world. Thereby, they are bound by the information exchanged with what they consider to be objects. The cartesian cut or, in modern terminology, the interface mediating this exchange, is again a construction. It serves as a “scaffolding,” an intermediate construction capable of providing the necessary conceptual means.

An attempt is made to formalize the interface, in particular the quantum interface and quantum measurements, by a symbolic information exchange. A principle of conservation of information is reviewed and a measure of information flux through the interface is proposed.

We cope with the question of why observers usually experience irreversibility in measurement processes if the evolution is reversible, i.e., one-to-one. And why should there be any meaningful concept of classical information if there is merely quantum information to begin with? We take the position here that the concept of irreversible measurement is no deep principle but originates in the practical inability to reconstruct a quantum state of the object.

Many issues raised apply also to the quantum’s natural double, virtual reality.

An experiment is proposed to test the conjecture that the self is transcendent.

## Reality construction by “knowables”

Otto Rössler, in a thoughtful book [1], has pointed to the significance of object-observer interfaces, a topic which had also been investigated in other contexts (cf., among others, refs. [2, 3, 4, 5, 6, 7, 8]). By taking up this theme, the following investigation is on the epistemology of interfaces, in particular of quantum interfaces. The informal notions of “cartesian cut” and “interface” are formalized. They are then applied to observations of quantum and virtual reality systems.

A generic interface is presented here as any means of communication or information exchange between some “observer” and some observed “object.” The “observer” as well as the “object” are subsystems of some larger, all-encompassing system called “universe.”

Generic interfaces are totally symmetric. There is no principal, *a priori* reason to call one subsystem “observer” and the other subsystem “object.” The denomination is arbitrary. Consequently, “observer” and “object” may switch identities.

Take, for example, an impenetrable curtain separating two parts of the same room. Two parties — call them Alice and John — are merely allowed to communicate by sliding through papers below the curtain. Alice, receiving the memos emanating from John’s side of the curtain, thereby effectively constructs a “picture” or representation of John and *vice versa*.

The cartesian cut spoils this total symmetry and arbitrariness. It defines a distinction between “observer” and “object” beyond doubt. In our example, one agent — say Alice — becomes the observer while the other agent becomes the observed object. That, however, may be a very arbitrary convention which not necessarily reflects the configuration properly.

A cartesian cut may presuppose a certain sense of “rationality,” or even “consciousness” on the “observer’s” side. We shall assume that some observer or agent exists which, endowed with rational intelligence, draws conclusions on the basis of certain premises, in particular the agent’s *state of knowledge*, or “knowables” to (re)construct “reality.” Thereby, we may imagine the agent as some kind of robot, some mechanistic or algorithmic entity. (From now on, “observer” and “agent” will be used as synonyms.) Note that the agent’s state of knowledge may not necessarily coincide with a complete description of the observed system, nor may the agent be in the possession of a complete description of its own side of the cut. Indeed, it is not unreasonable to

speculate that certain things, although knowable “from the outside” of the observer-object system, are principally unknowable to an intrinsic observer [7].

Although we shall come back to this issue later, the notion of “consciousness” will not be reviewed here. We shall neither speculate exactly what “consciousness” is, nor what may be the necessary and sufficient conditions for an agent to be ascribed “consciousness”. Let it suffice to refer to two proposed tests of consciousness by Turing and Greenberger [9].

With regards to the type of symbols exchanged, we shall differentiate between two classes: classical symbols, and quantum symbols. The cartesian cuts mediating classical and quantum symbols will be called “classical” or “quantum” (cartesian) cuts, respectively.

## Formalization of the cartesian cut

The task of formalizing the heuristic notions of “interface” and “cartesian cut” is, at least to some extent, analogous to the formalization of the informal notion of “computation” and “algorithm” by recursive function theory via the Church-Turing thesis.

In what follows, the informal notions of interface and cartesian cut will be formalized by symbolic exchange; i.e., by the mutual communication of symbols of a formal alphabet. In this model, an object and an observer alphabet will be associated with the observed object and with the observer, respectively.

Let there be an *object alphabet*  $\mathcal{S}$  with symbols  $s \in \mathcal{S}$  associated with the outcomes or “message” of an experiment possible results. Let there be an *observer alphabet*  $\mathcal{T}$  with symbols  $t \in \mathcal{T}$  associated with the possible inputs or “questions” an observer can ask.

At this point we would like to keep the observer and object alphabets as general as possible, allowing also for quantum bits to be transferred. Such quantum bits, however, have no direct operational meaning, since they cannot be completely specified. Only classical bits have a (at least in principle) unambiguous meaning, since they can be completely specified, copied and measured. We shall define an interface next.

- An interface  $I$  is an entity forming the common boundary between two

parts of a system, as well as a means of information exchange between those parts.

- By convention, one part of the of the system is called “observer” and the other part “object.”
- Information between the observer and the object via the interface is exchanged by symbols. The corresponding functional representation of the interface is a map  $I : \mathcal{T} \mapsto \mathcal{S}$ , where  $\mathcal{T}$  and  $\mathcal{S}$  are the observer and the object alphabets, respectively. Any such information exchange is called “*measurement*.”
- The interface ist *total* in the sense that the observer receives *all* symbols emanating from the object. (However, the object needs not receive all symbols emanating from the observer.)
- Types of interface include purely classical, quasi-classical, and purely quantum interfaces.
  - Classical scenario I: A classical interface is an interface defined in a classical system, for which the symbols in  $\mathcal{S}$  and  $\mathcal{T}$  are classical states encodable by classical bits “0” and “1” corresponding to “**true**” and “**false**,” respectively. This kind of binary code alphabet corresponds to yes-no outcomes to dichotomic questions; experimental physics in-a-nutshell. An example for a dichotomic outcome associated with is “there is a click in a counter” or “there is no click in a counter,” respectively.
  - Quasi-classical scenario II: a quasi-classical interface is an interface defined in a quantum system, whereby the symbols in  $\mathcal{S}$  and  $\mathcal{T}$  are classical states encoded by classical bits. This is the picture most commonly used for measurements in quantum mechanics.
  - Quantum scenario III: A quantum interface is an interface defined in a quantized system. In general, the quantum symbols in  $\mathcal{S}$  and  $\mathcal{T}$  are quantum states.

Informally, in a measurement, the object “feels” the observer’s question (in  $\mathcal{T}$ ) and responds with an answer (in  $\mathcal{S}$ ) which is felt by the observer (cf. Fig. 1).

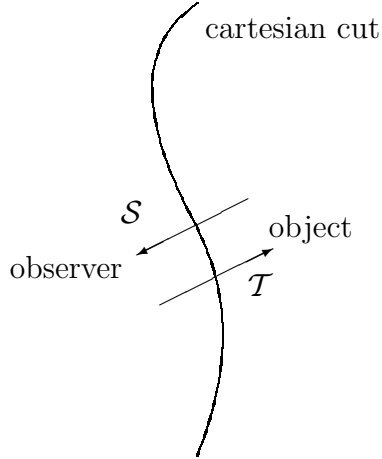


Figure 1: An interface as a cartesian cut between observer and object. The information flow across the interface is formalized by symbols.

The reader is encouraged to view the interface not as a static entity but as a *dynamic* one, through which information is constantly piped back and forth the observer and the object and the resulting time flow may also be viewed as the dynamic evolution of the system as a whole. In what follows it is important to stress that we shall restrict our attention to cases for which the interface is total; i.e., the observer receives *all* symbols emanating from the object.

## One-to-one quantum state evolution and “haunted” measurements

On a microphysical scale, we do not wish to restrict quantum object symbols to classical states. The concept pursued here is rather that of the quantum scenario III: a uniform quantum system with unitary, and thus reversible, one-to-one evolution. Any process within the entire system evolves according to a reversible law represented by a unitary time evolution  $U^{-1} = U^\dagger$ . As a result, the interface map  $I$  is one-to-one; i.e., it is a bijection.

Stated pointedly, we take it for granted that the wave function of the en-

tire system—including the observer and the observed object separated by the cartesian cut or interface—evolves one-to-one. Thus, in principle, previous states can be reconstructed by proper reversible manipulations.

In this scenario, what is called “measurement” is merely an exchange of quantum information. In particular, the observer can “undo” a measurement by proper input of quantum information via the quantum interface. In such a case, no information, no knowledge about the object’s state can remain on the observer’s side of the cut; all information has to be “recycled” completely in order to be able to restore the wave function of the object entirely in its previous form.

Experiments of the above form have been suggested [10] and performed under the name “haunted measurement” and “quantum eraser” [11]. These matters are very similar to the opening, closing and reopening of Schrödinger’s catalogue of expectation values [12, p. 823]: At least up to a certain magnitude of complexity, any measurement can be “undone” by a proper reconstruction of the wave-function. A necessary condition for this to happen is that *all* information about the original measurement is lost. In Schrödinger’s terms, the prediction catalog (the wave function) can be opened only at one particular page. We may close the prediction catalog before reading this page. Then we can open the prediction catalog at another, complementary, page again. By no way we can open the prediction catalog at one page, read and (irreversible) memorize the page, close it; then open it at another, complementary, page. (Two non-complementary pages which correspond to two co-measurable observables can be read simultaneously.)

## Where exactly is the interface located?

The interface has been introduced here as a scaffolding, an auxiliary construction to model the information exchange between the observer and the observed object. One could quite justifiably ask (and this question *has* indeed been asked by Professor Bryce deWitt), “where exactly *is* the interface in a concrete experiment, such as a spin state measurement in a Stern-Gerlach apparatus?”

We take the position here that the location of the interface very much depends on the physical proposition which is tested and on the conventions

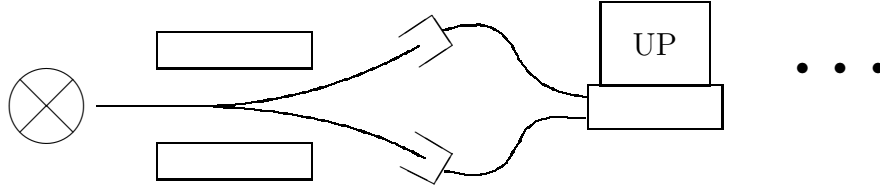


Figure 2: Where exactly is the interface located?

assumed. Let us take, for example, a statement like

“the electron spin in the  $z$ -direction is up.”

In the case of a Stern-Gerlach device, one could locate the interface at the apparatus itself. Then, the information passing through the interface is identified with the way the particle took.

One could also locate the interface at two detectors at the end of the beam paths. In this case, the information penetrating through the interface corresponds to which one of the two detectors (assumed lossless) clicks (cf. Fig. 2).

One could also situate the interface at the computer interface card registering this click, or at an experimenter who presumably monitors the event (cf. Wigner’s friend [13]), or at the persons of the research group to whom the experimenter reports, to their scientific peers, and so on.

Since there is no material or real substrate which could be uniquely identified with the interface, in principle it could be associated with or located at anything which is affected by the state of the object. The only difference is the reconstructibility of the object’s previous state (cf. below): the “more macroscopic” (i.e., many-to-one) the interface becomes, the more difficult it becomes to reconstruct the original state of the object.

## From one-to-one to many-to-one

If the quantum evolution is reversible, how come that observers usually experience irreversibility in measurement processes? We take the position here that *the concept of irreversible measurement is no deep principle but merely*

*originates in the practical inability to reconstruct a quantum state of the object.*

Restriction to classical state or information exchange across the quantum interface—the quasi-classical scenario II—effectively implements the standard quantum description of the measurement process by a classical measurement apparatus: there exists a clear distinction between the “internal quantum box,” the quantum object—with unitary, reversible, one-to-one internal evolution—and the classical symbols emanating from it. Such a reduction from the quantum to the classical world is accompanied by a loss of internal information “carried with the quantum state.” This effectively induces a many-to-one transition associated with the measurement process, often referred to as “wave function collapse.” In such a case, one and the same object symbol could have resulted from many different quantum states, thereby giving raise to irreversibility and entropy increase.

But also in the case of a uniform one-to-one evolution (scenario I), just as in classical statistical physics, reconstruction greatly depends on the possibility to “keep track” of all the information flow directed at and emanating from the object. If this flow is great and spreads quickly with respect to the capabilities of the experimenter, and if the reverse flow of information from the observer to the object through the interface cannot be suitably controlled [14, 15] then the chances for reconstruction are low.

This is particularly true if the interface is not total: in such a case, information flows off the object to regions which are (maybe permanently) outside of the observer’s control.

The possibility to reconstruct a particular state may widely vary with technological capabilities which often boil down to financial commitments. Thus, irreversibility of quantum measurements by interfaces appears as a gradual concept, depending on conventions and practical necessities, and not as a principal property of the quantum.

In terms of coding theory, the quantum object code is sent to the interface but is not properly interpreted by the observer. Indeed, the observer might only be able to understand a “higher,” macroscopic level of physical description, which subsumes several distinct microstates under one macro-symbol (cf. below). As a result, such macro-symbols are no unique encoding of the object symbols. Thus effectively the interface map  $I$  becomes many-to-one.

This also elucidates the question why there should be any meaningful concept of classical information if there is merely quantum information to begin



with: in such a scenario, classical information appears as an effective entity on higher, intermediate levels of description. Yet, the most fundamental level is quantum information.

## Do conscious observers “unthink”?

Because of the one-to-one evolution, a necessary condition for reconstruction of the object wave function is the complete restoration of the observer wave function as well. That is, the observer’s state is restored to its previous form, and no knowledge, no trace whatsoever can be left behind. An observer would not even know that a “measurement” has taken place. This is hard to accept, in particular if one assumes that observers have consciousness which are detached entities from and not mere functions of the quantum brain. Thus, in the latter case, one might be convinced that conscious observers “unthink” the measurement results in the process of complete restoration of the wave function. In the latter case, consciousness might “carry away” the measurement result via a process distinct from the quantum brain. (Cf. Wigner’s friend [13].)

But even in this second, dualistic, scenario, the conscious observer, after reconstruction of the wave function, would have no direct proof of the “previously measured fact,” although subsequent measurements might confirm his allegations. This amounts to a proposal of an experiment involving a conscious observer (*not* merely a rational agent) and a quantized object. The experiment tests the metaphysical claim that consciousness exists beyond matter [16]. As sketched above, the experiment involves four steps.

- Step I: The conscious observer measures some quantum observable on the quantized object which occurs irreducibly random according to the axioms of quantum theory. As a consequence, the observer “is aware of” the measurement result and ascribes to it an “element of physical reality” [17].
- Step II: The original quantum state of the quantized object is reconstructed. Thereby, all physical information about the measurement result is lost. This is also true for the brain of the conscious observer. Let us assume that the observer “is still aware of” the measurement result. In this case, the observer ascribes to it an “element of metaphysical reality.”

- Step III: The observer guesses or predicts the outcome of the measurement despite the fact that no empirical evidence about the outcome of the previous measurement exists.
- Step IV: The measurement is “re-done” and the actual measurement result is compared with the conscious observer’s prediction in step III. If the prediction and the actual outcome do not coincide, the hypothesis of a consciousness beyond matter is falsified.

As an analogy, one might think of a player in a virtual reality environment. Although at the observation level of the virtual reality, the measurement is undone, the player himself “knows” what has been there before. This knowledge, however, has been passed on to another interface which is not immanent with respect to the virtual reality. That is, it cannot be defined by intrinsic (endo-) means. Therefore, it can be called a *transcendent interface* with respect to the virtual reality. However, if we start with the real universe of the player, then the same interface becomes intrinsically definable. The hierarchical structure of meta-worlds has been the subject of conceptual and visual art [18, 19, 20] and literature [21].

## Parallels in statistical physics: from reversibility to irreversibility

The issue of “emergence” of irreversibility from reversible laws is an old one and subject of scientific debate at least since Boltzmann’s time [22]. We shall shortly review an explanation in terms of the emergence of many-to-one (irreversible) evolution relative to a “higher” macroscopic level of description from one-to-one (reversible) evolution at a more fundamental microscopic “complete” level of description. These considerations are based on the work of Jaynes [23, 24], Katz [25] and Hobson [26], among others. See Buček *et al.* [27] for a detailed review with applications.

In this framework, the many-to-one and thus irreversible evolution is a simple consequence of the fact that many different microstates, i.e., states on the fundamental “complete” level of physical description, are mapped onto a *single* macroscopic state (cf. Fig. 3). Thereby, knowledge about the microphysical state is lost; making impossible the later reconstruction of the microphysical state from the macroscopic one. (In the example drawn in

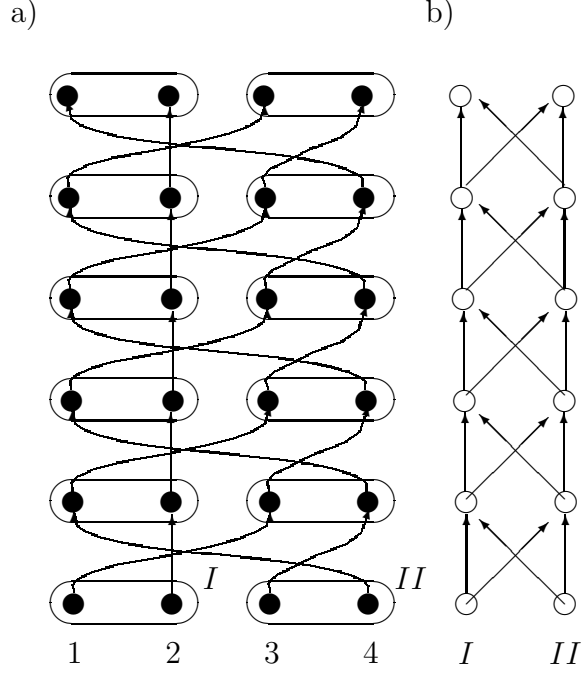


Figure 3: Full circles represents the “complete” level of description, open circles or the corresponding ovals represent intermediate or macroscopic levels of description. The microphysical states 1, 2 are mapped onto  $I$  and 3, 4 are mapped onto  $II$  by a many-to-one mapping. a) The temporal evolution in terms of the microstates is one-to-one; b) The evolution with respect to the macro-states is irreversible.

Fig. 3, observation of the “macrostate”  $II$  could mean that the system is either in microstate 1 or 2.) on some intermediate, “higher” level of physical description, whereas it remains reversible on the complete description level.

Here, just as in the quantum interface case, irreversibility in statistical physics is a gradual concept, very much depending on the observation level, which depends on conventions and practical necessities. Yet again, in principle the underlying complete level of description is one-to-one. As a consequence, this would for example make possible the reconstruction of the Library of Alexandria if one takes into account all smoky emanations thereof. The task of “reversing the gear,” of reconstructing the past and constructing

a different future, is thus not entirely absurd. Yet fortunately or unfortunately, for all practical purposes it remains impossible.

## Principle of information conservation

In another scenario (closely related to scenario I), classical information is a primary entity. The quantum is obtained as an effective theory to represent the state of knowledge, the “knowables,” of the observer about the object [28, 29, 30]. Thereby, quantum information appears as a derived theoretical entity, very much in the spirit of Schrödinger’s perception of the wave function as a catalogue of expectation values (cf. above).

The following circular definitions are assumed.

- An elementary object carries one bit of (classical) information [28].
- $n$  elementary objects carry  $n$  bits of (classical) information [28]. The information content present in the physical system *is exhausted* by the  $n$  bits given; nothing more can be gained by any perceivable procedure.
- Throughout temporal evolution, the amount of (classical) information measured in bits is conserved.

One immediate consequence seems a certain kind of irreducible randomness associated with requesting from an elementary object information which has not been previously encoded therein. We may, for instance, think of an elementary object as an electron which has been prepared in spin state “up” in some direction. If the electron’s spin state is measured in another direction, this must give rise to randomness since the particle “is not supposed to know” about this property. Yet, we may argue that in such a case the particle might respond with no answer at all, and not with the type of irreducible randomness which, as we know from the computer sciences [31, 32], is such a preciously expensive quality.

One way to avoid this problem is to assume that the apparent randomness does not originate from the object but is a property of the interface: the object always responds to the question it has been prepared for to answer; but the interface “translates” the observer’s question into the appropriate form suitable for the object. In this process, indeterminism comes in.

As a result of the assumption of the temporal conservation of information, the evolution of the system has to be one-to-one and, for finite systems, a permutation.

Another consequence of the conservation of information is the possibility to define continuity equations. In analogy to magnetostatics or thermodynamics we may represent the information flow by a vector which gives the amount of information passing per unit area and per unit time through a surface element at right angles to the flow. We call this the *information flow density*  $\mathbf{j}$ . The amount of information flowing across a small area  $\Delta A$  in a unit time is

$$\mathbf{j} \cdot \mathbf{n} \Delta A,$$

where  $\mathbf{n}$  is the unit vector normal to  $\Delta A$ . The information flow density is related to the average flow velocity  $v$  of information. In particular, the information flow density associated with an elementary object of velocity  $v$  per unit time is given by  $\mathbf{j} = \rho v$  bits per second, where  $\rho$  stands for the information density (measured in bits/ $m^3$ ). For  $N$  elementary objects per unit volume carrying one bit each,

$$\mathbf{j} = Nvi.$$

Here,  $i$  denotes the elementary quantity of information measured in bit units. The information flow  $I$  is the total amount of information passing per unit time through any surface  $A$ ; i.e.,

$$I = \int_A \mathbf{j} \cdot \mathbf{n} dA.$$

We have assumed that the cut is on a closed surface  $\mathcal{A}_c$  surrounding the object. The conservation law of information requires the following continuity equation to be valid:

$$\int_{\mathcal{A}_c} \mathbf{j} \cdot \mathbf{n} dA = -\frac{d}{dt}(\text{Information inside})$$

or, by defining an information density  $\rho$  and applying Gauss' law,

$$\nabla \cdot \mathbf{j} = -\frac{d\rho}{dt}.$$

To give a quantitative account of the present ability to reconstruct the quantum wave function of single photons, we analyze the “quantum eraser”

paper by Herzog, Kwiat, Weinfurter and Zeilinger [11]. The authors report an extension of their apparatus of  $x = 0.13$  m, which amounts to an information passing through a sphere of radius  $x$  of

$$I_{\text{qe}} = 4\pi x^2 ci = 6 \times 10^7 \text{bits/second}.$$

Here,  $\mathbf{j} = ci$  ( $c$  stands for the velocity of light in vacuum) has been assumed. At this rate the reconstruction of the photon wave function has been conceivable.

We propose to consider  $I$  as a measure for wave function reconstruction. In general,  $I$  will be astronomically high because of the astronomical numbers of elementary objects involved. Yet, the associated diffusion velocity  $v$  may be considerably lower than  $c$ .

Let us finally come back to the question, *“why should there be any meaningful concept of classical information if there is merely quantum information to begin with?”* A tentative answer in the spirit of this approach would be that *“quantum information is merely a concept derived from the necessity to formalize modes of thinking about the state of knowledge of a classical observer about a classical object. Although the interface is purely classical, it appears to the observer as if it were purely quantum or quasi-classical.”*

## Virtual reality as a quantum double

Just as quantum systems, virtual reality universes can have a one-to-one evolution. We shall shortly review reversible automata [33, 34] which are characterized by the following properties:

- a finite set  $S$  of states,
- a finite input alphabet  $I$ ,
- a finite output alphabet  $O$ ,
- temporal evolution function  $\delta : S \times I \rightarrow S$ ,
- output function  $\lambda : S \times I \rightarrow O$ .

$S \backslash I$	$\delta$			$\lambda$		
	1	2	3	1	2	3
$s_1$	$s_1$	$s_1$	$s_2$	1	2	2
$s_2$	$s_2$	$s_2$	$s_1$	1	3	3

Table 1: Transition and output table of a reversible automaton with two states  $S = \{s_1, s_2\}$  and three input/output symbols  $I = \{1, 2, 3\}$ . Neither its transition nor its output function is one-to-one.

The combined transition and output function  $U$  is reversible and thus corresponds to a permutation:

$$U : (s, i) \rightarrow (\delta(s, i), \lambda(s, i)), \quad (1)$$

with  $s \in S$  and  $i \in I$ . Note that neither  $\delta$  nor  $\lambda$  needs to be a bijection.

As an example, take the perturbation matrix

$$U = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 & 0 & 0 \end{pmatrix}.$$

It can be realized by a reversible automaton which is represented in Table 1. Neither its evolution function nor its transition function is one-to-one, since for example  $\delta(s_1, 3) = \delta(s_2, 1) = s_2$  and  $\lambda(s_1, 2) = \lambda(s_1, 3) = 2$ . Its flow diagram throughout five evolution steps is depicted in Figure 3, where the microstates 1, 2, 3, 4 are identified by  $(s_1, 1)$ ,  $(s_1, 2)$ ,  $(s_2, 1)$  and  $(s_2, 2)$ , respectively.

## Metaphysical speculations

Although the contemporaries always attempt to canonize their relative status of knowledge about the physical world, from a broader historical perspective

this appears sentimental at best and ridiculous at worst. The type of natural sciences which emerged from the Enlightenment is in a permanent scientific revolution. As a result, scientific wisdom is always transitory. Science is and needs to be in constant change.

So, what about the quantum? Quantum mechanics challenges the conventional rational understanding in the following ways:

- by allowing for randomness of single events, which collectively obey quantum statistical predictions;
- by the feature of complementarity; i.e., the mutual exclusiveness of the measurement of certain observables termed complementary. Complementarity results in a non-classical, non-distributive and thus non-boolean event structures;
- by non-standard probabilities which are based on non-classical, non-boolean event structures. These quantum probabilities cannot be properly composed from its proper parts, giving rise to the so-called “contextuality.”

I believe that, just as so many other formalisms before, also quantum theory will eventually give way to a more comprehensive understanding of fundamental physics, although at the moment it appears almost heretic to pretend that there is something “beyond the quantum”. Exactly how this progressive theory beyond the quantum will look like, nobody presently can say [35]. (Otherwise, it would not be beyond anymore, but there would be another theory lurking beyond the beyond.) In view of the quantum challenges outlined before, it may be well worth speculating that the revolution will drastically change our perception of the world.

It may well be that epistemic issues such as the ones reviewed here will play an important role therein. I believe that the careful analysis of conventions which are taken for granted and are never mentioned in standard presentations of the quantum and relativity theory [36] will clarify some misconceptions.

Are quantum-like and relativity-like theories consequences of the modes we use to think about and construct our world? Do they not tell us more about our projections than about an elusive reality?



Of course, physical constants such as Planck’s constant or the velocity of light *are* physical input. But the *structural form* of the theories might be conventional.

Let me also state that one-to-one evolution is a sort of “Borgesian” nightmare, a hermetic prison: the time evolution is a constant permutation of one and the same “message” which always remains the same but expresses itself through different forms. Information is neither created nor discarded but remains constant at all times. The implicit time symmetry spoils the very notion of “progress” or “achievement,” since what is a valuable output is purely determined by the subjective meaning the observer associates with it and is devoid of any syntactic relevance. In such a scenario, any gain in knowledge remains a merely subjective impression of ignorant observers.

## References

- [1] Otto E. Rössler. *Endophysics. The World as an Interface*. World Scientific, Singapore, 1998. With a foreword by Peter Weibel.
- [2] R. J. Boskovich. *De spacio et tempore, ut a nobis cognoscuntur*. Vienna, 1755. English translation in [37].
- [3] T. Toffoli. The role of the observer in uniform systems. In G. Klir, editor, *Applied General Systems Research*. Plenum Press, New York, London, 1978.
- [4] Karl Svozil. Connections between deviations from lorentz transformation and relativistic energy-momentum relation. *Europhysics Letters*, 2:83–85, 1986. excerpts from [38].
- [5] Karl Svozil. Operational perception of space-time coordinates in a quantum medium. *Il Nuovo Cimento*, 96B:127–139, 1986.
- [6] Otto E. Rössler. Endophysics. In John L. Casti and A. Karlquist, editors, *Real Brains, Artificial Minds*, page 25. North-Holland, New York, 1987.
- [7] Karl Svozil. *Randomness & Undecidability in Physics*. World Scientific, Singapore, 1993.

- [8] Harald Atmanspacher and G. Dalenoort, editors. *Inside Versus Outside*, Berlin, 1994. Springer.
- [9] Daniel M. Greenberger. Private communication in the Viennese coffee-house “Landmann.” Recollection from the author’s memory: Indication of consciousness is associated with the breaking of a super-rule, such described in the Genesis, 1986.
- [10] Daniel B. Greenberger and A. YaSin. “Haunted” measurements in quantum theory. *Foundation of Physics*, 19(6):679–704, 1989.
- [11] Thomas J. Herzog, Paul G. Kwiat, Harald Weinfurter, and Anton Zeilinger. Complementarity and the quantum eraser. *Physical Review Letters*, 75(17):3034–3037, 1995.
- [12] Erwin Schrödinger. Die gegenwärtige Situation in der Quantenmechanik. *Naturwissenschaften*, 23:807–812, 823–828, 844–849, 1935. English translation in [39] and [40, pp. 152-167].
- [13] Eugene P. Wigner. Remarks on the mind-body question. In I. J. Good, editor, *The Scientist Speculates*, pages 284–302. Heinemann and Basic Books, London and New York, 1961. Reprinted in [40, pp. 168-181].
- [14] W. H. Zurek. Decoherence and the transition from the quantum to classical. *Physics Today*, pages 36–44, October 1991.
- [15] W. H. Zurek. Decoherence, einselection, and the existential interpretation (the rough guide). *Phil.Trans.Roy.Soc.Lond.*, A356:1793–1820, 1998. e-print at <http://xxx.lanl.gov/abs/quant-ph/9805065>.
- [16] John C. Eccles. The mind-brain problem revisited: The microsite hypothesis. In John C. Eccles and O. Creutzfeldt, editors, *The Principles of Design and Operation of the Brain*, page 549. Springer, Berlin, 1990.
- [17] Albert Einstein, Boris Podolsky, and Nathan Rosen. Can quantum-mechanical description of physical reality be considered complete? *Physical Review*, 47:777–780, 1935. Reprinted in [40, pages. 138-141].
- [18] Peter Weibel. various video-installations.

- [19] Total Recall. Movie, Director Paul Verhoeven, USA. 113 minutes. English, budget of US\$65 million, 1990.
- [20] The Matrix. Movie, Directors Andy Wachowski and Larry Wachowski, USA. English, budget of US\$80 million, 1999.
- [21] Daniel F. Galouye. *Simulacron 3*. 1964.
- [22] Jean Bricmont. Science of chaos or chaos in science? *Annals of the New York Academy of Sciences*, 775:131–176, 1996. also reprinted in [41] and as e-print <http://xxx.lanl.gov/abs/chao-dyn/9603009>.
- [23] Edwin Thompson Jaynes. Information theory and statistical mechanics. In K. W. Ford, editor, *Statistical Physics 3 (Brandeis Summer Institute 1962)*. Benjamin, New York, 1962.
- [24] Edwin Thompson Jaynes. Probability theory: The logic of science. URL <http://bayes.wustl.edu/etj/prob.html>, 1985.
- [25] A. Katz. *Principles of Statistical Mechanics. The Information Theory Approach*. W. H. Freeman and Company, San Francisco, 1967.
- [26] A. Hobson. *Concepts in Statistical Mechanics*. Gordon and Breach, New York, 1971.
- [27] V. Bužek, G. Drobny, R. Derka, G. Adam, and H. Wiedemann. Quantum state reconstruction from incomplete data. *Chaos, Solitons & Fractals*, 10(6):981–1074, 1999.
- [28] Anton Zeilinger. A foundational principle for quantum mechanics. *Foundations of Physics*, 29(4):631–643, 1999.
- [29] Časlav Brukner and Anton Zeilinger. Malus’ law and quantum information. *Acta Physica Slovaca*, 49(4):647–652, 1999.
- [30] Časlav Brukner and Anton Zeilinger. Operationally invariant information in quantum mechanics. *Physical Review Letters*, 83(17):3354–3357, 1999.
- [31] Gregory J. Chaitin. *The Unknowable*. Springer-Verlag, Singapore, 1999.

- [32] Cristian Calude. Algorithmic information theory. lecture notes. 1992.
- [33] Karl Svozil. The Church-Turing thesis as a guiding principle for physics. In Cristian S. Calude, John Casti, and Michael J. Ninneen, editors, *Unconventional Models of Computation*, pages 371–385, Singapore, 1998. Springer.
- [34] Karl Svozil. One-to-one. *Complexity*, 42(1):25–29, 1998.
- [35] Imre Lakatos. *Philosophical Papers. 1. The Methodology of Scientific Research Programmes*. Cambridge University Press, Cambridge, 1978.
- [36] Karl Svozil. Relativizing relativity.  
e-print <http://xxx.lanl.gov/abs/physics/9809025>, 1999.
- [37] R. J. Boskovich. De spacio et tempore, ut a nobis cognoscuntur. In J. M. Child, editor, *A Theory of Natural Philosophy*, pages 203–205. Open Court (1922) and MIT Press, Cambridge, MA, 1966.
- [38] Karl Svozil. On the setting of scales for space and time in arbitrary quantized media. *Lawrence Berkeley Laboratory preprint*, LBL-16097, May 1983.
- [39] J. D. Trimmer. The present situation in quantum mechanics: a translation of Schrödinger’s “cat paradox”. *Proc. Am. Phil. Soc.*, 124:323–338, 1980. Reprinted in [40, pp. 152-167].
- [40] John Archibald Wheeler and Wojciech Hubert Zurek. *Quantum Theory and Measurement*. Princeton University Press, Princeton, 1983.
- [41] Jean Bricmont. Science of chaos or chaos in science? In Paur R. Gross, Norman Levitt, and Martin W. Lewis, editors, *Flight from Science and Reason*. John Hopkins University Press, 1997.